Electroweak Physics with SoLID and a future positron beam at JLab

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Eur. Phys. J. A manuscript No. (will be inserted by the editor) https://arxiv.org/abs/2007.15081

An experimental program with high duty-cycle polarized and unpolarized positron beams at Jefferson Lab

Eur. Phys. J. A manuscript No. (will be inserted by the editor) https://arxiv.org/abs/2103.12555

Accessing weak neutral-current coupling g_{AA}^{eq} using positron and electron beams at Jefferson Lab

(A New Proposal to Jefferson Lab PAC-49)

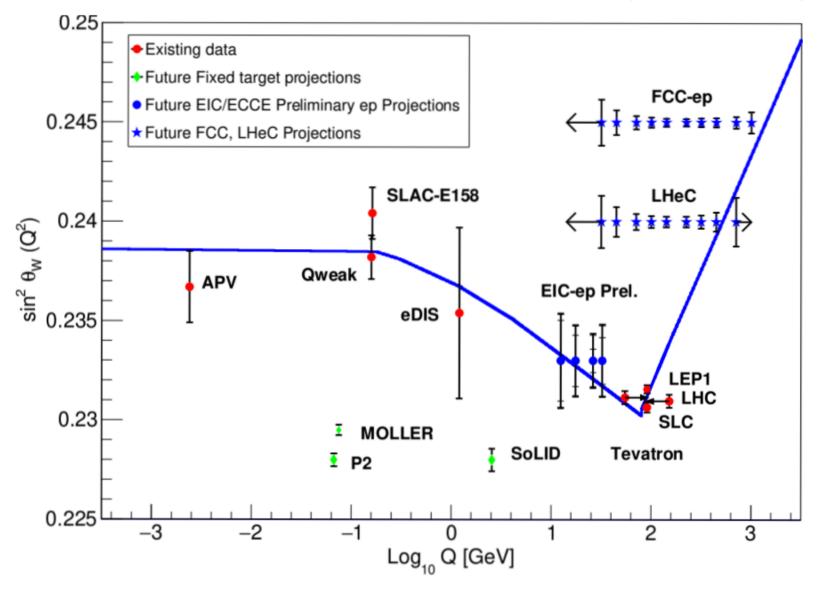
Measurement of the Asymmetry $A_d^{e^+e^-}$ between $e^+-^2{\rm H}$ and $e^--^2{\rm H}$ Deep Inelastic Scattering Using SoLID and PEPPo at JLab

May 24, 2021

-(pending approval from JLab Hall A and SoLID Collaborations)



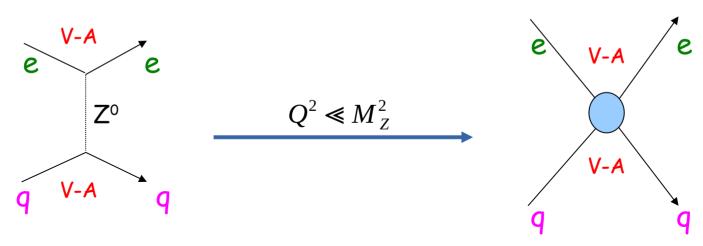
The Landscape of Electroweak Physics Study



EIC projections from arxiv.org/2204.07557 [hep-ph] LHeC projection (60GeV x 7 TeV, ~1000fb⁻¹) from EPJC 80 (2020) 9, 831 arxiv.org/2007.11799; FCC-ep projections: priv. comm. D. Britzger

(points with uncertainties comparable to or smaller than Qweak are shown, full range shown as arrows),

Neutral-Current Effective Couplings in (Low Energy) Electron Scattering



$$L_{NC}^{lq} = \frac{G_F}{\sqrt{2}} \sum_{q} \left[C_{0q} \overline{l} \gamma^{\mu} l \overline{q} \gamma_{\mu} q + C_{1q} \overline{e} \gamma^{\mu} \gamma_5 l \overline{q} \gamma_{\mu} q + C_{2q} \overline{e} \gamma^{\mu} e \overline{q} \gamma_{\mu} \gamma_5 q + C_{3q} \overline{l} \gamma^{\mu} \gamma_5 l \overline{q} \gamma_{\mu} \gamma_5 q \right]$$

VV (identical to γ) AV, VA (parity-violating)

$$C_{1u} = 2g_A^e g_V^u = -\frac{1}{2} + \frac{4}{3}\sin^2(\theta_W)$$
 $C_{2u} = 2g_V^e g_A^u = -\frac{1}{2} + 2\sin^2(\theta_W)$ $C_{3u} = -2g_A^e g_A^u = \frac{1}{2}$

$$C_{1d} = 2g_A^e g_V^d = \frac{1}{2} - \frac{2}{3}\sin^2(\theta_W)$$

$$C_{2u} = 2 g_V^e g_A^u = -\frac{1}{2} + 2 \sin^2(\theta_W)$$

$$C_{1d} = 2g_A^e g_V^d = \frac{1}{2} - \frac{2}{3}\sin^2(\theta_W) \qquad C_{2d} = 2g_V^e g_A^d = \frac{1}{2} - 2\sin^2(\theta_W) \qquad C_{3d} = -2g_A^e g_A^d = -\frac{1}{2}$$

$$C_{3u} = -2g_A^e g_A^u = \frac{1}{2}$$

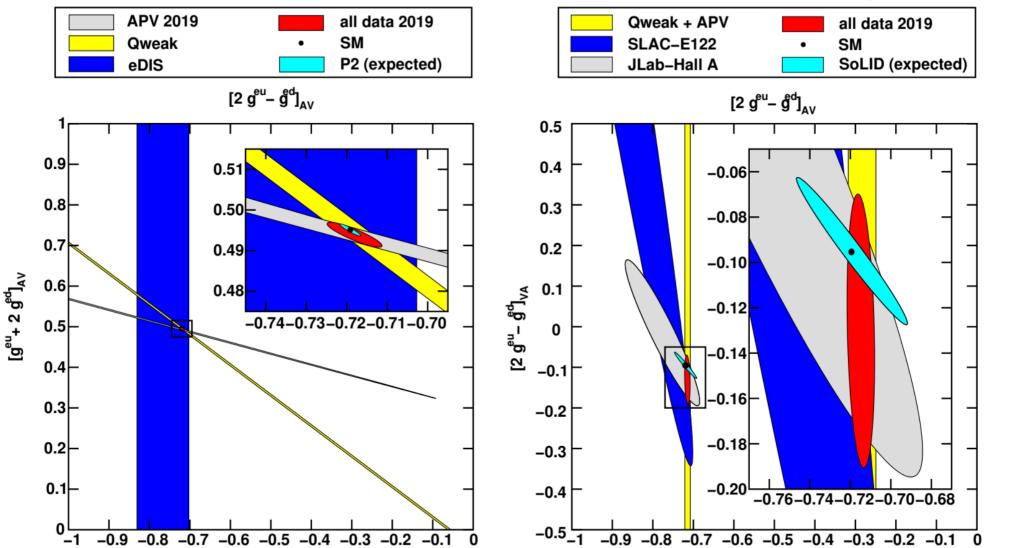
$$C_{3d} = -2 g_A^e g_A^d = -\frac{1}{2}$$

- A new set of notation $g_{AV,VA,AA}^{eq}$ introduced in 2013 <u>Erler&Su, Prog. Part. Nucl. Phys. 71</u>, 119 (2013)
- Example: In PVES, we can measure C_{1,2}

Current Knowledge on C_{1q}, C_{2q}

all are 68% C.L. limit

https://arxiv.org/abs/2103.12555



CERN for muon: $2C_{3u}^{\mu q} - C_{3d}^{\mu q} = 1.57 \pm 0.38$

Argento et al., PLB120B, 245 (1983)



In the Parton Model (low energy)

$$A_{RL}^{e^{\pm .}} = \frac{\sigma_{R}^{e^{\pm .}} - \sigma_{L}^{e^{\pm .}}}{\sigma_{R}^{e^{\pm .}} + \sigma_{L}^{e^{\pm .}}}$$

$$\left(A_{RL}^{e^{\pm \cdot}} = -A_{LR}^{e^{\pm \cdot}}\right)$$

$$A_{RL}^{e^{\pm \cdot}} = \frac{\sigma_R^{e^{\pm \cdot}} - \sigma_L^{e^{\pm \cdot}}}{\sigma_R^{e^{\pm \cdot}} + \sigma_L^{e^{\pm \cdot}}} \qquad A_d = |\lambda| (108 \, ppm) \, Q^2 \big[(2 \, C_{1u} - C_{1d}) + Y(y) (2 \, C_{2u} - C_{2d}) \, R_V(x) \big] \\ (A_{RL}^{e^{\pm \cdot}} = -A_{LR}^{e^{\pm \cdot}}) \qquad \text{beam polarization} \qquad Y(y) = \frac{1 - (1 - y)^2}{1 + (1 - y)^2} \qquad R_V(x) = \frac{u_V(x) + d_V(x)}{u(x) + \overline{u}(x) + \overline{d}(x)}$$

(indicates spin flip of quarks)

In the Parton Model (low energy)

$$A_{RL}^{e^{\pm .}} = \frac{\sigma_{R}^{e^{\pm .}} - \sigma_{L}^{e^{\pm .}}}{\sigma_{R}^{e^{\pm .}} + \sigma_{L}^{e^{\pm .}}}$$

$$\left(A_{RL}^{e^{\pm \cdot}} = -A_{LR}^{e^{\pm \cdot}}\right)$$

$$A_{RL}^{e^{+} e^{-}} = \frac{\sigma_{R}^{e^{+}} - \sigma_{L}^{e^{-}}}{\sigma_{R}^{e^{+}} + \sigma_{L}^{e^{-}}}$$

$$\left(A_{RL}^{e^+e^-} \neq -A_{LR}^{e^+e^-}\right)$$

$$A_{RR}^{e^{+} \cdot e^{-}} = \frac{\sigma_{R}^{e^{+}} - \sigma_{R}^{e^{-}}}{\sigma_{R}^{e^{+}} + \sigma_{R}^{e^{-}}}$$

$$\left(A_{RR}^{e^+e^-} \neq -A_{LL}^{e^+e^-}\right)$$

$$A_{unpol}^{e^{+}e^{-}} = \frac{\sigma^{e^{+}} - \sigma^{e^{-}}}{\sigma^{e^{+}} + \sigma^{e^{-}}}$$

$$A_{RL}^{e^{\pm}} = \frac{\sigma_R^{e^{\pm}} - \sigma_L^{e^{\pm}}}{\sigma_R^{e^{\pm}} + \sigma_L^{e^{\pm}}} \qquad A_d = |\lambda| (108 \, ppm) \, Q^2 \big[\big(2 \, C_{1u} - C_{1d} \big) + Y \, (y) \big(2 \, C_{2u} - C_{2d} \big) \, R_V(x) \big]$$

$$(A_{RL}^{e^{\pm}} = -A_{LR}^{e^{\pm}}) \qquad \text{beam polarization} \qquad Y(y) = \frac{1 - (1 - y)^2}{1 + (1 - y)^2} \qquad R_V(x) = \frac{u_V(x) + d_V(x)}{u(x) + \bar{d}(x)}$$

$$Y(y) = \frac{1 - (1 - y)^{2}}{1 + (1 - y)^{2}} \qquad R_{V}(x) = \frac{u_{V}(x) + d_{V}(x)}{u(x) + \overline{u}(x) + d(x) + \overline{d}(x)}$$

(indicates spin flip of quarks)

$$A_{RL}^{e^{+}\cdot e^{-}} = \frac{\sigma_{R}^{e^{+}} - \sigma_{L}^{e^{-}}}{\sigma_{R}^{e^{+}\cdot e^{-}}}$$

$$A_{RL,d}^{e^{+}\cdot e^{-}} = \frac{\sigma_{R}^{e^{+}} - \sigma_{L}^{e^{-}}}{\sigma_{R}^{e^{+}\cdot e^{-}}}$$

$$A_{RL,d}^{e^{+}\cdot e^{-}} = (108 \ ppm)Q^{2} Y(y) R_{V}(x) [|\lambda|(2C_{2u} - C_{2d}) - (2C_{3u} - C_{3d})]$$

$$A_{RL,d}^{e^{+}\cdot e^{-}} = A_{RL,d}^{e^{+}\cdot e^{-}}$$
(flie 13.1 for LD)

(flip $|\lambda|$ for LR)

"B" in CERN measurement

"B" in CERN r
$$A_{RR}^{e^+ \cdot e^-} = \frac{\sigma_R^{e^+} - \sigma_R^{e^-}}{\sigma_R^{e^+} + \sigma_R^{e^-}} \qquad A_{RR,d}^{e^+ \cdot e^-} = (108 \, ppm) \, Q^2 \Big[|\lambda| \big(2 \, C_{1u} - C_{1d} \big) - Y \, (y) \, R_V(x) \big(2 \, C_{3u} - C_{3d} \big) \Big]$$

(flip $|\lambda|$ for LL)

(no polarization needed!)

$$A_{unpol}^{e^{+}e^{-}} = \frac{\sigma^{e^{+}} - \sigma^{e^{-}}}{\sigma^{e^{+}} + \sigma^{e^{-}}} \qquad A_{d}^{e^{+}e^{-}} = -(108 \ ppm)Q^{2}Y(y)R_{V}(x)(2C_{3u} - C_{3d})$$

"direct" access to $2C_{3u}$ - C_{3d}

e⁺e⁻ for Structure Function Study

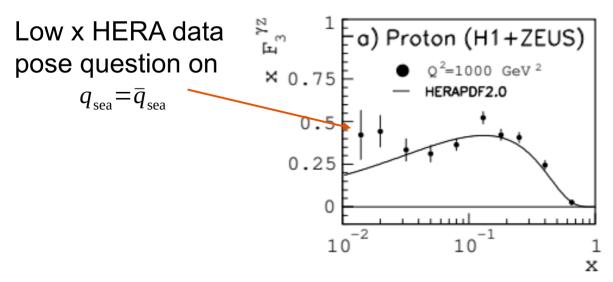
$$A_{\text{unpol}}^{e^{+}e^{-}} = \frac{G_{F}Q^{2}}{2\sqrt{2}\pi\alpha} \frac{g_{A}^{e}}{2} Y(y) \frac{F_{3}^{\gamma Z}}{F_{1}^{\gamma}}$$

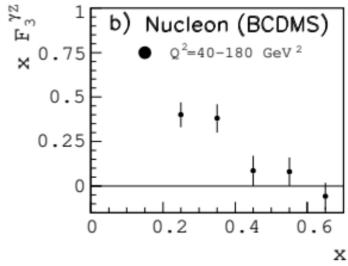
(in Apv, $F_3^{\gamma Z}$ is suppressed by g_V^e)

In the parton model:

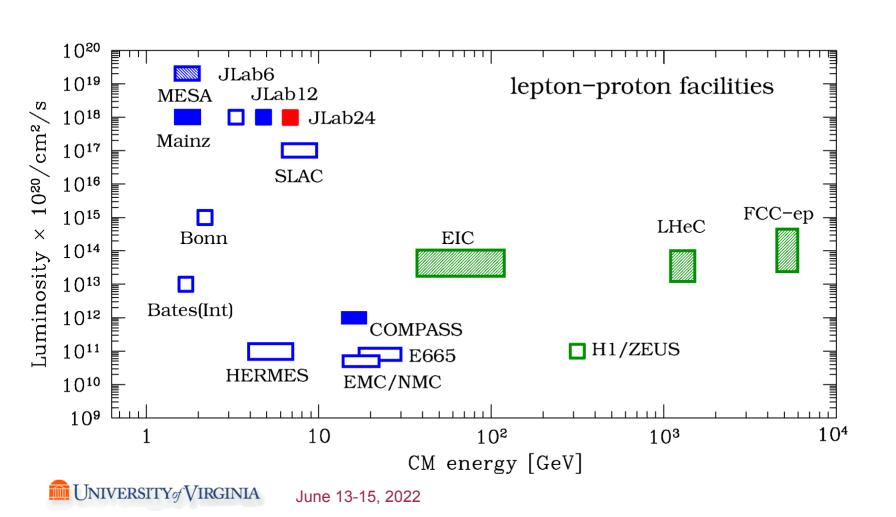
$$F_1^{\gamma}(x,Q^2) = 1/2 \sum Q_q^2 [q + \overline{q}]$$

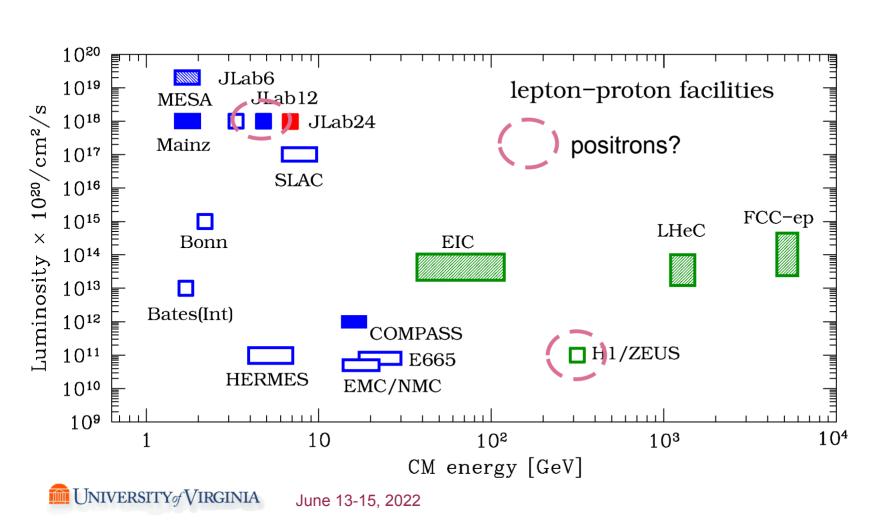
$$F_3^{\gamma Z}(x,Q^2) = 2 \sum g_A^q[q - \overline{q}]$$

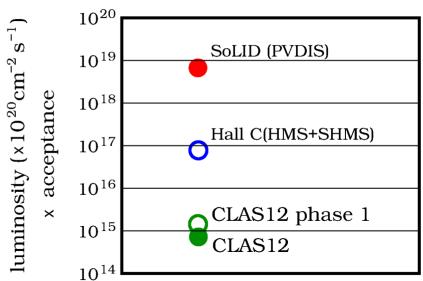


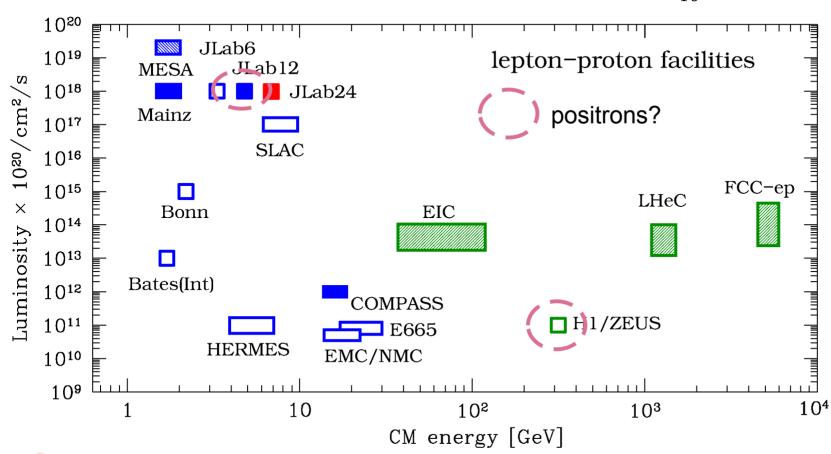


By measuring $A_{p,d}^{e+e^-}$ we can access $F_3^{\gamma Z}(x,Q^2)$





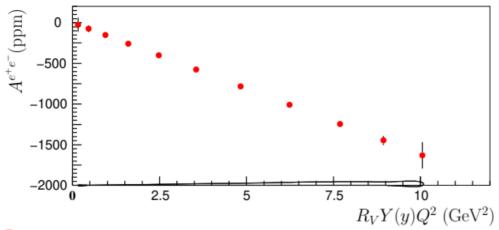




- ◆SoLID PVDIS configuration + 40cm LD2
- ◆PEPPo: 3uA unpolarized beam
- ◆Reverse magnet polarity of SoLID for positron detection

if we consider only statistics and assume A=0 at Q²=0:



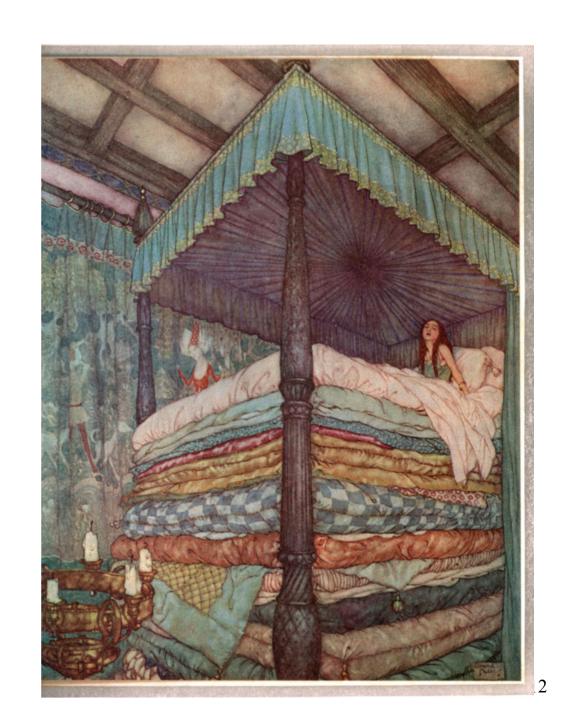


$$A_d^{e^{+}e^{-}} = -(108 \ ppm)Q^2 Y R_V (2C_{3u} - C_{3d})$$

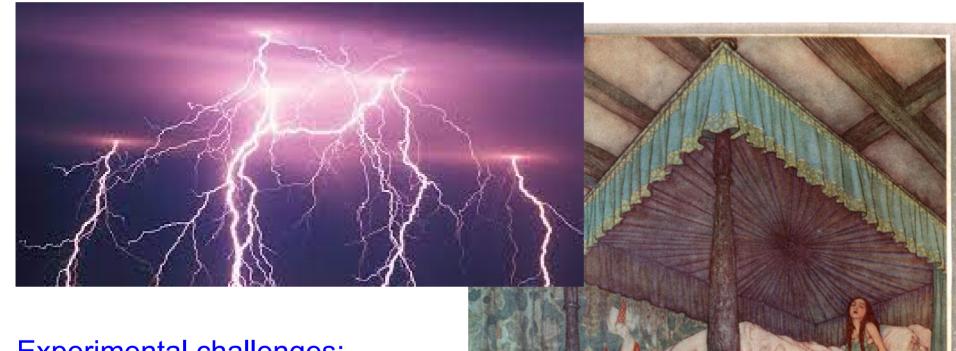
$$\to 1.5 \pm 0.007$$

Experimental challenges:

beam energy difference beam luminosity difference beam position difference charged pion background pair production background magnet and detector stability



Theoretical challenges: - higher-order QED corrections



Experimental challenges:

beam energy difference beam luminosity difference beam position difference charged pion background pair production background magnet and detector stability

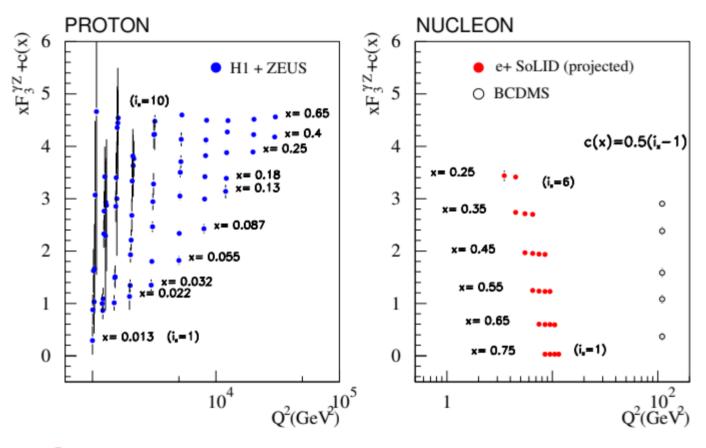


PR12-21-006 Lepton Charge Asymmetry

- 104 PAC days
- positron beam 3uA unpolarized
- beam control (1E-4 beam energy, ? beam position, "fast switch")

$$\Delta (2C_{3u} - C_{3d})_{\text{total}} = \pm 0.053 (\exp) \pm 0.009 (1\% \text{ QED})$$

+0.000 - 0.035 (HT, CJ15) $\approx \pm 0.060$



PAC49 report:

Issues: The PAC is pleased to see such an interesting and far-reaching proposal.
.....At the same time, the requirements on the accelerator and theory are both daunting.

Summary: ... leads us to defer the proposal in its present form.

Summary

- By comparing e- vs. e+ DIS cross section, we can form lepton-charge asymmetry that is directly proportional to a new set of eq EW NC coupling: C_{3q} or g_{AA} .
- So far, challenges in both experimental and theoretical systematic effects are both daunting and require further work to make the measurement compelling.
- A phased-approach can also be considered: study TPE in DIS using 11
 GeV and EW physics with a 22 GeV (e+) beam.

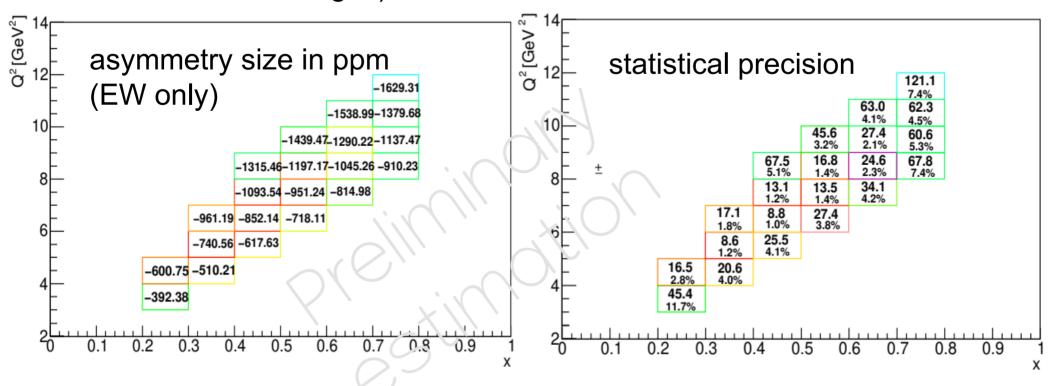
Backup



Summary of Challenges and Why They Exist?

- With a positron beam, the best physics impact comes from comparison between
 e+ and e- scattering, rather than measuring the same observable (e.g. Apv) as
 electrons
- If positron vs. electron comparison is our goal, then all systematic effect related to the beam need to be controlled to high precision
- Frequent ("weekly") and fast switch between e+ and e- beams is required to control differences in beam and run conditions → impact on positron beam design.
- Measurements where signal is tiny (EW physics) will be extremely difficult
- Particle background effects on the detector, trigger, and DAQ system.
- There is no well established calculation for TPE (QED NLO) in DIS. All previous (SLAC) data indicated zero but with poor precision;
- HERA data provided only slight constraint on QED NLO in DIS
- We could consider a "phase" approach: study DIS TPE first, then EW physics

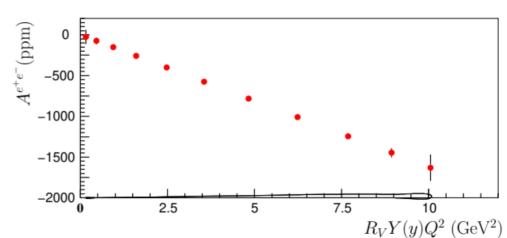
What can we do with 80 days of 3uA beam on a 40cm LD2 target? (in absence of all challenges):



if we consider only statistics and assume A=0 at Q²=0:

$$A_d^{e^+ \cdot e^-} = -(108 \ ppm)Q^2 Y R_V (2C_{3u} - C_{3d})$$

$$\rightarrow 1.5 \pm 0.007$$

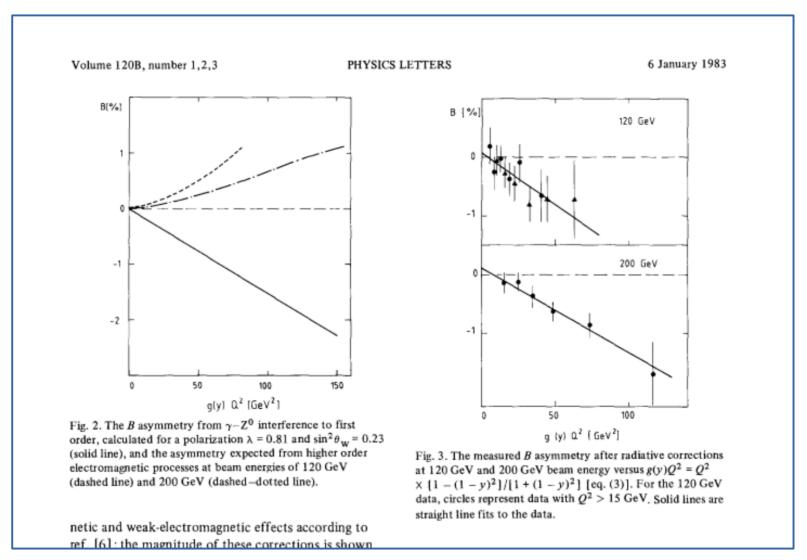




Past Experiment – BCDMS

1983 CERN, using polarized μ + vs. μ - beams:

$$2C_{3u}^{\mu q}-C_{3d}^{\mu q}=1.57\pm0.38$$



a measurement for the electron is highly desired



Experimental Challenges

QED higher order (scaled by $1/5) \rightarrow \Delta A_{OED}$

