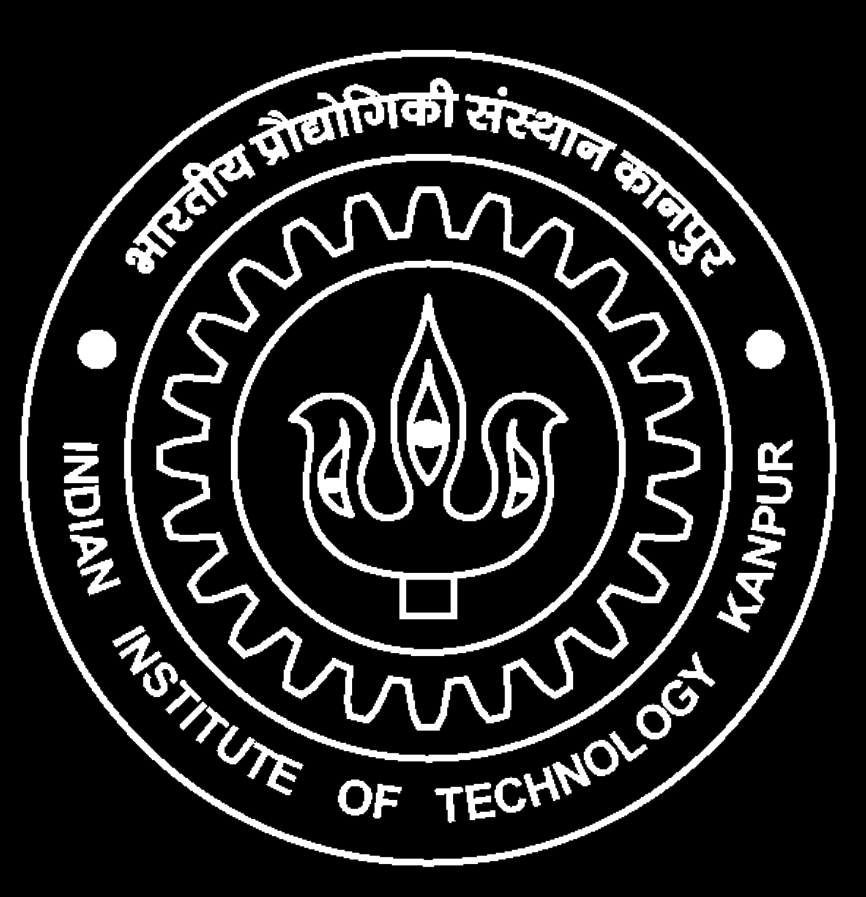




$\cos 2\phi_t$ azimuthal asymmetry in back-to-back J/ψ -jet production in $e p \rightarrow e J/\psi Jet X$ at the EIC

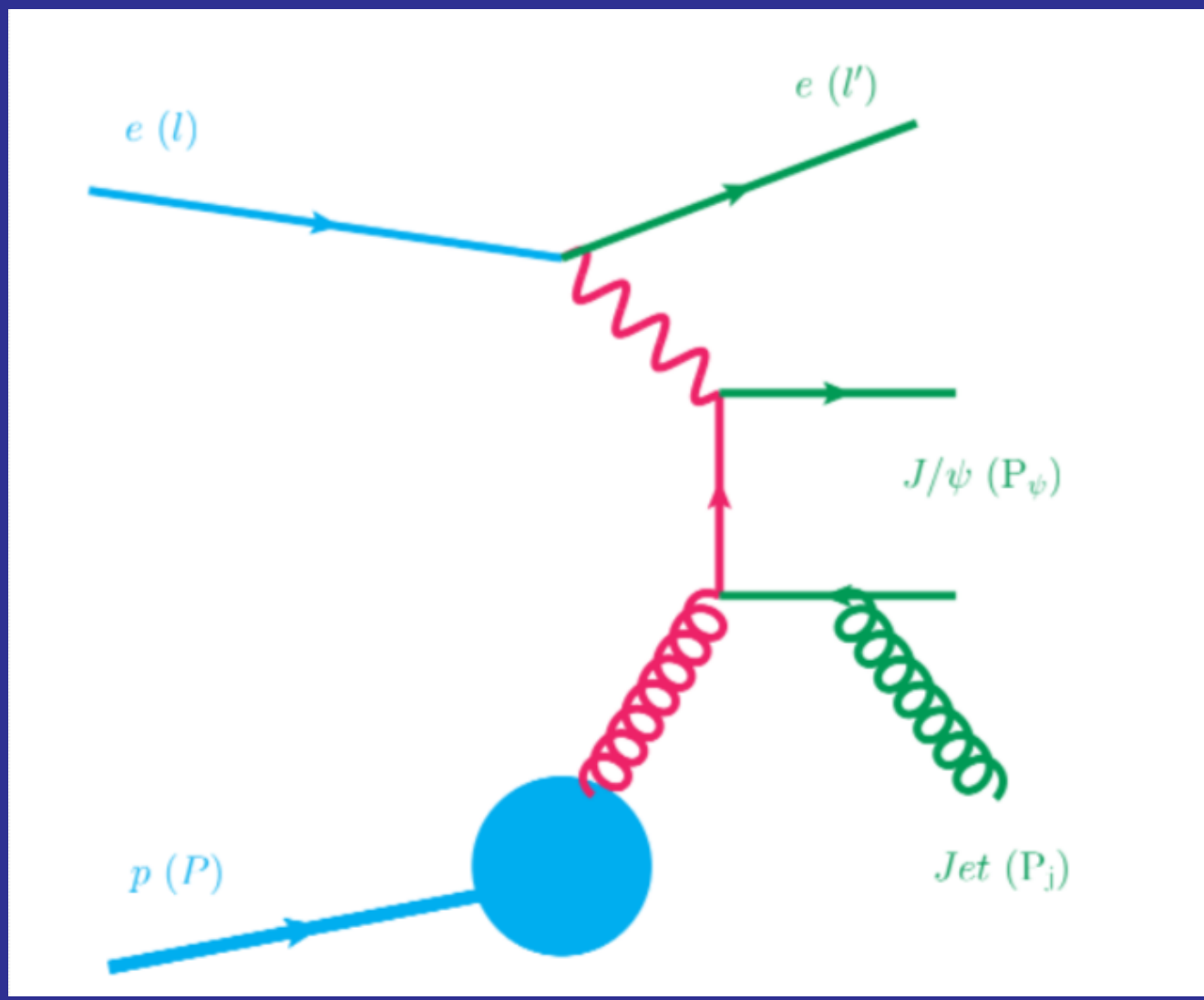
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Introduction

- We present the $\cos 2\phi_t$ azimuthal asymmetry in $e p \rightarrow e J/\psi Jet X$, where the J/ψ -jet pair is almost back-to-back in the transverse plane, within the framework of the generalized parton model(GPM).
- We use non-relativistic QCD(NRQCD) to calculate the J/ψ production amplitude and incorporate both color singlet(CS) and color octet(CO) contributions to the asymmetry.
- We estimate the asymmetry using different parameterizations of the gluon TMDs in the kinematics that can be accessed at the future electron-ion collider (EIC) and also investigate the impact of transverse momentum dependent (TMD) evolution on the asymmetry.

Calculation



$$e^-(l) + p(P) \rightarrow e^-(l') + J/\psi(P_\psi) + Jet(P_j) + X$$

In small-x $\rightarrow \gamma^* + g \rightarrow c\bar{c}(^{2S+1}L_J^{(1,8)}) + g$
D'Alesio (2019)

virtual photon and Proton along $\pm z$ axis

Leptonic plane \Rightarrow measuring azimuthal angles

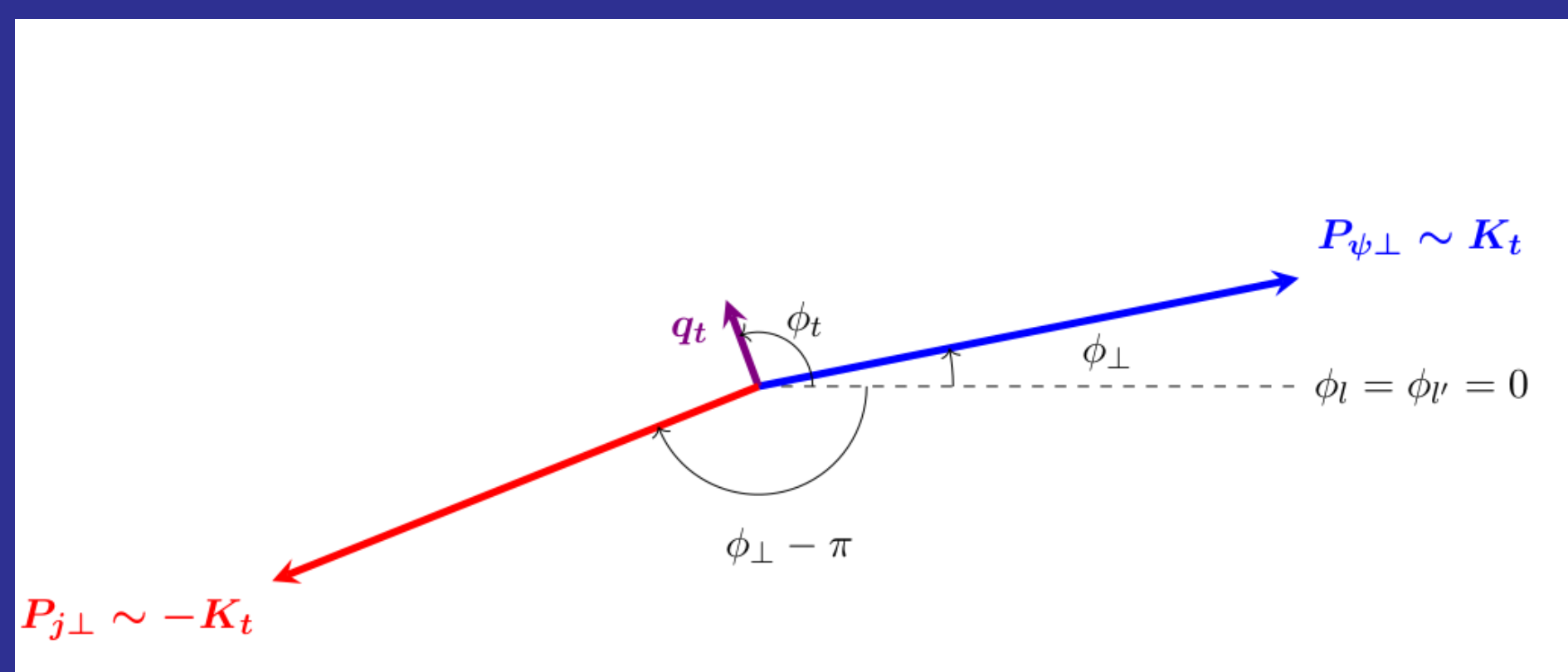
$$Q^2 = -q^2, \quad s = (P+l)^2, \quad x_B = \frac{Q^2}{2P \cdot q},$$

$$y = \frac{P \cdot q}{P \cdot l}, \quad z = \frac{P \cdot P_\psi}{P \cdot q}.$$

back to back scattering $\Rightarrow |\mathbf{q}_t|^2 \ll |\mathbf{K}_t|^2 \sim M_\psi^2$
 \Rightarrow TMD factorization.

ϕ_t and ϕ_\perp

$$\mathbf{q}_t \equiv \mathbf{P}_{\psi\perp} + \mathbf{P}_{j\perp}, \quad \mathbf{K}_t \equiv \frac{\mathbf{P}_{\psi\perp} - \mathbf{P}_{j\perp}}{2}$$



small transverse momentum allows us to factorize the TMDs

J/ψ production

$$d\sigma^{ab \rightarrow J/\psi} = \sum_n d\sigma[ab \rightarrow c\bar{c}(^{2S+1}L_J^{(1,8)})]$$

$$\langle 0 | \mathcal{O}^{J/\psi}(^{2S+1}L_J^{(1,8)}) | 0 \rangle$$

$d\sigma[ab \rightarrow c\bar{c}(^{2S+1}L_J^{(1,8)})]$: High energy perturbative part.

$\langle 0 | \mathcal{O}^{J/\psi}(^{2S+1}L_J^{(1,8)}) | 0 \rangle$: Non perturbative Long Distance Matrix Element. $c\bar{c}(^{2S+1}L_J^{(1,8)})$ to J/ψ

Bodwin, Braaten, Lepage (1994)

Differential cross-section and $\cos 2\phi_t$ asymmetry

$$\frac{d\sigma}{dz dy dx_B d^2\mathbf{q}_t d^2\mathbf{K}_t} = d\sigma^U + d\sigma^T$$

$$\frac{d\sigma}{dz dy dx_B d^2\mathbf{q}_t d^2\mathbf{K}_t} = \frac{1}{(2\pi)^4} \frac{1}{16sz(1-z)Q^4} \left\{ (\mathbb{A}_0 + \mathbb{A}_1 \cos \phi_\perp + \mathbb{A}_2 \cos 2\phi_\perp) f_1^g(x, \mathbf{q}_t^2) + \frac{\mathbf{q}_t^2}{M_P^2} h_1^{\perp g}(x, \mathbf{q}_t^2) (\mathbb{B}_0 \cos 2\phi_t + \mathbb{B}_1 \cos(2\phi_t - \phi_\perp) + \mathbb{B}_2 \cos 2(\phi_t - \phi_\perp) + \mathbb{B}_3 \cos(2\phi_t - 3\phi_\perp) + \mathbb{B}_4 \cos(2\phi_t - 4\phi_\perp)) \right\}$$

$$\langle \cos 2\phi_t \rangle \equiv A^{\cos 2\phi_t} = 2 \frac{\int d\phi_t d\phi_\perp \cos 2\phi_t d\sigma(\phi_t, \phi_\perp)}{\int d\phi_t d\phi_\perp d\sigma(\phi_t, \phi_\perp)}$$

$$\langle \cos 2\phi_t \rangle \equiv A^{\cos 2\phi_t} = \frac{\int \mathbf{q}_t d\mathbf{q}_t \frac{\mathbf{q}_t^2}{M_P^2} \mathbb{B}_0 h_1^{\perp g}(x, \mathbf{q}_t^2)}{\int \mathbf{q}_t d\mathbf{q}_t \mathbb{A}_0 f_1^g(x, \mathbf{q}_t^2)}$$

$$U = \frac{2 * |\mathbb{B}_0|}{\mathbb{A}_0} \bigg|_{\frac{\mathbf{q}_t^2}{2M_P^2} |h_1^{\perp g}(x, \mathbf{q}_t^2)| = f_1^g(x, \mathbf{q}_t^2)}$$

D'Alesio (2019)

TMD parameterizations

Gaussian parameterization : Transverse momentum dependent part is gaussian in nature. C. Pisano, et. al. (2012)

Spectator model : It is a numerical model to describe TMDs in small-x region. P. Tael., et.al. (2020)

TMD Evolution: In TMD evolution the TMDs are solutions of Collins Soper Sterman evolution equation. D. Boer (2020)

Results and Discussion

