# How the 20+ GeV SIDIS/TMD can encounter lattice QCD?

The Next Generation of 3D Imaging Jefferson Lab, Jul. 7-8, 2022

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#### What can lattice QCD calculate?

# Short answer: almost everything for leading-power TMDs ③



#### TMDs from global analyses

Semi-inclusive deep inelastic scattering:  $l + p \longrightarrow l + h(P_h) + X$ 

$$\frac{d\sigma^W}{dxdydz_h d^2 \mathbf{P}_{hT}} \sim \int d^2 \mathbf{b}_T \ e^{i\mathbf{b}_T \cdot \mathbf{P}_{hT}/z}$$

×
$$f_{i/p}(x, \mathbf{b}_T, Q, Q^2) D_{h/i}(z_h, \mathbf{b}_T, Q, Q^2)$$



Kang, Prokudin, Sun and Yuan, Phys. Rev. D 93, 014009 (2016)

$$f_{i/p}(x, \mathbf{b}_T, \mu, \zeta) = f_{i/p}^{\text{pert}}(x, b^*(b_T), \mu, \zeta)$$
Phys. Rev. D 93, 014009 (2016)
$$\times \left(\frac{\zeta}{Q_0^2}\right)^{g_K(b_T)/2} \xrightarrow{f_{i/p}^{\text{NP}}(x, b_T)} \xrightarrow{f_{i/p}^{\text{NP}}(x, b_T)} \text{Intrinsic TMD}$$

#### TMDs from global analyses

Results of the TMDs and Collins-Soper kernel:



Bacchetta, Bertone, Bissolotti, et al., MAP Collaboration, 2206.07598

## TMDs from lattice QCD

Quasi-TMDs in Large-Momentum Effective Theory:

(Naive) Quasi-TMD  $\checkmark$  NLO Perturbative matching  $\checkmark$  $\frac{\tilde{f}_{i/p}^{[s]}(x, \mathbf{b}_{T}, \mu, \tilde{P}^{z})}{\sqrt{S_{r}^{q}(b_{T}, \mu)}} = C(\mu, x\tilde{P}^{z}) \exp\left[\frac{1}{2}K(\mu, b_{T})\ln\frac{(2x\tilde{P}^{z})^{2}}{\zeta}\right] \times f_{i/p}^{[s]}(x, \mathbf{b}_{T}, \mu, \zeta) \left\{1 + \mathcal{O}\left[\frac{1}{(x\tilde{P}^{z}b_{T})^{2}}, \frac{\Lambda_{QCD}^{2}}{(x\tilde{P}^{z})^{2}}\right]\right\}$ See Ebert, Schindler, Stewart and YZ, JHEP 04 (2022) 178

and references therein.

**Reduced soft function**

Ji, Liu and Liu, NPB 955 (2020), PLB 811 (2020).

The Lorentz invariant approach by Musch, Hägler, Engelhardt, Negele and Schäfer has mainly been applied to calculate the ratios of TMD moments.

#### **TMDs from lattice QCD**

$$\frac{\tilde{f}_{i/p}^{[s]}(x, \mathbf{b}_T, \mu, \tilde{P}^z)}{\sqrt{S_r^q(b_T, \mu)}} = C(\mu, x \tilde{P}^z) \exp\left[\frac{1}{2}K(\mu, b_T) \ln\frac{(2x\tilde{P}^z)^2}{\zeta}\right] \times f_{i/p}^{[s]}(x, \mathbf{b}_T, \mu, \zeta) \left\{1 + \mathcal{O}\left[\frac{1}{(x\tilde{P}^z b_T)^2}, \frac{\Lambda_{\text{QCD}}^2}{(x\tilde{P}^z)^2}\right]\right\}$$

#### Matching coefficient:

- Independent of spin;
- No quark-gluon or flavor mixing, which makes gluon calculation much easier.

#### **TMDs from lattice QCD**

$$\frac{\tilde{f}_{i/p}^{[s]}(x, \mathbf{b}_T, \mu, \tilde{P}^z)}{\sqrt{S_r^q(b_T, \mu)}} = C(\mu, x \tilde{P}^z) \exp\left[\frac{1}{2}K(\mu, b_T)\ln\frac{(2x\tilde{P}^z)^2}{\zeta}\right] \times f_{i/p}^{[s]}(x, \mathbf{b}_T, \mu, \zeta) \left\{1 + \mathcal{O}\left[\frac{1}{(x\tilde{P}^z b_T)^2}, \frac{\Lambda_{\text{QCD}}^2}{(x\tilde{P}^z)^2}\right]\right\}$$

\* Collins-Soper kernel;

$$K(\mu, b_T) = \frac{d}{d \ln \tilde{P}^z} \ln \frac{\tilde{f}_{i/p}^{[s]}(x, \mathbf{b}_T, \mu, \tilde{P}^z)}{C(\mu, x \tilde{P}^z)}$$

\* Flavor separation;

$$\frac{f_{i/p}^{[s]}(x, \mathbf{b}_T)}{f_{j/p}^{[s']}(x, \mathbf{b}_T)} = \frac{\tilde{f}_{i/p}^{[s]}(x, \mathbf{b}_T)}{\tilde{f}_{j/p}^{[s']}(x, \mathbf{b}_T)}$$

\* Spin-dependence, e.g., Sivers function;

\* Full TMD kinematic dependence.

\* Twist-3 PDFs from small  $b_T$  expansion of TMDs.

YONG ZHAO, 07/08/2022

#### **Collins Soper kernel**



#### Soft function



#### **Important systematics:**

Matching beyond leading order, Fourier transform, power corrections at finite P<sup>z</sup>.

## Conclusion

- The method for calculating all the leading-power TMDs is complete;
- Lattice results for the Collins-Soper kernel and soft function are promising, but systematics need to be under control;
- What to expect in the short term?
  - Improved Collins-Soper kernel and soft function results;
  - Ratios of spin-dependent TMDs;
  - Ratios of quark TMDs of different flavors (challenge in the calculation of disconnected lattice diagrams);
- What to expect in the long term?
  - x and  $b_T$  dependence of intrinsic TMDs;
  - Twist-3 PDFs.

#### **Backup slides**

Data used by the MAP collaboration in 2206.07598



Bacchetta, Bertone, Bissolotti, et al., MAP Collaboration, 2206.07598

#### **Backup slides**



*i*, *j* (including spinor indices) remain intact



 $\propto \delta_{ij}$  Can mix with singlet channel and with gluons

$$b^2 = - b_z^2 - b_T^2 < b_T^2 \sim 1/\Lambda_{\rm QCD}^2$$

Hard particles cannot propagate that far!