

# Search for a Nonzero Strange Form Factor of the Proton at $2.5 \text{ (GeV/c)}^2$

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Hall C Meeting 1/13/23

$\Psi$  Indiana University



Jefferson Lab

The logo for Jefferson Lab, featuring the text "Jefferson Lab" in black, with a red swoosh underline that starts under "Jefferson" and ends under "Lab".

# Search for a Nonzero Strange Form Factor of the Proton at $2.5 \text{ (GeV/c)}^2$

## Collaborators

- Bogdan Wojtsekhowski (*spokesperson-contact*)

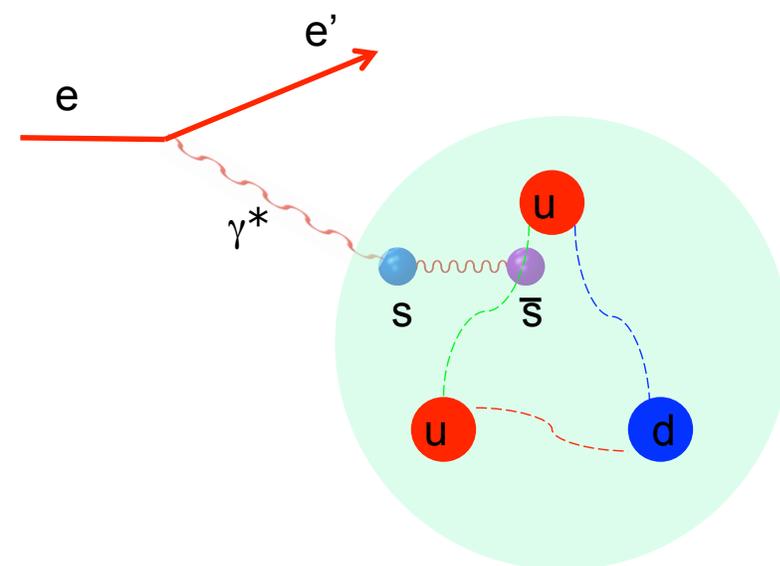


- Rakitha Beminiwattha



LOUISIANA TECH  
UNIVERSITY

- Kent Paschke



# Charge symmetry and the nucleon form factors

## Charge Symmetry

$$G_E^p = \frac{2}{3} G_E^{u,p} - \frac{1}{3} G_E^{d,p}$$

$$G_E^n = \frac{2}{3} G_E^{u,n} - \frac{1}{3} G_E^{d,n}$$

Charge symmetry is assumed for the form factors,  $G_E^{u,p} = G_E^{d,n}$ , etc. and used to find the flavor separated form-factors, measuring  $G_{E,M}^{p,n}$  to find  $G_{E,M}^{u,d}$

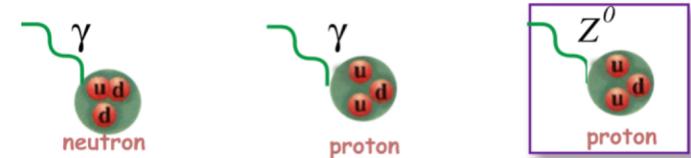
$$G_E^p = \frac{2}{3} G_E^{u,p} - \frac{1}{3} G_E^{d,p} - \frac{1}{3} G_E^s$$

$$G_E^n = \frac{2}{3} G_E^{u,n} - \frac{1}{3} G_E^{d,n} - \frac{1}{3} G_E^s$$

But this can be broken! One way is to have a non-zero strange form-factor, which breaks the "2 equations and 2 unknowns" system

The weak form factor provides a third linear combination:

$$G_E^{p,Z} = \left(1 - \frac{8}{3} \sin^2 \theta_W\right) G_E^{u,p} + \left(-1 + \frac{4}{3} \sin^2 \theta_W\right) G_E^{d,p} + \left(-1 + \frac{4}{3} \sin^2 \theta_W\right) G_E^s$$



A strange quark form factor would be indistinguishable from a broken charge symmetry in u,d flavors

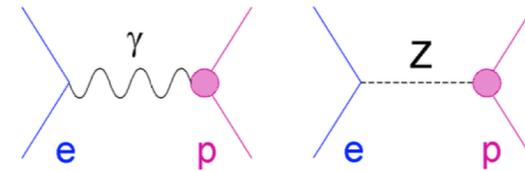
$$\delta G_E^u \equiv G_E^{u,p} - G_E^{d,n}$$

$$\delta G_E^d \equiv G_E^{d,p} - G_E^{u,n}$$

*So, more generally: this experiment tests the assumption of charge symmetry which is crucial to the flavor decomposition of the form factors*

# Strangeness form factors

## Polarized electron beam elastic e-p scattering



$$A_{PV} = -\frac{G_F Q^2}{4\pi\alpha\sqrt{2}} \cdot \left[ (1 - 4\sin^2\theta_W) - \frac{\epsilon G_E^p G_E^n + \tau G_M^p G_M^n}{\epsilon(G_E^p)^2 + \tau(G_M^p)^2} - \frac{\epsilon G_E^p G_E^s + \tau G_M^p G_M^s}{\epsilon(G_E^p)^2 + \tau(G_M^p)^2} \right. \\ \left. + \epsilon'(1 - 4\sin^2\theta_W) \frac{G_M^p G_A^{Zp}}{\epsilon(G_E^p)^2 + \tau(G_M^p)^2} \right]$$

$$A_{PV} = 150 \text{ ppm at } \theta = 15.5^\circ, Q^2 = 2.5 \text{ GeV}^2 \text{ (for sFF} = 0)$$

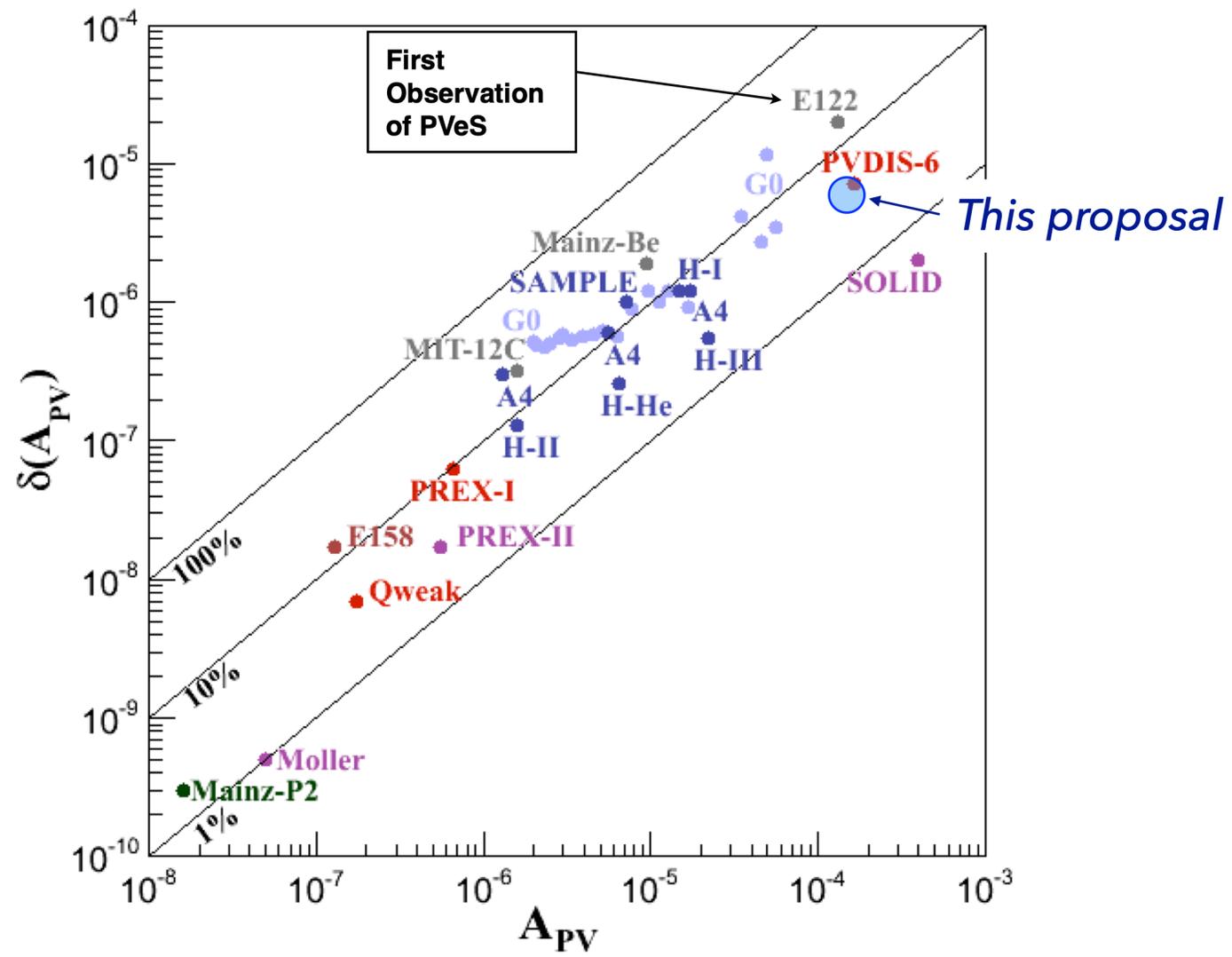
$$A_{PV} = (-226 \text{ ppm}) * [ 0.075 + 0.542 - 6.43 * ( G_M^s + 0.32 G_E^s ) + 0.038 ]$$

$Q_w$

EMFF

axial

strange form-factors



# Breakdown of U/D scaling at larger $Q^2$

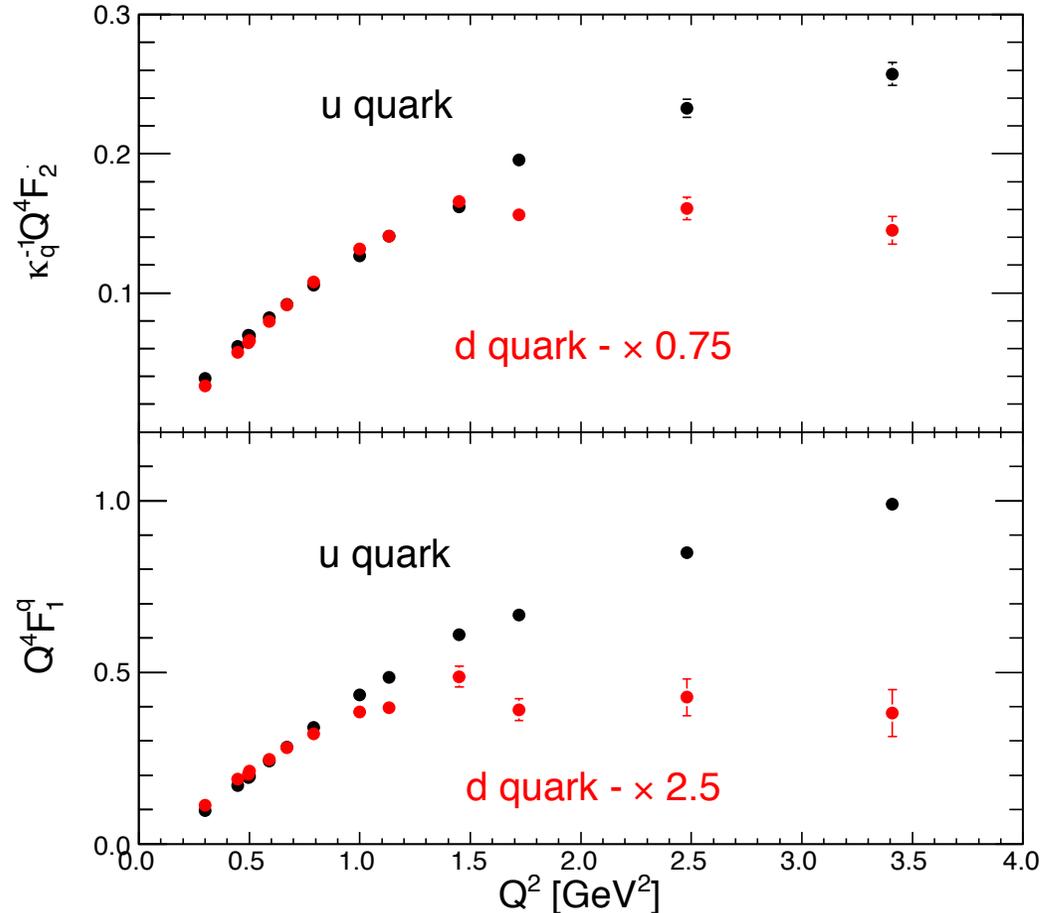
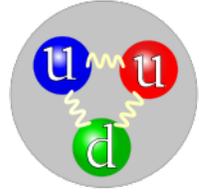


FIG. 3: The  $Q^2$ -dependence for the  $u$ - and  $d$ -contributions to the proton form factors (multiplied by  $Q^4$ ). The data points are explained in the text.

## Flavor decomposition of the elastic nucleon electromagnetic form factors

G.D. Cates,<sup>1</sup> C.W. de Jager,<sup>2</sup> S. Riordan,<sup>3</sup> and B. Wojtsekhowski<sup>2,\*</sup>

<sup>1</sup>University of Virginia, Charlottesville, VA 22903

<sup>2</sup>Thomas Jefferson National Accelerator Facility, Newport News, VA 23606

<sup>3</sup>University of Massachusetts, Amherst, MA 01003

(Dated: March 6, 2011)

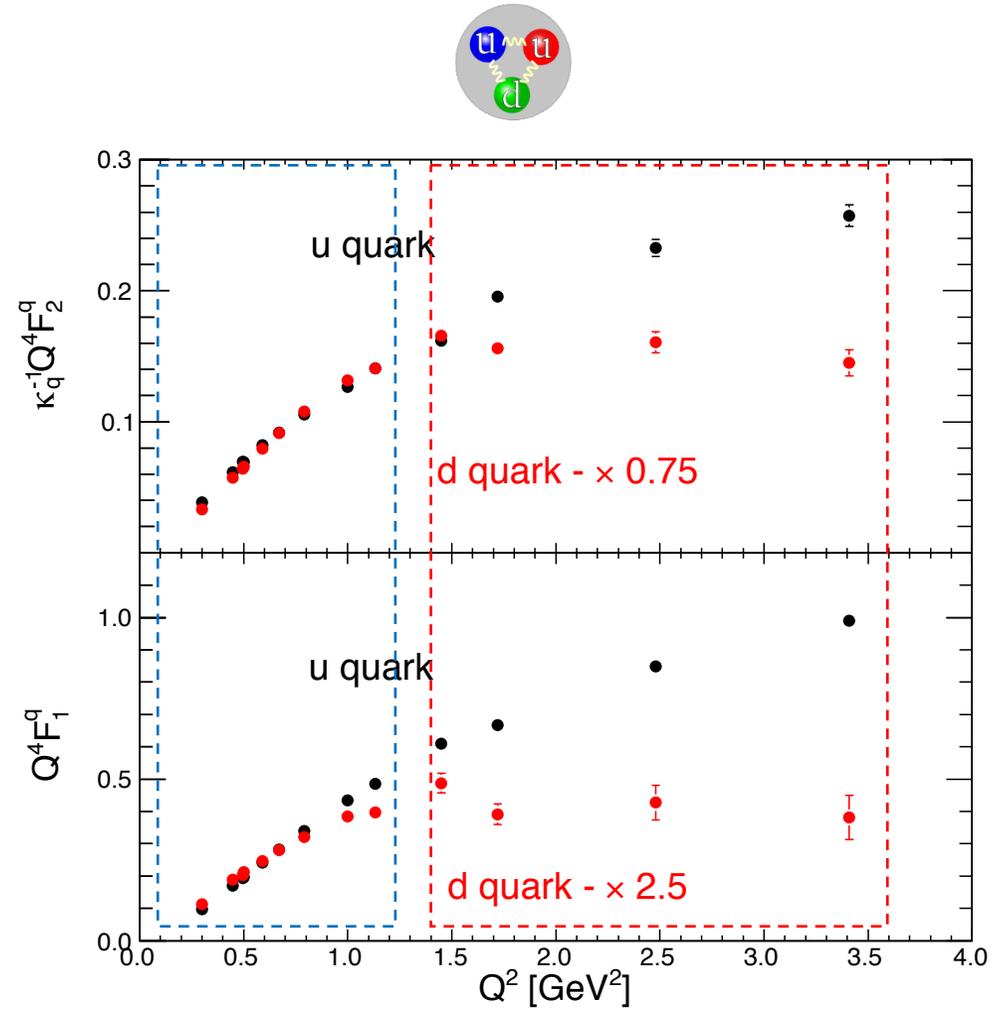
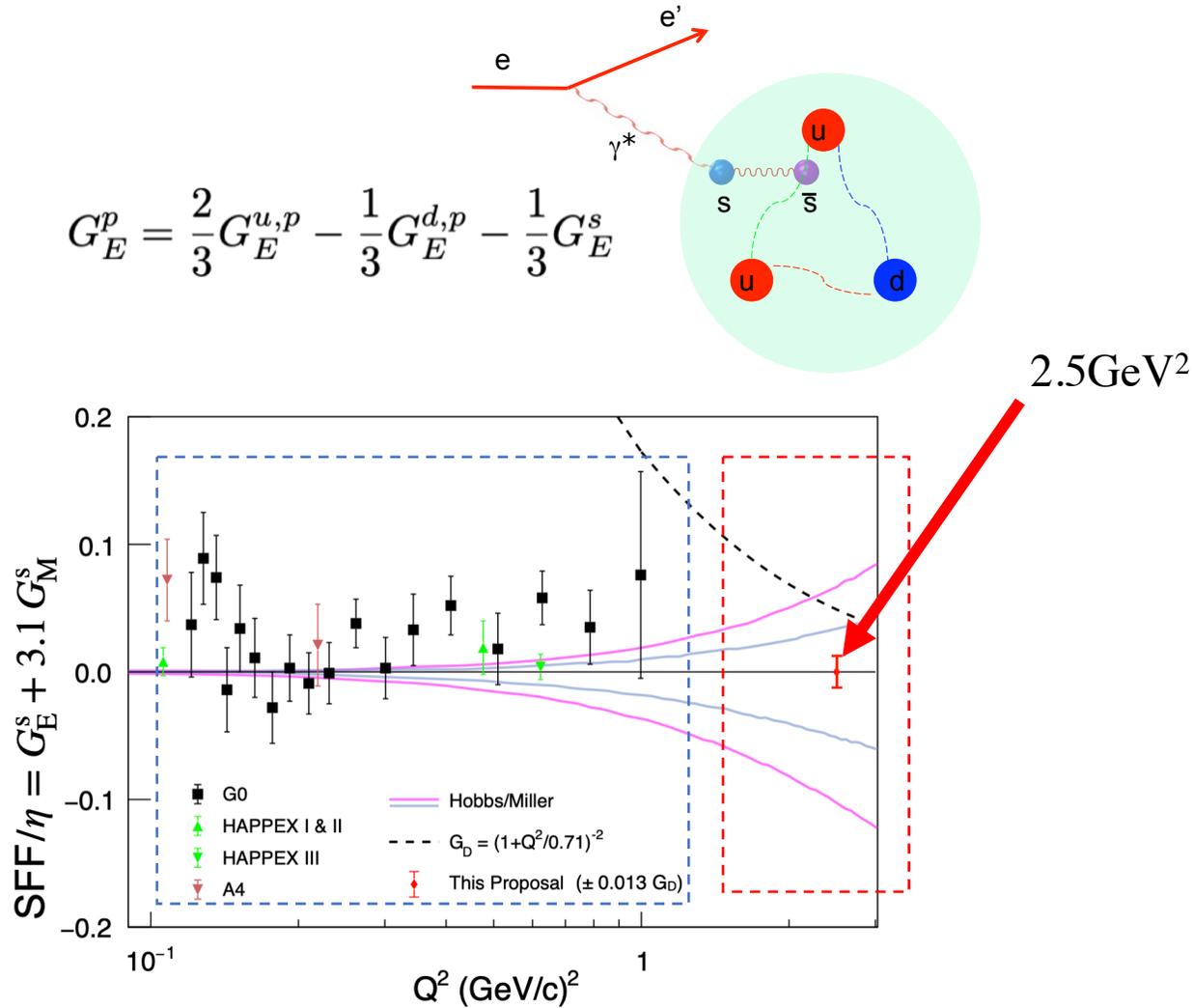
The  $u$ - and  $d$ -quark contributions to the elastic nucleon electromagnetic form factors have been determined using experimental data on  $G_E^n$ ,  $G_M^n$ ,  $G_E^p$ , and  $G_M^p$ . Such a flavor separation of the form factors became possible up to  $3.4 \text{ GeV}^2$  with recent data on  $G_E^n$  from Hall A at JLab. At a negative four-momentum transfer squared  $Q^2$  above  $1 \text{ GeV}^2$ , for both the  $u$ - and  $d$ -quark components, the ratio of the Pauli form factor to the Dirac form factor,  $F_2/F_1$ , was found to be almost constant, and for each of  $F_2$  and  $F_1$  individually, the  $d$ -quark component drops continuously with increasing  $Q^2$ .

- Why is there a breakdown of U/D scaling at  $> 1 \text{ GeV}^2$
- Diquark?
  - “In the framework of Dyson-Schwinger equation calculations, the reduction of the ratios  $F_1^d/F_1^u$  and  $F_2^d/F_2^u$  at high  $Q^2$  is related to diquark degrees of freedom”
- Unless there’s something strange going on....

$$F_1 = \frac{G_E + \tau G_M}{1 + \tau} \quad F_2 = -\frac{G_E - G_M}{1 + \tau}$$

$$\tau = Q^2/4M^2$$

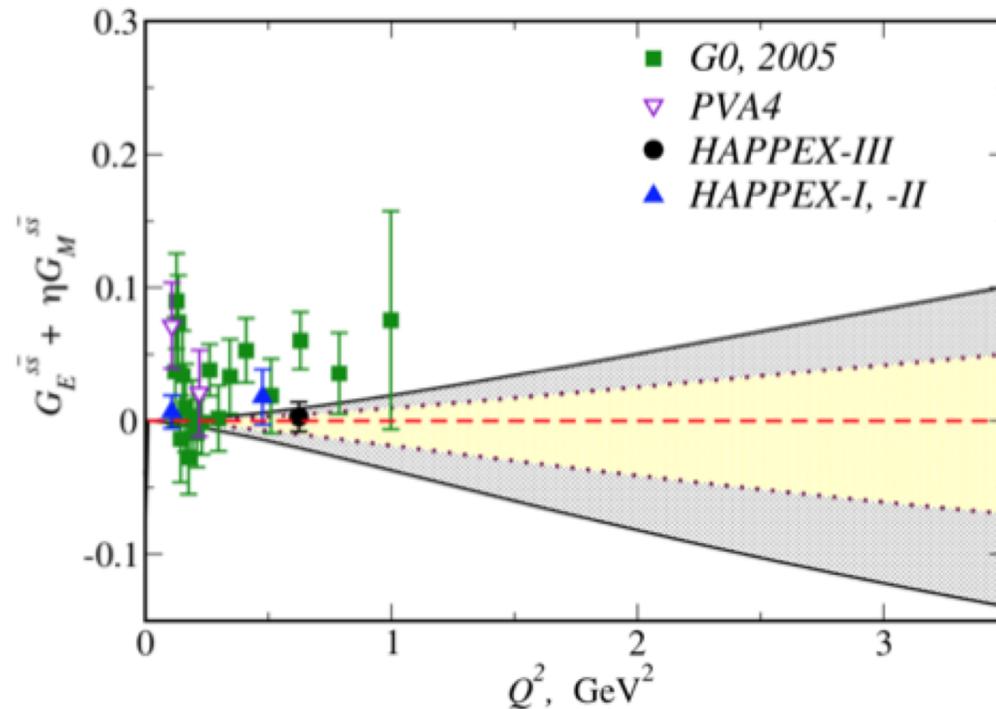
# Strange FF zero at high $Q^2$ ?



- How will SFF impact flavor decomposition of the nucleon FF at higher  $Q^2$ ?

# Strange form-factor predictions

T.Hobbs & J.Miller, 2018



Conclusion: sFF small (but non-zero) at low  $Q^2$ , but quite reasonable to think they may grow relatively large at large  $Q^2$

$G_D = 0.0477$  at  $2.5 \text{ GeV}^2$   
uncertainty here ranges from (0.036,-0.051)

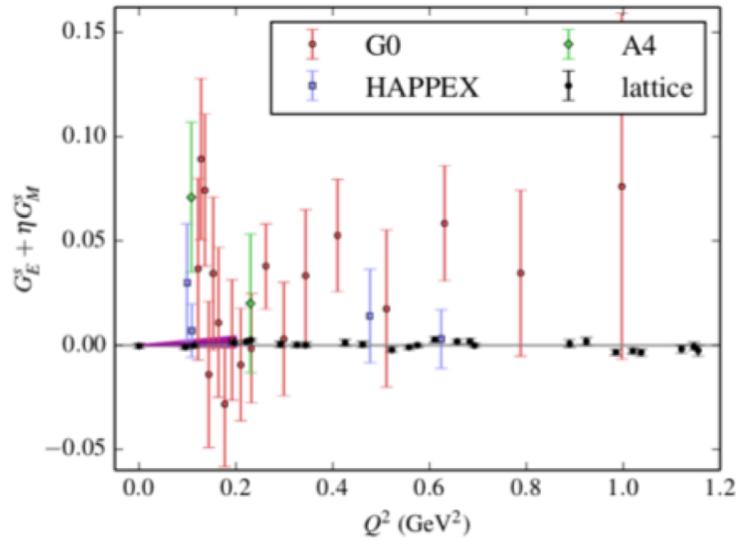
$G_s/G_D \sim 1$  is not excluded

Follows work from *Phys.Rev.C* 91 (2015) 3, 035205  
(LFWF to tie DIS and elastic measurements in a simple model)

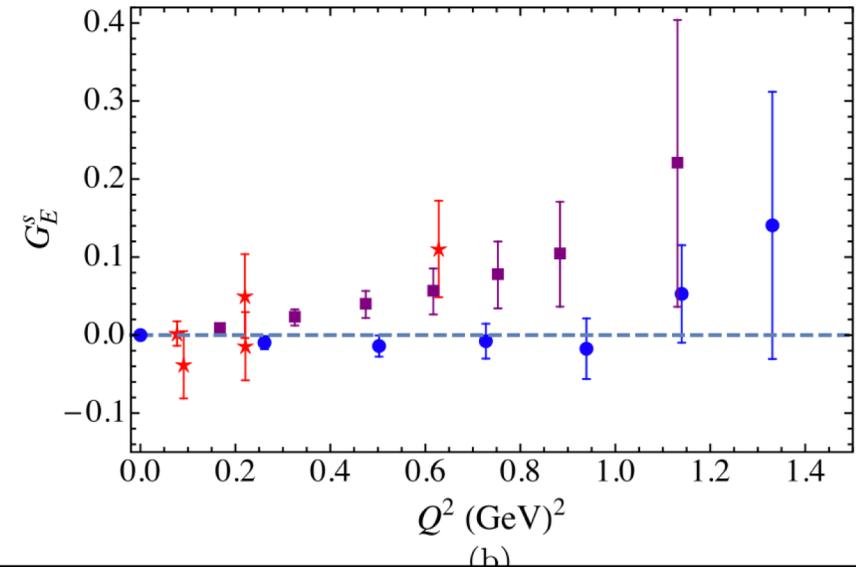
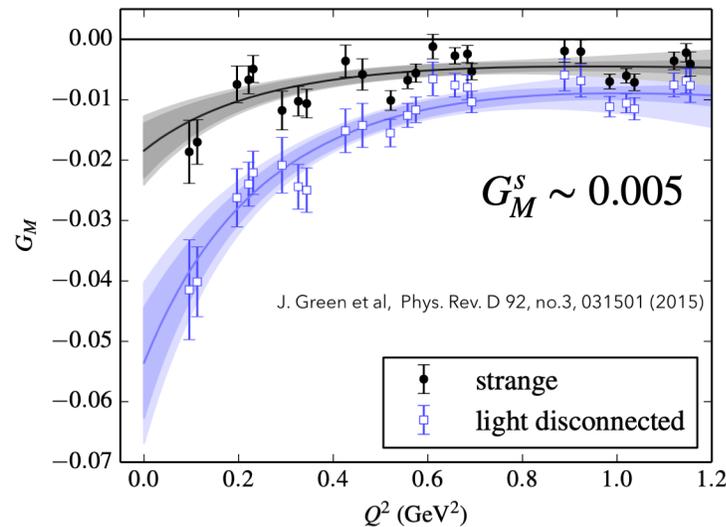
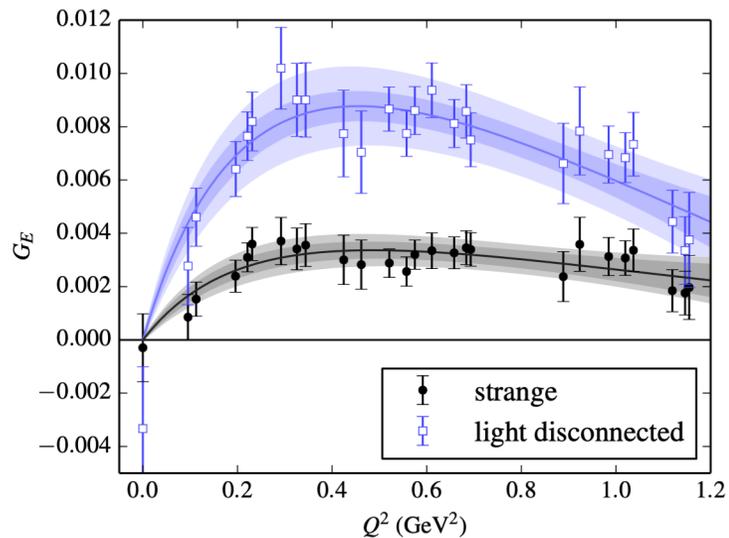
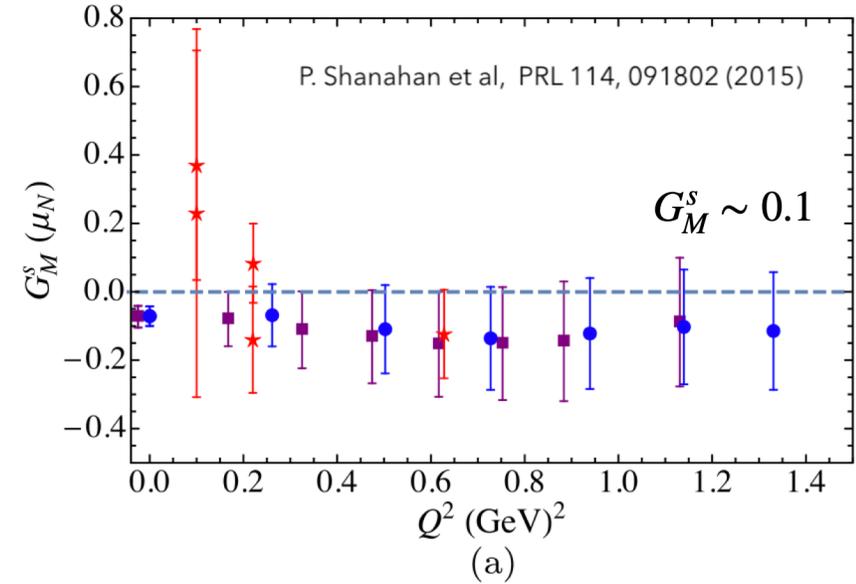
Tim Hobbs and Jerry Miller have both joined the collaboration

# Strange form-factors on the lattice

J.Green etal, 2015

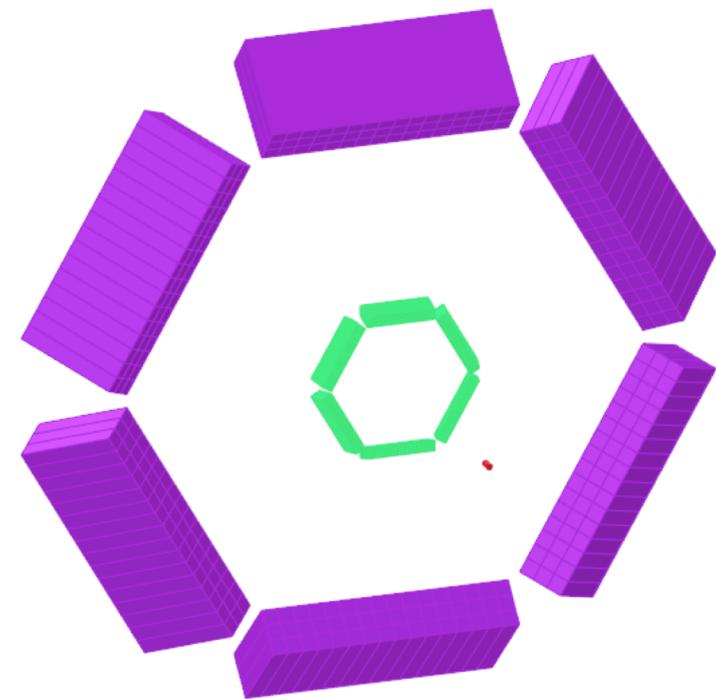
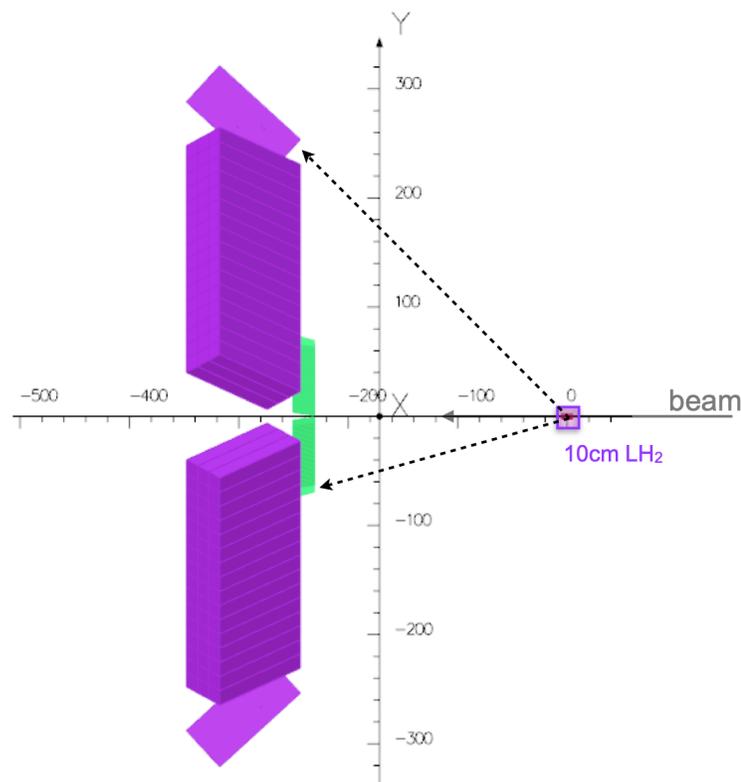
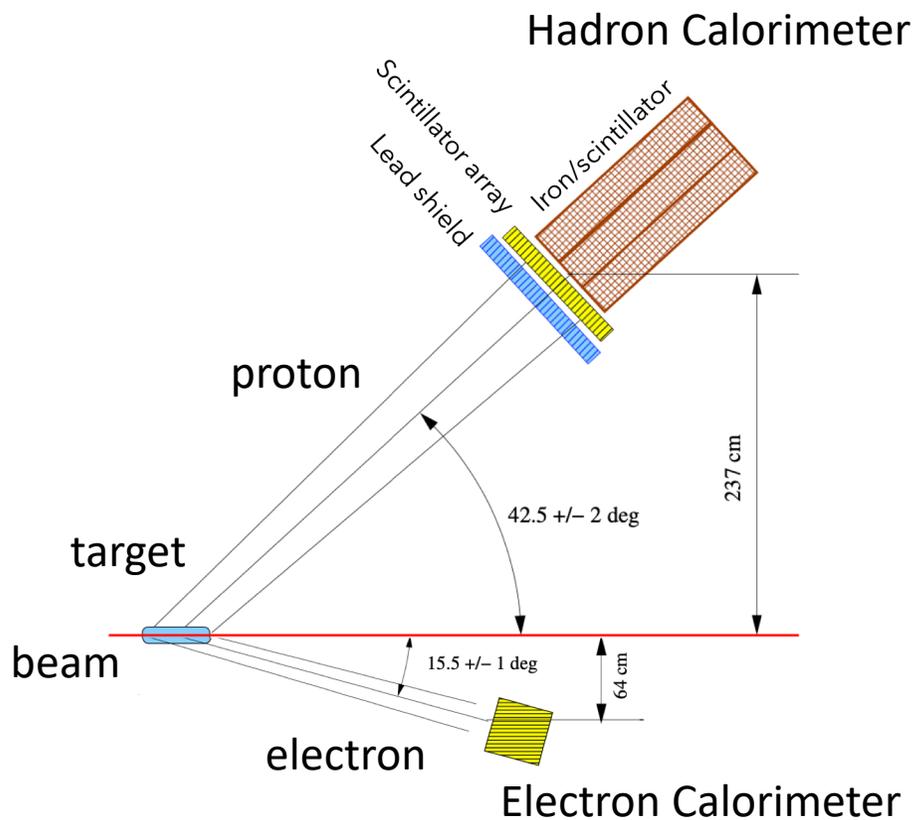


Some lattice calculations predict non-zero central values that would be visible with the proposed precision



# Experimental concept

- Elastic kinematics between electron and proton
- Full azimuthal coverage,  $\sim 42$  msr
- High resolution calorimeter for electron arm
- Angular correlation e-p
- 6.6 GeV beam
- Scattered electron at 15.5 degrees
- Scattered proton at 42.4 degrees
- 10 cm LH<sub>2</sub> target,  $60 \mu\text{A}$ ,  $\mathcal{L} = 1.6 \times 10^{38} \text{ cm}^{-2}/\text{s}$



# Detector System

## HCAL - hadron calorimeter

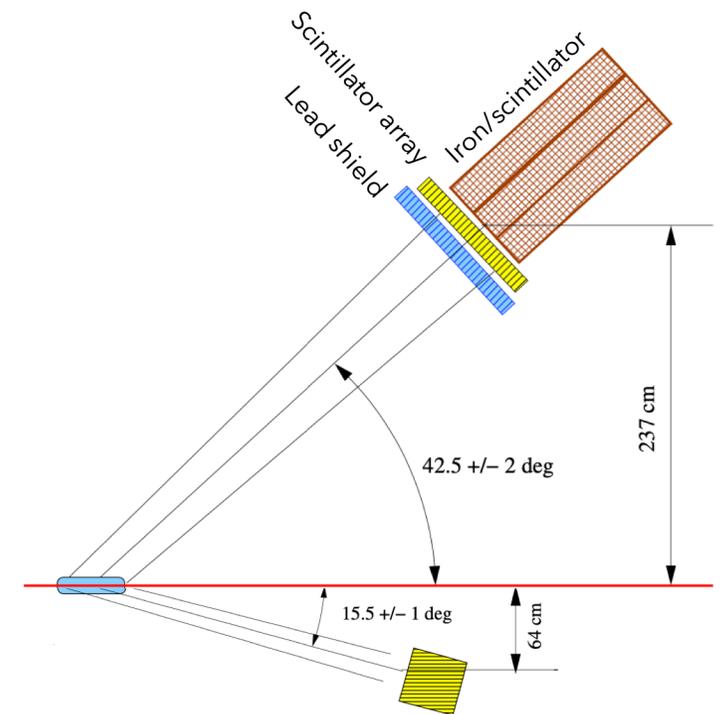
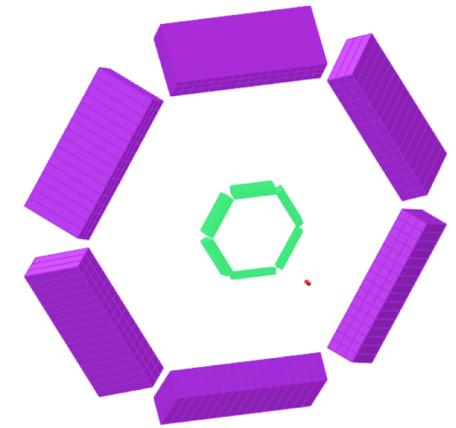
- Reassembled from detector elements from the SBS HCAL
- 288 blocks, each  $15.5 \times 15.5 \times 100 \text{ cm}^3$
- iron/scintillator sandwich with wavelength shifting fiber readout

## ECAL - electron calorimeter

- Reassembled from detector elements from the NPS calorimeter
- 1000 blocks, each  $2 \times 2 \times 20 \text{ cm}^3$
- $\text{PbWO}_4$  scintillator

## Scintillator array

- Used for improved position resolution in front of HCAL
- Not used to form trigger
- 7200 blocks, each  $3 \times 3 \times 10 \text{ cm}^3$
- Lead shield in front (thickness to be optimized) to reduce photon load



# Calorimeters reusing components

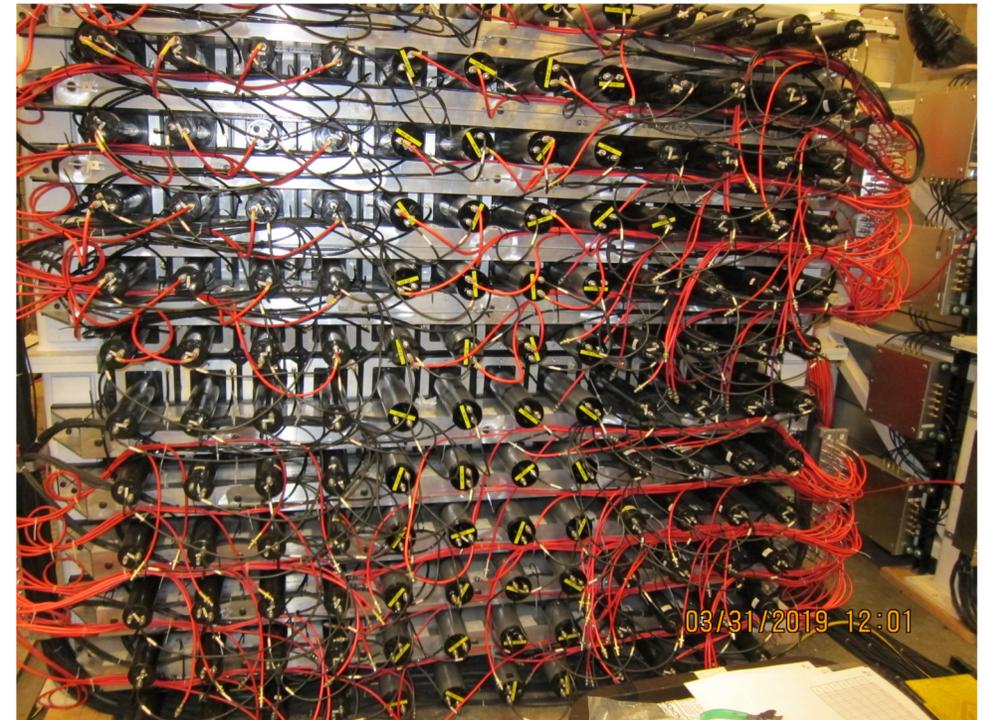
## NPS electromagnetic calorimeter

- 1080 PBWO<sub>4</sub> scintillators, PMTs + bases
- will run in future NPS experiment

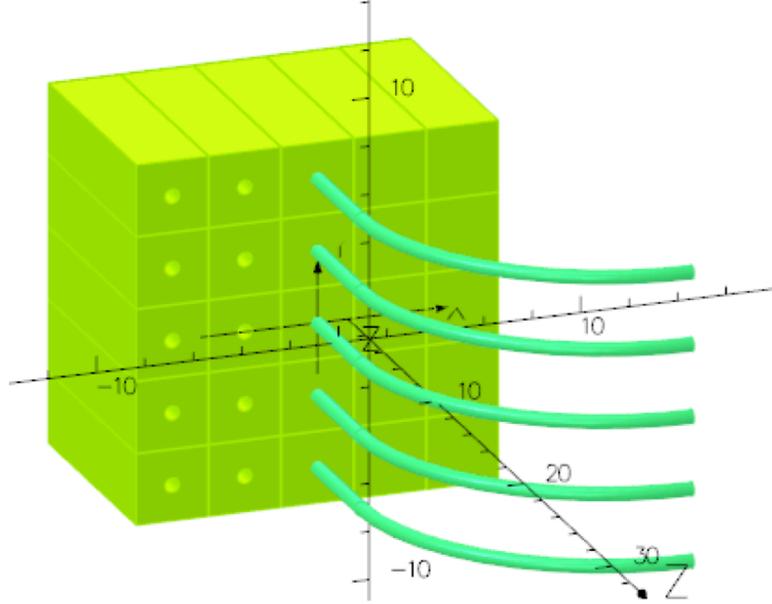


## SBS hadronic calorimeter

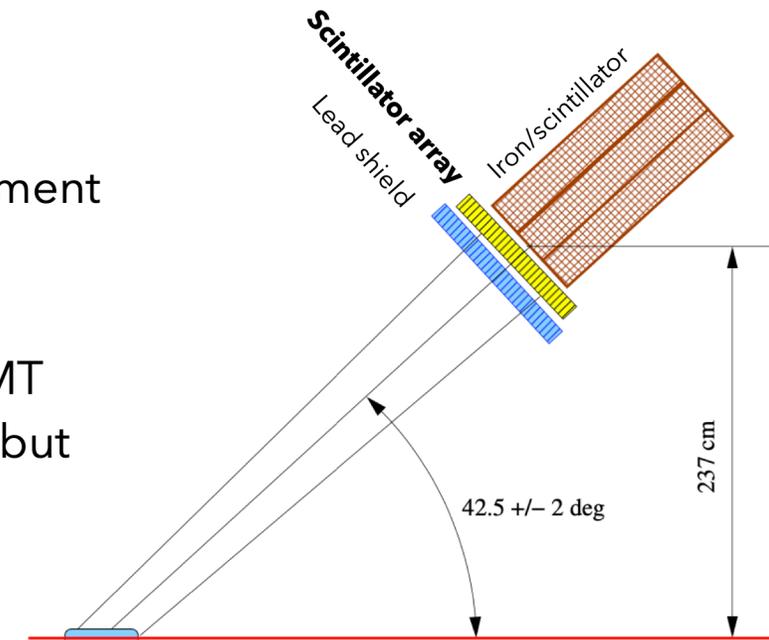
- 288 iron/scintillator detectors, PMTs + bases
- Already in use with SBS



# Scintillator Array



- New detector, must be built for this experiment
- Extruded plastic scintillator block
- Readout with wavelength-shifting fiber
- Each fiber read by pixel on multi-anode PMT
- Originally proposed for 2x2 cross-section, but 3x3 provides sufficient resolution
- 7200 blocks, each 3 x 3 x 10 cm<sup>3</sup>



Design matches scintillator array built for GEP

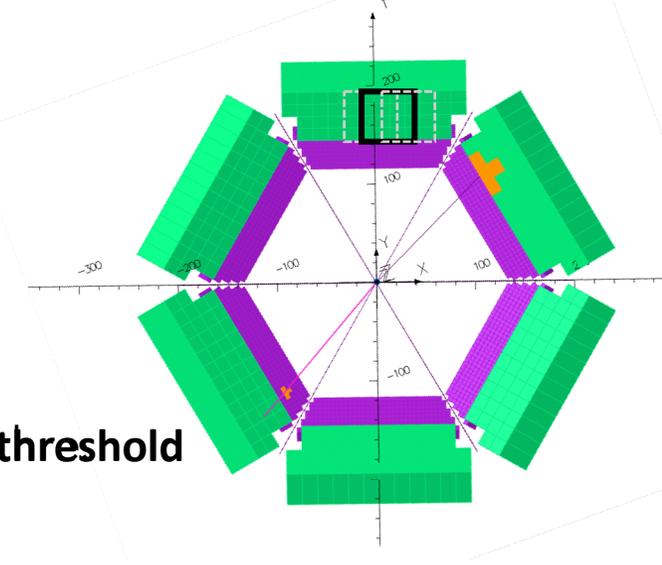
- 2400 elements, 0.5 x 4 x 50 cm<sup>3</sup>
- Already built, will run soon



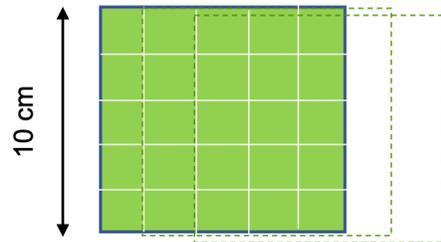
# Triggering and Analysis

Grouping into "subsystems" for energy threshold and coincidence triggering

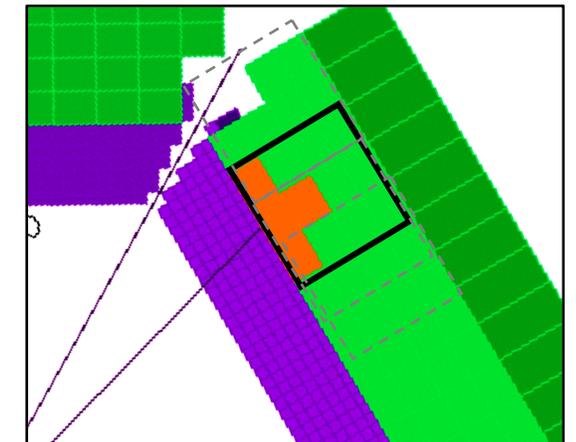
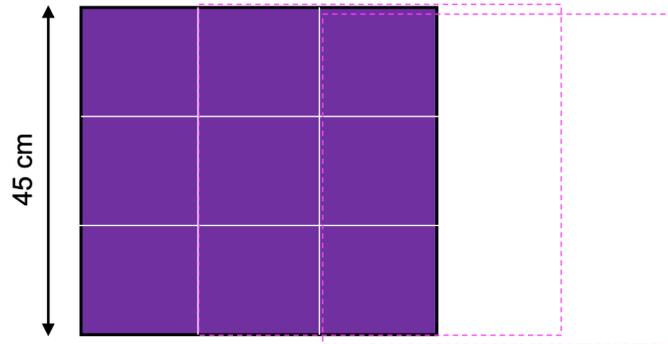
- each polar column of detectors, overlapping with neighbors
- sum amplitude with conservative coincidence timing window
- compare to conservative energy threshold
- trigger when complementary (ECAL and HCAL) subsystems are both above **threshold**



Electron subsystems



Proton subsystems



Subsystems adjacent crystals

- 1000  $\text{PbWO}_4$  crystals
- $2 \times 2 \times 20 \text{ cm}^3$
- 5x5 grouping for subsystem
- 200 overlapping subsystems

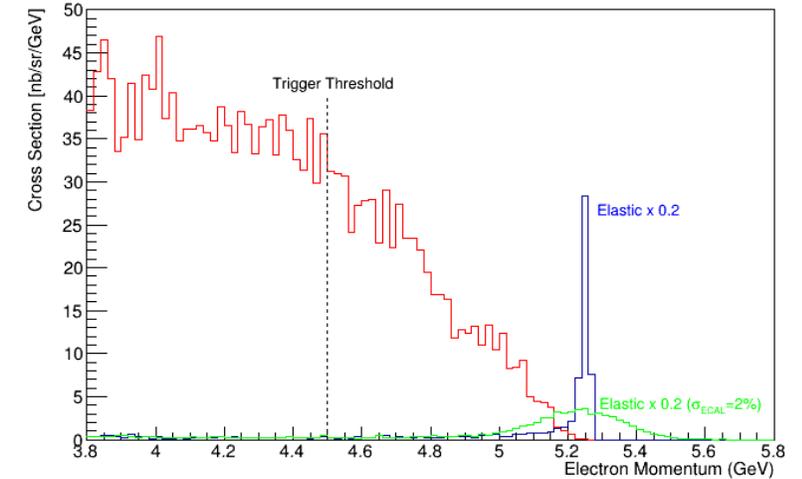
- 288 iron/scintillators
- $15.5 \times 15.5 \times 100 \text{ cm}^3$
- 3x3 grouping for subsystem
- 96 overlapping subsystems

# Rates and Precision

Beam and target: 60 uA on 10 cm LH<sub>2</sub> => luminosity is  $1.6 \times 10^{38}$  cm<sup>-2</sup>/s

## Trigger (online)

- Elastic 37 kHz signal in full detector
- Inelastic (pion production) coincidence trigger rate ~10 kHz
- Accidental coincidence rate < 0.2 kHz
  - ~60 kHz total singles rate in ECAL > 5 GeV energy threshold
  - ~1.2MHz total singles rate in HCAL > 50 MeV energy threshold
- Temporal coincidence cut 40ns
- ~50 kHz total coincidence trigger rate



## Offline analysis

- clustering, scintillator array to improve geometric cuts, tighter acceptance and ECAL cut, 4ns timing
- Accepted elastic signal reduced to 14 kHz - production statistics
- Inelastic (pion production) <0.5%, accidentals < $1 \times 10^{-5}$  due to higher E cut, angular precision

Beam polarization 85%

30 days runtime → Raw asymmetry statistical precision  $\delta(A_{\text{raw}}) \sim 5$  ppm

→  $A_{\text{PV}} = -150 \pm 6.2$  ppm



# Error budget

$A_{PV}$  expected to be -150 ppm (without Strange FF)

quantity	value	contributed uncertainty
Beam polarization	$85\% \pm 1.5\%$	1.8%
Beam energy	$6.6 + / - 0.003 \text{ GeV}$	0.1%
Scattering angle	$15.5^\circ \pm 0.03^\circ$	0.4%
Beam asymmetries	$<100 \text{ nm}, <10 \text{ ppm}$	0.2%
Backgrounds	$< 0.5\%$	0.5%
$G_E^n / G_{\text{Dipole}}$	$0.41 \pm 0.04$	0.6%
$G_E^p / G_{\text{Dipole}}$	$0.75 \pm 0.02$	0.5%
$G_{Mn} / G_{\text{Dipole}}$	$1.01 \pm 0.02$	1.7%
$G_M^p / G_{\text{Dipole}}$	$1.08 \pm 0.01$	0.9%
$G_A^{Zp} / G_{\text{Dipole}}$	$-0.15 \pm 0.02$	0.9%
Total systematic uncertainty:		3.0%

(Polarimetry precision better than 1% has been achieved for multiple experiments)

or 4.5 ppm

Statistical precision for  $A_{PV}$ : 6.2 ppm (4.1%)

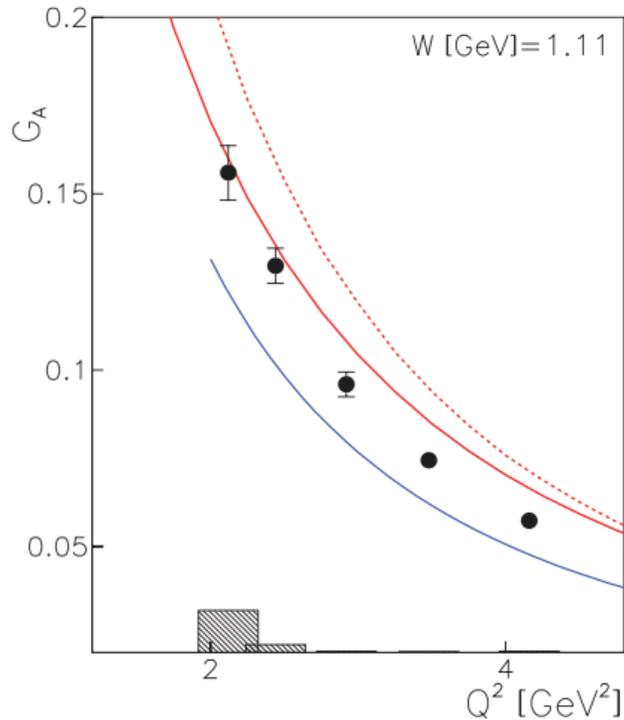
- With 30 days of production running

There is also an uncertainty from radiative correction, is small except for a dominant "anapole" piece.

If the anapole uncertainty is not improved, this would contribute at additional 4.1 ppm (2.7%) uncertainty

# Axial FF Contributions

$$A_A = (1 - 4 \sin^2 \theta_W) \epsilon' G_M^p \tilde{G}_A$$



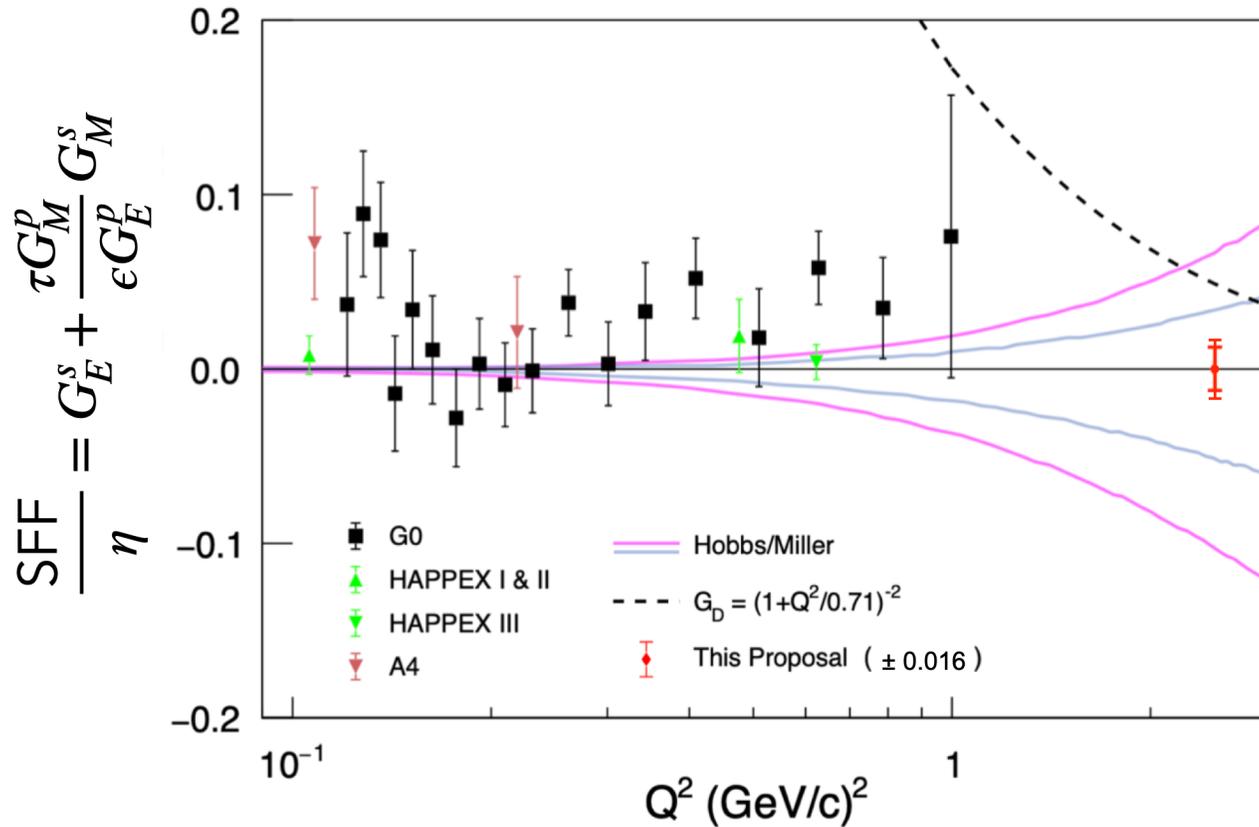
K. Park *et al.* [CLAS Collaboration],  
Phys. Rev. C **85**, 035208 (2012).

- Axial form factor parameterization  $G_A^p = 0.15$  at  $Q^2 = 2.5 \text{ GeV}^2$   
*C. Chen, C. S. Fischer, C. D. Roberts, and J. Segovia, Form factors of the nucleon axial current, Physics Letters B 815, 136150 (2021)*
- Confirmed with pion photoproduction measurements  
K. Park *et al.* [CLAS Collaboration], Phys. Rev. C **85**, 035208 (2012).
  - ( $\sim 15\%$  interpretation uncertainty)  
I.V. Anikin, V.M. Braun, and N. Offen, Phys.Rev.D 94 (2016) 3, 034011.
- How uncertain is this measurement because of it?
  - **Axial term  $\sim 6\%$  of  $A_{\text{PV}}$**
  - $\sim 15\%$  uncertainty, so estimate **1% relative uncertainty** on the 4% statistical measurement

# Projected result

$$\delta A_{PV} = \pm 6.2 \text{ (stat)} \pm 4.5 \text{ (syst)}$$

$$\delta (G_E^s + 3.1G_M^s) = \pm 0.013 \text{ (stat)} \pm 0.010 \text{ (syst)} = 0.016 \text{ (total)}$$



If  $G_M^s = 0$ ,  $\delta G_E^s \sim 0.016$ , (about 34% of  $G_D$ )

If  $G_E^s = 0$ ,  $\delta G_M^s \sim 0.0052$ , (about 11% of  $G_D$ )

The proposed measurement is especially sensitive to  $G_M^s$

The proposed error bar reaches the range of lattice predictions, and the empirically unknown range is much larger.

# Summary

## *Search for a Nonzero Strange Form Factor of the Proton at $2.5 \text{ (GeV/c)}^2$*

- 10+ years after the last sFF searches were performed, a new experiment is proposed for much higher  $Q^2$ , motivated by interest in flavor decomposition of electromagnetic form factors
- Measurement of  $A_{pv}$  in e-p scattering at high  $Q^2 = 2.5 \text{ GeV}^2$  with projected accuracy of 11% of the dipole value
- The proposed error bar is in the range possibly suggested by lattice predictions, and significantly inside the range from the simple extrapolation from previous data
- Coincidence-Parity measurement with highly segmented calorimeters
- Goal: Confirm whether Strange Form Factor is zero or non-zero in this kinematic region
- Result informs the breakdown of U/D scaling
- These results will be crucial to support the interpretation of the nucleon form-factors as constraints on GPDs
- Allows for flavor decomposition at high  $Q^2$

