

# Measuring the Neutron Spin Asymmetry $A_1^n$ in the Valence Quark Region in Hall C at Jefferson Lab

#### Melanie Cardona (Rehfuss)

Temple University

On Behalf of the E12-06-110 Collaboration



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### Nucleon spin structure: current status





→ Today: large uncertainties! in  $\Delta G$   $\Delta \Sigma = \int (\Delta u + \Delta \bar{u} + \Delta d + \Delta \bar{d} + \Delta s + \Delta \bar{s}) dx \sim 30\%$   $\Delta G = \int dx \Delta g \sim 20\%, L_q \sim ??, L_g \sim ??$ Little known about quark OAM ( $L_q$ )

Quark spin seems to play a smaller role in the nucleon spin decomposition than predicted by the CQM, which expected  $\Delta\Sigma \sim 75\%$ ,  $L_q \sim 25\%$ 

LQCD & high-x physics can help! Due to the non-perturbative nature of QCD, making absolute predictions of nucleon spin structure is generally difficult, but ...

 $\rightarrow$  LQCD can compute  $L_q = J_q - \Delta \Sigma_q$ ,  $J_g$  (@ physical  $\pi$  mass!)





### $A_1^n$ @ high- x: a key observable for spin structure





- Free of sea effects ( $q\bar{q}$  pairs and hard gluons)
- Spin is assumed to be carried by the valence quarks

0.9

0.7E

0.6

0.5E

 $(\underline{b} + \underline{d})/(\underline{b}\Delta)$ 

0.3

→ A poorly-explored region due to low rates at high x (need high luminosity, Hall C's 12 GeV-era polarized <sup>3</sup>He target reached  $2x10^{36}$  cm<sup>-2</sup>s<sup>-1</sup>!)

• Which models will our data agree with? How much of a role does  $L_q$  play in forming the nucleon spin?





**Polarized** and sea quark PDFs for  $Q^2 = 10 \text{ GeV}^2$  from the NNPDFpol1.1 parameterization

See Nocera ER, et al. Nucl. Phys. B887:276 (2014).

### Accessing spin structure: polarized DIS cross sections

 $\nu, Q^2$ 

U: 
$$\frac{d^2\sigma}{d\Omega dE'} = \frac{\alpha^2}{4E^2 \sin^4\left(\frac{\theta}{2}\right)} \left(\frac{2}{M}F_1(x,Q^2)\sin^2\left(\frac{\theta}{2}\right) + \frac{1}{\nu}F_2(x,Q^2)\cos^2\frac{\theta}{2}\right)$$

P: 
$$\frac{d^2\sigma}{d\Omega dE'}(\downarrow \uparrow \uparrow \uparrow \uparrow) = \frac{4\alpha^2 E'}{MQ^2 \nu E} \left[ (E + E'\cos\theta)g_1(x,Q^2) - \frac{Q^2}{\nu}g_2(x,Q^2) \right] = \Delta\sigma$$

$$\frac{d^2\sigma}{d\Omega dE'}(\downarrow \Rightarrow -\uparrow \Rightarrow) = \frac{4\alpha^2 \sin\theta E'^2}{MQ^2 \nu^2 E} \left[\nu g_1(x,Q^2) + 2Eg_2(x,Q^2)\right] = \Delta \sigma_\perp$$

 $Q^2 = 4EE' \sin^2(\theta/2)$ 4-momentum transfer $\nu = E - E'$ Energy transfer $W = M^2 + 2M\nu - Q^2$ Final state hadronic mass $\theta$ Scattering angle $x = Q^2/2M\nu$ Quark fractional momentum

Quark Parton Model:

•  $F_1(x) = \frac{1}{2} \Sigma e_i^2 [q_i^{\uparrow}(x) + q_i^{\downarrow}(x)]$ where  $q_i(x) = q_i^{\uparrow}(x) + q_i^{\downarrow}(x)$  is the probability of finding a quark q of flavor i with momentum fraction x

• 
$$g_1(x) = \frac{1}{2} \Sigma e_i^2 [q_i^{\uparrow}(x) - q_i^{\downarrow}(x)]$$
  
where  $\Delta q_i(x) = q_i^{\uparrow}(x) - q_i^{\downarrow}(x)$  is the sum  
over the helicity distribution for a quark q of flavor  
*i* with momentum fraction x

#### Hadrons

nucleo

W

•  $g_2(x)$  describes the transverse spin structure of the nucleon, which vanishes in the QPM (quark-gluon correlations)

# What is $A_1$ ? the virtual photon-nucleon asymmetry



Our wide Q<sup>2</sup> range (over 10 GeV<sup>2</sup>) will allow for further study of A<sub>1</sub>'s Q<sup>2</sup> –dependence @ a given x value in the valence region

• We need a transverse and longitudinal component to reconstruct the asymmetry along the virtual photon direction:

$$A_{\parallel} = \frac{\sigma^{\downarrow\uparrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\downarrow\uparrow} + \sigma^{\uparrow\uparrow\uparrow}} \quad \text{and} \quad A_{\perp} = \frac{\sigma^{\downarrow\Rightarrow} - \sigma^{\uparrow\Rightarrow}}{\sigma^{\downarrow\Rightarrow} + \sigma^{\uparrow\Rightarrow}}$$
$$\rightarrow \quad A_{1} = \frac{A_{\parallel}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$$

- $\sigma^{\downarrow\uparrow}(\sigma^{\uparrow\uparrow})$  is the cross section for a **longitudinally polarized target** with the electron spin aligned antiparallel (parallel) to the target spin
- $\sigma^{\downarrow\Rightarrow}(\sigma^{\uparrow\Rightarrow})$  is the cross section for a **transversely polarized target** with the electron spin aligned antiparallel (parallel) to the beam direction
- $\eta, \xi$ , and *d* are kinematic factors, and *D* depends on the ratio of the longitudinal and transverse virtual-photon absorption cross sections  $R = \sigma_L / \sigma_T$

# **Experimental Setup**

#### **Spectrometers:**

- High Momentum Spectrometer (HMS)
- Super High Momentum Spectrometer (SHMS)

#### **Electron Beam:**

- 1-pass @ 2.2 GeV (elastic, Δ(1232))
- 5-pass @ 10.5 GeV (DIS)
- Beam polarization: 85%
   (<3% uncertainty according to Moller)</li>
- Circular beam raster with  $\sim 2.0$  mm radius
- < 50 ppm avg. charge asymmetry

#### **Polarized <sup>3</sup>He Target**

- <sup>3</sup>He production cell (40 cm)
- $\sim 50\%$  in-beam polarization
- 30 uA beam current
- 3% uncertainty in polarimetry







 $A_1^n$  production began on Jan. 12<sup>th</sup>, 2020 and ended on March 13<sup>th</sup>, 2020

#### Data Quality: Helicity-sorted charge, live-time, and beam trip cuts



#### Data Quality: Helicity-sorted charge, live-time, and beam trip cuts



correct for any biasing

#### Sign Convention for $A_{\parallel}$ : <sup>3</sup>He Elastic Asymmetries



#### Sign Convention for $A_{\perp}$ : <sup>3</sup>He $\Delta(1232)$ Asymmetries



#### SHMS & HMS Parallel RAW Asymmetries



#### SHMS & HMS Perpendicular RAW Asymmetries



# Summary

#### E12-06-110 is a high-impact experiment on nucleon spin-structure

- The measurements of  $A_1^n$  at high x allow us to test fundamental predictions of the nucleon spin structure
- Combined with precision proton data, the high-precision neutron data will allow us to extract polarized-to-unpolarized quark PDF ratios distributions (Δq) and spin-flavor distributions (Δu/u) and (Δd/d)

The results will help answer questions like,
How much of a role does L<sub>q</sub> play?
(to what degree are the quarks' spin aligned parallel to the nucleon spin?)

**Thanks for listening!** 



Analysis Flowchart

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#### PhD Candidates

#### Spokespeople













#### SHMS & HMS Asymmetries

➤ LT<sup>±</sup>, Q<sup>±</sup>, IHWP-corrected, Wien-Flip corrected (2/18/20, beginning w/ SHMS 10355 & HMS 3163)

> 
$$x_{BJ}$$
 bins = 0.35, 0.40, 0.45, 0.50, 0.55, 0.60, 0.65, 0.70, 0.75  
all  $\pm$  0.025

>  $x_{BI}$  bins = 0.80 (0.775, 0.83), 0.86 (0.83, 0.89), 0.92 (0.89, 0.95), 0.98 (0.95, 1.0)

$$A = \frac{\sum_{i} \left(\frac{A_{i}}{(\Delta A_{i})^{2}}\right)}{\sum_{i} \left(\frac{1}{(\Delta A_{i})^{2}}\right)} \qquad \Delta A = \sqrt{\frac{1}{\sum_{i} \left(\frac{1}{(\Delta A_{i})^{2}}\right)}} \qquad i \sim \operatorname{run} \#$$
$$\Delta A_{i} = \sqrt{\frac{1 - A_{i}^{2}}{N}} \qquad \Delta A_{phys,i} = \frac{\Delta A_{i}}{P_{b}P_{t}f_{N_{2}}} \qquad f_{N_{2}} = 0.92 \qquad \text{(for now)}$$

# Extracting $g_{1,2}/F_1 \& A_1, A_2$

$$\frac{\boldsymbol{g}_{1}^{^{3}H\boldsymbol{e}}}{\boldsymbol{F}_{1}^{^{3}H\boldsymbol{e}}} = \left(\frac{1}{\mathrm{d}'}\right) \left(\mathrm{A}_{\parallel} + \tan\left(\frac{\theta}{2}\right) A_{\perp}\right)$$

$$\frac{g_2^{^{3}He}}{F_1^{^{3}He}} = \left(\frac{y}{2d'}\right) \left(-A_{\parallel} + \left(\frac{E - E'\cos(\theta)}{E'\sin(\theta)}\right)A_{\perp}\right)$$

$$A_1 = \frac{1}{D(1+\eta\xi)} A_{\parallel} - \frac{\eta}{d(1+\eta\xi)} A_{\perp}$$
$$A_2 = \frac{\xi}{D(1+\eta\xi)} A_{\parallel} + \frac{1}{d(1+\eta\xi)} A_{\perp}$$

 $A_{\parallel} \& A_{\perp}$  are the electron **physics** double-spin asymmetries

#### Electron Beam Energy E = 10.38 GeV (fixed)

$$D = \frac{E - \epsilon E'}{E(1 + \epsilon R)}$$

$$\epsilon = \frac{1}{1 + 2\left(1 + \frac{\nu^2}{Q^2}\right) \tan^2\left(\frac{\theta}{2}\right)}$$

$$\eta = \frac{\epsilon \sqrt{Q^2}}{E - E'\epsilon} \quad \xi = \eta(1 + \epsilon)/2\epsilon$$

$$\nu = E - E' \qquad y = \nu/E$$

$$d = D\sqrt{\frac{2\epsilon}{1 + \epsilon}} \qquad R(x, Q^2) = \frac{\sigma_L}{\sigma_T} (1998)$$

$$d' = \frac{(1 - \epsilon)(2 - y)}{y(1 + \epsilon R)}$$

### Nuclear Corrections & Quark Flavor Decomposition

•  $A_1^n$  is ultimately extracted from  $A_1^{^{3}He}$  as

$$A_1^n = \frac{F_2^{^{3}He} \left[ A_1^{^{3}He} - 2\left(\frac{F_2^p}{F_2^{^{3}He}}\right) P_p A_1^p \left(1 - \frac{0.014}{2P_p}\right) \right]}{P_n F_2^n (1 + \frac{0.056}{P_n})}$$

where  $P_n = 0.86^{+0.036}_{-0.02}$  and  $P_p = -0.028^{+0.009}_{-0.004}$  are the effective nucleon polarizations of the neutron and proton inside <sup>3</sup>He

• Combining neutron  $g_1/F_1$  data with measurements on the proton allows a flavor decomposition to separate the polarized-to-unpolarized-PDF ratios for up and down quarks:

# $A_1^p$ : 3-parameter, $Q^2$ – independent fit

 $A_1^p = (0.041 \pm 0.008) + (1.442 \pm 0.081)x + (-0.599 \pm 0.163)x^2$ 



Figure D.1: Our fit to world  $A_1^p$  data. The error bars on the data are the in-quadrature sum of their statistical and systematic errors. The yellow band indicates the error on the fit.

David Flay Thesis, Appendix D.2

### Kinematics and Beam Time

Kine	Spec	$E_b$	$E_p$	$\theta$	beam time	
		GeV	GeV	(0)	(hours)	
$\Delta(1232)$	SHMS	2.183	-1.79736	8.5	4.0	
Elastic	SHMS	2.183	-2.12860	8.5	8.0	

Kine	Spec	$E_b$	$E_p$	$\theta$	$ e^{-}$ production	$n \mid e^+ \text{ pro}$	d.   Tot. Time			
		GeV	GeV	(0)	(hours)	(hour	s) (hours)			
DIS										
3	HMS	10.5	2.90	30.0	88.0	0.0	88.0			
4	HMS	10.5	3.50	30.0	511.0	0.0	511.0			
B	SHM	S 10.5	3.40	30.0	511.0	4.0	515.0			
$\mathbf{C}$	SHM	$S \mid 10.5$	2.60	30.0	88.0	4.0	92.0			

- Longitudinal asymmetry of elastic scattering and transverse asymmetry of  $\Delta(1232)$  are used to check  $P_b P_t$  (sign convention)
- SHMS (B) and HMS (4) used to determine physics asymmetry of <sup>3</sup>He at high x, high  $Q^2$
- SHMS (C) and HMS (3) used to cover medium x with high  $Q^2$  to improve upon 6 GeV results

### **Projected Results**



# $A_1^n$ and $\Delta q/q$ predictions as $x \to 1$

SU(6) spin-flavor symmetric neutron wavefunction polarized in  $+\hat{z}$  direction, with diquark states  $(qq)_{S,S_z}$ :

$$|n^{\uparrow}\rangle = \frac{1}{\sqrt{2}} |d^{\uparrow}(ud)_{00}\rangle + \frac{1}{\sqrt{18}} |d^{\uparrow}(ud)_{10}\rangle - \frac{1}{3} |d^{\downarrow}(ud)_{11}\rangle - \frac{1}{3} |u^{\uparrow}(dd)_{10}\rangle - \frac{\sqrt{2}}{3} |u^{\downarrow}(dd)_{11}\rangle$$

$$\frac{\frac{F_{2}}^{n}}{F_{2}^{p}} \frac{d}{u} \frac{\Delta d}{\Delta u} \frac{\Delta u}{u} \frac{\Delta d}{d} A_{1}^{n} A_{1}^{p}}{\frac{\Delta d}{d} A_{1}^{n} A_{1}^{p}}$$

$$\frac{DSE-1 \ 0.49 \ 0.28 \ -0.11 \ 0.65 \ -0.26 \ 0.17 \ 0.59}{DSE-2 \ 0.41 \ 0.18 \ -0.07 \ 0.88 \ -0.33 \ 0.34 \ 0.88}$$

$$\frac{0_{|ud|}^{+}}{14 \ 0 \ 0 \ 1 \ 0 \ 1 \ 1}$$

$$\frac{1}{NJL} \ 0.43 \ 0.20 \ -0.06 \ 0.80 \ -0.25 \ 0.35 \ 0.77$$

$$\frac{SU(6) \ \frac{2}{3} \ \frac{1}{2} \ -\frac{1}{4} \ \frac{2}{3} \ -\frac{1}{3} \ 0 \ \frac{5}{9}}{CQM \ \frac{1}{4} \ 0 \ 0 \ 1 \ -\frac{1}{3} \ 1 \ 1 \ 1 \ 1}$$

Craig D. Roberts, et al: 10.1016/j.physletb.2013.09.038

#### Previous Results



## <sup>3</sup>He Performance Evolution



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# <sup>3</sup>He Target Polarization throughout $A_1^n$ Production Running



 $90^{\circ} \rightarrow \bot$  to  $e^{-}$  beam  $180^{\circ} \rightarrow \parallel$  to  $e^{-}$  beam

Plot courtesy of Junhao Chen, College of William & Mary

## Target Field/Spin Direction

<b>Target Holding Field Direction</b>	<sup>3</sup> He Spin Direction
+X Beam RIGHT (90°)	Beam LEFT
-X Beam LEFT (270°)	Beam RIGHT
+Z DOWNSTREAM (0°)	UPSTREAM
-Z UPSTREAM (180°)	<b>DOWNSTREAM</b>

The target was always pumped in the low-energy state (<sup>3</sup>He spin is **opposite of the holding field**) during data-taking

### Hall C Spectrometers & Detectors

**HMS detectors** 

**SHMS detectors** 



Spectrometer	Central Momentum (GeV/c)	Momentum Acceptance	Momentum Resolution	Scattering Angle	Solid Angle Acceptance (msr)	Horizontal Acceptance (mrad)	Vertical Acceptance (mrad)
HMS	0.5 - 7.5	(-9%, 9%)	0.02%	12.5° - 90°	8.1	±32	±85
SHMS	2.0 - 11.0	(-10%, 22%)	0.03% - 0.08%	5.5° - 40°	> 4.0	<u>±</u> 24	$\pm 40$