The GPD program with CLAS12: highlights and perspectives

CLAS collaboration

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A set of distributions encoding the nucleon structure



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2 / 20

The generalized parton distributions and the nucleon

At leading twist there are 8 GPDs:

- 4 chiral-even GPDs: H, E, \tilde{H} and \tilde{E} .
- 4 chiral-odd GPDs: H_T , E_T , \tilde{H}_T and \tilde{E}_T .

Using the GPDs, we can determine the total angular momentum of quarks in the nucleon.

$$\int_{-1}^{1} x \left[H^f(x,\xi,0) + E^f(x,\xi,0) \right] dx = J^f \qquad \forall \xi \ .$$

By Fourier transform of the GPD H at $\xi=0$ (need extrapolation), we obtain the distribution in the transverse plane of the partons as a function of their longitudinal momentum.



Several GPDs but a universality principle

GPDs are universal: the same GPDs parameterize DVCS, DVMP, TCS, DDVCS,...





Deep virtual meson production has an additionnal non-perturbative part: The meson. Although there is an additionnal "unknown", it can conveniently be used for flavor separation or study of specific GPDs.

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GPD @ CLAS12

June 13th 2022 4 / 20

A spectrometer to get them all at once

The deep exclusive processes are:

- associated with a low cross section => 10^{35} s⁻¹cm⁻².
- diversifed with respect to their final state => RICH, LTCC, ToF.
- ... exclusive => Large acceptance with Central and Forward detectors.

CLAS12 spectrometer can simulatenously look at all these processes. Collecting data since 2018 with a 85%-longitudinally polarized beam on:

- unpolarized protons @ 6.6/7.5 and 10.6 GeV.
- unpolarized deuterium @ 10.4 GeV.
- ... and today with longitudinally polarized protons and neutrons @ 10.6 GeV.





Image: A matrix

All detectors required for exclusive processes

A 10.6-GeV electron beam scatters off a 5-cm LH₂ target. The beam is ${\sim}85\%$ longitudinally polarized. All CLAS12 detectors are necessary to reconstruct all particles of a DVCS final state:

- The electron is going through Cerenkov detector, drift chambers and electromagnetic calorimeter.
- The photon is either detected in a sampling calorimeter or a small PbWO₄-calorimeter close to the beamline.
- The recoil proton goes in the Silicon and Micromegas detector.



DVCS and GPDs

•
$$Q^2 = -q^2 = -(k - k')^2$$
.
• $x_B = \frac{Q^2}{2p \cdot q}$

- x average longitudinal momentum fraction carried by the active quark.
- $\xi \sim \frac{x_B}{2-x_B}$ the longitudinal momentum transfer.
- $t = (p' p)^2$ squared momentum transfer to the nucleon.

The GPDs enter the DVCS amplitude through a complex integral. This integral is called a *Compton form factor* (CFF).

$$\mathfrak{H}_{++}(\xi,t) = \int_{-1}^{1} H(x,\xi,t) \left(\frac{1}{\xi-x-i\epsilon} - \frac{1}{\xi+x-i\epsilon} \right) dx \; .$$

At leading order, $Im \mathfrak{H}(\xi, t) \propto (H(\xi, \xi, t) - H(-\xi, \xi, t)).$

7/20

- 3

Photon electroproduction

We use leptons beam to generate the γ^{\ast} in the initial state... not without consequences.

Indeed, experimentally we measure the cross section of the process $ep \to ep\gamma$ and not strictly $\gamma^* p \to \gamma p$.

Photon electroproduction and GPDs

The interference term allows to access the phase of the DVCS amplitude, *i.e* allows to isolate imaginary and real parts of CFFs.

$$\begin{split} c^{DVCS}_{0,UU} &\sim \quad 4(1-x_B) \left(\mathfrak{H}\mathfrak{H}^* + \widetilde{\mathfrak{H}}\widetilde{\mathfrak{H}}^* \right) \,, \\ c^{\mathfrak{I}}_{1,UU} &\sim \quad F_1 \; Re\mathfrak{H} + \xi(F_1 + F_2) \; Re\widetilde{\mathfrak{H}} \,, \\ s^{\mathfrak{I}}_{1,LU} &\sim \quad F_1 \; Im \mathfrak{H} + \xi(F_1 + F_2) \; Im \widetilde{\mathfrak{H}} \,, \\ s^{\mathfrak{I}}_{1,LU} &\sim \quad F_1 \; Im \widetilde{\mathfrak{H}} + \xi(F_1 + F_2) Im \widetilde{\mathfrak{H}} \,, \end{split}$$

In the present talk, we will be interested in deriving the beam-spin asymetry defined as:

$$A = \frac{\Delta^4 c}{d^4 \sigma}$$

Figure: Unpolarized and beam-helicity cross sections at $Q^2=2.3 \text{ GeV}^2$, $x_B=0.36$, t=-0.3 GeV² (Hall A).

GPD @ CLAS12

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DVCS: Unpolarized protons from LH2 at 10.6 GeV

•
$$A_{LU} \propto Im \left[F_1 \mathcal{H} + \xi \left(F_1 + F_2\right) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}\right]$$

- Performed by G. Christiaens, M. Defurne, D. Sokhan.
- Analysis note approved and article under Ad-Hoc review... Submission to PRL.
- Tested against ANNs fitted on 6 GeV data, new data significantly improved our knowledge of BSA.
- Comparisons with KM15 and VGG/GK models are performed as well.

DVCS: Unpolarized protons/neutrons from LD2 at 10.2 GeV

•
$$A_{LU} \propto Im \left[F_1 \mathcal{H} + \xi \left(F_1 + F_2\right) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}\right]$$

- Performed by A. Hobart and S. Niccolai.
- Analysis note under review.
- Measuring the proton allows to check for nuclear effect on proton structure.
- With $F_2 >> F_1$, the neutron is more sensitive to E than the proton.

Top: pDVCS BSA integrated over all phase space. Bottom: nDVCS BSA integrated over all phase space.

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June 13th 2022 11 / 20

DVCS: ongoing analysis on Unpolarized protons from LH2

- Cross sections must be measured to completely characterize DVCS.
- S. Lee and R. Milner extracting cross section with 10.6 GeV data, shown @ DNP22.
- Normalization is being carefully checked.
- Data collected @ 6.5 and 7.5 GeV to separate DVCS and interference terms.
- J. Tan and FX Girod currently finalizing the BSA analysis shown @ DNP22.

DVCS: Data taking on longitudinally polarized target

- $A_{UL} \propto Im \left[F_1 \tilde{\mathcal{H}} + \xi \left(F_1 + F_2 \right) \left(\mathcal{H} + x_B/2 \mathcal{E} \right) \xi \frac{t}{4M^2} F_2 \tilde{\mathcal{E}} \right]$
- First electrons sent for RGC on saturday evening to Hall B.
- Will continuously run until March 2023.
- Both NH₃ and ND₃ longitudinally polarized.
- NH₃ polarized up to 90% last week in the Hall.

Perspectives: DVCS with positron beam

- Conditionally approved by the PAC.
- The most straightforward approach to separate DVCS and I.

$$\sigma(\lambda, \pm e) = \left| \mathfrak{T}^{BH} \right|^2 + \left| \mathfrak{T}^{DVCS} \right|^2 \mp \mathfrak{I} ,$$

•
$$A_{UU}^C \propto Re\left[F_1\mathcal{H} + \xi (F_1 + F_2)\tilde{\mathcal{H}} - \frac{t}{4M^2}F_2\mathcal{E}\right]$$

• The predictions for the charge asymmetry differ by much from one model to another.

Figure: Bayesian reweighting with CLAS12 prediction. Dutrieux *et al*, EPJA 57, 8, 250.

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TCS: First measurement ever with CLAS12

- TCS is a conjuguate process of DVCS.
- Two observables were extracted.
- $A_{FB} = \frac{\sigma(\theta,\phi) \sigma(180 \theta,\phi + 180)}{\sigma(\theta,\phi) + \sigma(180 \theta,\phi + 180)}$

→ 0.6

0.5

- A_{FB} and $A_{\odot U}$ gives Re and Im of $F_1 \mathcal{H} + \xi (F_1 + F_2) \tilde{\mathcal{H}} \frac{t}{4M^2} F_2 \mathcal{E}$
- Published in PRL by P. Chatagnon and S. Niccolai (127, 262501).

-t (GeV²)

15 / 20

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Perspectives: Double DVCS @ CLAS12

- Probing the GPD outside the diagonal x = ξ unlike DVCS and TCS.
- DVCS and TCS limits of DVCS when Q^2 and Q'^2 at 0.
- But cross section much lower by a factor 300 than DVCS.
- LOI12-16-004 for a HL-CLAS12 (10³⁷).
- Predictions for 100 days run.

GPD @ CLAS12

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π^0 -electroproduction

The π^0 -electroproduction cross section can be written as the sum of structure functions corresponding to response to virtual photon polarization.

$$\frac{d^{4}\sigma}{dtd\phi dQ^{2}dx_{B}} = \frac{1}{2\pi}\Gamma_{\gamma^{*}}(Q^{2}, x_{B}, E_{e}) \Big[\frac{d\sigma_{T}}{dt} + \epsilon \frac{d\sigma_{L}}{dt} + \epsilon \frac{d\sigma_{TT}}{dt}\cos(2\phi) + \lambda_{e}\sqrt{2\epsilon(1-\epsilon)}\frac{d\sigma_{TL}}{dt}\sin(\phi) + \sqrt{2\epsilon(1+\epsilon)}\frac{d\sigma_{TL}}{dt}\cos(\phi)\Big],$$

Assuming the factorization for transversely polarized photon:

$$\frac{d\sigma_{T}}{dt} = \frac{4\pi\alpha}{2k'} \frac{\mu_{\pi}^{2}}{Q^{8}} \left[(1-\xi^{2}) |\mathcal{H}_{T}|^{2} - \frac{t'}{8m^{2}} |2\widetilde{\mathcal{H}}_{T} + \mathcal{E}_{T}|^{2} \right],$$
(2)
$$\frac{d\sigma_{TT}}{dt} = \frac{4\pi\alpha}{2k'} \frac{\mu_{\pi}^{2}}{Q^{8}} \frac{t'}{16m^{2}} |2\widetilde{\mathcal{H}}_{T} + \mathcal{E}_{T}|^{2},$$

π^{0} -electroproduction: BSA and cross sections

- With the 10.6 GeV beam, new data at large Q² but also Rosenbluth separation possible when compared to 6 GeV data.
- BSA analysis note under review submitted by A. Kim *et al.*
- Active work on cross-section extraction By R. Johnston shown @ DNP22.

Conclusion

- The GPD experimental program with CLAS12 is one of the richest ever, thanks to the large acceptance of the spectrometer and the high luminosity.
- Data was collected at multiple energies, on proton and neutron. And a longitudinally polarized target is in place as I speak.
- Several analyses under review, becoming as many publications in the months to come.
- Many ongoing analyses on DVMP in addition to π⁰, including φ-electroproduction by P. Moran on proton data and N. Ramasubramanian on deuterium data.
- First TCS ever measurements published!
- A positron beam would complete the set of DVCS observables collected by CLAS12, and unambiguously constrain the real part of CFF.
- HL-CLAS12 would give us the opportunity to measure double DVCS for the first time, and constrain the GPDs away from the diagonal x=ξ.

19 / 20

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