

# Kalman Filter for track reconstruction in BONuS12

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HUGS

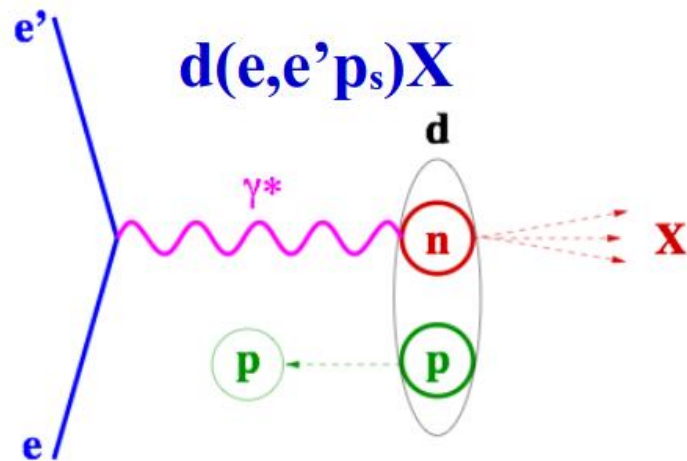
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## Quasi-Free Neutron Structure at Large $x_B$

Parton Distribution Functions (PDF) :

- Provide information on the partons longitudinal momentum distributions.
- Measurable via Deep Inelastic Scattering (DIS).
- For nucleons, the unpolarized DIS cross section is parametrized by two PDFs:  $F_{1,2}(x)$ .



$$\frac{F_{2n}}{F_{2p}} \approx \frac{1 + 4d/u}{4 + d/u} \rightarrow \frac{d}{u} = \frac{4F_{2n}/F_{2p} - 1}{4 - 4F_{2n}/F_{2p}}$$

## Tagged-proton nDVCS

General Parton Distribution Functions (GPD) :

- Mapping out simultaneously the space and momentum components of quarks and gluons.
- Measurable via Deeply Virtual Compton Scattering (DVCS).
- At JLab there are four GPDs accessible:

$$H(x, \xi, t), \tilde{H}(x, \xi, t), E(x, \xi, t), \tilde{E}(x, \xi, t)$$

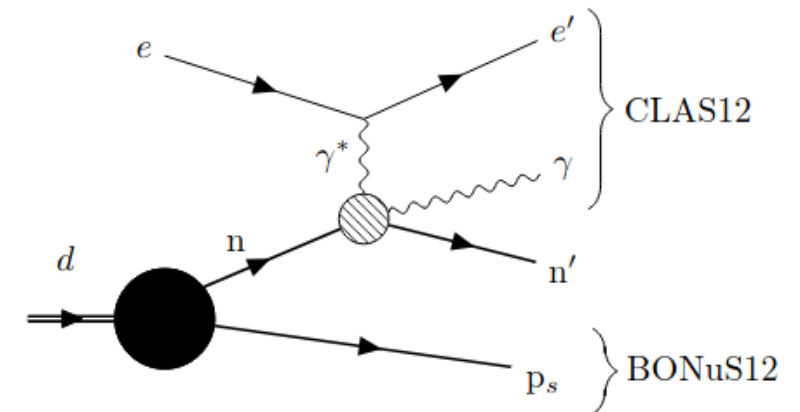


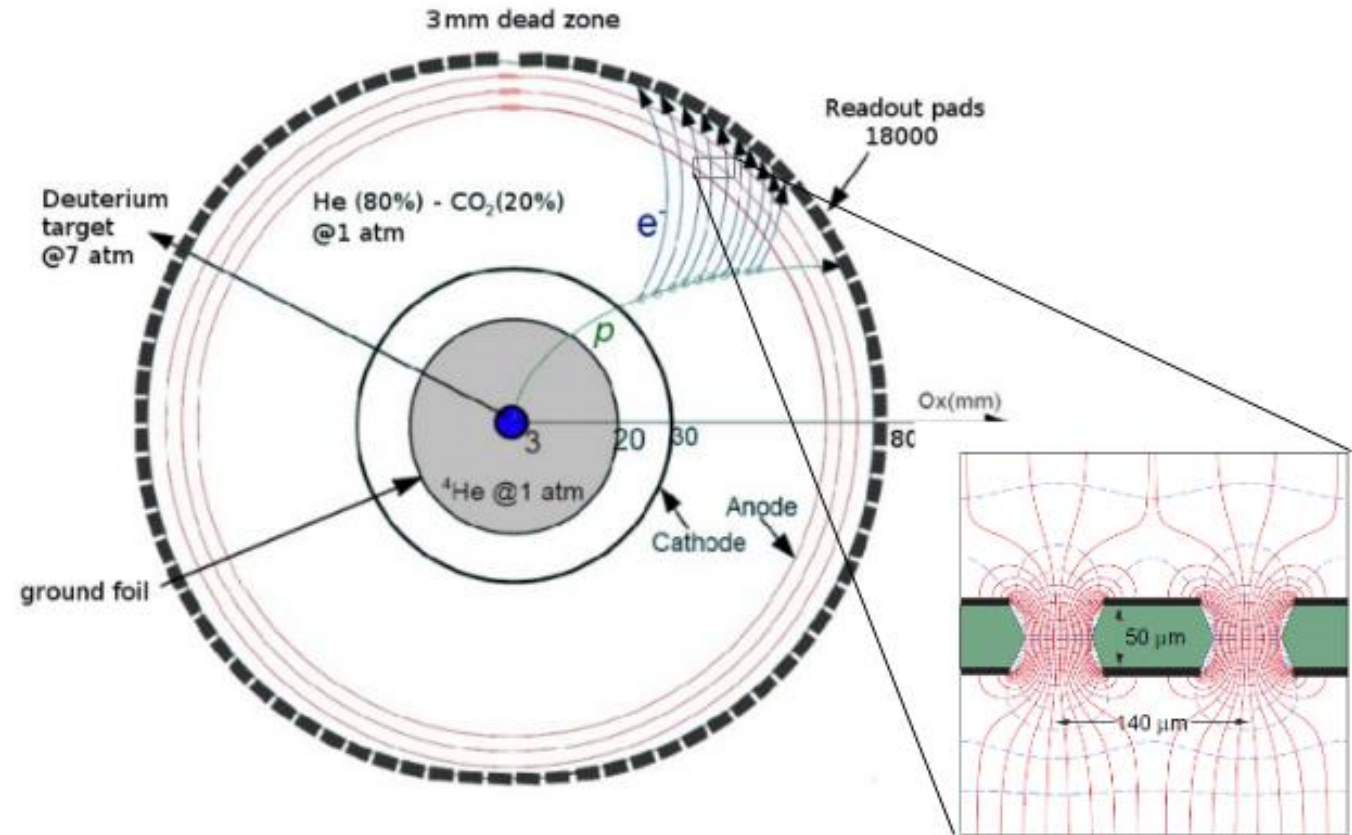
Figure 1.4: Proton-tagged neutron DVCS diagram in deuterium.



# Introduction: RTPC Detector

- Design :
  - Nearly 100% azimuthal coverage
  - 400 mm long , 160 mm  $\varnothing$ .
  - 6 mm diameter target : 5 atm D.
  - 30 mm radius of cathode foil (4  $\mu\text{m}$  thick).
  - 40 mm drift region : 80%  $^4\text{He}$  and 20%  $\text{CO}_2$ .
  - 3 GEMs layers, gain of 100/layer
  - 17280 readout elements.
- Work principle:
  - Charged particle ionizes the gas atoms.
  - Under EM field, released electrons follow their drift paths.
  - Amplifications via the 3 GEM layers.
  - Use energy loss for particle identification.

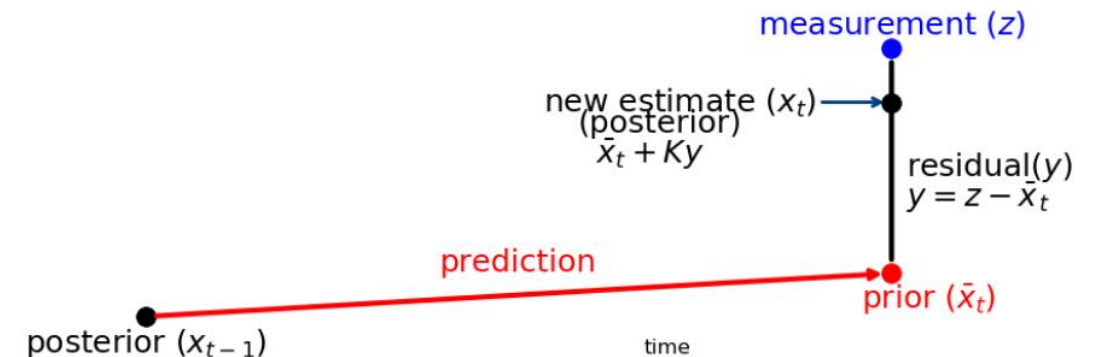
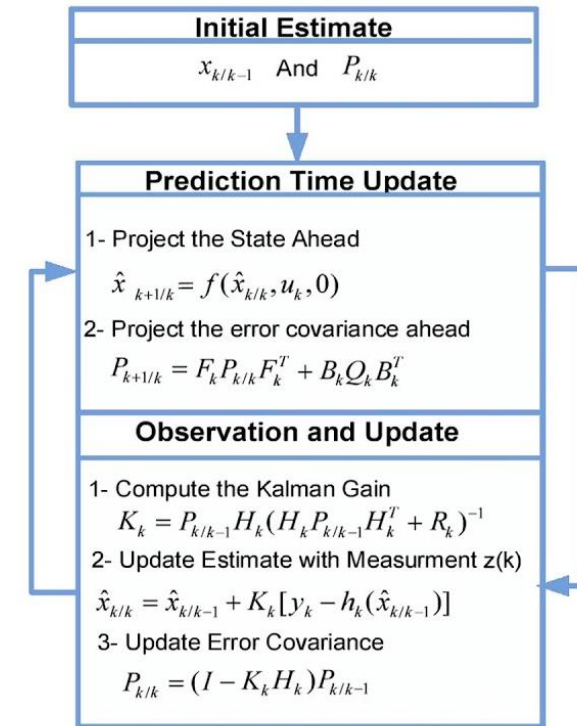
$$\frac{dE}{dx} = \frac{\sum_i \frac{ADC_i}{G_i}}{path}$$





# Theory of the Kalman Filter

- Principle :
  - The Kalman filter is a set of mathematical equations that provides an efficient computational (recursive) means to estimate the state of a process.
  - The algorithm works by a two-phase process :
    - The prediction phase : the Kalman filter produces estimates of the current state variables, along with their uncertainties.
    - The update phase : next measurement is observed, these estimates are updated.
- Components :
  - Measurement vector (site) : describe the particle in detector coordinate system (z).
  - State vector : describe the particle (x).
  - Propagator : describe the motion of the system, propagate the state vector (f).
  - Matrix Propagator : Jacobian matrix of the propagator (F).
- Kalman Filter is useful for take in account the energy loss and need smaller matrix inversion.





The equation of motion of a charged particle in a magnetic field is :

$$\frac{d^2 \vec{r}}{ds^2} = \kappa \frac{q}{p} \left( \frac{d\vec{r}}{ds} \times \vec{B} \right)$$

Suppose the magnetic is uniform, and we assume its direction is parallel with the z-axis of the coordinate system. In that case, the trajectory of the charged particle can be solved analytically, which is a helix.

In the BONuS12 experiment, we must consider the energy loss and the non-uniformity of the magnetic field. So, there is no more analytic solution, we must use numerical methods to solve equation of motion.

The state vector for BONuS12 is :

$$x = (x, y, z, p_x, p_y, p_z)^T$$

We solve this equation with a Runge-Kutta 4 order algorithm :

$$y_{n+1} = y_n + \frac{1}{6}h(k_1 + 2k_2 + 2k_3 + k_4)$$

$$k_1 = f(t_n, y_n)$$

$$k_2 = f(t_n + \frac{h}{2}, y_n + h\frac{k_1}{2})$$

$$k_3 = f(t_n + \frac{h}{2}, y_n + h\frac{k_2}{2})$$

$$k_4 = f(t_n + h, y_n + hk_3)$$

After every step, we take into account the energy loss. We need a small step size to calculate the energy loss correctly.

The Jacobian of the propagator is computed with the numerical derivative method :

$$f'(x) = \frac{f(x+h) - f(x-h)}{2h}$$



## Energy loss

The Bethe-Bloch formula, which is modified taking into account various corrections :

$$\frac{dE}{dx} = 2\pi r_e^2 m c^2 n_{el} \frac{z^2}{\beta^2} \left[ \ln \left( \frac{2mc^2 \beta^2 \gamma^2 T_{up}}{I^2} \right) - \beta^2 \left( 1 + \frac{T_{up}}{T_{max}} \right) - \delta - \frac{2C_e}{Z} + S + F \right]$$

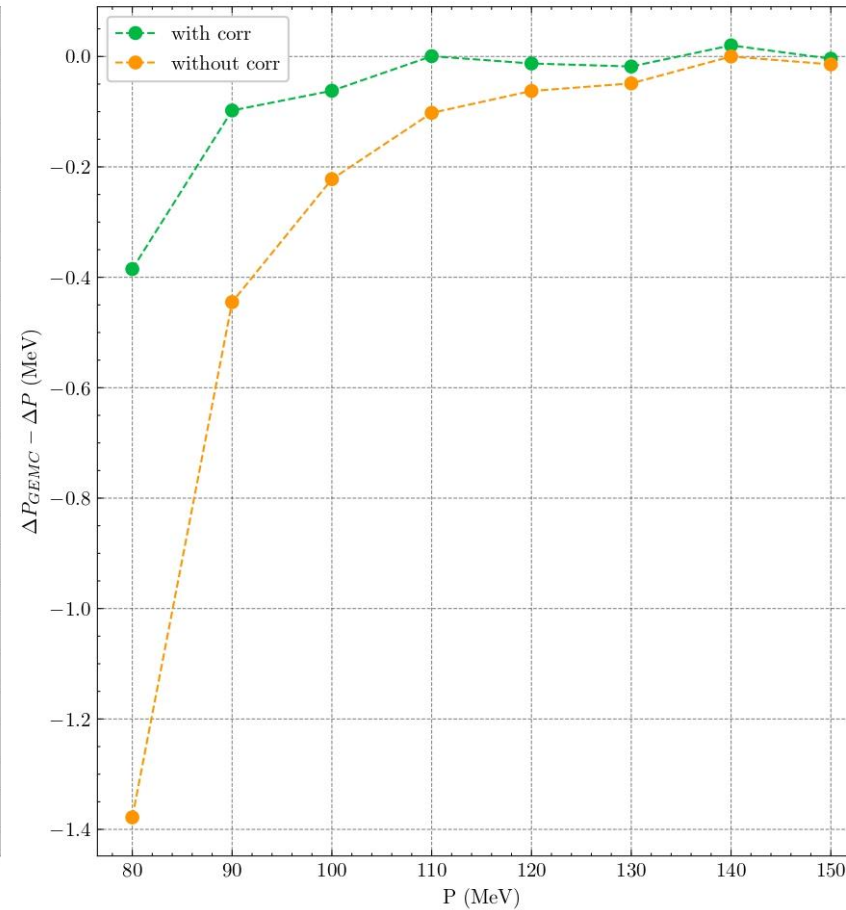
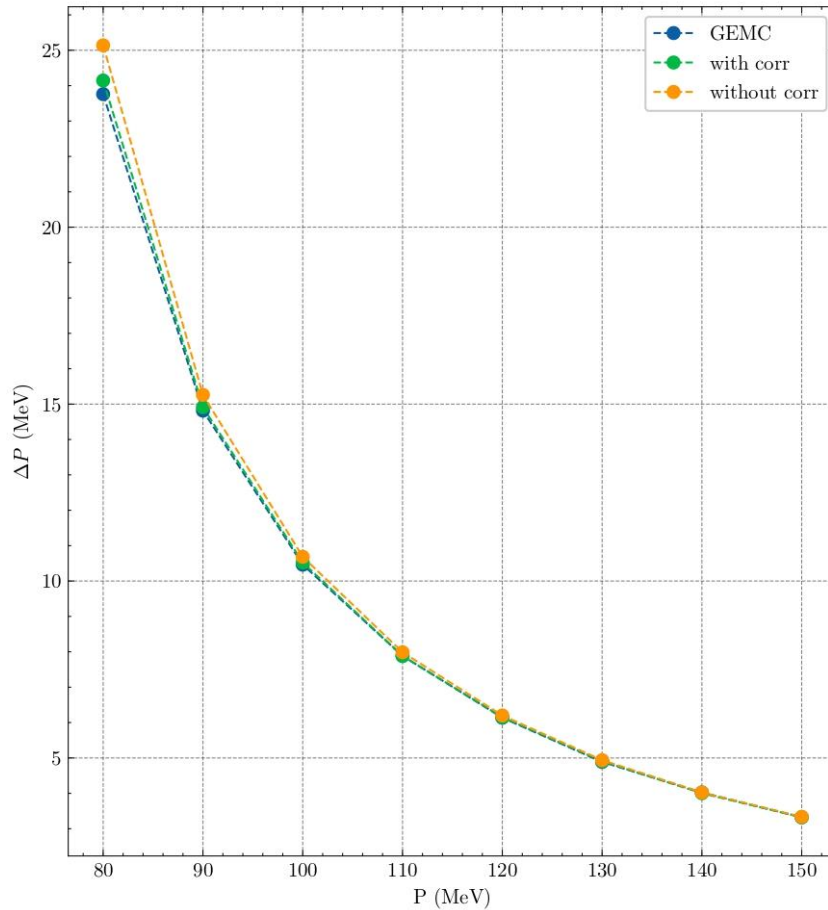
- Standard **Bethe-Bloch** formula.
- **Shell Correction** : is the so-called shell correction term which accounts for the fact of interaction of atomic electrons with atomic nucleus. This term more visible at low energies and for heavy atoms. **Most important correction**.
- **Density Correction** : is a correction term which considers the reduction in energy loss due to the so-called density effect. This becomes important at high energies because media have a tendency to become polarized as the incident particle velocity increases.
- **High Order Corrections** : Mott correction term, Finite size correction term, Barkas correction, Bloch correction, Spin Correction.

Use GEANT4 algorithm rewrite in Java



# Energy loss

$\Delta P$  for proton in BONuS12 with uniform B



Compute energy loss with and without correction :

- Difference in momentum between vertex and GEM
- Important at low momentum.

Mostly due to shell correction





- Initialization :
  - State vector come from classic fit.
  - Error covariance matrix come from observed resolution.
- Measurement errors :
  - R for measurement points. It's weight by the ADC.
  - Specific one for beamline vertex where errors on z and phi are big.
- System error Q :
  - Multiple scattering (very low).
  - Energy loss fluctuation.

$$\Delta r = 70/r \text{ mm}$$

$$\Delta \phi = 1^\circ$$

$$\Delta z = 2 \text{ mm}$$

$$R = \begin{bmatrix} \Delta r^2 & 0 & 0 \\ 0 & \Delta \phi^2 & 0 \\ 0 & 0 & \Delta z^2 \end{bmatrix}$$

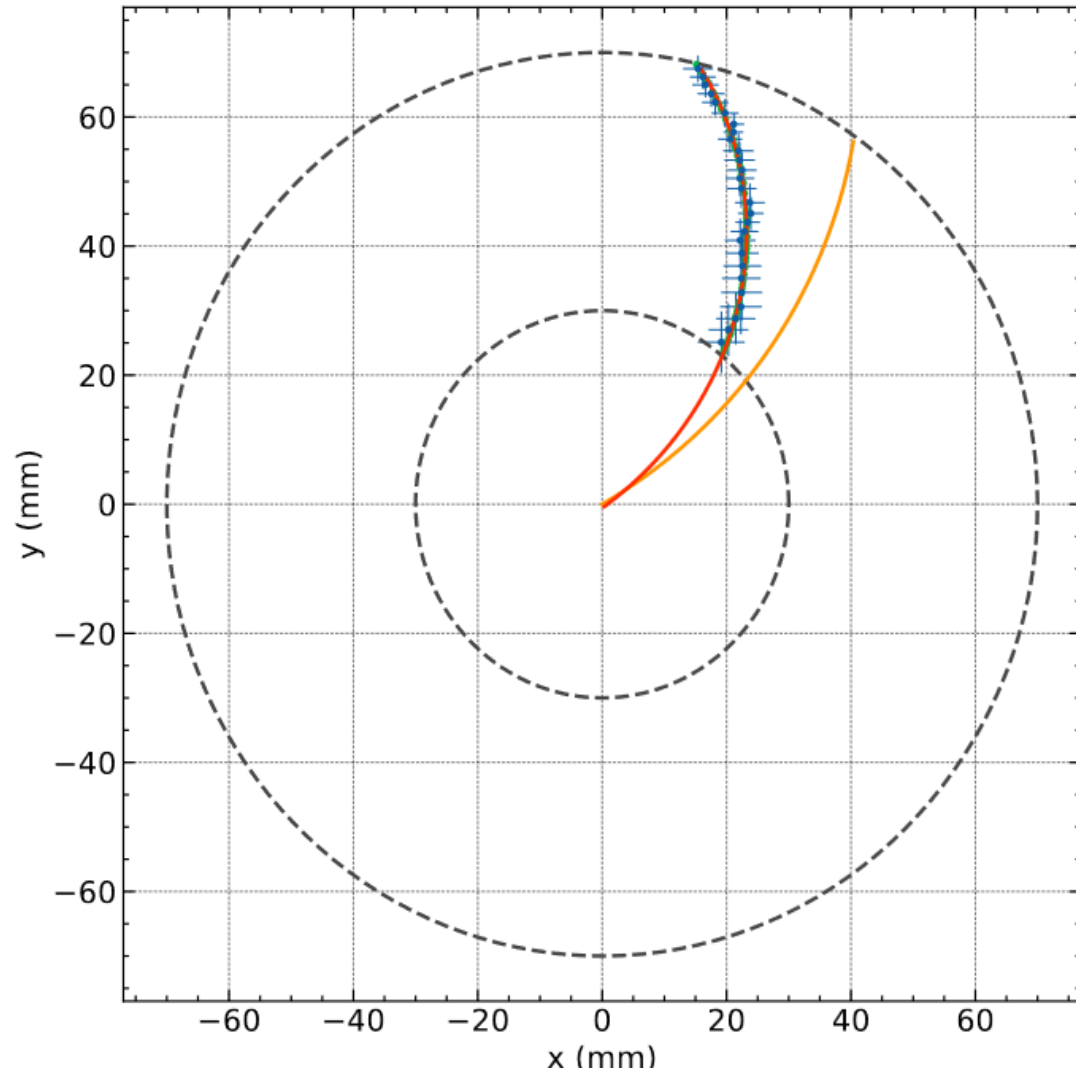
$$Q = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & g(\Omega^2, px) & 0 & 0 \\ 0 & 0 & 0 & 0 & g(\Omega^2, py) & 0 \\ 0 & 0 & 0 & 0 & 0 & g(\Omega^2, pz) \end{bmatrix}$$

$$\Omega^2 = 2\pi r_e^2 m_e c^2 N_{el} \frac{Z_h^2}{\beta^2} T_{max} s \left( 1 - \frac{\beta^2}{2} \frac{T_c}{T_{max}} \right),$$



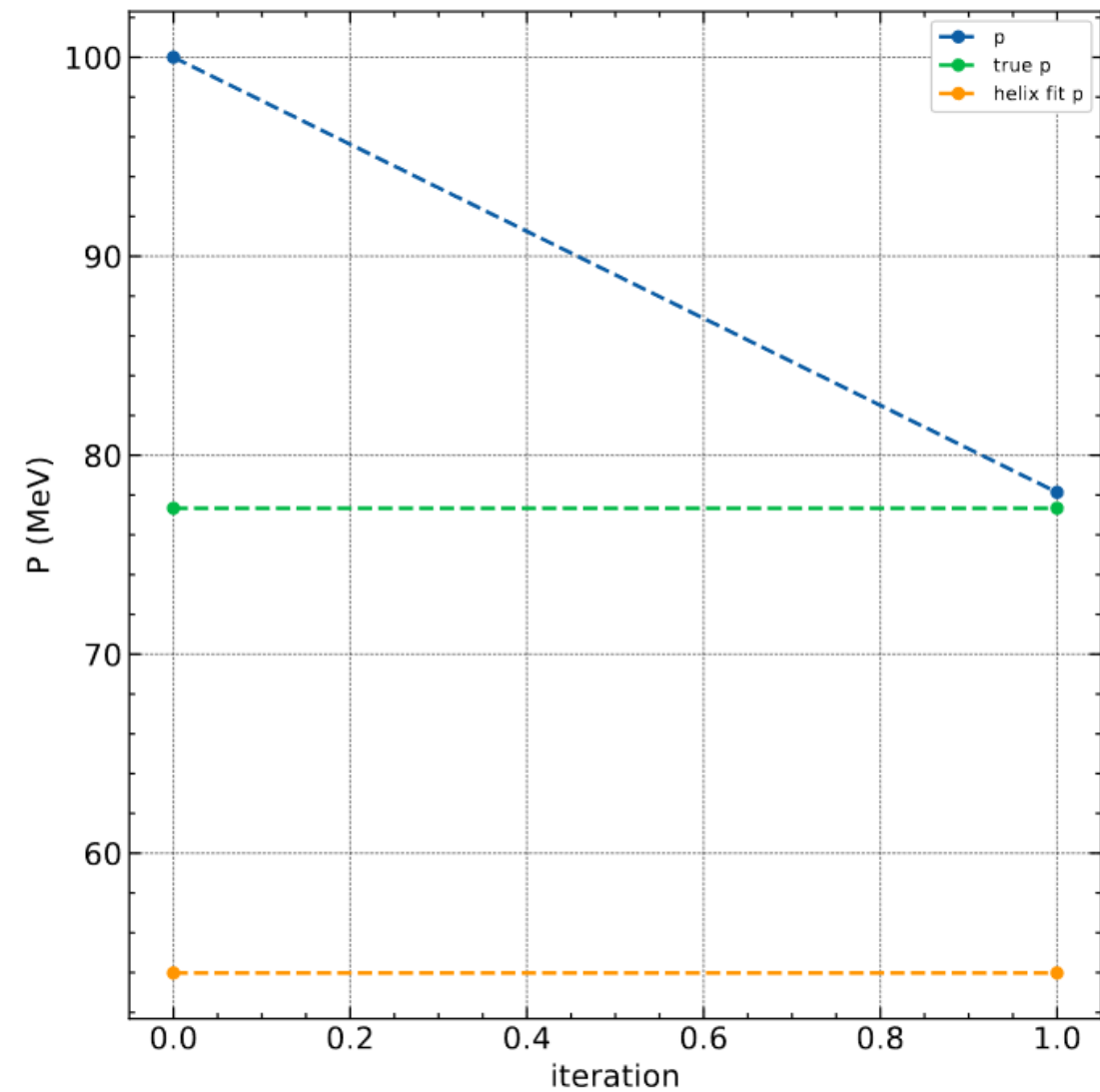


## Result for one track



Initial track

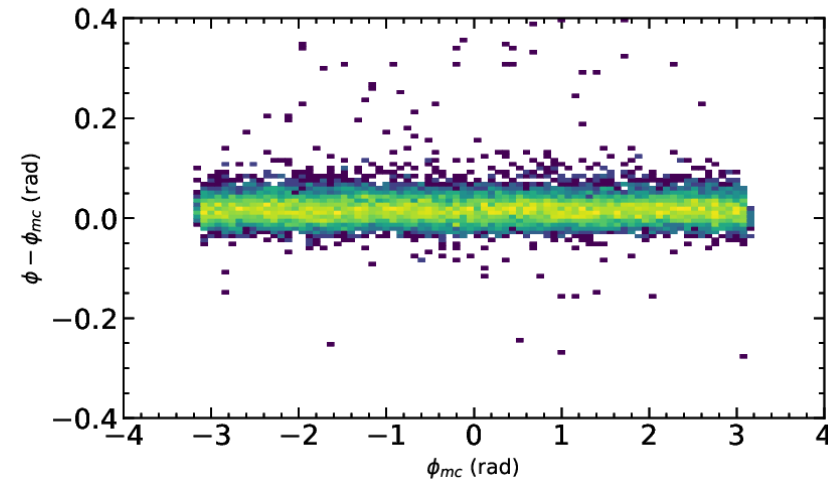
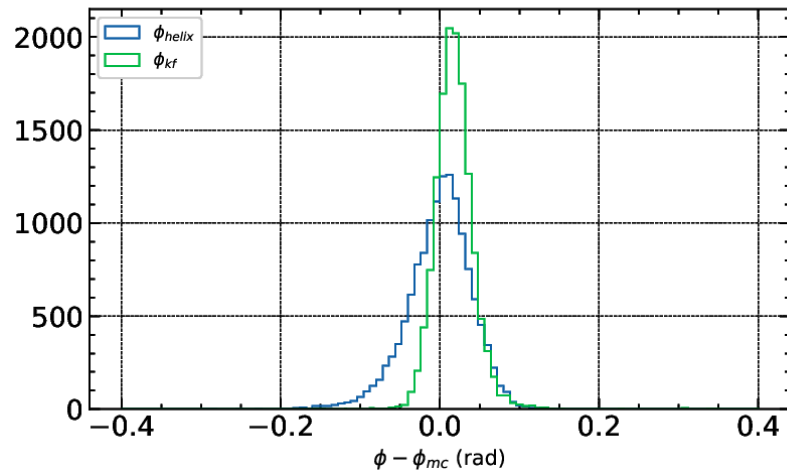
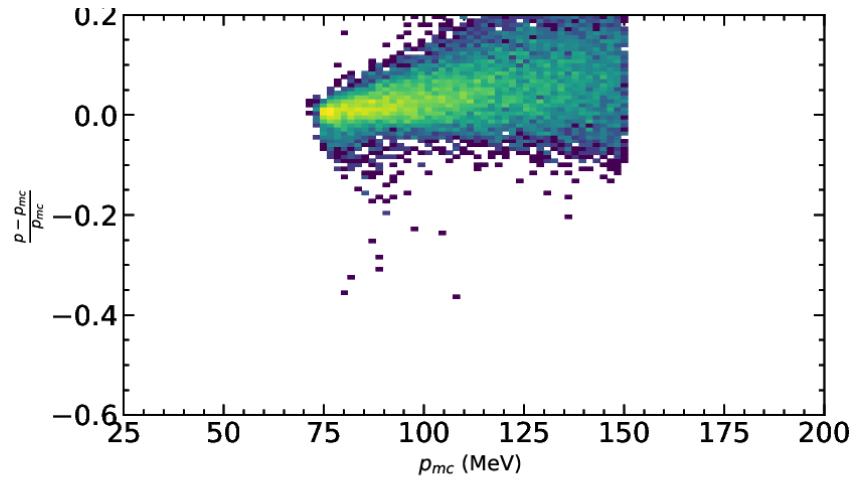
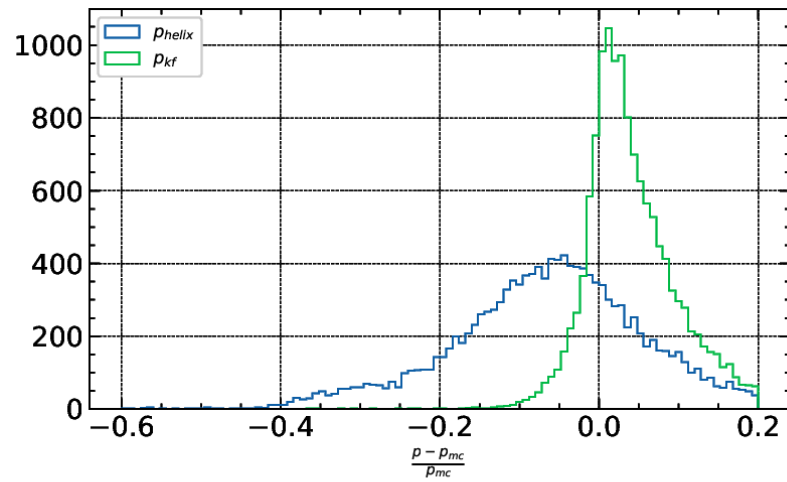
Final track





## Result for simulation

Data from simulation : 70 to 150 MeV/c proton (full range)

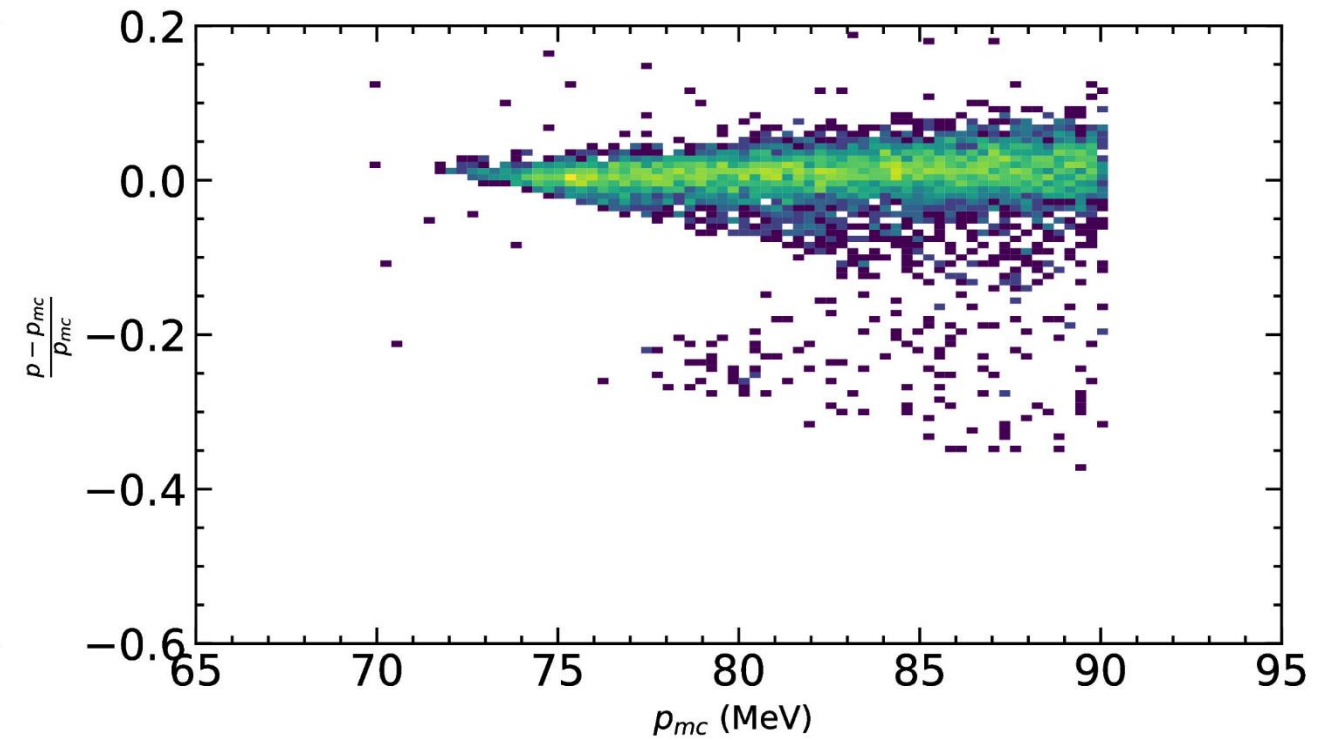
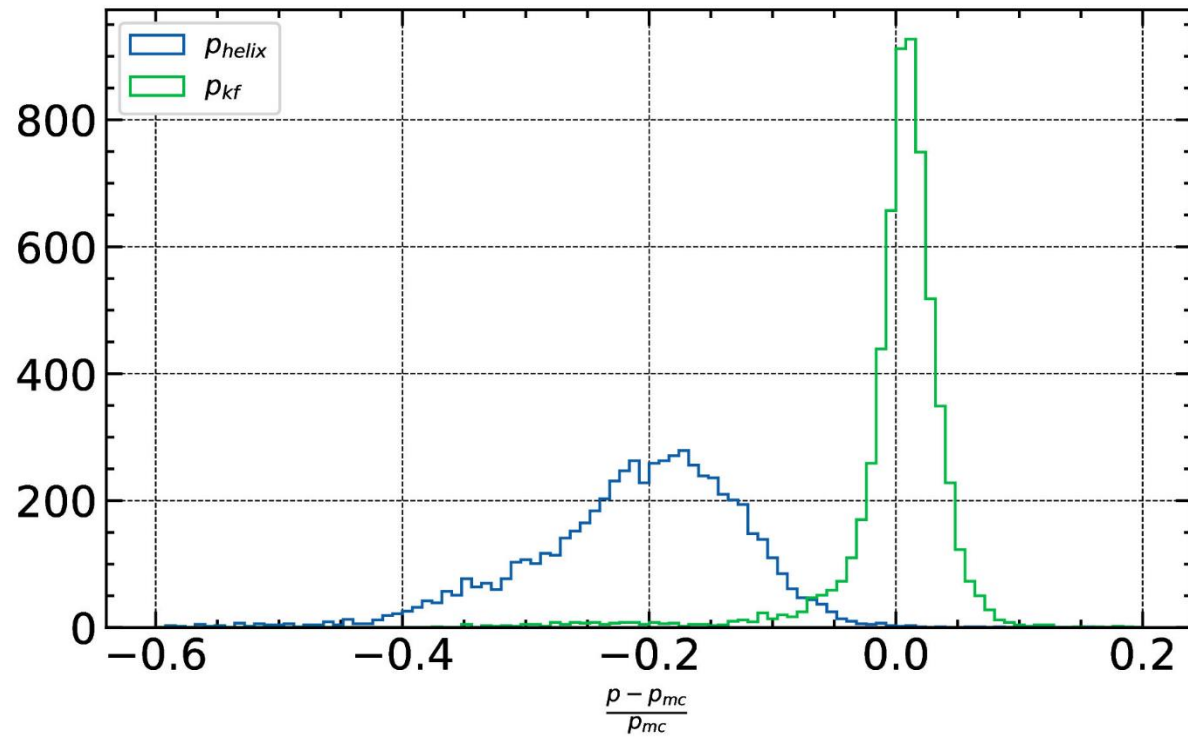


Work well on low momentum,  
still some issues with higher  
momentum.



## Result for simulation

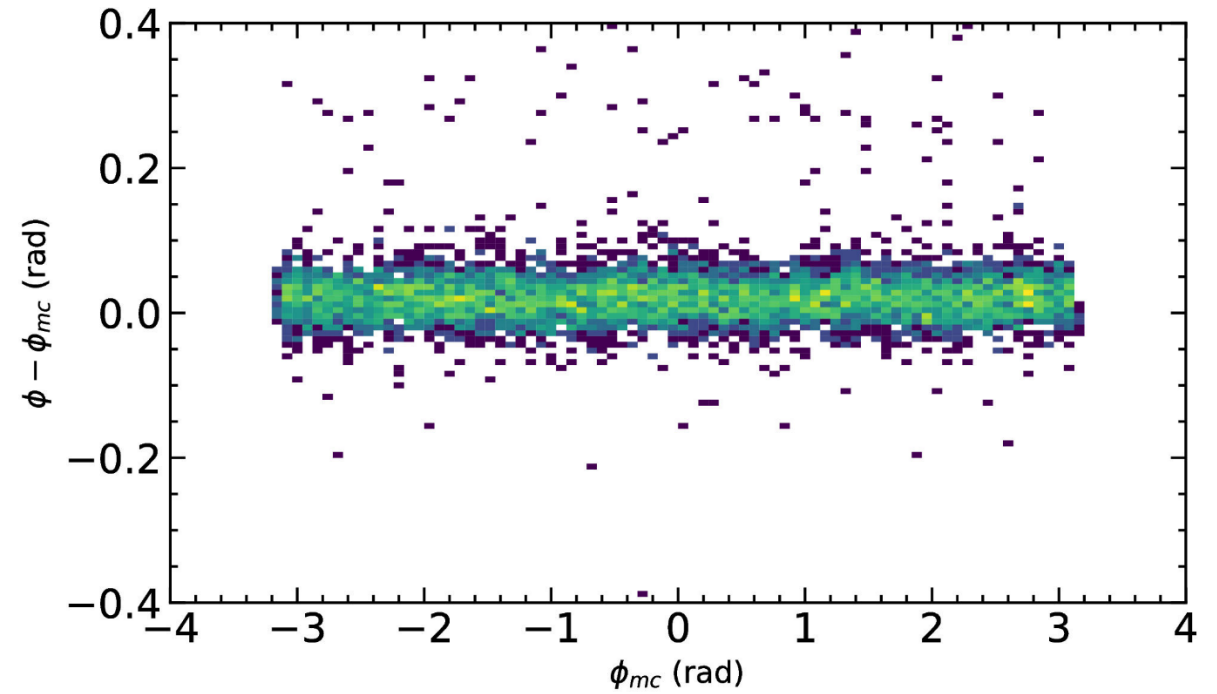
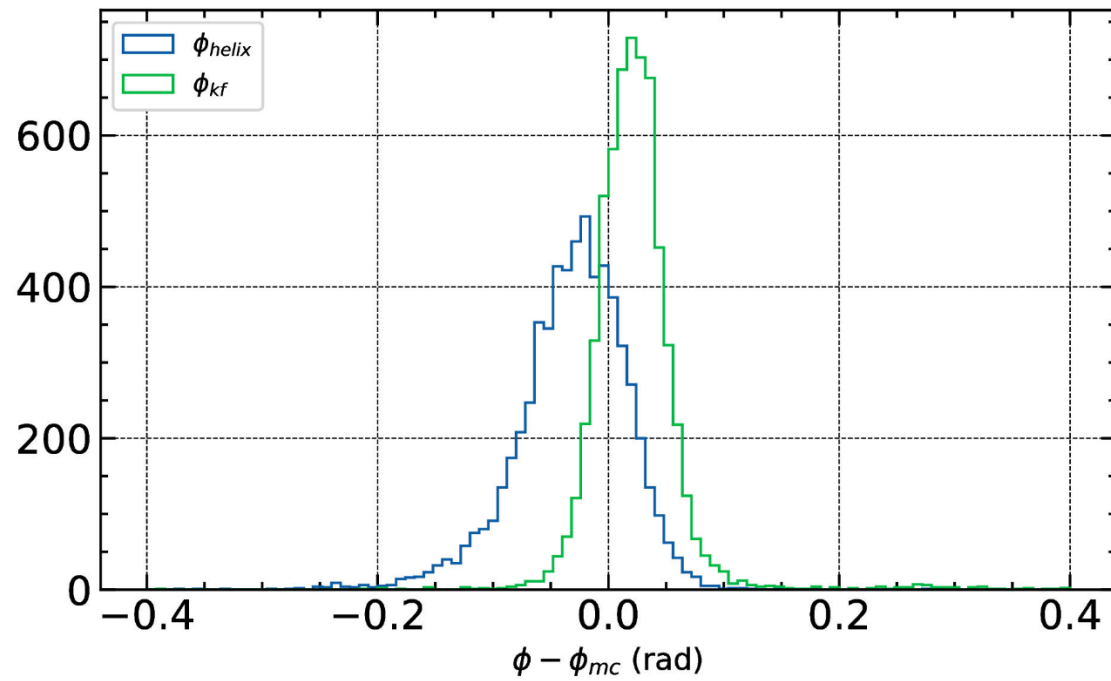
Data from simulation : 70 to 90 MeV/c proton





## Result for simulation

Data from simulation : 70 to 90 MeV/c proton





- Summary :
  - Kalman Filter can find the momentum for low-energy protons.
  - Use the Kalman Filter to compute a more precise  $dE_{dx}$ .
- Perspectives :
  - Test and fine-tune the parameters of the Kalman Filter on radiative elastic scattering data (need lower momentum to test).
  - Cut hit which are too far from the tracks.