

Correlations in partonic and hadronic interactions

Duke University (NC), March 7-12, 2022

**Generalized parton
distributions at**

COMPASS

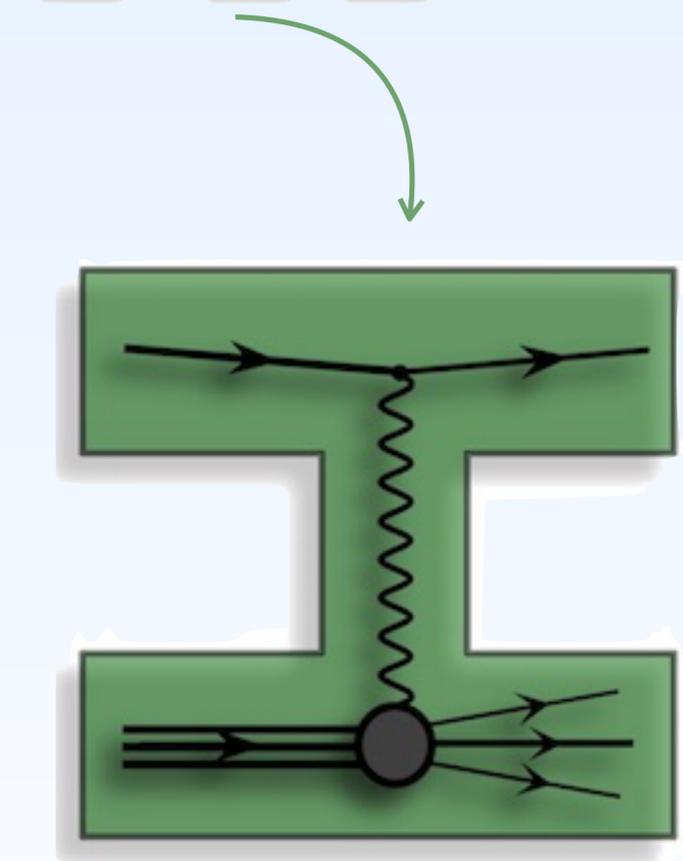
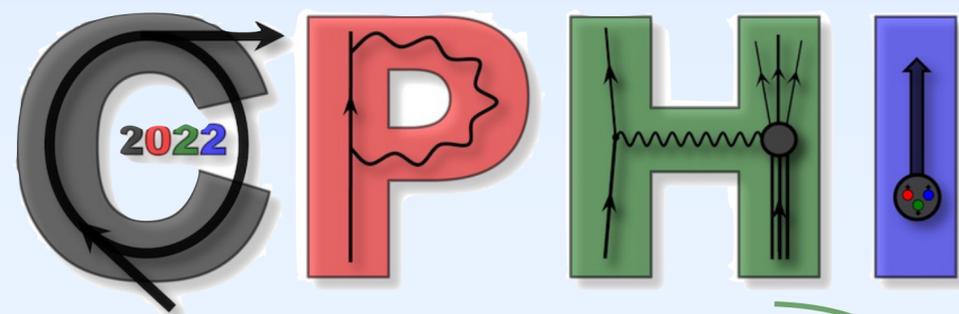


March 8, 2022

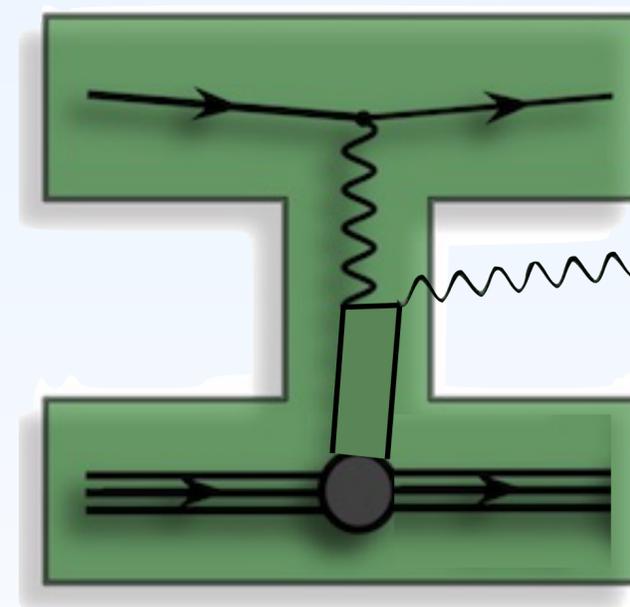
Caroline Riedl

for the COMPASS Collaboration





inclusive or semi-inclusive
lepton-proton deep-inelastic
scattering



deeply virtual Compton
scattering (DVCS)

scattered beam lepton

real photon

proton stays intact

exclusive
measurement

Physics questions

How do quarks & gluons, and their dynamics, make up proton spin?

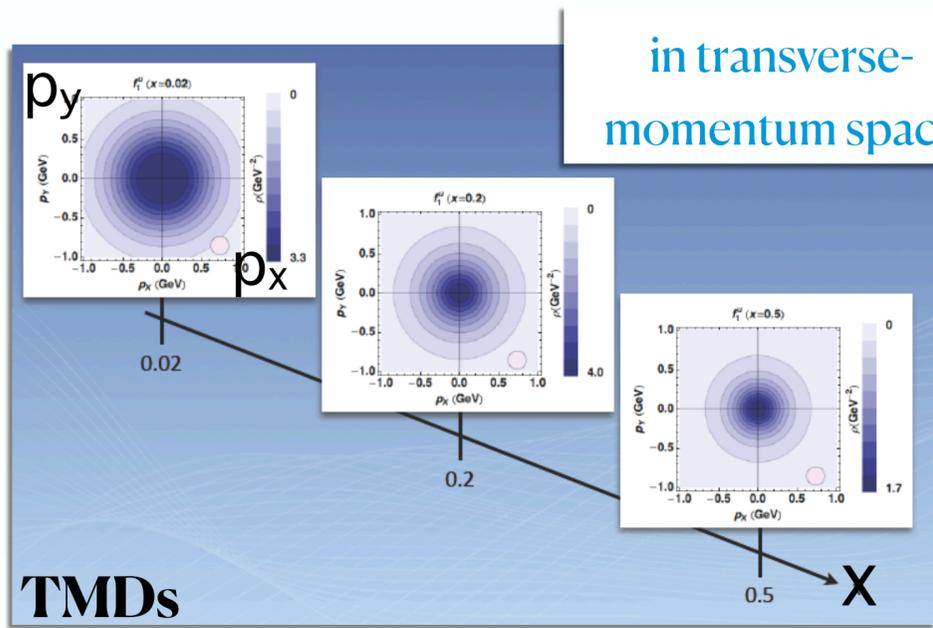
$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L_g$$

Spin puzzle

quark spin gluon spin orbital angular momenta of quarks and gluons

How is the proton spin correlated with the motion of quarks/gluons?

in transverse-momentum space



TMDs

X

Nucleon tomography

Transverse Momentum Dependent PDFs

Generalized Parton Distributions

TMDs
 $f(x, k_T)$

GPDs
 $H(x, \xi, t)$

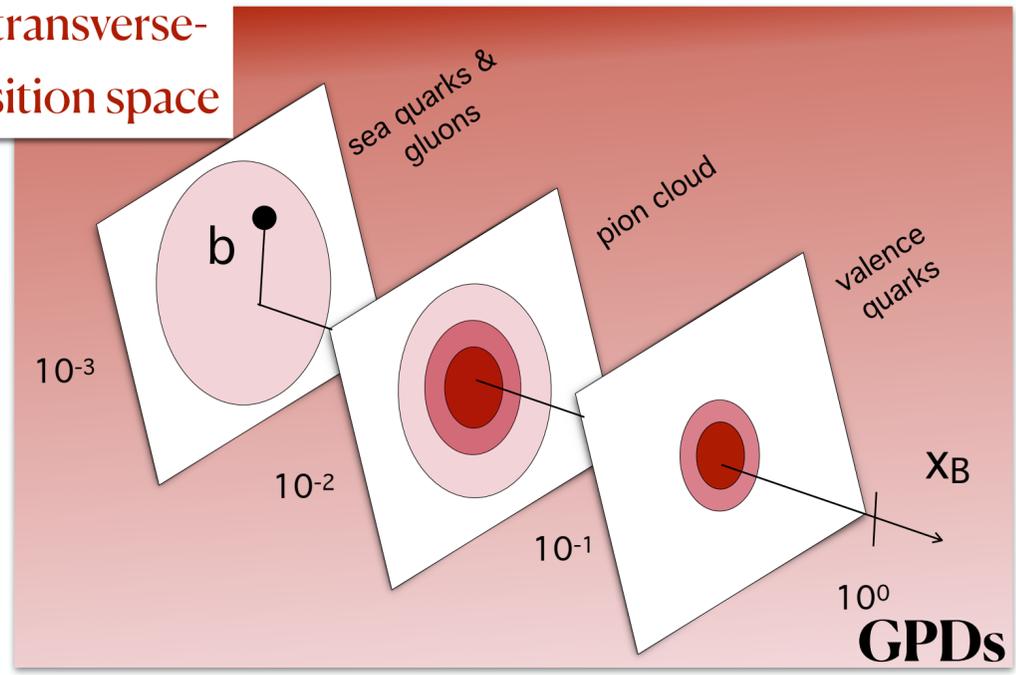
k_T -integration

$\xi=0, t=0$

(collinear) PDFs $q(x)$, 1D

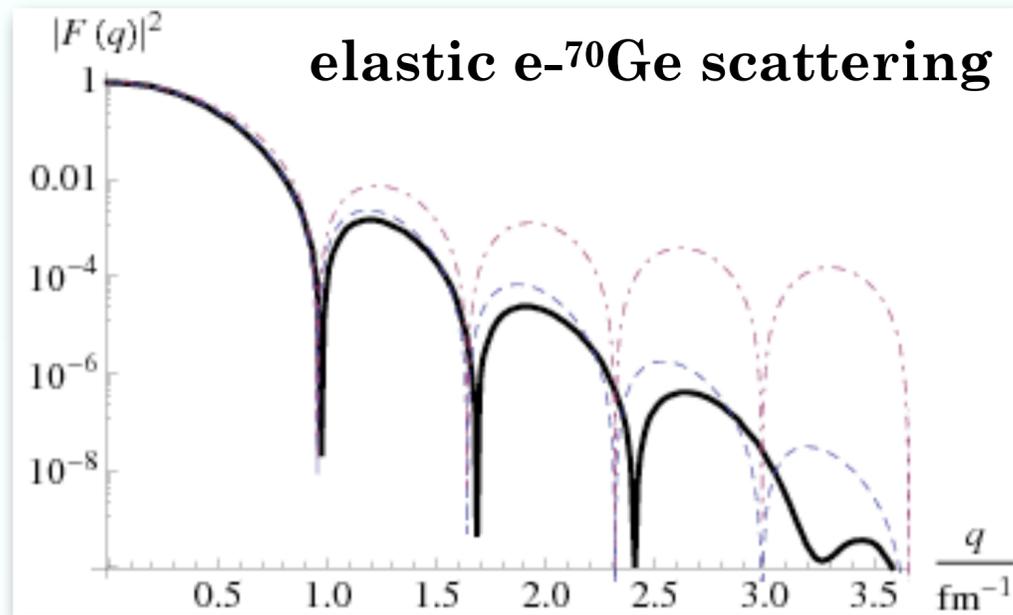
How does the proton spin influence the spatial distribution of partons?

in transverse-position space

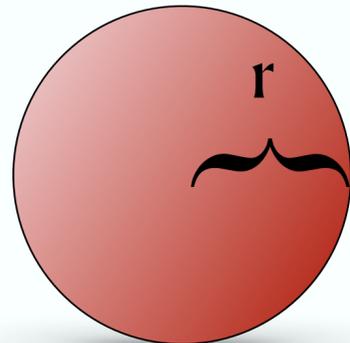


GPDs

Generalized Parton Distributions



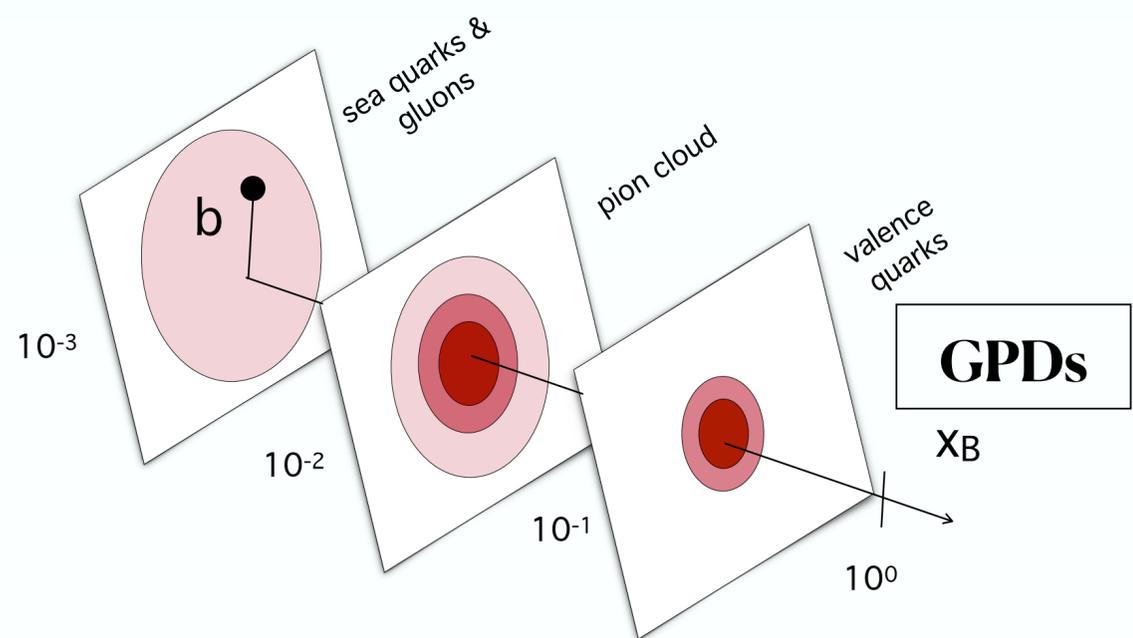
Form factors *from elastic scattering*



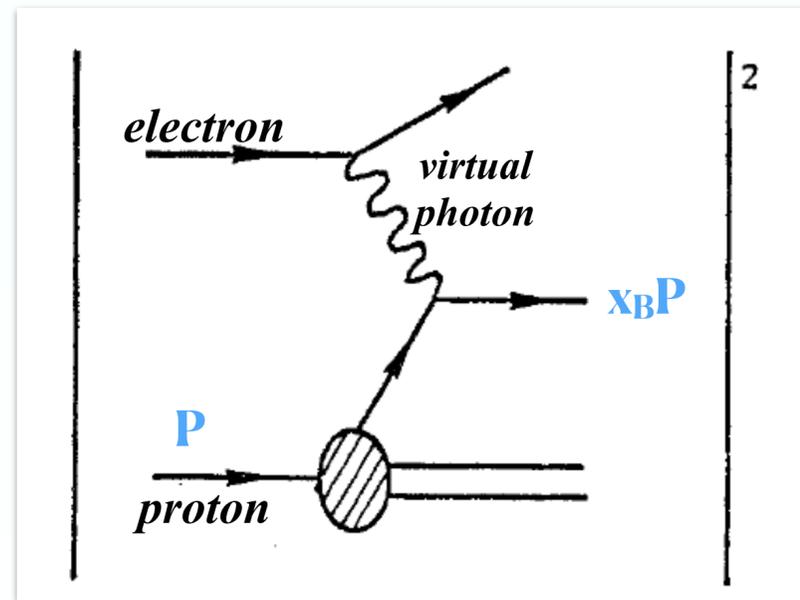
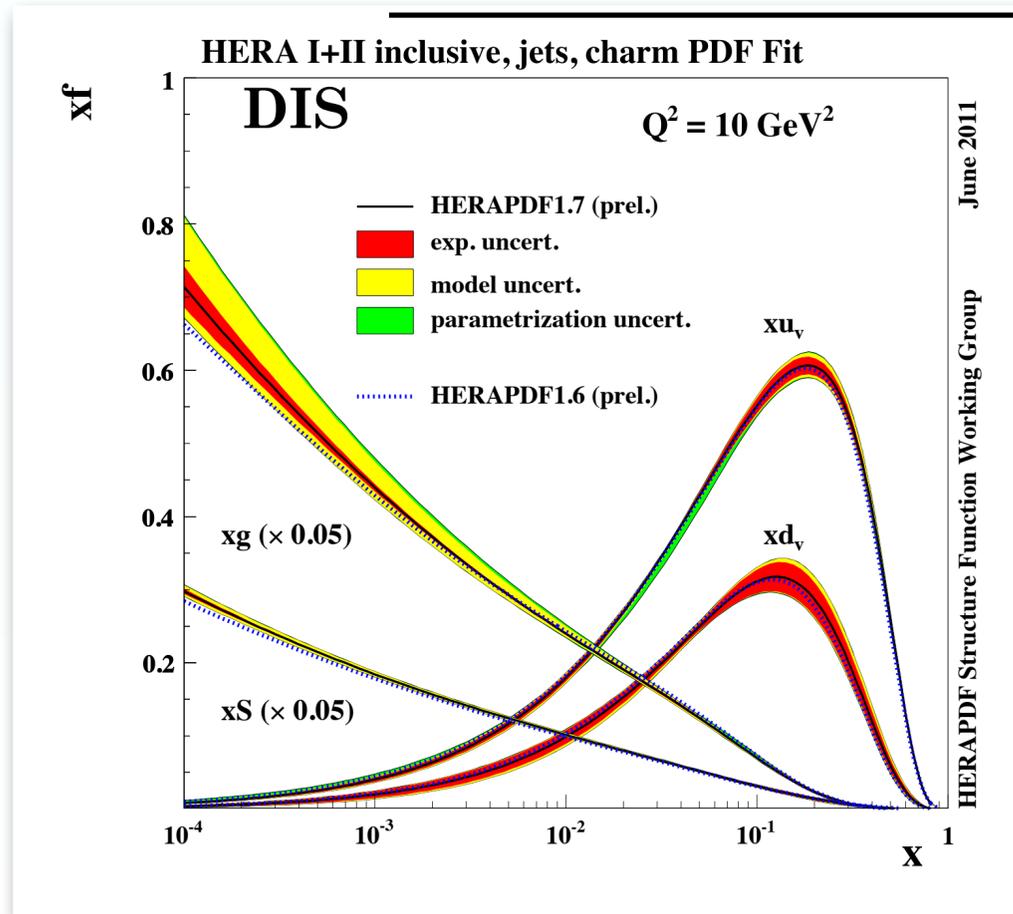
$$\int_{-1}^1 H^q(x, \xi, t) dx = F^q(t)$$

transverse parton positions

longitudinal parton momenta



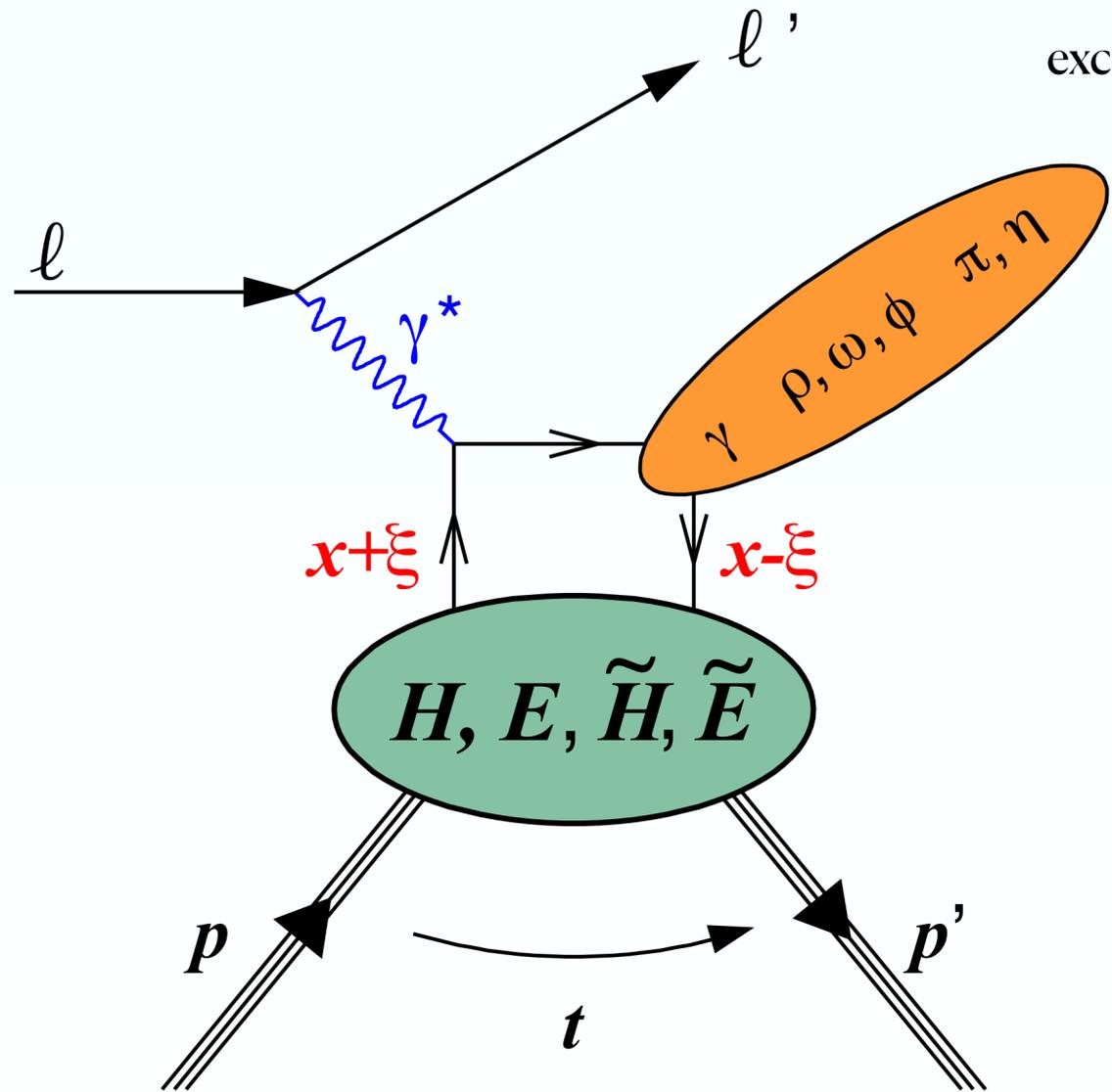
GPDs



PDFs *from deep-inelastic scattering*

$$H^q(x, 0, 0) = q(x)$$

Hard exclusive reactions



exclusive measurement = detection of entire final state
(or assumed to be known)



Standard channels to access
generalized parton
distributions are DVCS &
DVMP

$$\ell p \rightarrow \ell p \gamma$$

Deeply Virtual Compton
Scattering (DVCS)

$$\ell p \rightarrow \ell p M$$

Deeply Virtual
Meson Production (DVMP)

- x, ξ : longitudinal momentum fractions of probed quark
- **skewness** $\xi \approx x_B / (2 - x_B)$ in Bjorken limit
(Q^2 large & x_B, t fixed)
- **average momentum x : mute variable,**
not accessible in DVCS & DVMP. Is not x -Bjorken!
- **t : squared 4-momentum transfer to target**

Experimental access to GPDs

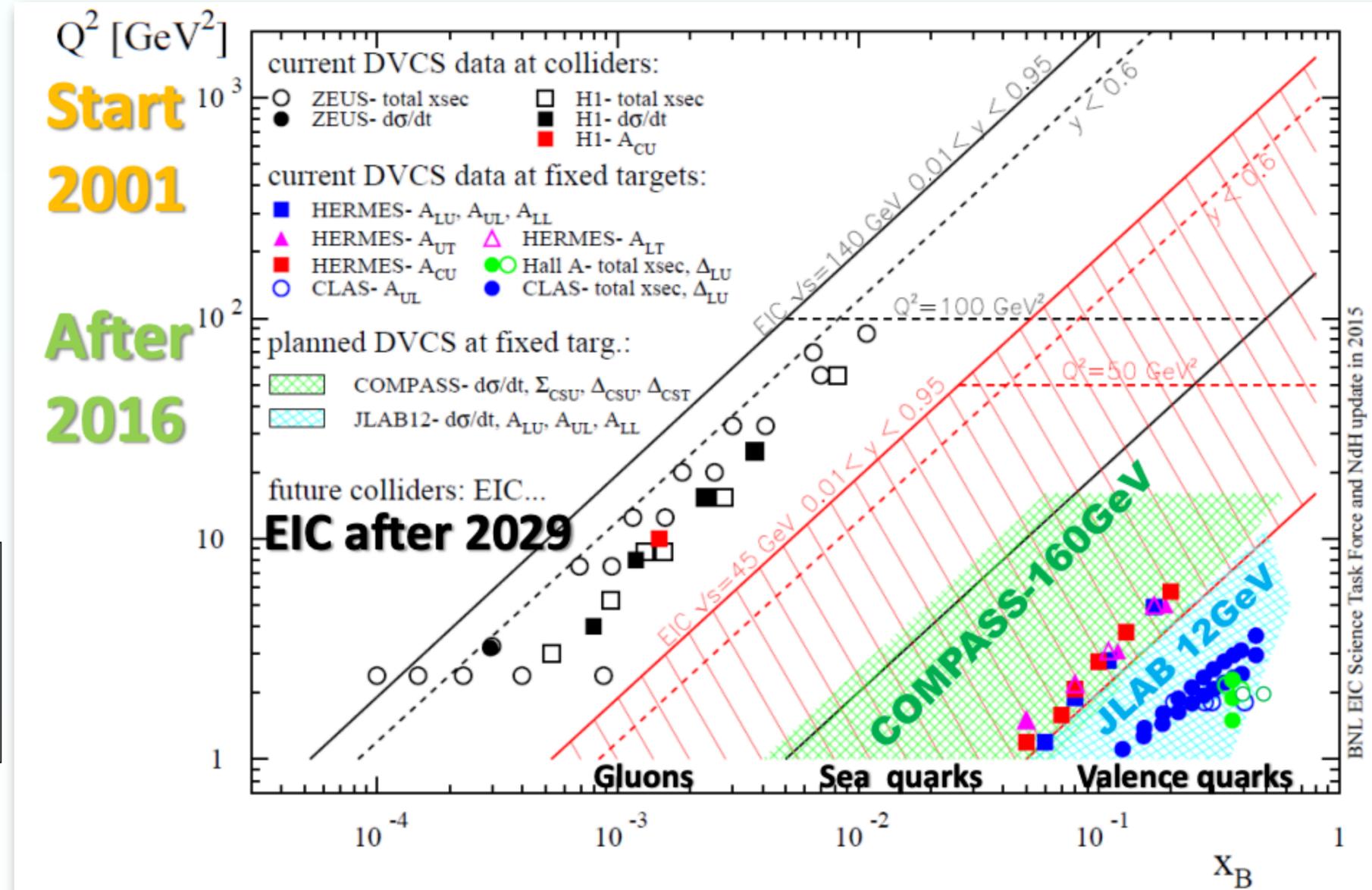
From HERMES & JLab-6 & HERA to COMPASS & JLab12 & RHIC to the EIC

- Different exclusive final-state particles allow to probe different GPDs
- **4 chiral-even GPDs (conserve quark helicity)**
- 4 chiral-odd GPDs (flip quark helicity)

| GPDs | flips nucleon helicity | conserves nucleon helicity |
|-----------------------------------|-------------------------------|-------------------------------|
| does not depend on quark helicity | E | H |
| depends on quark helicity | \tilde{E} | \tilde{H} |

| | | |
|----------|----------------------------------|----------------------------|
| $F_1(x)$ | $J^P=1^-$ vector mesons | $J^P=1^-$ photon (DVCS) |
| $g_1(x)$ | $J^P=0^-$ pseudoscalar mesons | |

forward limit $\xi \rightarrow 0, t \rightarrow 0$



@leading twist for a spin-1/2 target

DVCS: Compton form factors (CFFs) \leftrightarrow GPDs

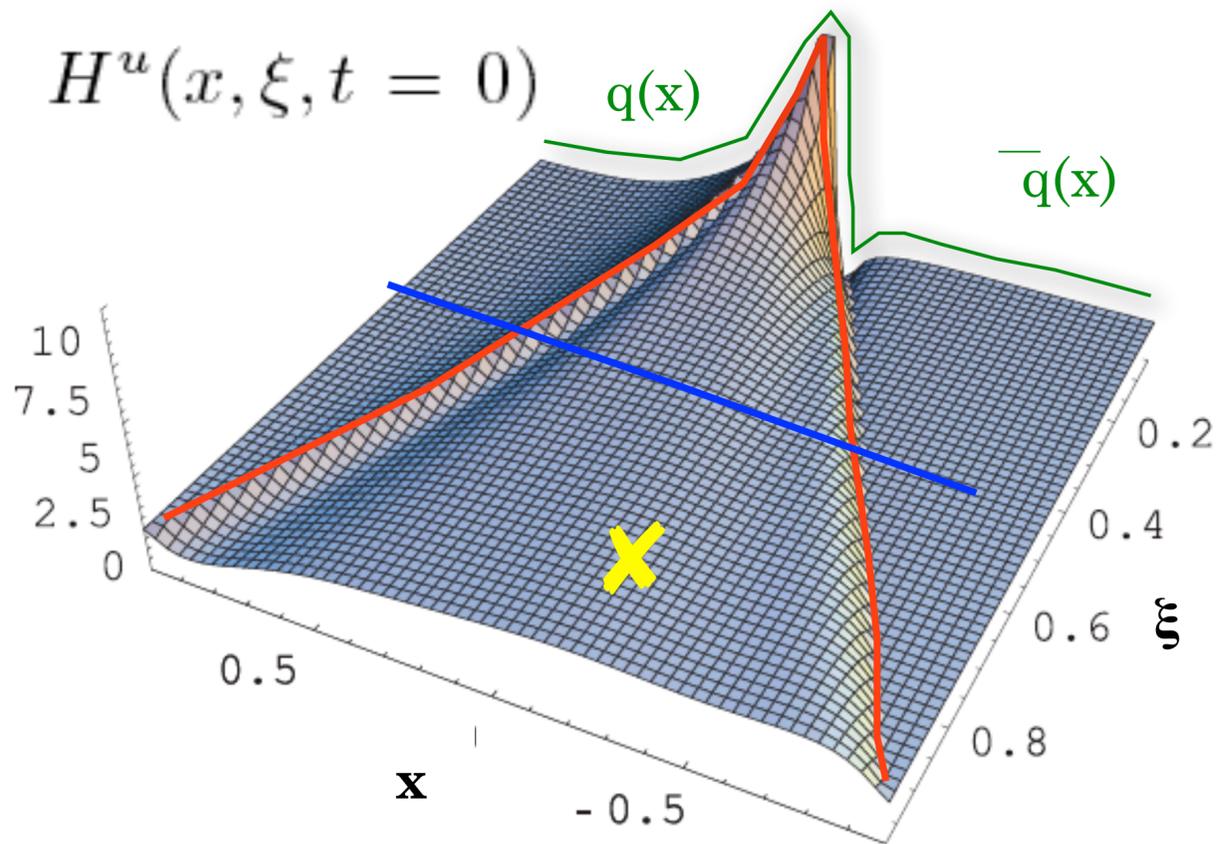
In DVCS, the experimentally accessed quantity is a (complex) **Compton Form Factor (CFF)**:

$$\mathcal{H}(\xi, t) = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i\pi H(\xi, \xi, t)$$

CFF **hard scattering kernel** \otimes **GPD**

$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx \boxed{C_q^\mp(\xi, x)} \boxed{F^q(x, \xi, t)}$$

assuming factorization $\dagger(Q^2 \text{ large \& } t \text{ small})$



Im(τ_{DVCS})
 $x = \xi$

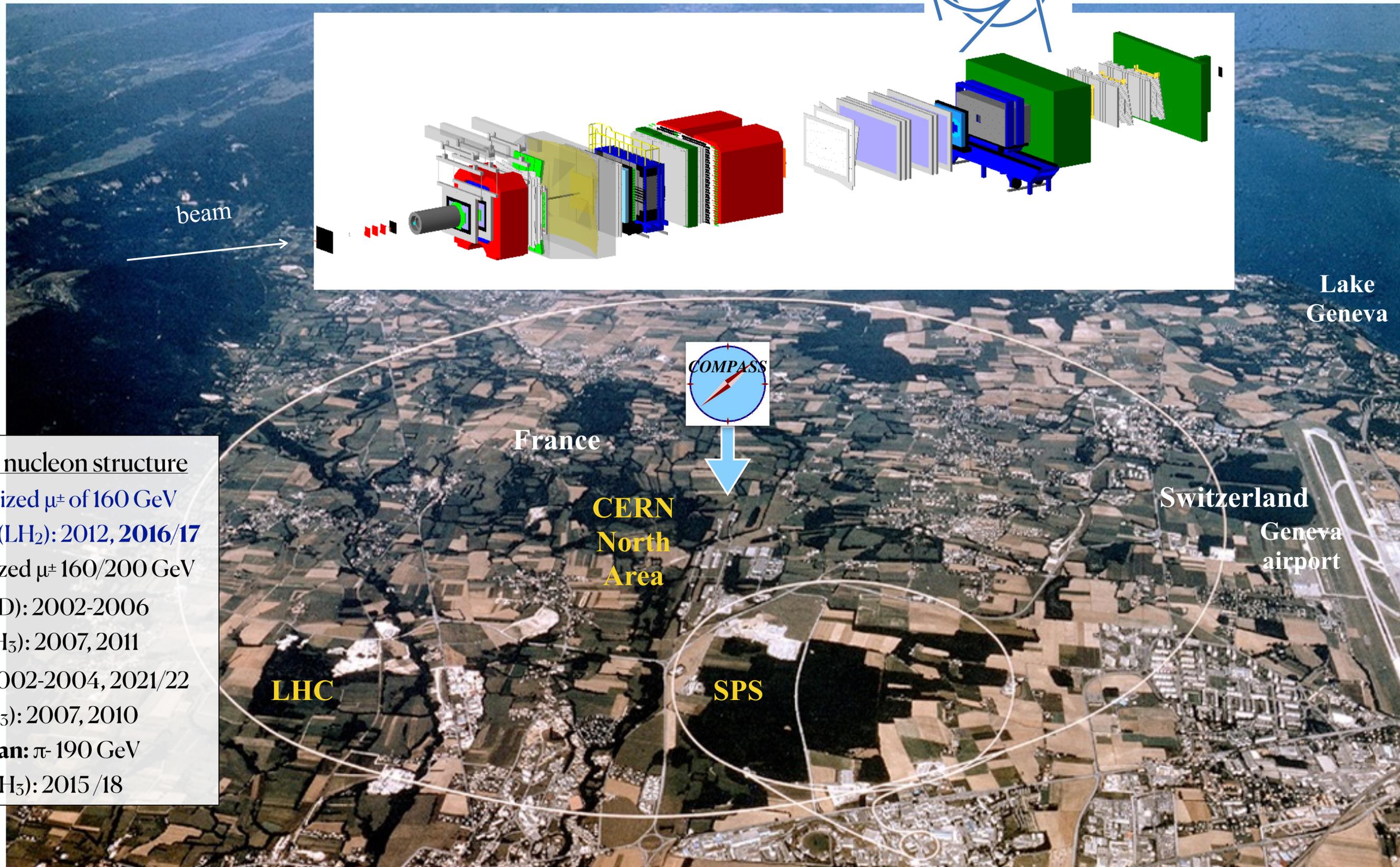
Re(τ_{DVCS})
integral over x

Im(τ_{DVCS})
 $|x| < \xi$
DDVCS

Dispersion relation with D-term $D(t)$: related to shear forces and radial distribution of pressure inside the nucleon

$$\text{Re}\mathcal{H}(\xi, t) = \mathcal{P} \int_{-1}^{+1} dx \frac{\text{Im}\mathcal{H}(x, t)}{x - \xi} + D(t)$$

COMPASS @



beam

Lake Geneva



France

CERN
North
Area

Switzerland
Geneva
airport

LHC

SPS

COMPASS nucleon structure

GPD: polarized μ^\pm of 160 GeV

unpol proton (LH_2): 2012, 2016/17

SIDIS: polarized μ^\pm 160/200 GeV

$d \rightarrow (^6\text{LiD})$: 2002-2006

$p \rightarrow (\text{NH}_3)$: 2007, 2011

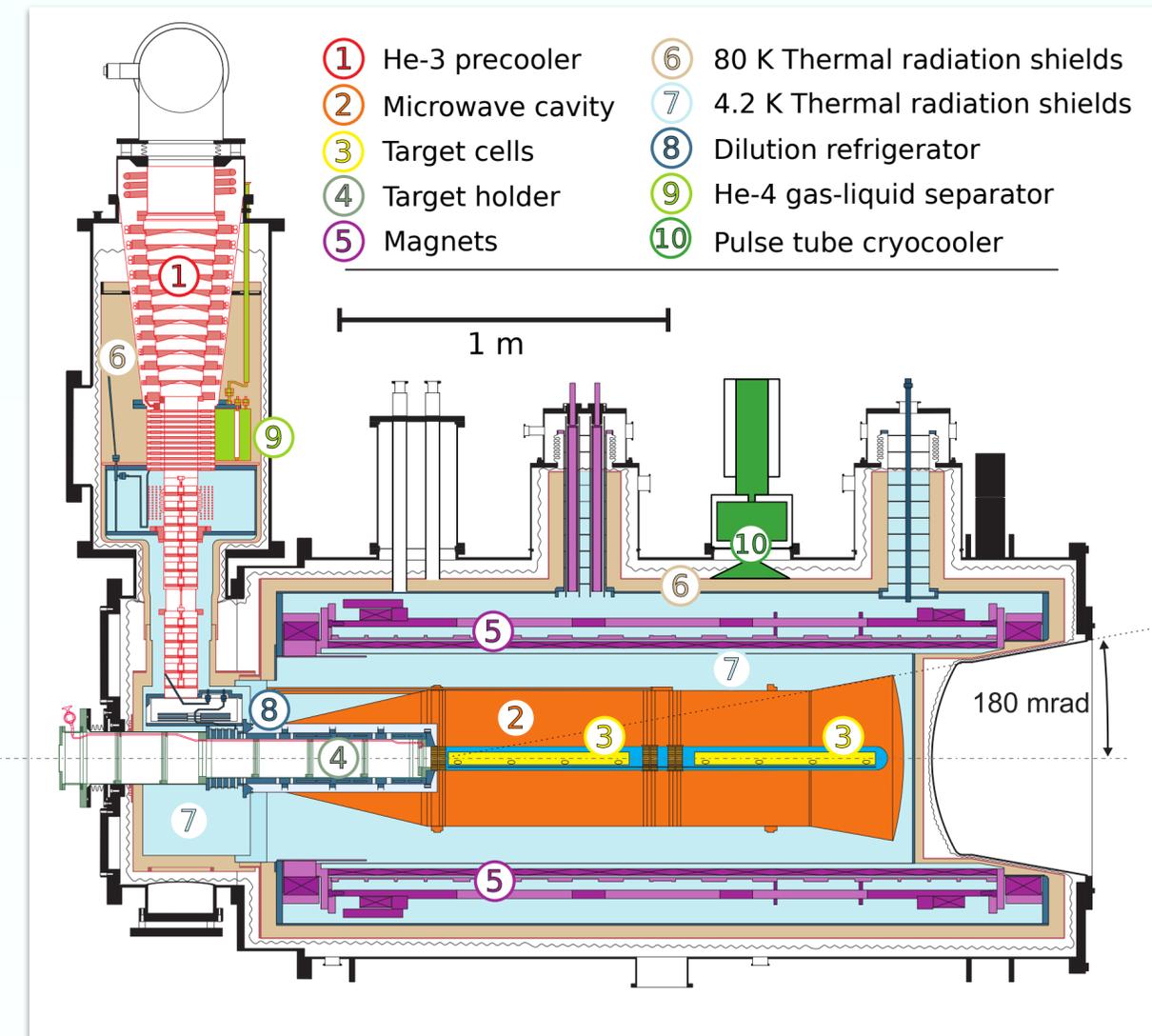
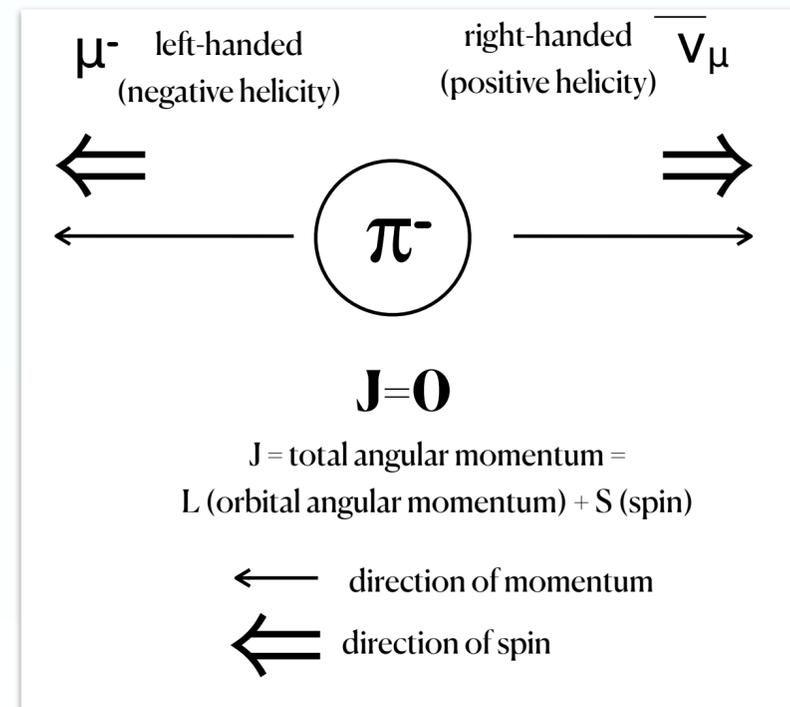
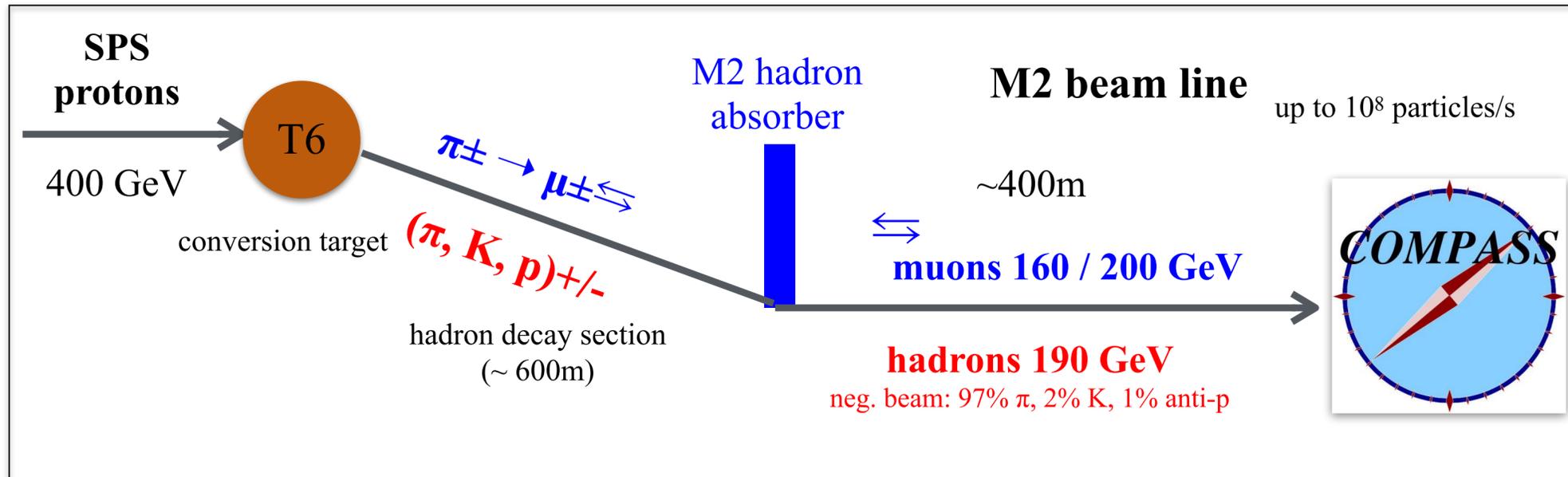
$d \uparrow (^6\text{LiD})$: 2002-2004, 2021/22

$p \uparrow (\text{NH}_3)$: 2007, 2010

Drell-Yan: π^- 190 GeV

$p \uparrow (\text{NH}_3)$: 2015/18

CERN SPS polarized muon beams and COMPASS polarized target



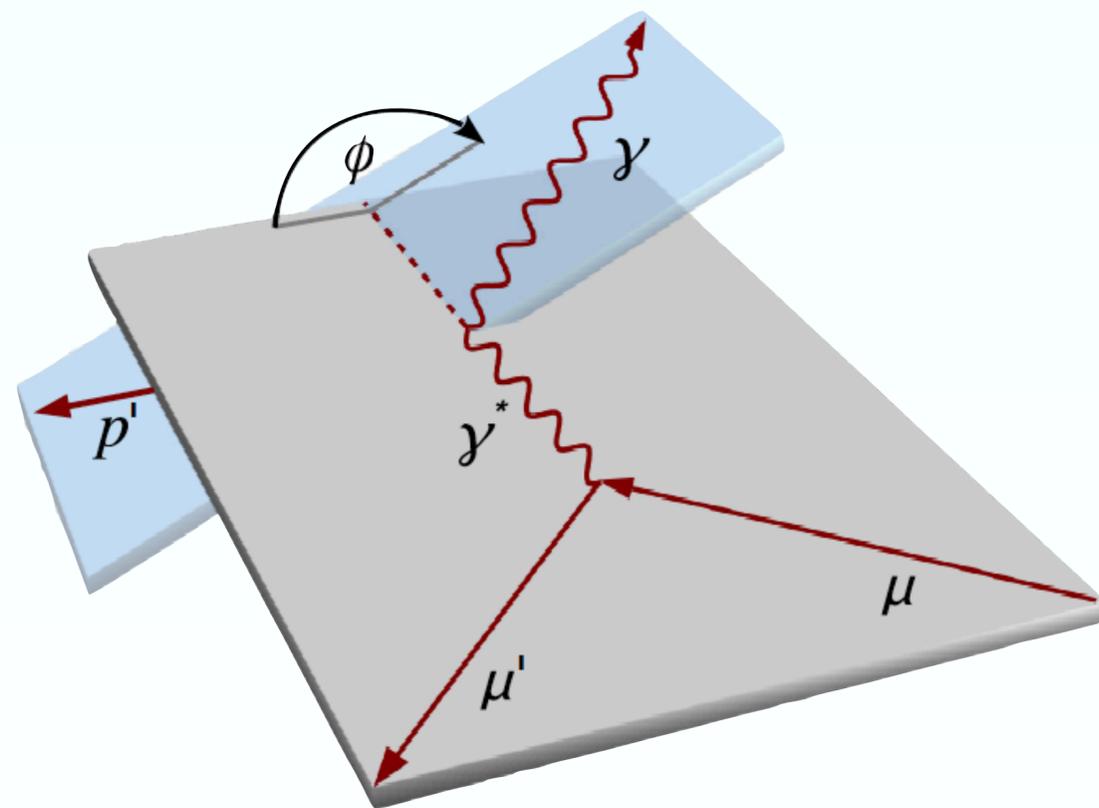
NH₃: ammonia beads, **⁶LiD**: deuterated lithium
dilution factor ~ 0.22 (NH₃), 0.5 (LiD)

- Polarization achieved by **Dynamic Nuclear Polarization (DNP)**
 - dilution refrigerator: ~60mK
 - dipole magnet (transverse): 0.5T
 - solenoid (longitudinal): 2.5T
 - microwave system

~ 80% for protons
~ 50% for deuterons

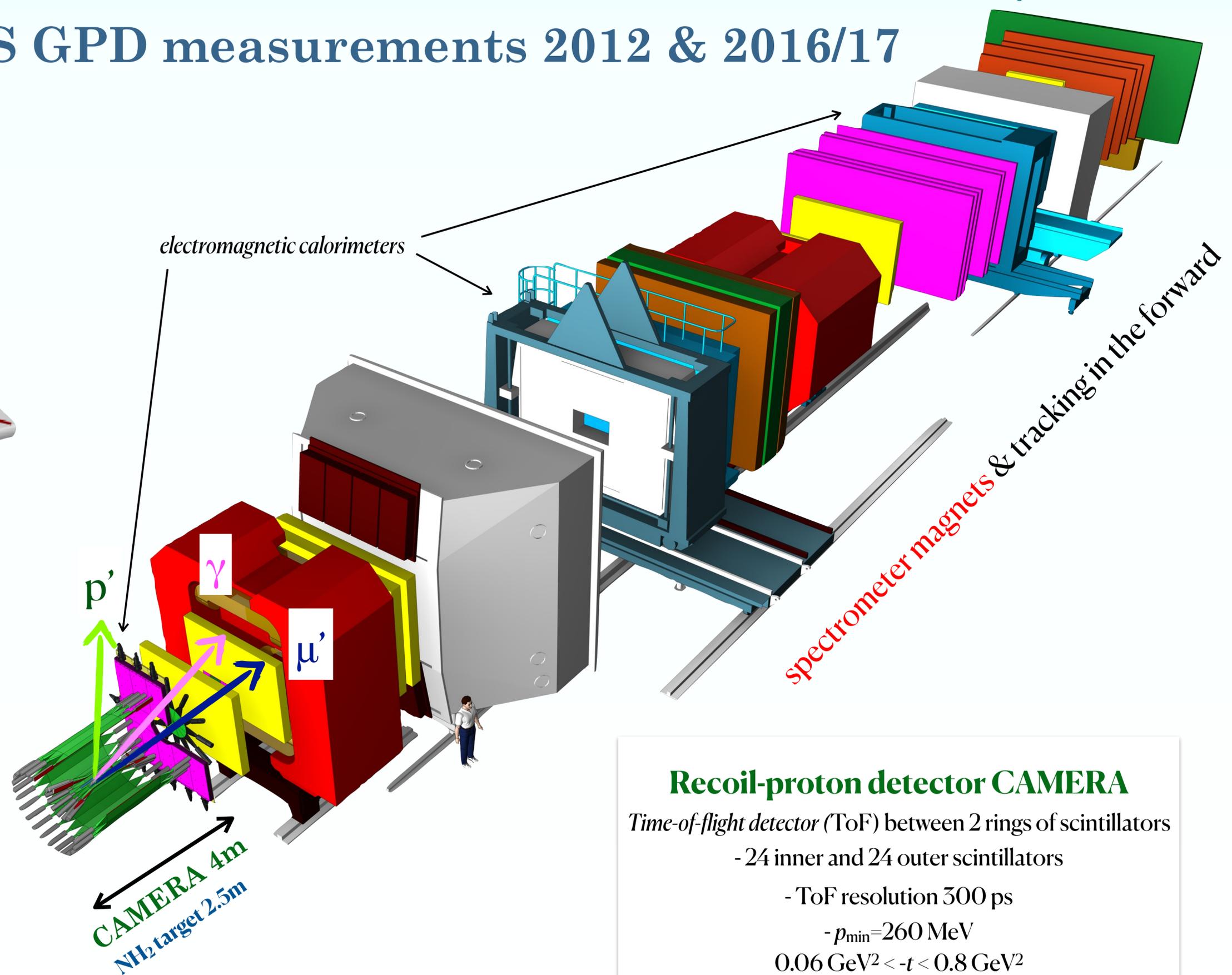
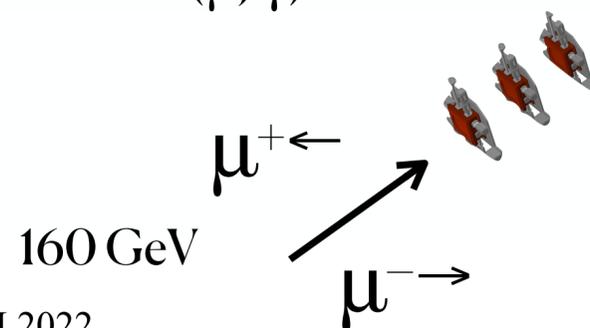
- Polarization determined with **Nuclear Magnetic Resonance (NMR)**

COMPASS GPD measurements 2012 & 2016/17



Azimuthal angle $\phi_{\gamma\gamma^*}$ between lepton-scattering $\mu\mu'$ plane and plane defined by virtual γ^* and real γ photons

Separate kinematic measurements in recoil-proton detector (p) and forward spectrometer (μ, γ)



Recoil-proton detector CAMERA

Time-of-flight detector (ToF) between 2 rings of scintillators

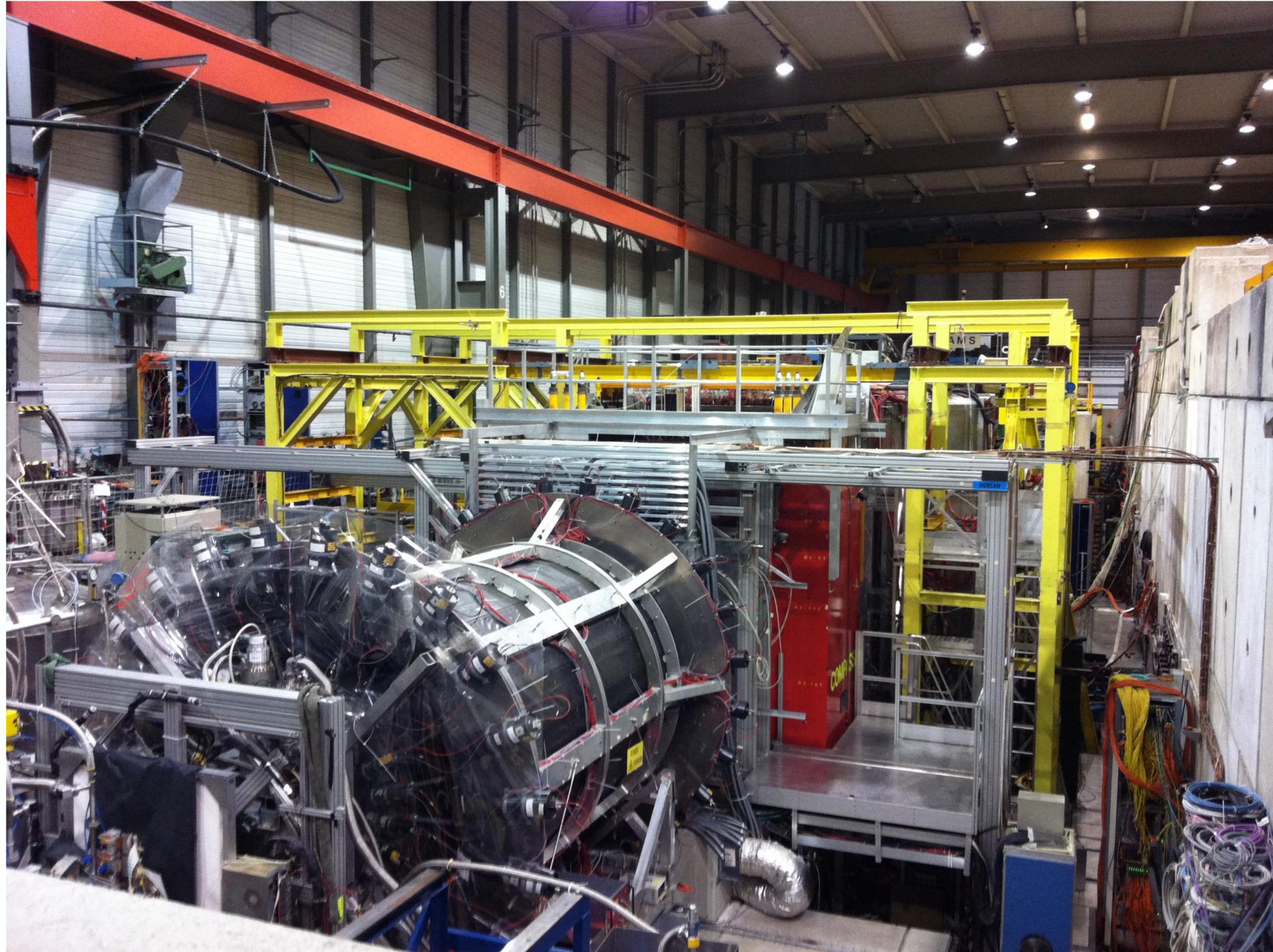
- 24 inner and 24 outer scintillators

- ToF resolution 300 ps

- $p_{\min} = 260$ MeV

$0.06 \text{ GeV}^2 < -t < 0.8 \text{ GeV}^2$

COMPASS GPD measurements 2012 & 2016/17

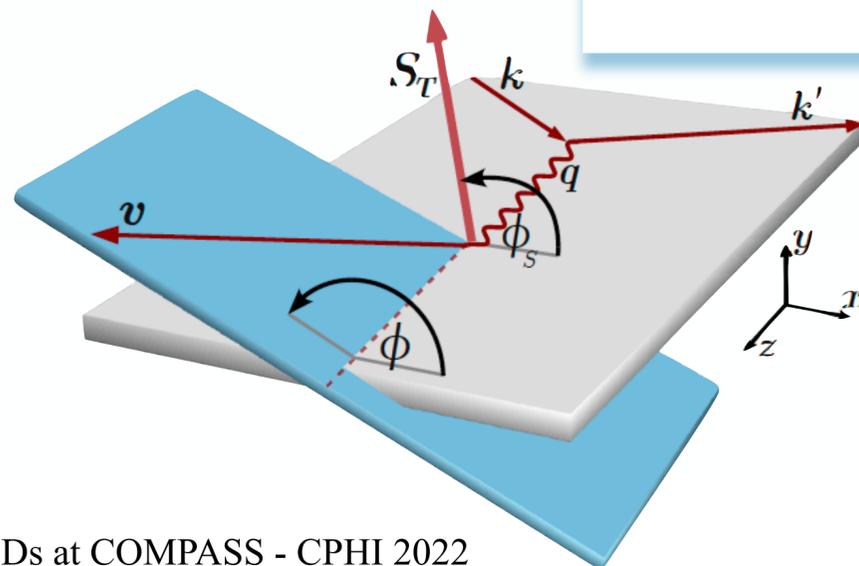
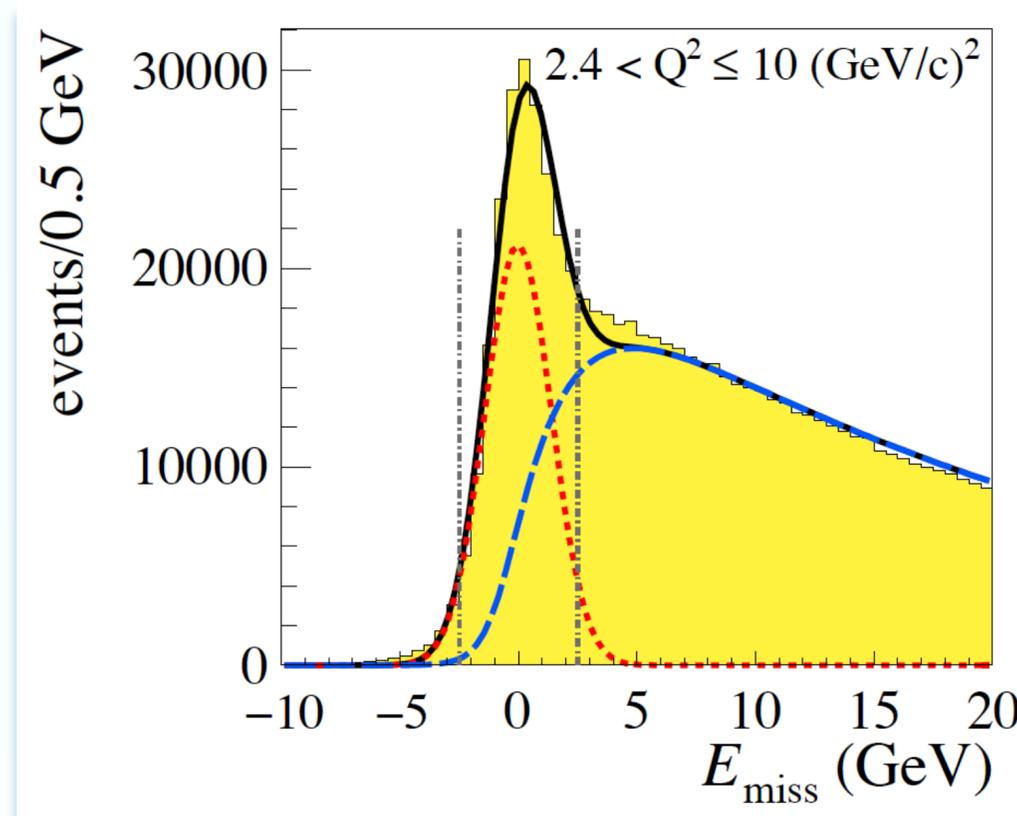


Selection of exclusive event sample

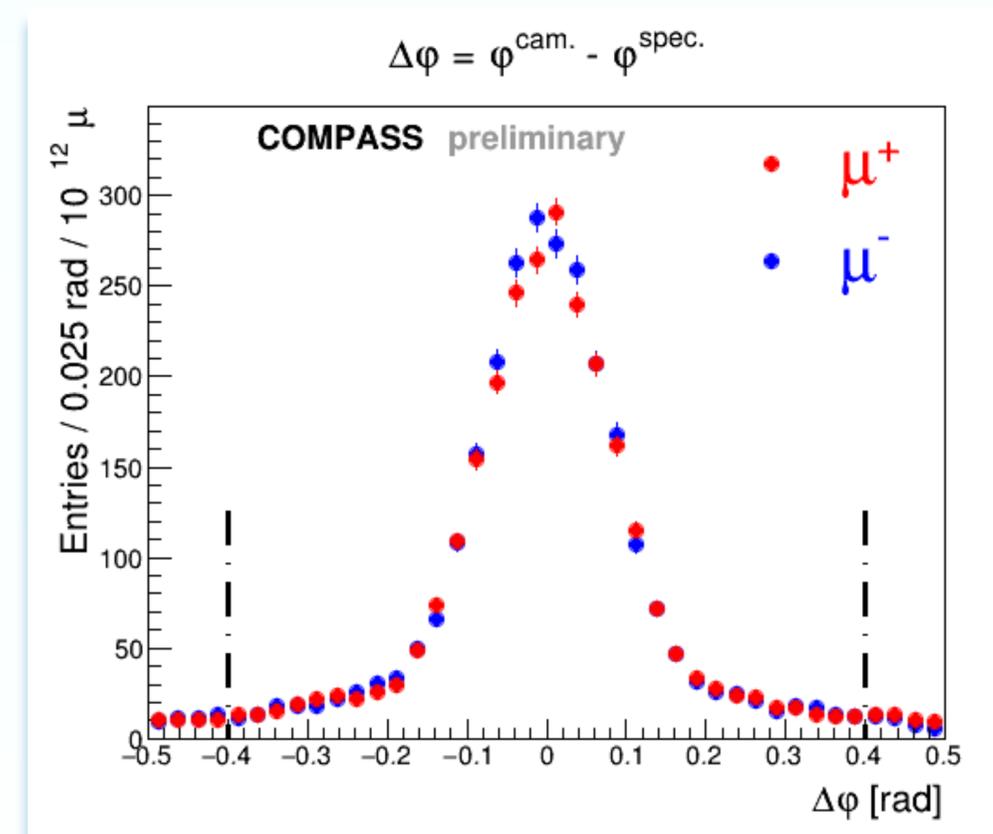
DVMP without recoil-proton detection:
missing energy technique assuming proton mass

DVCS with recoil-proton detector (RPD):
comparison of proton kinematics measured in
RPD vs. expected in spectrometer (from $\mu\gamma$)

Exclusive ρ^0 production
 $\mu p^{(\uparrow)} \rightarrow \mu p \rho^0$
with simulated **exclusive**
signal & **SIDIS** background



In case of transverse target
polarization: additional **azimuthal**
angle ϕ_s defined by direction of
transverse target-polarization vector



DVCS
 $\mu^\pm p \rightarrow \mu^\pm p \gamma$

+ kinematically complete event reconstruction via
kinematic event fitting

Access to CFFs at COMPASS

ϕ -modulation in cross section
(azimuthal asymmetry analysis)

DVCS $\sigma_{\gamma^* \gamma N} \sim \left| \begin{array}{c} \text{DVCS} \\ + \\ \text{Bethe-Heitler (BH)} \end{array} \right|^2$

$$= |\mathcal{T}_{\text{BH}}|^2 + \left(\mathcal{T}_{\text{DVCS}} \mathcal{T}_{\text{BH}}^* + \mathcal{T}_{\text{DVCS}}^* \mathcal{T}_{\text{BH}} \right) + |\mathcal{T}_{\text{DVCS}}|^2$$

Analysis of azimuthal modulations
(HERMES- and JLab-type) on DVCS on
the unpolarized proton in progress

dominant at small x_B
(remainder ~ 5% from KM / GK model)

The DVCS / Bethe-Heitler interference
term allows to disentangle
Re(τ_{DVCS}) and **Im(τ_{DVCS})**
magnitude and **phase**
of DVCS amplitude τ_{DVCS}

$$\mathcal{S}_{CS,U} \equiv d\sigma^{\leftarrow+} + d\sigma^{\leftarrow-}$$

$$\mathcal{D}_{CS,U} \equiv d\sigma^{\leftarrow+} - d\sigma^{\leftarrow-}$$

$$\mathcal{A}_{CS,U} \equiv \frac{d\sigma^{\leftarrow+} - d\sigma^{\leftarrow-}}{d\sigma^{\leftarrow+} + d\sigma^{\leftarrow-}} = \frac{\mathcal{D}_{CS,U}}{\mathcal{S}_{CS,U}}$$

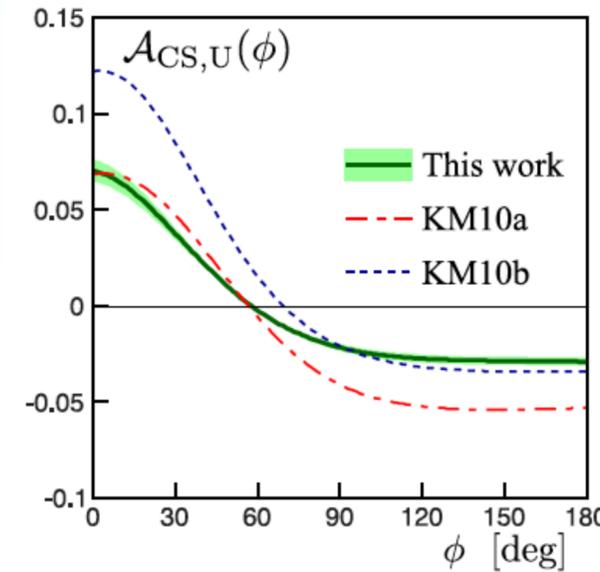
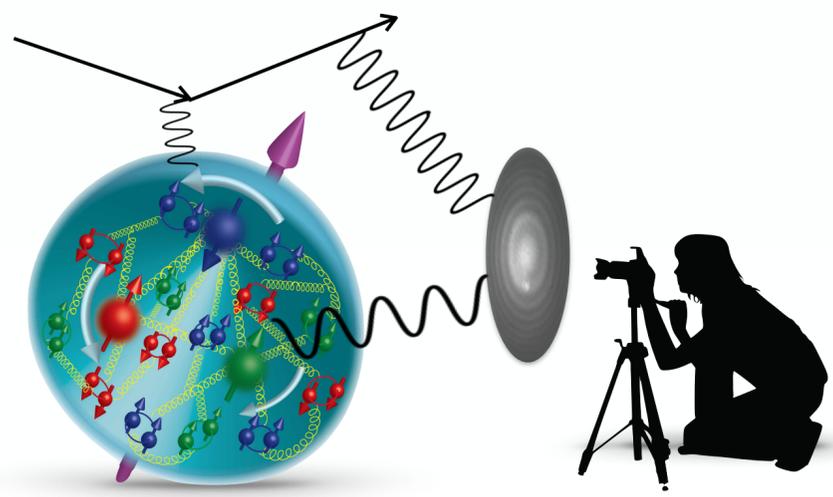
Spin-independent DVCS cross section \propto

$$4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2} \mathcal{E}\mathcal{E}^*$$

$$\text{Im} \left(F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E} \right)$$

$$\text{Re} \left(F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E} \right)$$

at leading twist (twist 2)
on the proton



Sign & magnitude of $\cos\phi$
amplitude for beam-charge
asymmetry? (changes sign
between HERMES and HERA)

Kroll, Moutarde, Sabatié, Eur. Phys. J. C
(2013) 73:2278

Test of GPD universality: use DVMP data
to constrain GPD parameters

Transverse imaging of the nucleon

DVCS cross section /
transverse imaging

Impact-parameter representation of parton distribution function:

$$q^f(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i\Delta_\perp \cdot \mathbf{b}_\perp} H^f(x, 0, -\Delta_\perp^2)$$

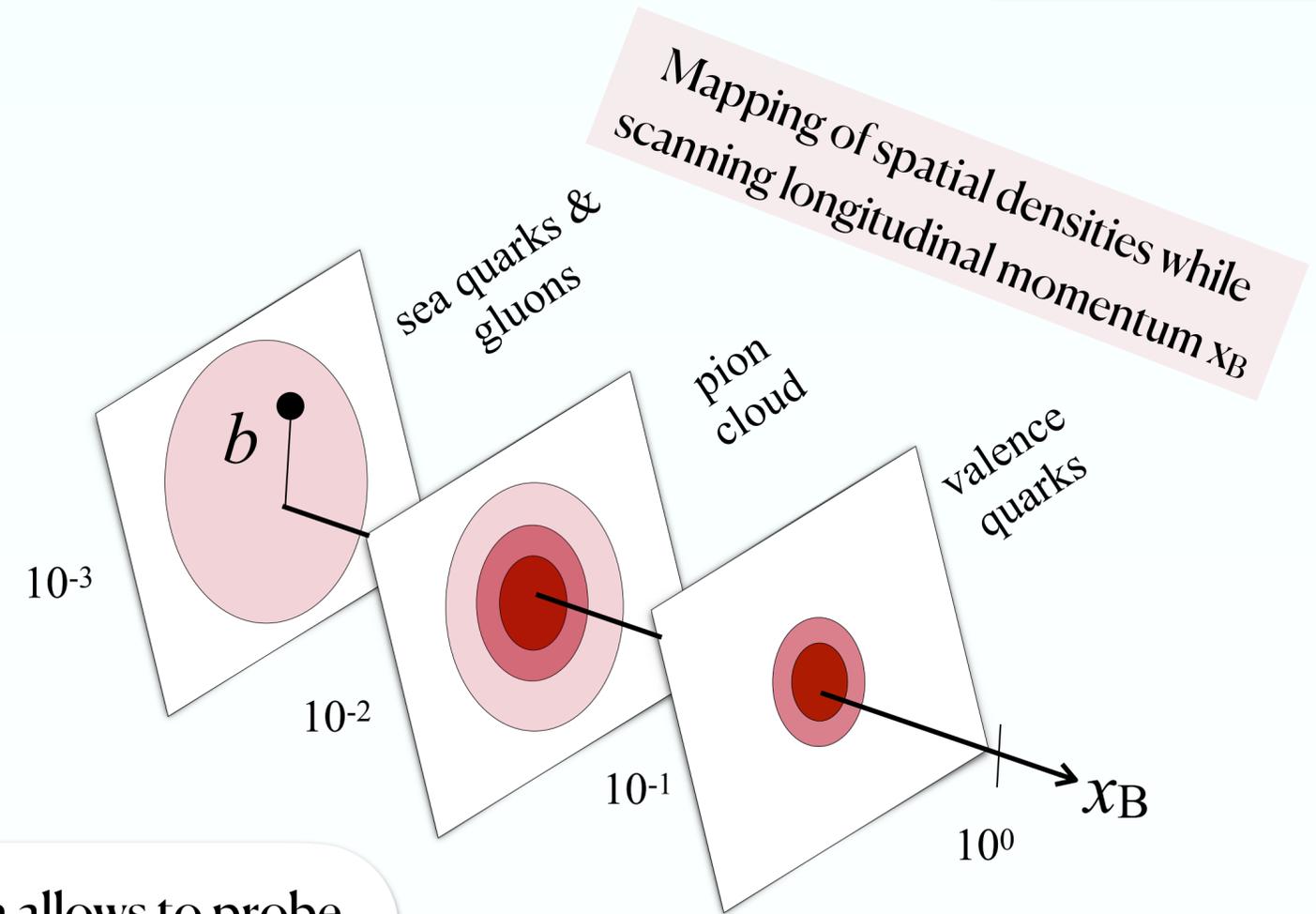
[Burkardt, Int. J. Mod. Phys. A18 (2003) 173]

“spatial parton density = Fourier transform of GPD”

\mathbf{b}_\perp is the impact parameter,

Δ_\perp is the difference of initial and final transverse momenta,

Δ_\perp^2 is related to the Mandelstam- t



The differential DVCS cross section allows to probe the **transverse extension of partons** in the nucleon:

$$\frac{d\sigma^{\text{DVCS}}}{dt} \propto e^{-b|t|}$$

$b = \text{“}t\text{-slope”} = \text{average impact parameter}$

3-dim “tomographic images” of the nucleon in longitudinal momentum and transverse position

Extraction of pure DVCS yield at COMPASS

DVCS cross section /
transverse imaging

$$|\mathcal{T}_{\text{BH}}|^2 + (\mathcal{T}_{\text{DVCS}}\mathcal{T}_{\text{BH}}^* + \mathcal{T}_{\text{DVCS}}^*\mathcal{T}_{\text{BH}}) + |\mathcal{T}_{\text{DVCS}}|^2$$

BH reference yield:

at small

$\langle x_B \rangle = 0.0085$

DVCS amplitude:

ϕ -modulations in cross section

at medium $\langle x_B \rangle = 0.0200$

Transverse imaging:

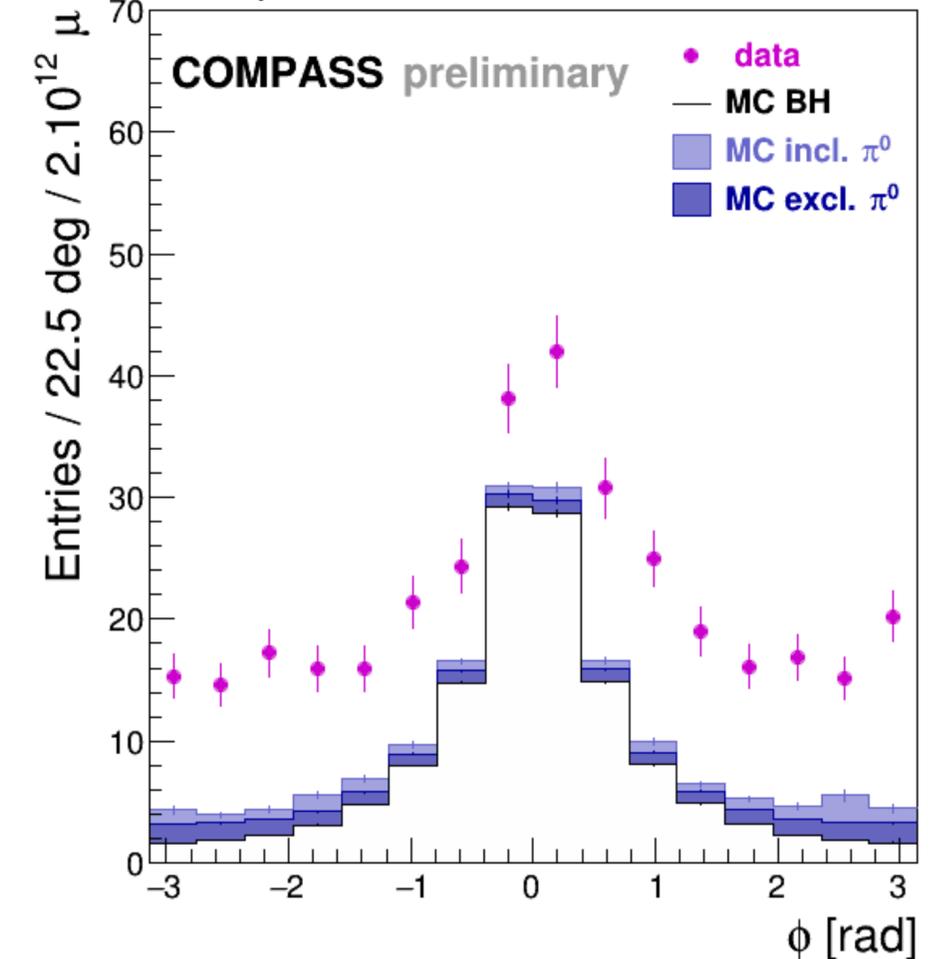
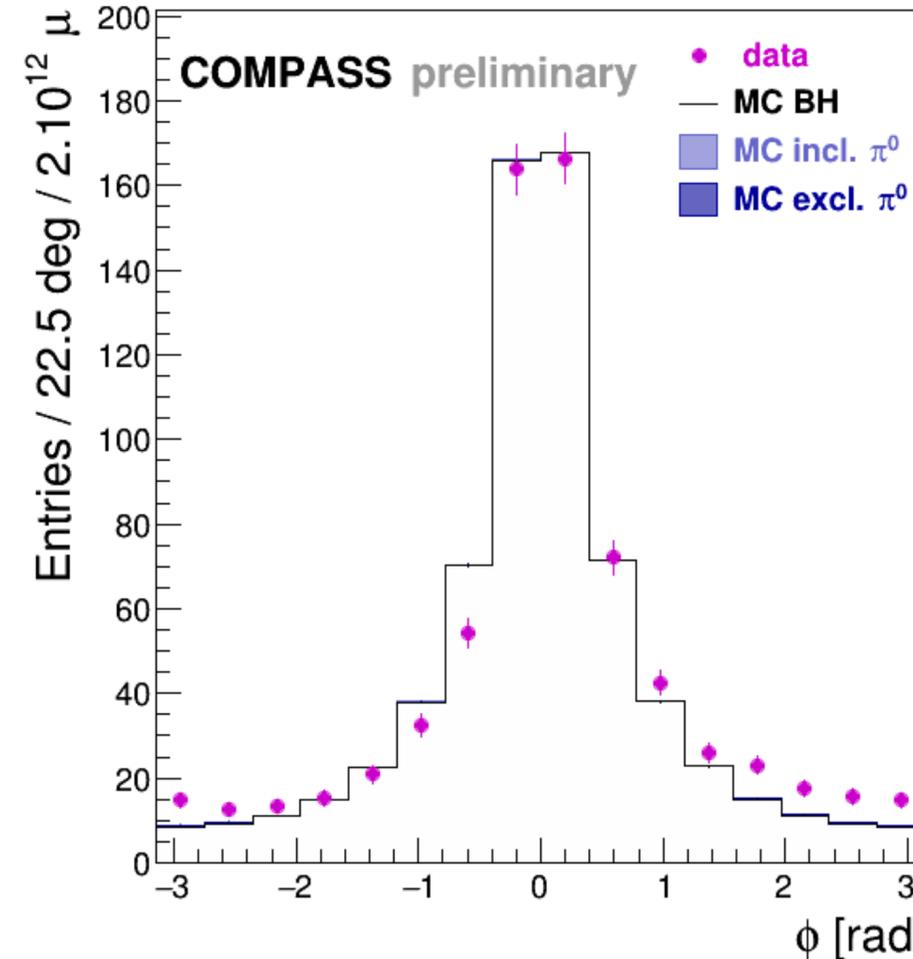
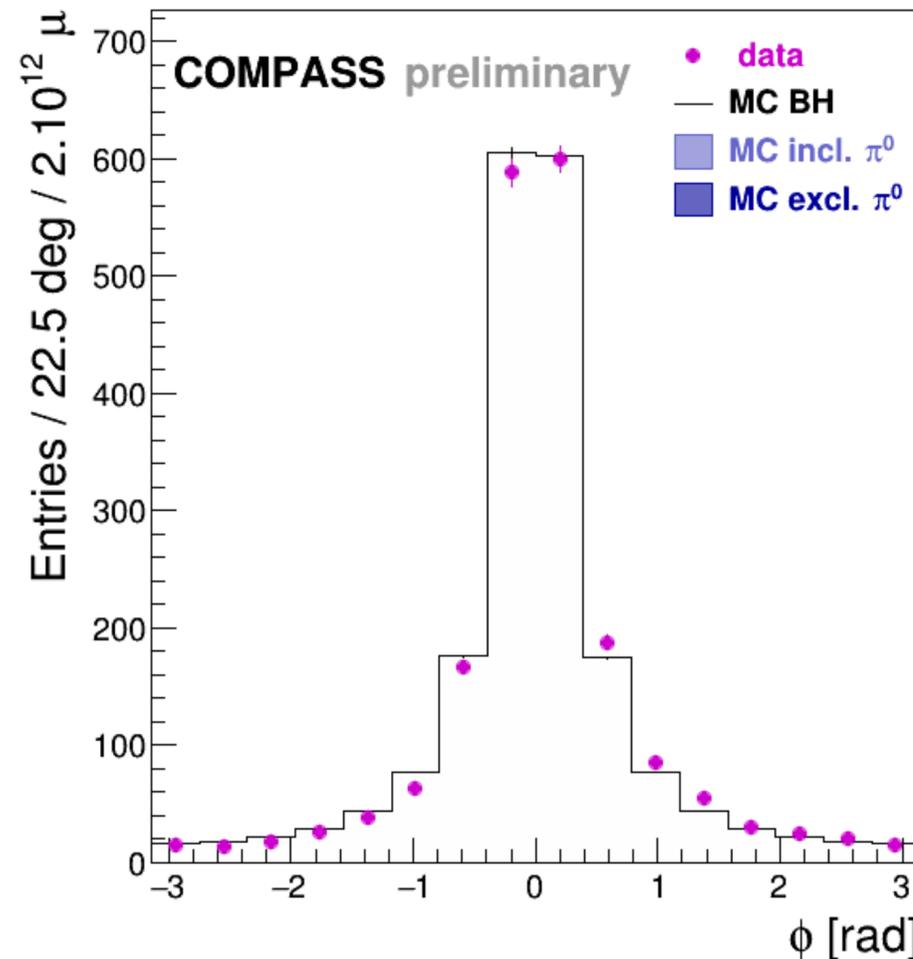
ϕ -integrated cross section

at medium $\langle x_B \rangle = 0.0630$

$80 < \nu \text{ [GeV]} < 144$

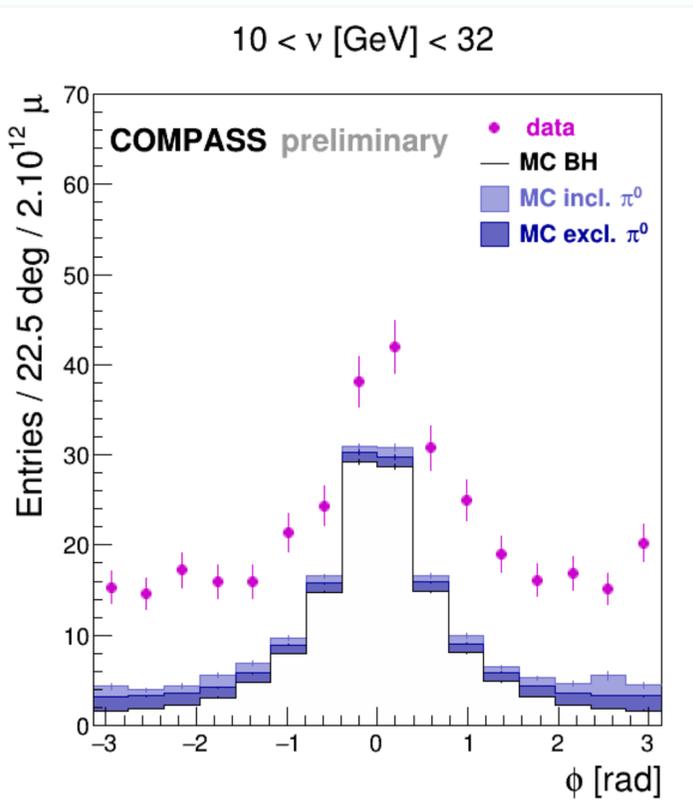
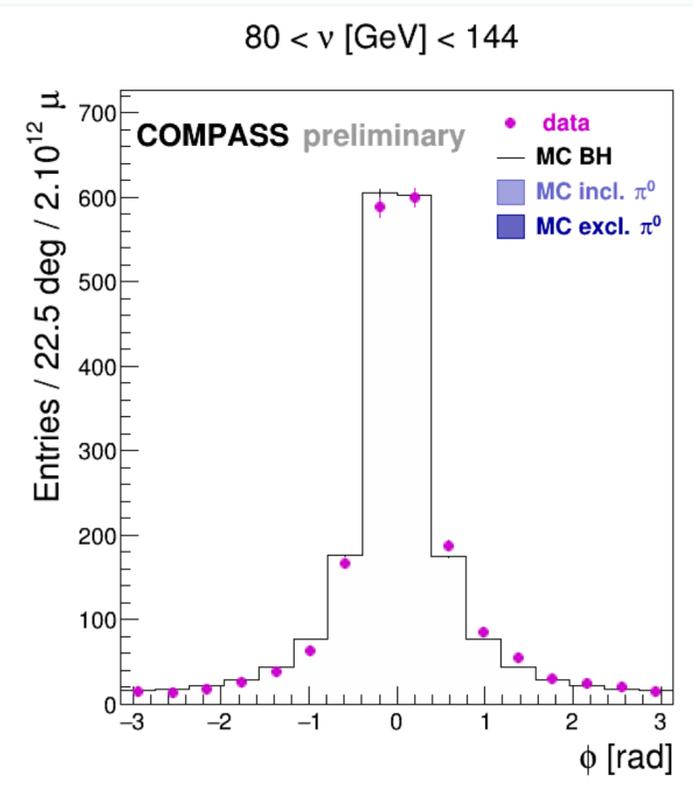
$32 < \nu \text{ [GeV]} < 80$

$10 < \nu \text{ [GeV]} < 32$



Extraction of pure DVCS yield at COMPASS

DVCS cross section /
transverse imaging



DVCS / BH: $\mu p \rightarrow \mu p \gamma$
 π^0 production: $\mu p \rightarrow \mu p \pi^0 \rightarrow \mu p \gamma(\gamma)(X)$ exclusive or SIDIS
 $\gamma\gamma$: “visible π^0 ”, γ : invisible π^0 background

Determine BH reference yield at low- $x_B \Leftrightarrow$ high- v : tune MC to data

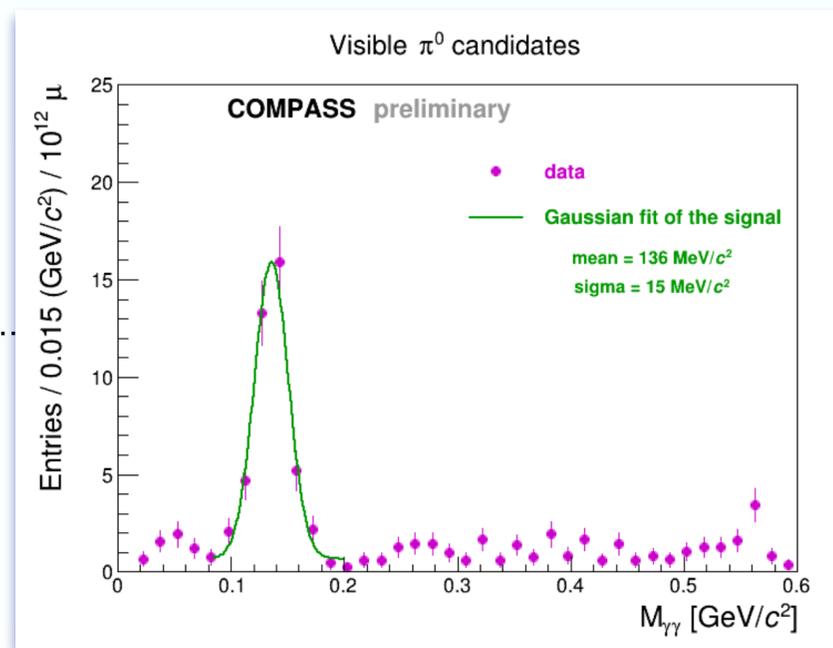
Subtract BH yield in high- $x_B \Leftrightarrow$ low- v bin

Subtract measured visible π^0 yield in high- $x_B \Leftrightarrow$ low- v bin

Estimate π^0 invisible background from MC:
 SIDIS 40% (LEPTO) + exclusive 60% (HEPGEN with GK model)
 with 10% uncertainty each

Remove invisible π^0 yield
 (invisible normalized to visible yield)

Pure DVCS yield



portion of the 2016 data = 2x 2012 data

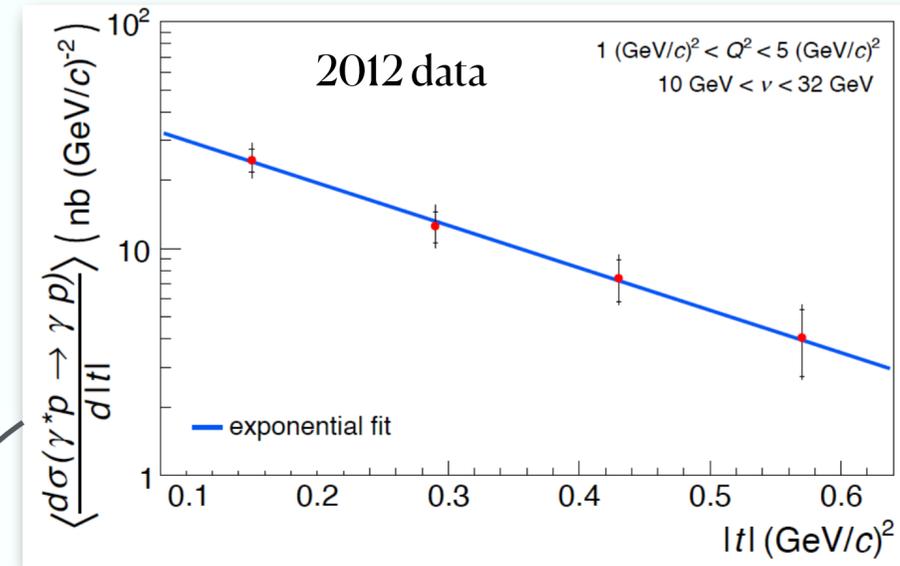
COMPASS ϕ -integrated DVCS cross section

DVCS cross section /
transverse imaging

- Relation between parton average squared transverse extension and slope

$$\langle r_{\perp}^2(x_{Bj}) \rangle \approx 2 \langle B(x_{Bj}) \rangle \hbar^2$$

- Sea-quark domain between gluons and valence-quarks
- Pilot run:** $\langle Q^2 \rangle = 1.8 \text{ GeV}^2$ & $\langle x_{Bj} \rangle = 0.056$
- 2016/17 run:** $\langle Q^2 \rangle = 1.8 \text{ GeV}^2$ & $\langle x_{Bj} \rangle = 0.063$ (ECal0 extension)
- So far, $\approx 30\%$ of COMPASS cross-section data analyzed.

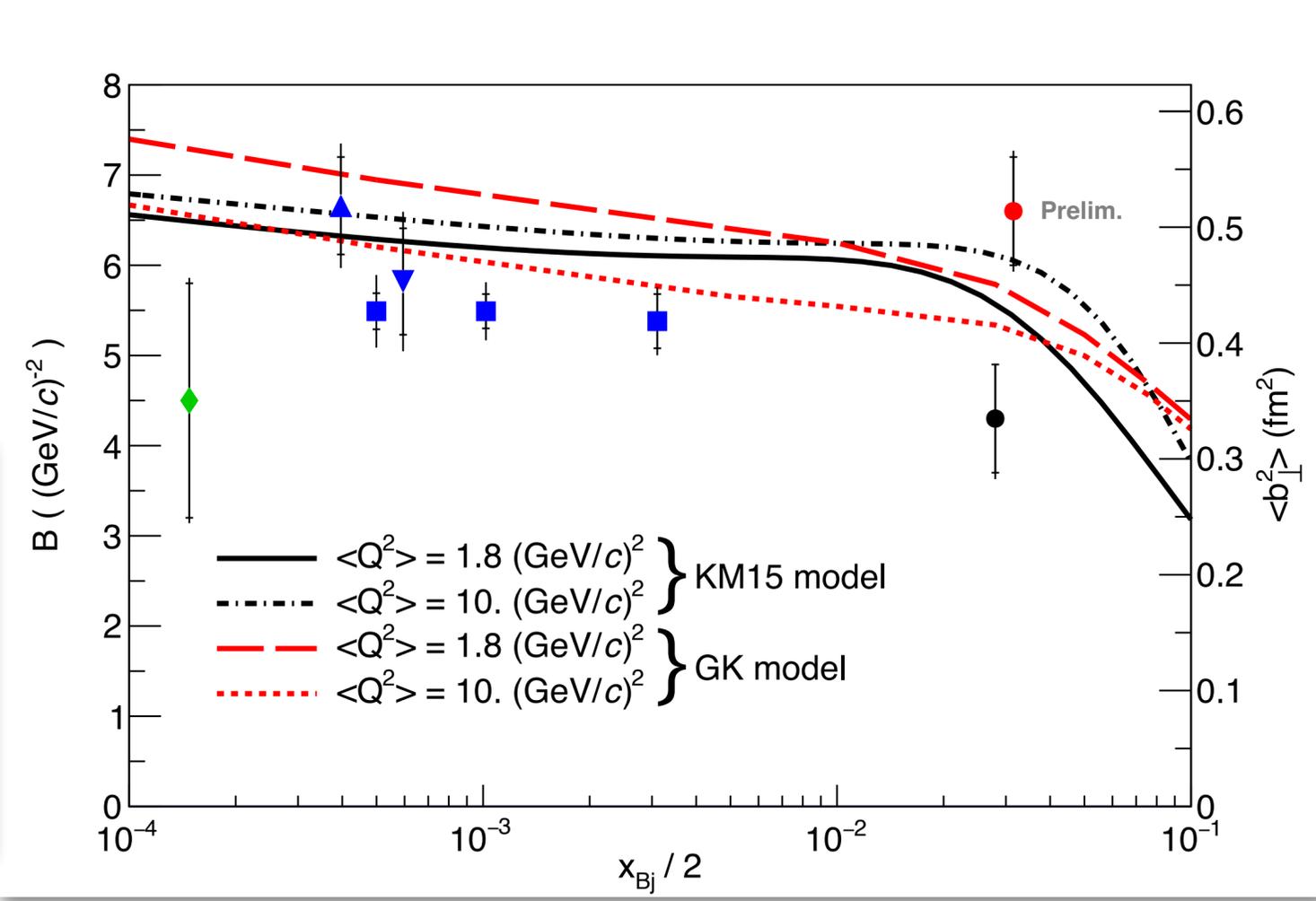


$$\frac{d\sigma^{\text{DVCS}}}{dt} \propto e^{-b|t|}$$

Differential DVCS cross section with b (or B) = “ t -slope” = average impact parameter

2016/17 ($\approx 23\%$)
2012

| | | |
|---|--|-----------------------------|
| ● | COMPASS: $\langle Q^2 \rangle = 1.8 \text{ (GeV/c)}^2$ | This Analysis |
| ● | COMPASS: $\langle Q^2 \rangle = 1.8 \text{ (GeV/c)}^2$ | Phys. Lett. B793 (2019) 188 |
| ◆ | ZEUS: $\langle Q^2 \rangle = 3.2 \text{ (GeV/c)}^2$ | JHEP 0905 (2009) 108 |
| ▲ | H1: $\langle Q^2 \rangle = 4.0 \text{ (GeV/c)}^2$ | } Eur. Phys. C44 (2005) 1 |
| ▼ | H1: $\langle Q^2 \rangle = 8.0 \text{ (GeV/c)}^2$ | |
| ■ | H1: $\langle Q^2 \rangle = 10. \text{ (GeV/c)}^2$ | Phys. Lett. B681 (2009) 391 |



[COMPASS PLB 793 (2019) 188] (2012 data)

Exclusive meson production

| | | |
|---|--|--|
| chiral-even GPDs at leading twist | | chiral-odd GPDs at higher twist |
| H, E | J ^P =1 ⁻ vector mesons | H _T , E _T , Ē _T =2H̄ _T +E _T |
| H̄ | J ^P =0 ⁻ pseudoscalar mesons | |
| longitudinally polarized virtual photon & vector meson) | | |

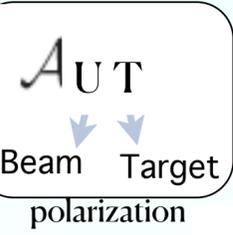
- Different mesons filter different quark flavors and have different sensitivity to gluon GPDs.

$$\begin{aligned}
 E^{\rho^0} &= \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u + \frac{1}{3} E^d + \frac{3}{4} E^g \right) \\
 E^\omega &= \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u - \frac{1}{3} E^d + \frac{3}{4} E^g \right) \\
 E^\phi &= -\frac{1}{3} E^s + \frac{1}{8} E^g
 \end{aligned}$$

- Deeply virtual meson production allows also access to higher-twist chiral-odd GPDs.
- Some GPDs are related to transverse-momentum dependent PDFs (TMDs):
 - **GPD E ↔ Sivers TMD**: involve switch of nucleon helicity $\sim \vec{S}_T \cdot (\hat{P} \times \vec{k}_T)$
 ⇒ sensitive to spin-orbit correlations... orbital angular momentum
 - **GPD H_T ↔ transversity TMD**: both are chiral-odd
 - **GPD Ē_T ↔ Boer-Mulders TMD**: T-odd (as is Sivers TMD)
- Helicity conservation in interactions of light quarks with gluons or photons
 ⇒ initial parton helicity flip needs compensation by **higher-twist meson wave functions**

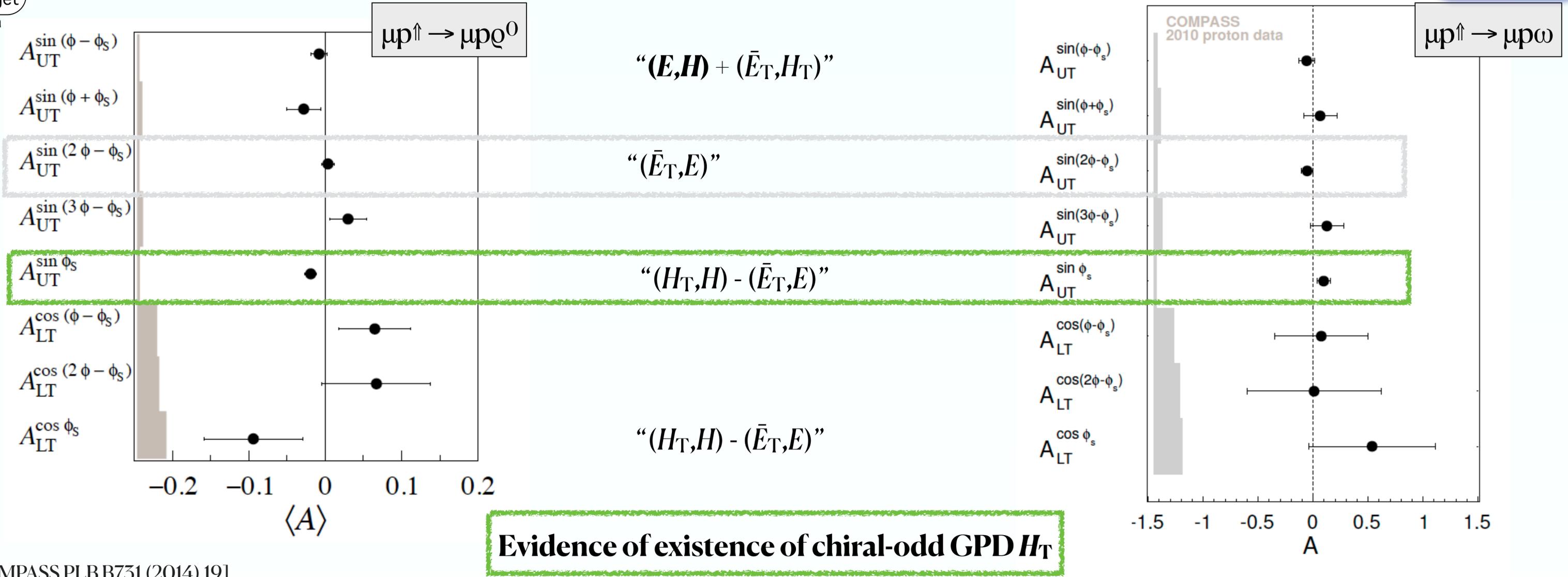
Goloskokov, Kroll, Eur. Phys. J. C 74, 2725 (2014)

We will now look at **transverse-spin asymmetries**, spin-density matrix elements (**SDMEs**), and the **cross section for exclusive π⁰ production**.



Transverse-target spin asymmetries in DVMP

A_{UT} / chiral-odd GPDs



[COMPASS PLB B731 (2014) 19]

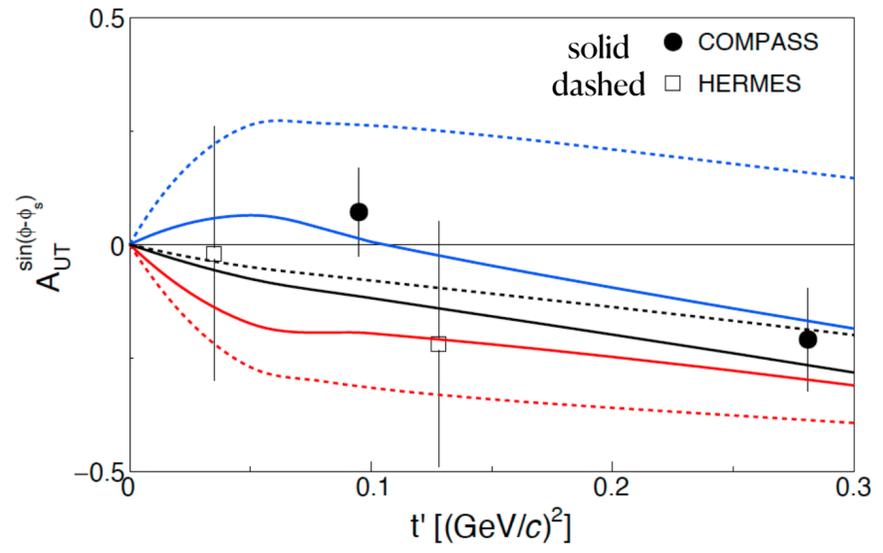
$$E^{\rho^0} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u + \frac{1}{3} E^d + \frac{3}{4} E^g \right)$$

$$E^\omega = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u - \frac{1}{3} E^d + \frac{3}{4} E^g \right)$$

$$E^\phi = -\frac{1}{3} E^s + \frac{1}{8} E^g$$

GPDE

- Cancellation effects for q^0 ...
- GPD flavor separation by measuring both meson channels...



[COMPASS NPB 915 (2017) 454]

Contribution of **pion pole** expected to be sizeable for ω and small for q^0

GK 2014

positive $\pi\omega$ form factor
no pion pole
negative $\pi\omega$ form factor

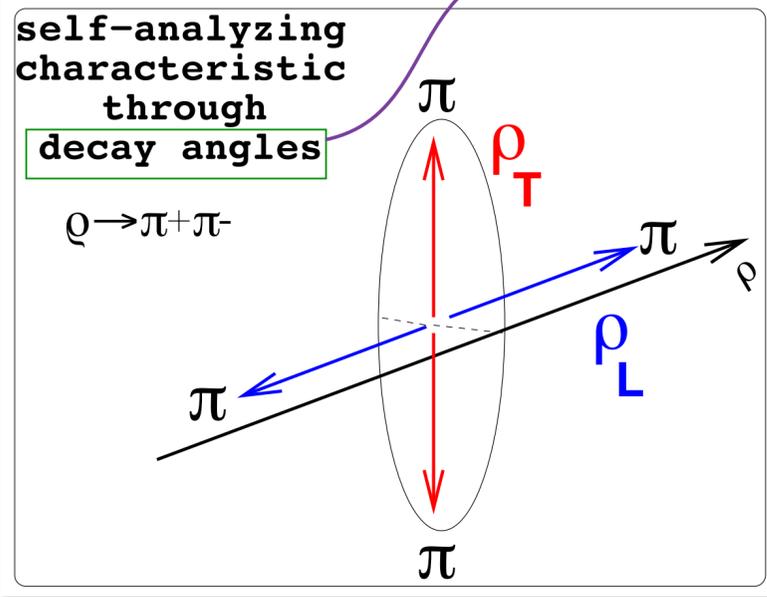
[COMPASS NPB 865 (2012) 1 (not shown, includes q on deuteron)]

Spin density matrix elements in $\mu p \rightarrow \mu p \text{VM}$

SDMEs / chiral-odd GPDs

$$\frac{d\sigma}{dx_B dQ^2 dt} W(x_B, Q^2, t, \phi, \phi_S, \varphi, \vartheta)$$

$W(\dots)$ parametrized by SDMEs

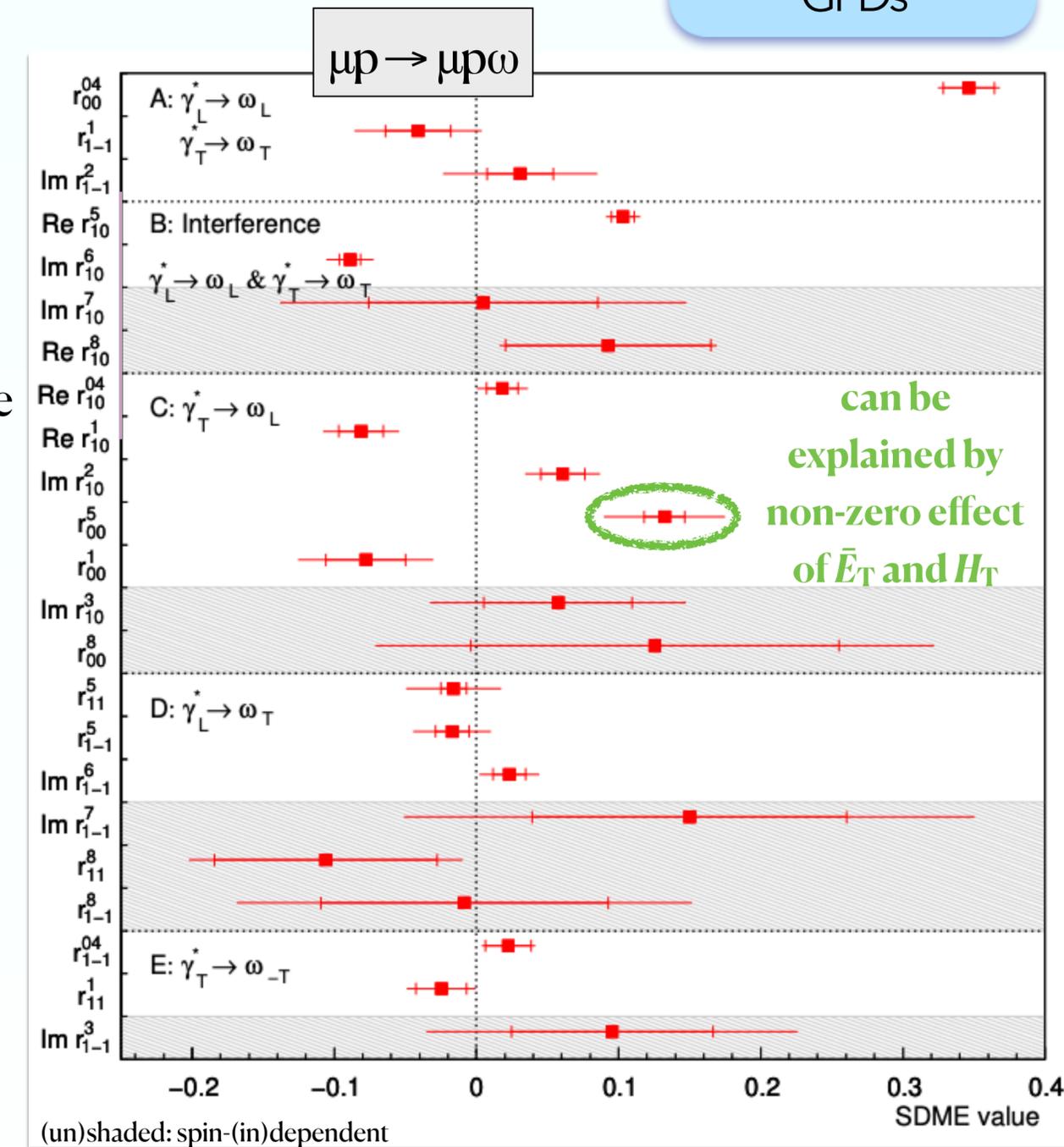


- Spin density matrix elements describe how the spin components of the virtual photon are transferred to the created vector meson
- Test of s-channel helicity conservation (SCHC), $\lambda_{\gamma^*} = \lambda_{\text{VM}}$, : only SDMEs of classes A&B are not restricted to =0 if SCHC. Observed: considerable SCHC in $\gamma^*_T \rightarrow \omega_L$ (class C)
- SDMEs measurements provide further constraints on GPD parameterizations beyond cross-section and spin-asymmetry measurements.

- Sensitivity to chiral-odd GPDs H_T and \bar{E}_T .

- Tests of hierarchy of helicity amplitudes
- Cross-section ratio R of longitudinal to transverse vector mesons,...

$$R' = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}}$$



[COMPASS EPJC (2021) 81126]

SDMEs $\mu p \rightarrow \mu p \rho^0$ to be published

$\mu p \rightarrow \mu p \pi^0$

Exclusive π^0 cross section at COMPASS

π^0 cross section /
chiral-odd GPDs

$$\frac{d^2\sigma_{\gamma^*p}}{dt d\phi} = \frac{1}{2\pi} \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cos(2\phi) \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) \frac{d\sigma_{LT}}{dt} \right]$$

L, T indices indicate polarization of virtual photon. Double index = interference

$$\left[(1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2 \right]$$

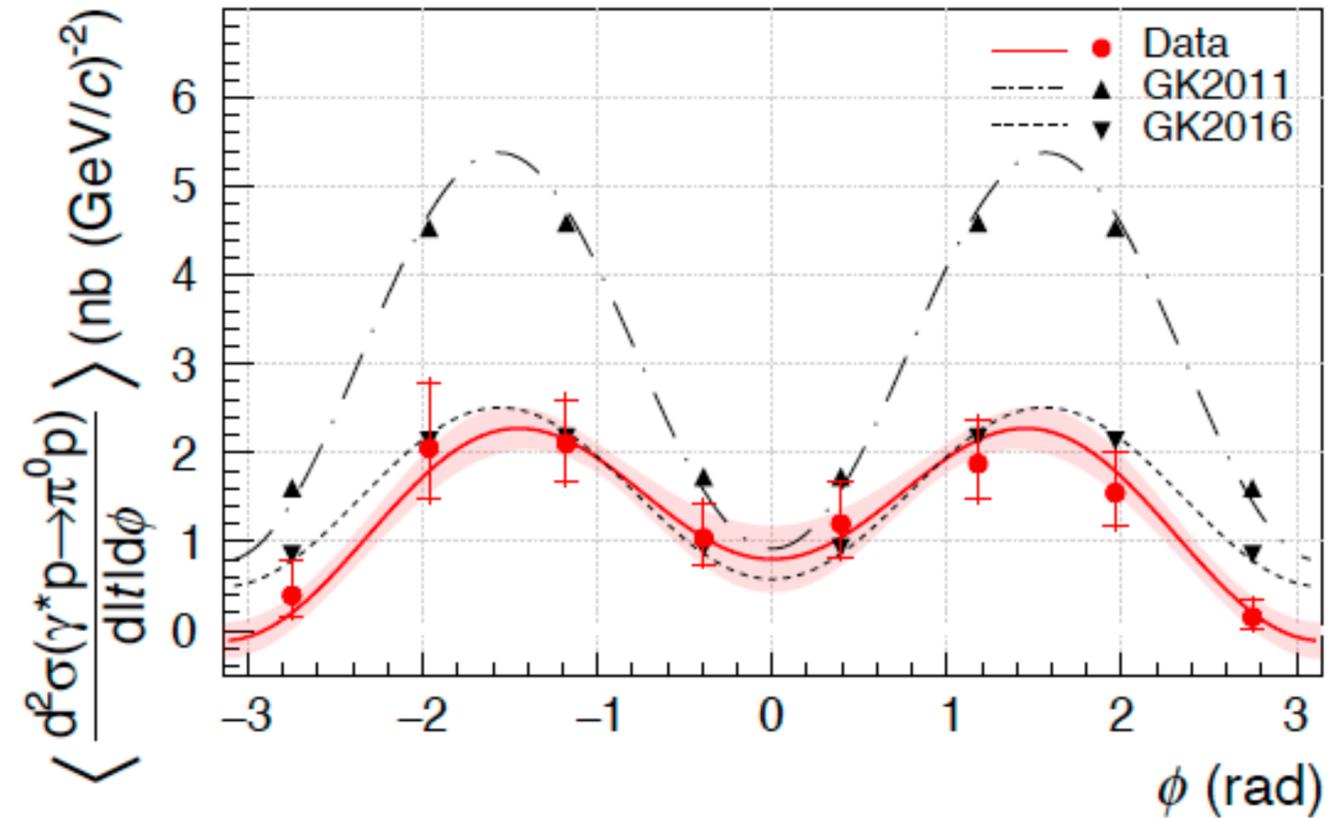
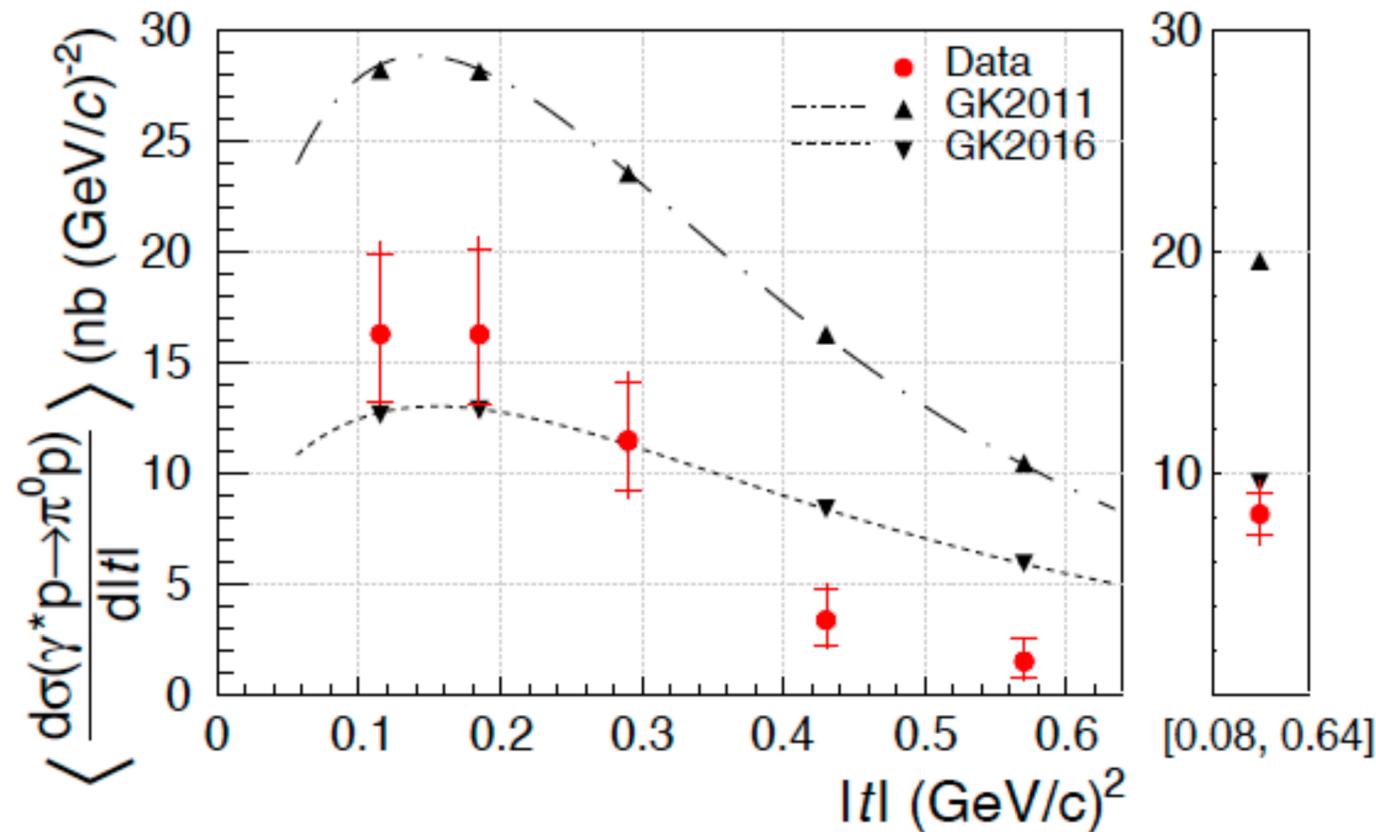
sensitivity to chiral-odd GPDs

$$t' |\langle \bar{E}_T \rangle|^2$$

$$\xi \sqrt{1 - \xi^2} \sqrt{-t'} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

Dip at small $|t|$ indicative of large effect by chiral-odd GPD \bar{E}_T

$$\left[(1 - \xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4M^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right]$$



Summary and outlook: GPDs at **COMPASS**



- 2012 GPD pilot run & 2016/17 GPD runs with recoil-proton detector
Transverse runs without recoil-proton detector
- **DVCS**: *transverse extension of partons in the proton*
 - t -slope of DVCS cross section in the kinematic domain between the other fixed-target experiments and HERA ep collider
 - Azimuthal asymmetry analysis ongoing
- **DVMP**: *input for GPD constraints, in particular chiral-odd GPDs*
 - Transverse target spin asymmetries for ρ^0 and ω vector mesons
 - SDMEs for ρ^0 and ω vector mesons
 - π^0 cross section
- More data are being analyzed.

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M.Gorzellik, PhD thesis, University of Freiburg (2018)

A. Sandacz and P. Sznajder, arXiv:1207.0333[hep-ph]

GK model:

S. V. Goloskokov and P. Kroll, Eur. Phys. J. C 42 (2005) 281

S. V. Goloskokov and P. Kroll, Eur. Phys. J. C 53 (2008) 367

S. V. Goloskokov and P. Kroll, Eur. Phys. J. C 65 (2010) 137

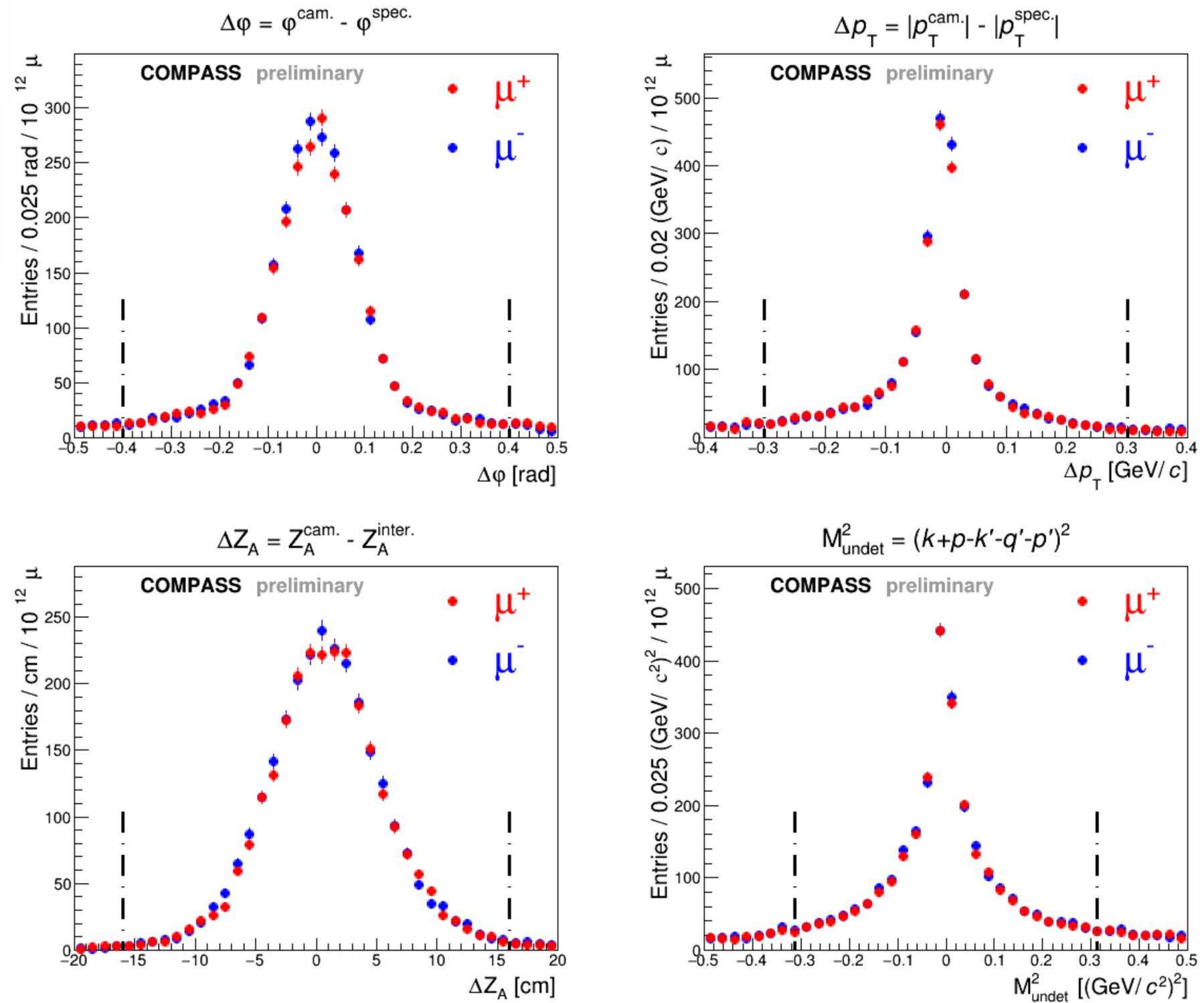
KM15 model:

K. Kumericki and D. Müller, Nucl. Phys. B 841 (2010) 1

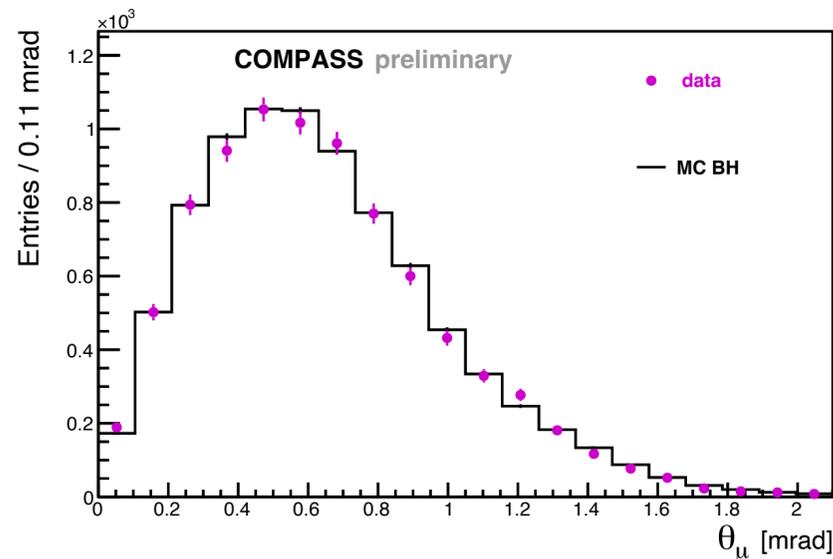
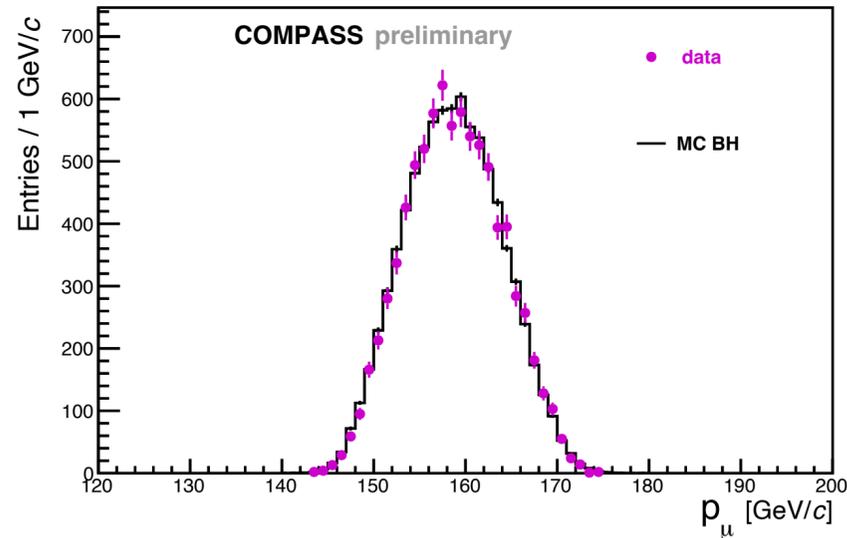
K. Kumericki and D. Müller, EPJ Web Conf. 112 (2016) 01012

Backup

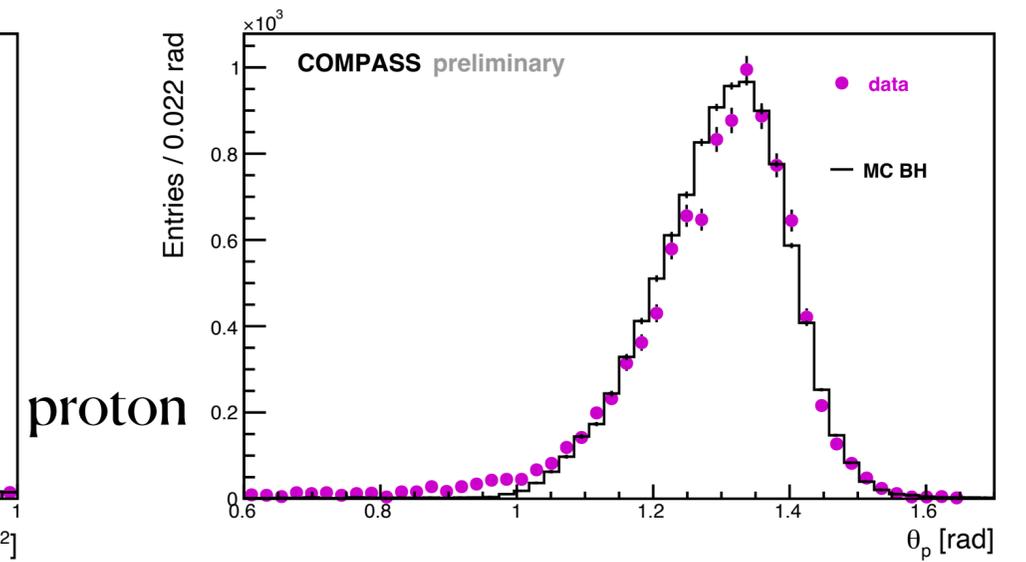
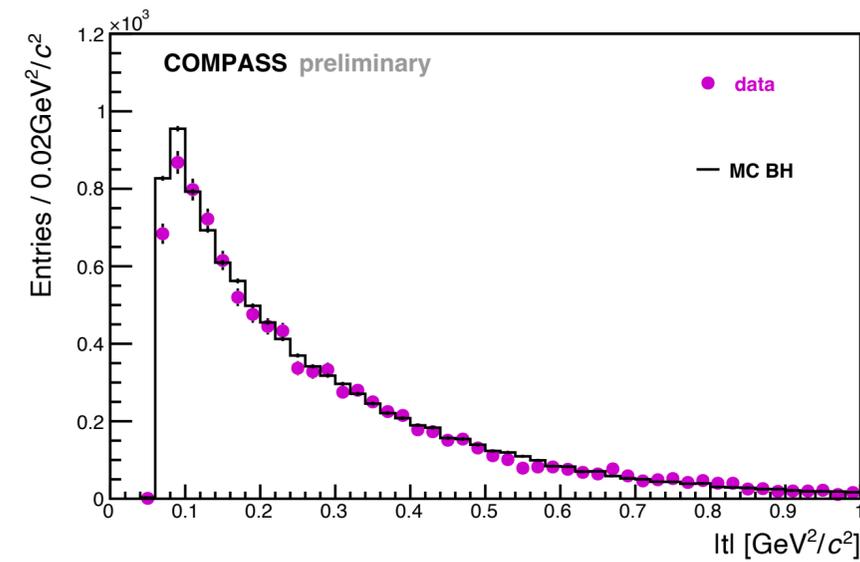
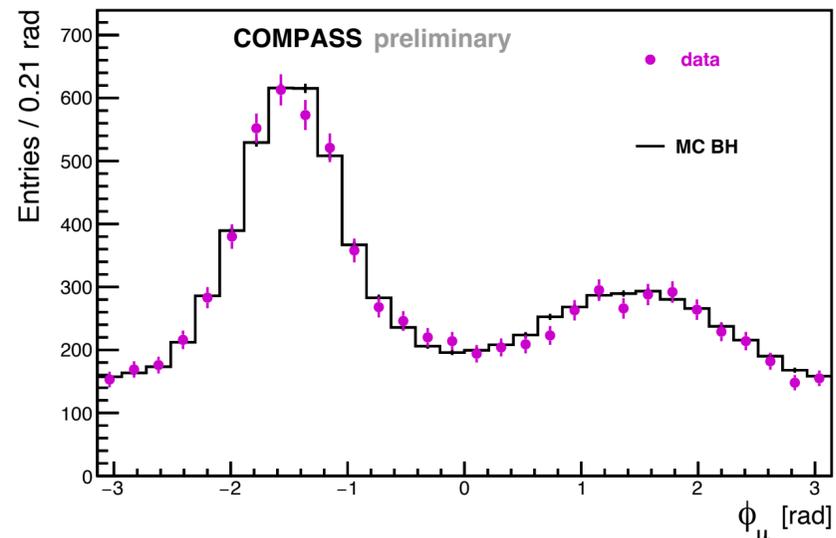
COMPASS DVCS 2016 data - 1D exclusivity cuts



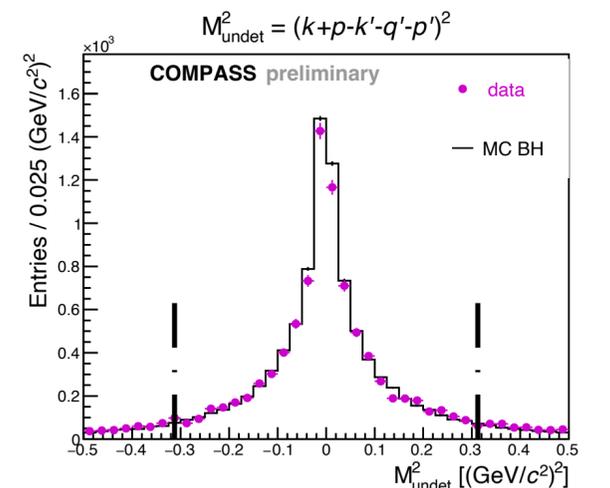
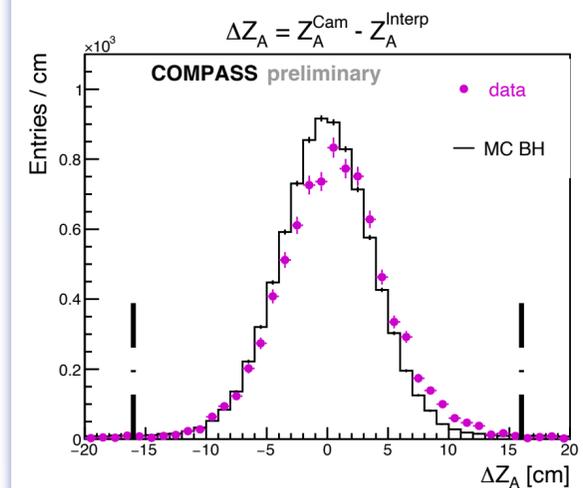
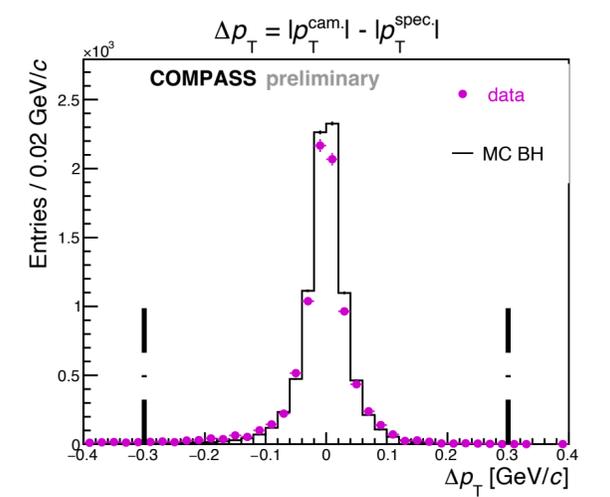
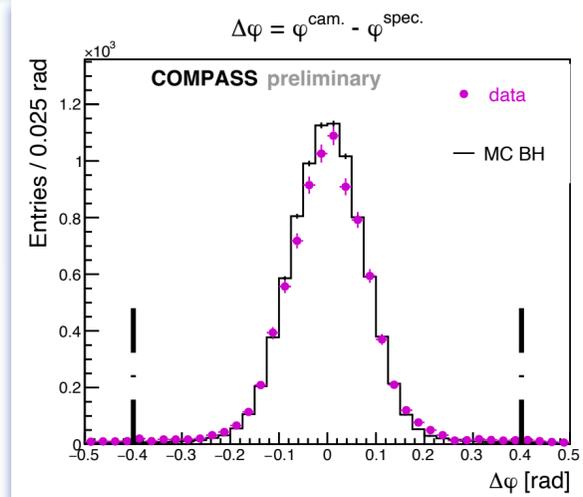
COMPASS DVCS 2016 data - BH data vs. MC, $80 < \nu < 144 \text{ GeV}$



(incoming muon)



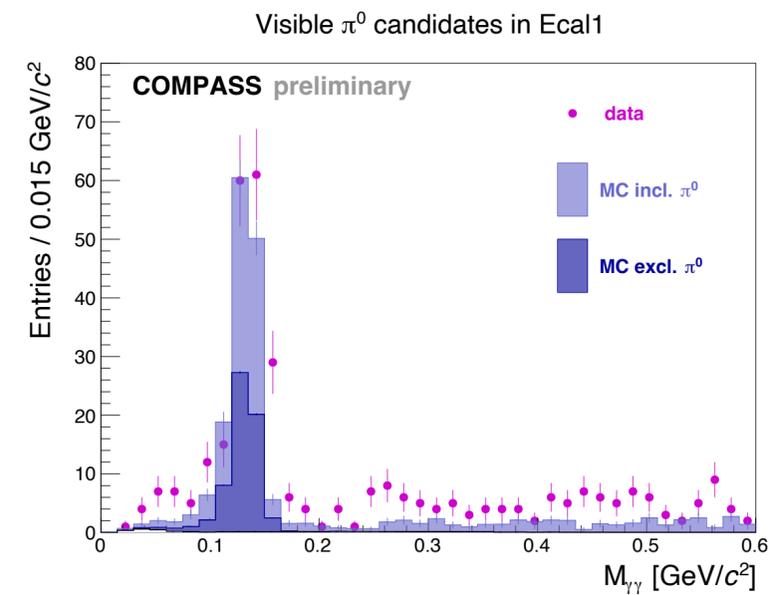
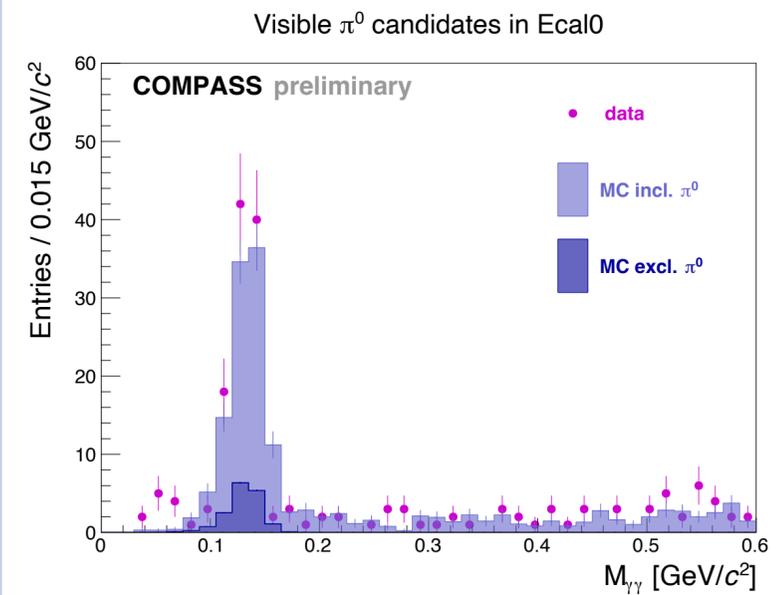
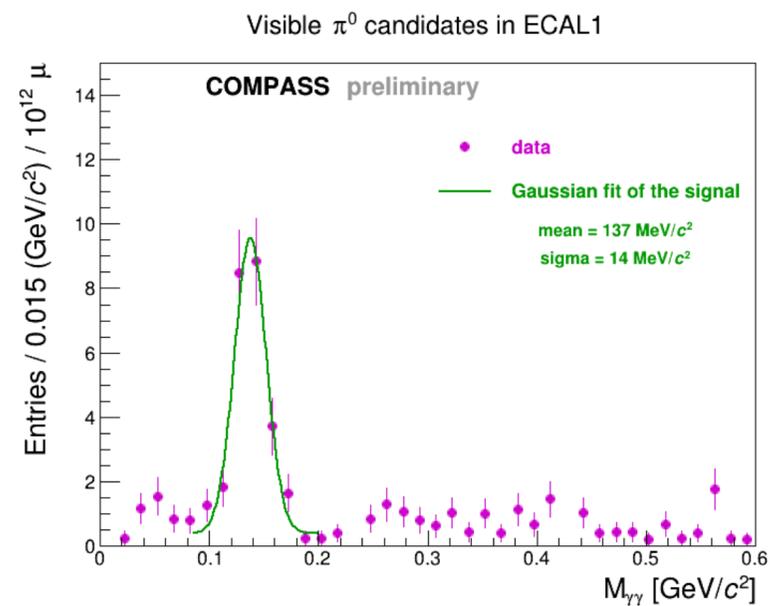
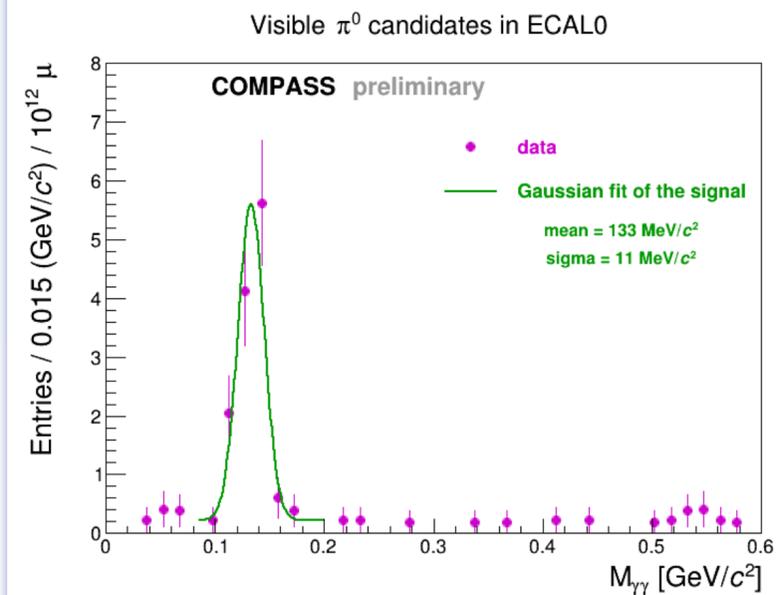
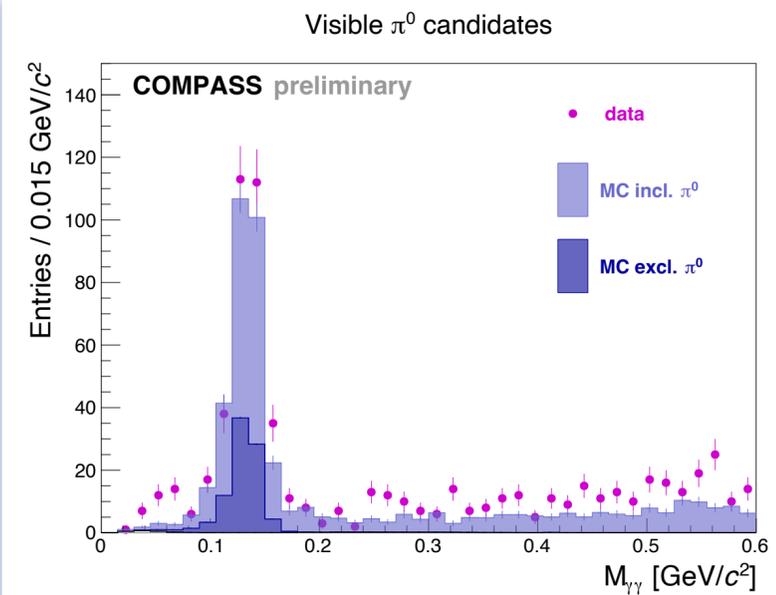
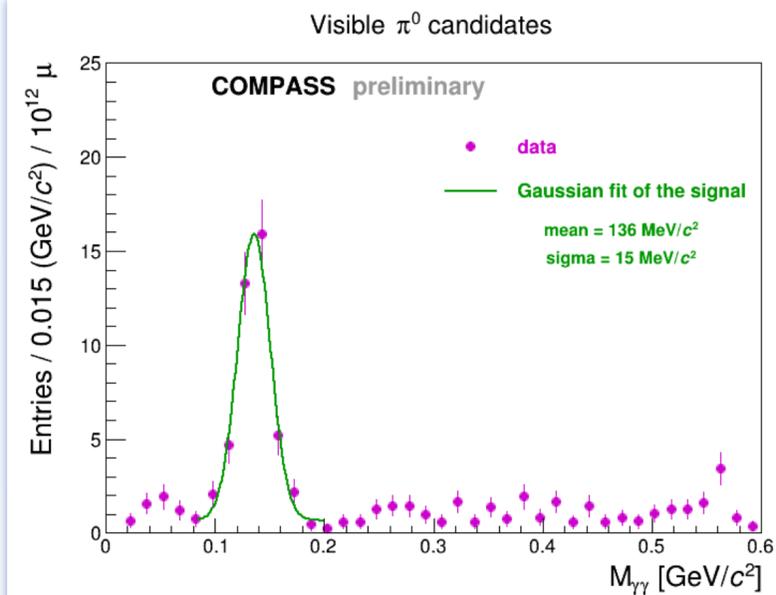
proton



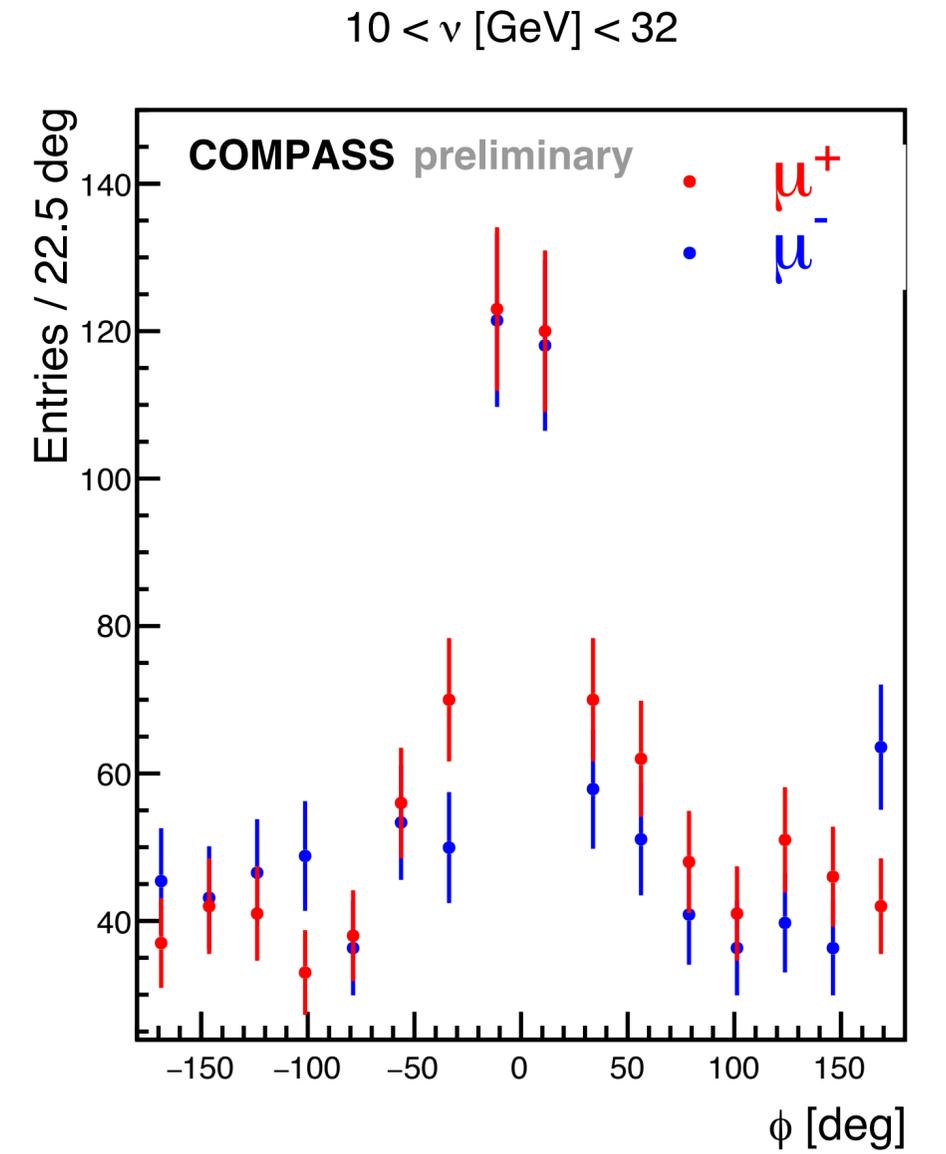
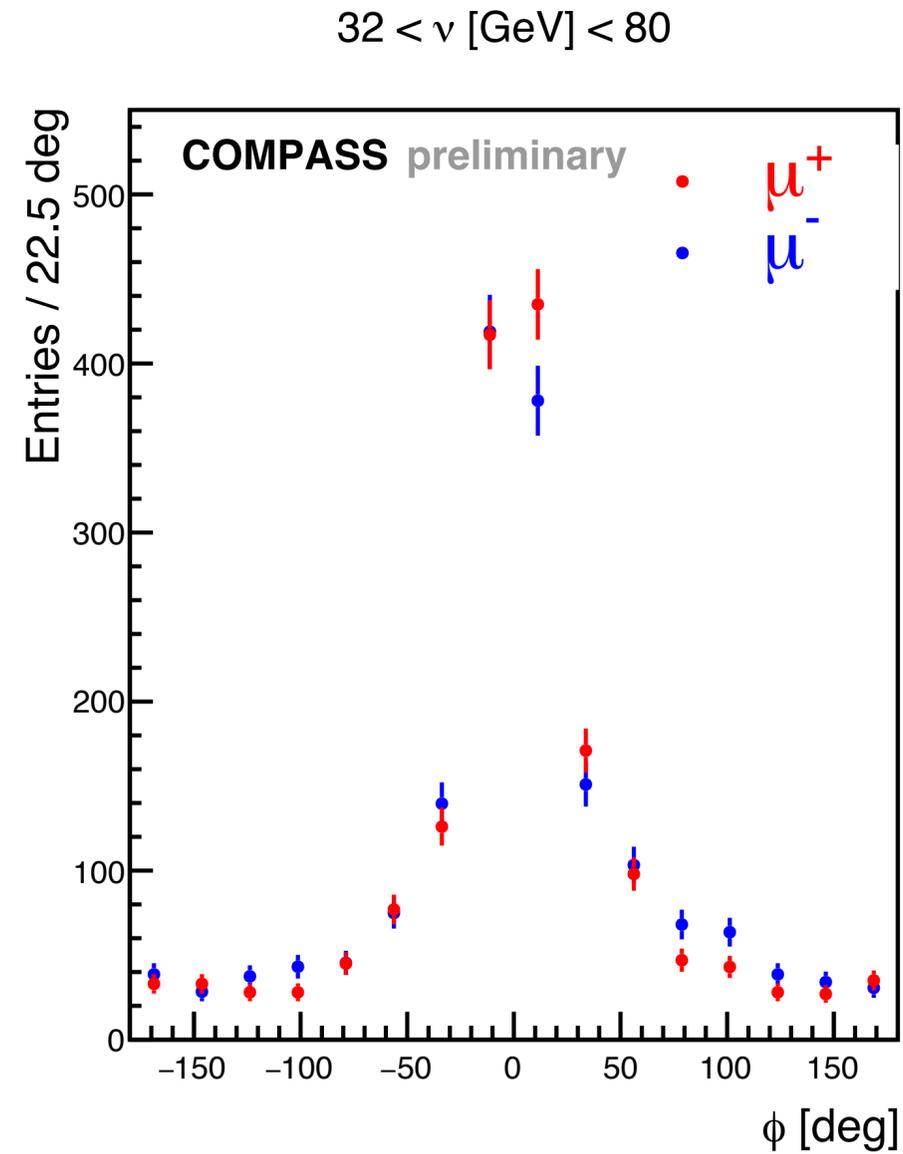
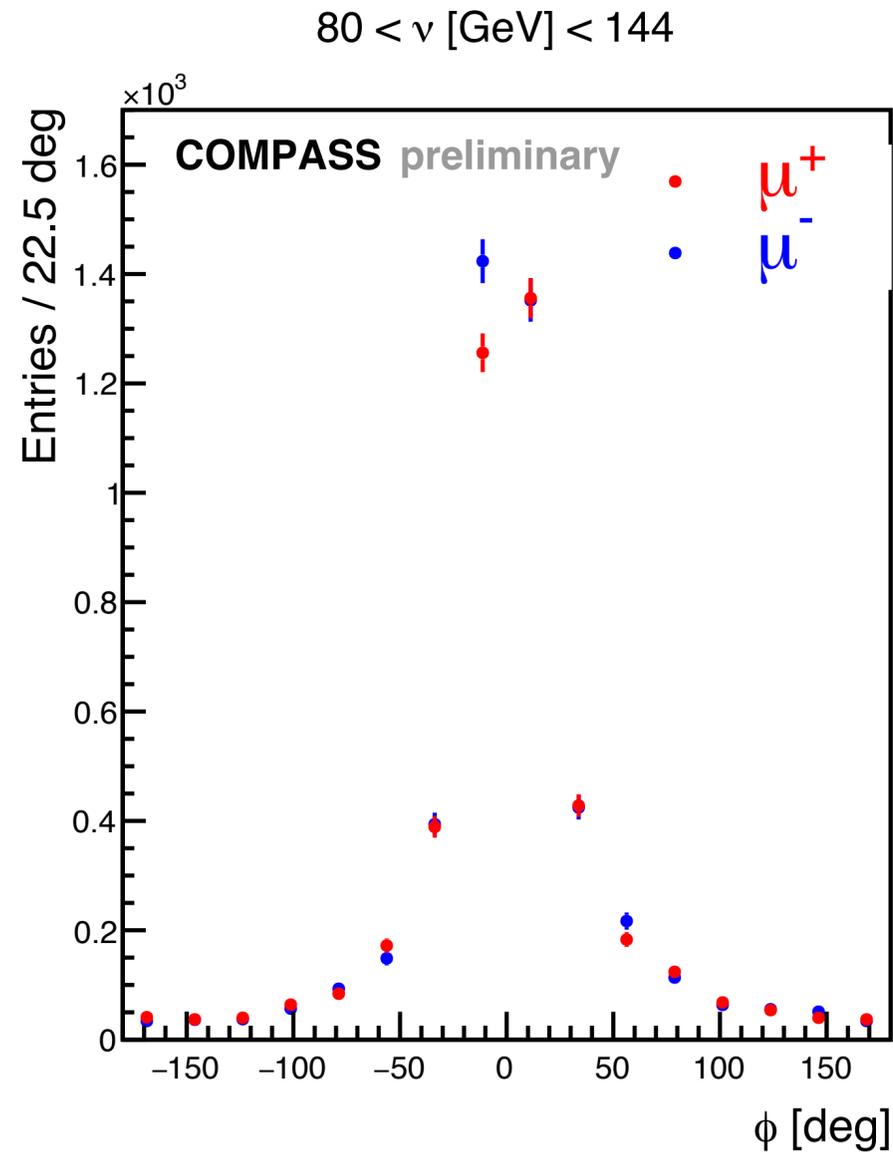
exclusive variables

COMPASS DVCS 2016 data - invariant mass of visible π^0

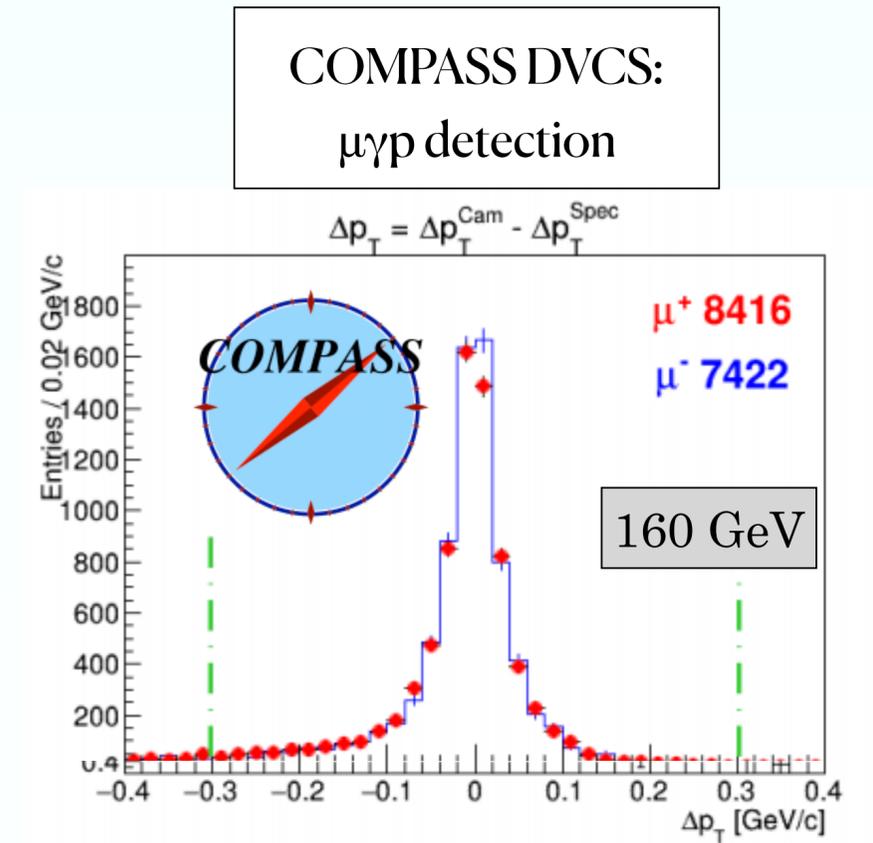
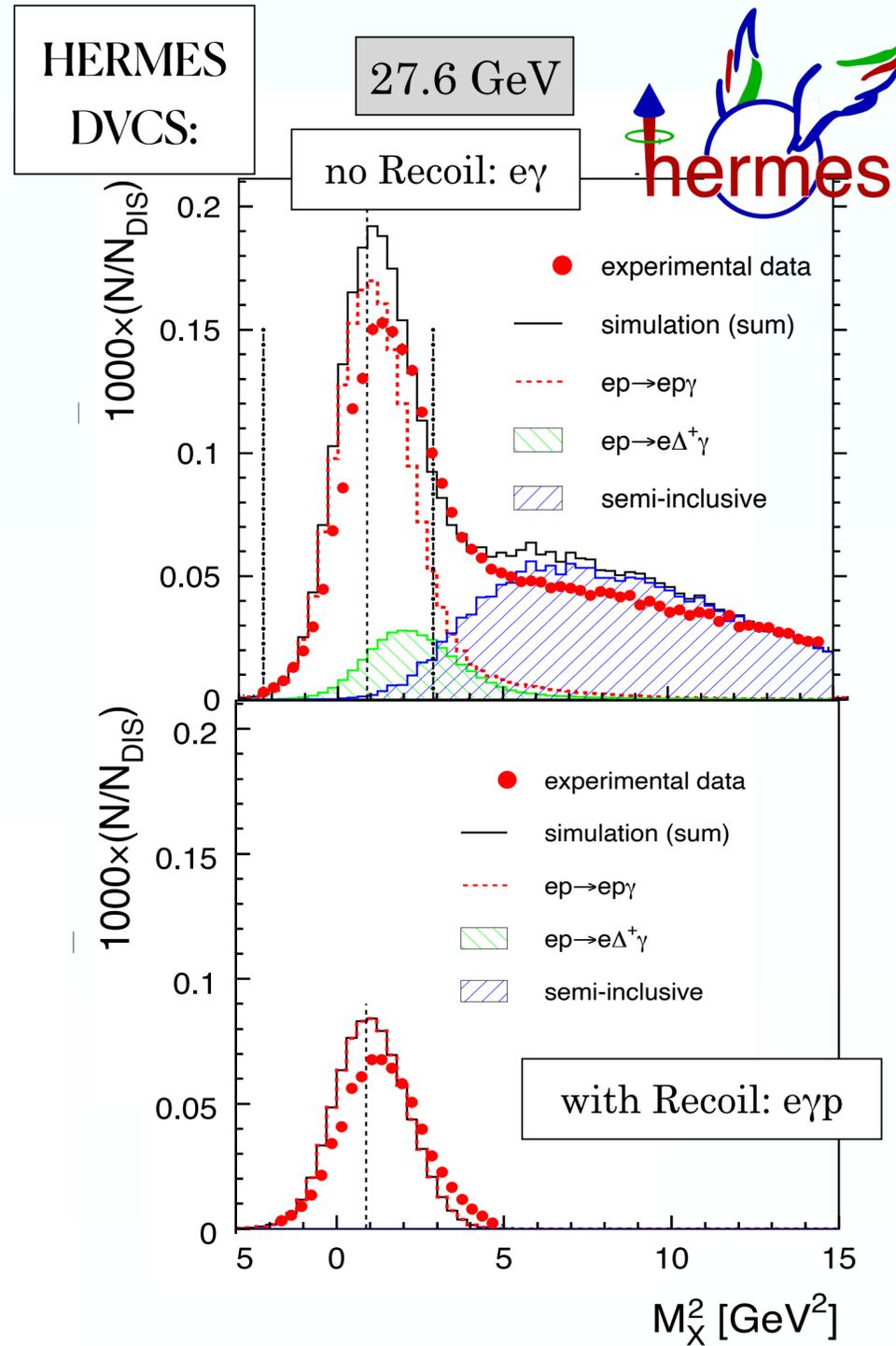
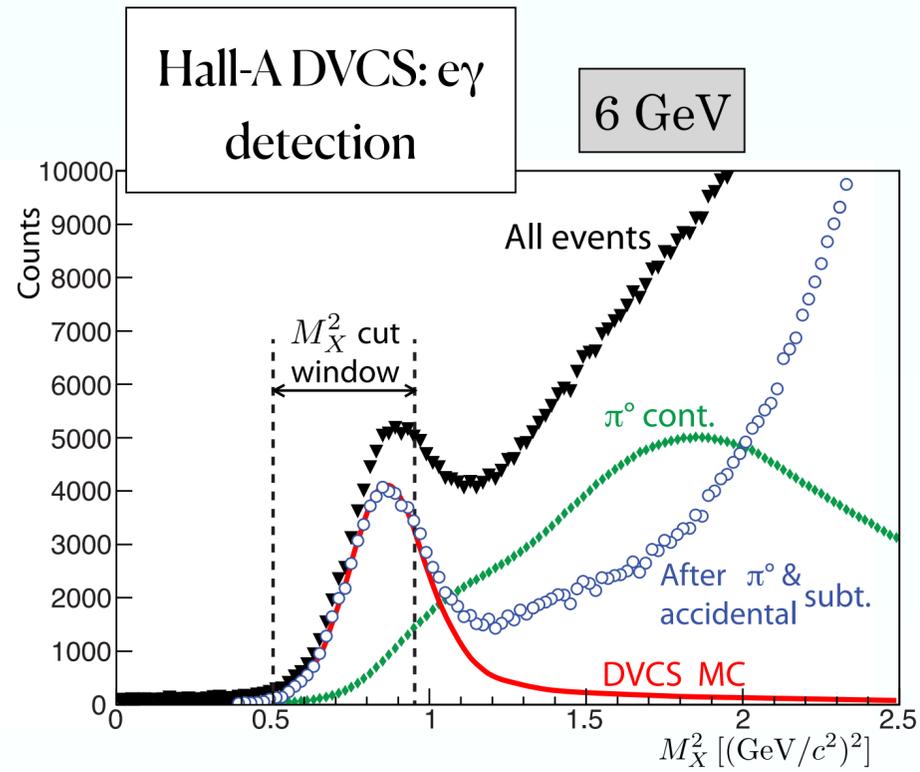
The visible π^0 contamination is removed in the ϕ -distribution



COMPASS DVCS 2016 data - $\phi_{\gamma\gamma^*}$ distributions

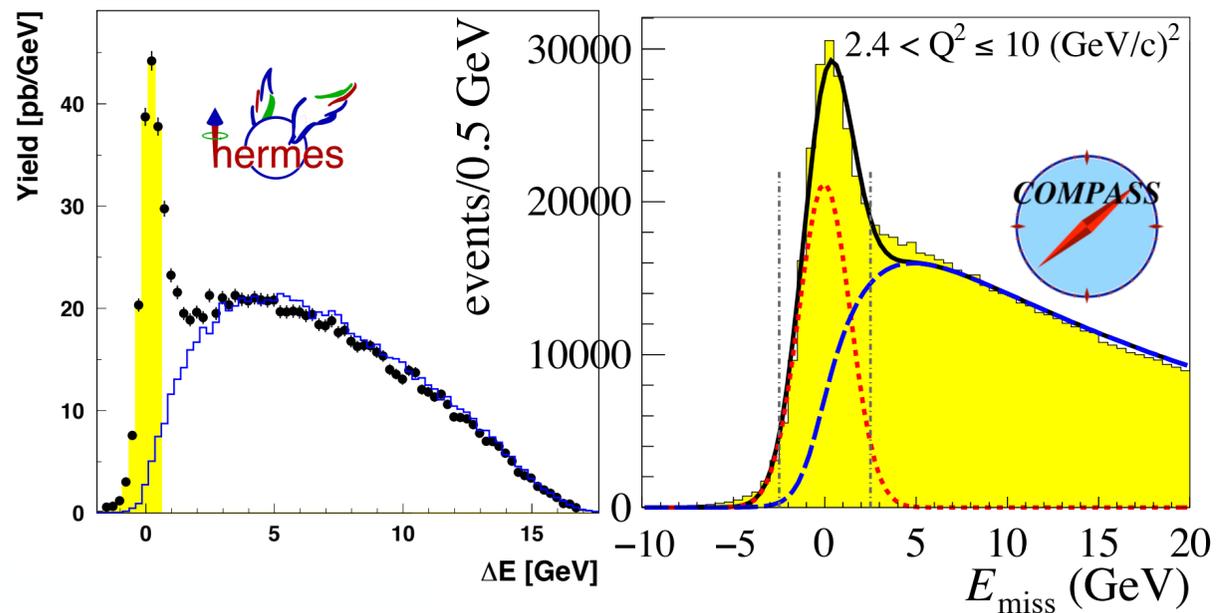


Selection of exclusive data sample



HERMES & COMPASS excl. ρ :

$\ell\pi-\pi^+$ detection



CLAS DVCS:

no Inner Calo: ep or $e\gamma p$

with Inner Calo: $e\gamma p$



H1/ZEUS DVCS:

$e\gamma$ + forward veto

ZEUS subsample: $e\gamma p$



Experimental access to CFFs at HERMES & JLab

$$\text{DVCS} \quad \sigma_{\gamma^* \gamma N} \sim \left| \begin{array}{c} \text{DVCS} \\ + \\ \text{Bethe-Heitler (BH)} \end{array} \right|^2$$

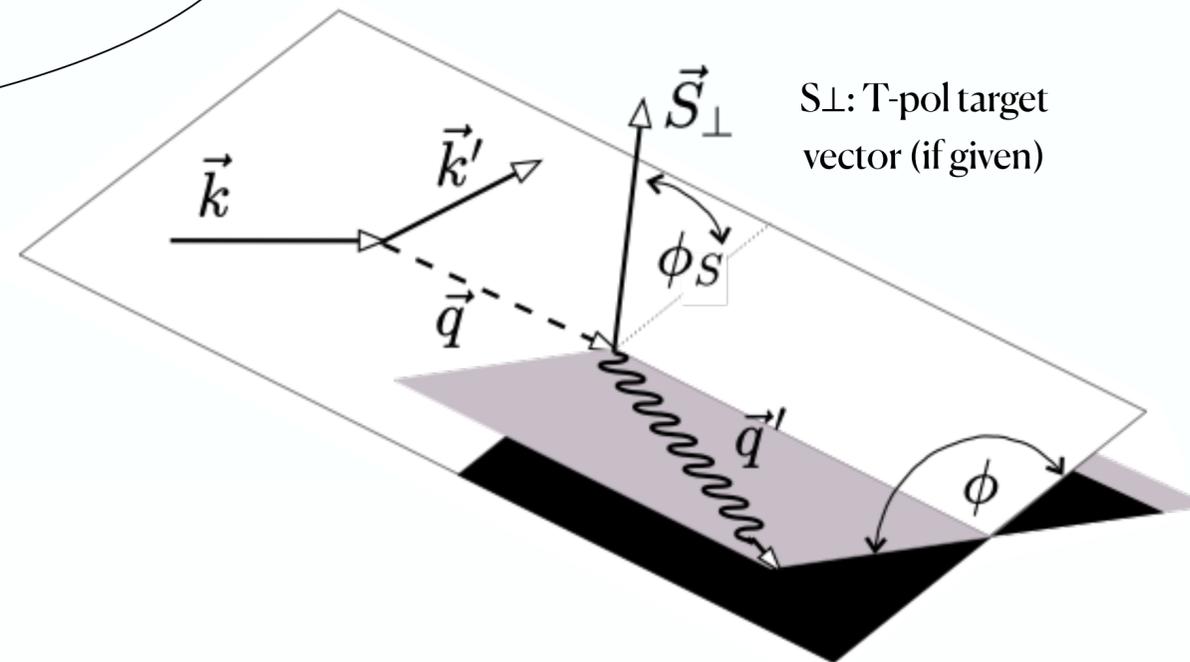
$$= |\mathcal{T}_{\text{BH}}|^2 + \boxed{(\mathcal{T}_{\text{DVCS}} \mathcal{T}_{\text{BH}}^* + \mathcal{T}_{\text{DVCS}}^* \mathcal{T}_{\text{BH}})} + |\mathcal{T}_{\text{DVCS}}|^2$$

The DVCS-Bethe_Heitler interference term helps to disentangle $\text{Re}(\tau_{\text{DVCS}})$ and $\text{Im}(\tau_{\text{DVCS}})$ magnitude and phase of DVCS amplitude τ_{DVCS}

Example: azimuthal beam-spin asymmetry

$$\mathcal{A}_{\text{LU}}(\phi)$$

\downarrow Beam \downarrow Target
 polarization

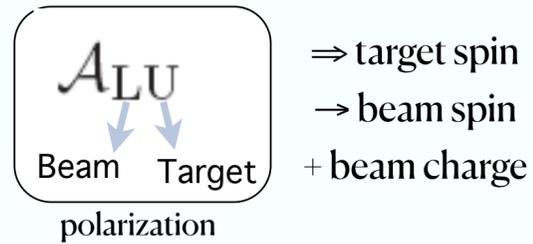


Lepton k with charge C_B & polarization P_B off nucleon
 Azimuthal angle ϕ between lepton scattering plane and plane spanned by virtual & real photons

$$\sigma(\phi; P_B, C_B) = \sigma_{\text{UU}}(\phi) \cdot [1 + P_B \mathcal{A}_{\text{LU}}^{\text{DVCS}}(\phi) + P_B C_B \mathcal{A}_{\text{LU}}^{\text{I}}(\phi) + C_B \mathcal{A}_C(\phi)]$$

Different experimental configurations (beam polarization, beam charge, target polarization, and their combinations) provide access to different parts or aspects of CFFs.

Different experimental configurations to map out CFFs



unpolarized target:

Best access

$$F_1 \mathcal{H} + \frac{x_B}{2 - x_B} (F_1 + F_2) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}$$

↓ dominant for the proton
↓ dominant for the neutron

$$\mathcal{A}_{LU}(\phi) \equiv \frac{d\sigma^{\rightarrow} - d\sigma^{\leftarrow}}{d\sigma^{\rightarrow} + d\sigma^{\leftarrow}}$$

Beam-helicity asymmetry
More Fourier coefficients accessible with 2 beam charges

$$\mathcal{A}_C(\phi) \equiv \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-}$$

Beam-charge asymmetry

Im(\mathcal{H})
Re(\mathcal{H})

Compton Form Factors (CFFs)

longitudinally polarized target:

$$\frac{x_B}{2 - x_B} (F_1 + F_2) \left(\mathcal{H} + \frac{x_B}{2} \mathcal{E} \right) + F_1 \tilde{\mathcal{H}} - \frac{x_B}{2 - x_B} \left(\frac{x_B}{2} F_1 + \frac{t}{4M^2} F_2 \right) \tilde{\mathcal{E}}$$

$\mathcal{A}_{UL}(\phi, e_\ell) \equiv$ Longitudinal target-spin asymmetry

$$\frac{[\sigma^{\leftarrow \Rightarrow}(\phi, e_\ell) + \sigma^{\rightarrow \Rightarrow}(\phi, e_\ell)] - [\sigma^{\leftarrow \Leftarrow}(\phi, e_\ell) + \sigma^{\rightarrow \Leftarrow}(\phi, e_\ell)]}{[\sigma^{\leftarrow \Rightarrow}(\phi, e_\ell) + \sigma^{\rightarrow \Rightarrow}(\phi, e_\ell)] + [\sigma^{\leftarrow \Leftarrow}(\phi, e_\ell) + \sigma^{\rightarrow \Leftarrow}(\phi, e_\ell)]}$$

analog: Double-spin (LL) asymmetry

transversely polarized target:

$$\frac{t}{4M^2} \left[(2 - x_B) F_1 \mathcal{E} - 4 \frac{1 - x_B}{2 - x_B} F_2 \mathcal{H} \right]$$

$$\mathcal{A}_{UT}^{\text{DVCS}}(\phi, \phi_S) \mathcal{A}_{UT}^{\text{I}}(\phi, \phi_S)$$

Transverse target-spin asymmetry

$$\mathcal{A}_{LT}^{\text{I}}(\phi, \phi_S) \mathcal{A}_{LT}^{\text{BH+DVCS}}(\phi, \phi_S)$$

Double-spin (LT) asymmetry

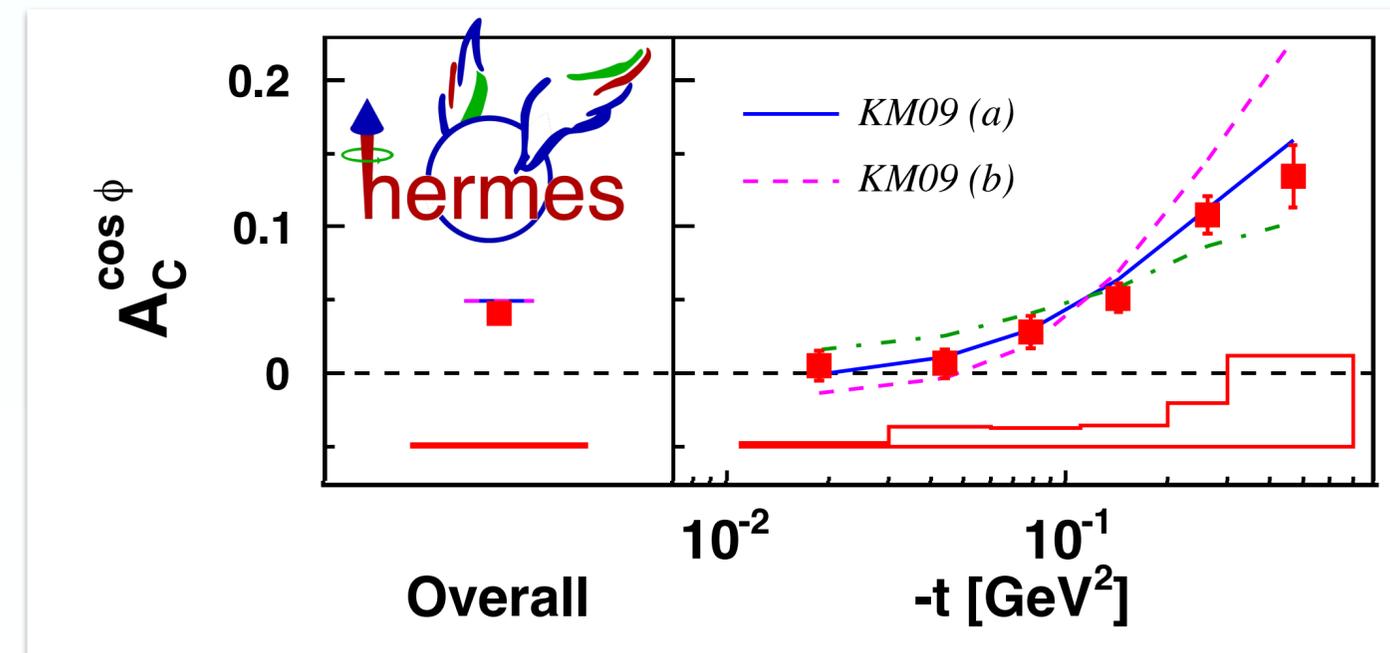
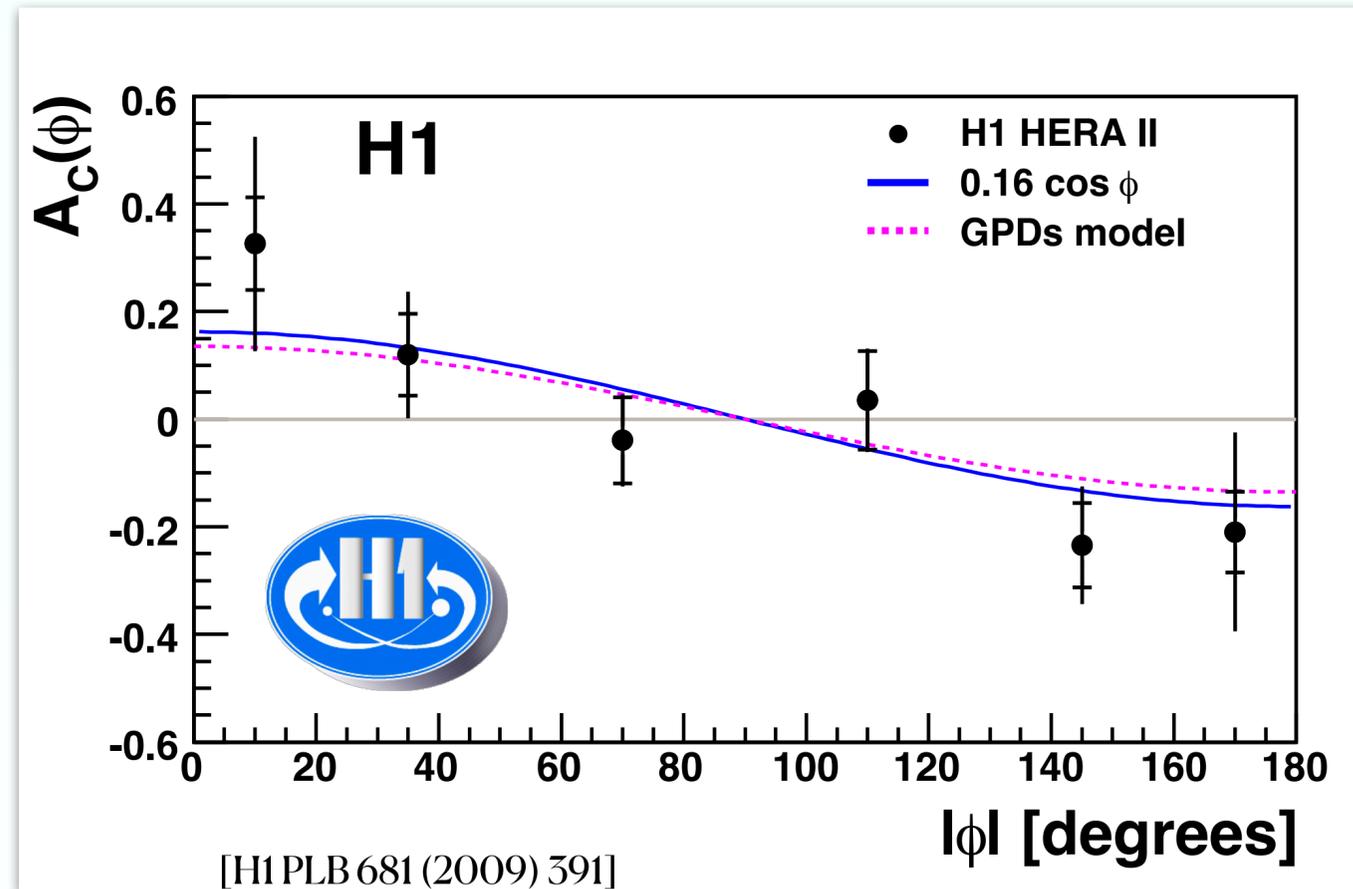
H1 & HERMES: DVCS beam-charge asymmetry

$\text{Re}(\tau_{\text{DVCS}}) > 0$
for HERA (small x)

$\text{Re}(\tau_{\text{DVCS}}) < 0$
for HERMES (larger x)

Where is the zero crossing? COMPASS
measurement at intermediate energy

- $\rho = \text{Re}(\tau_{\text{DVCS}}) / \text{Im}(\tau_{\text{DVCS}})$
 - $\rho = 0.20 \pm 0.05(\text{stat}) \pm 0.08(\text{sys})$
 - In good agreement with theoretical calculation (dispersion relation)
- H1@HERA/DESY: first and only measurement at collider
 - low $x_B = 10^{-4} \dots 10^{-2}$
 - $6.5 < Q^2 < 80 \text{ GeV}^2$
 - $30 < W < 140 \text{ GeV}$
 - $|t| < 1 \text{ GeV}^2$



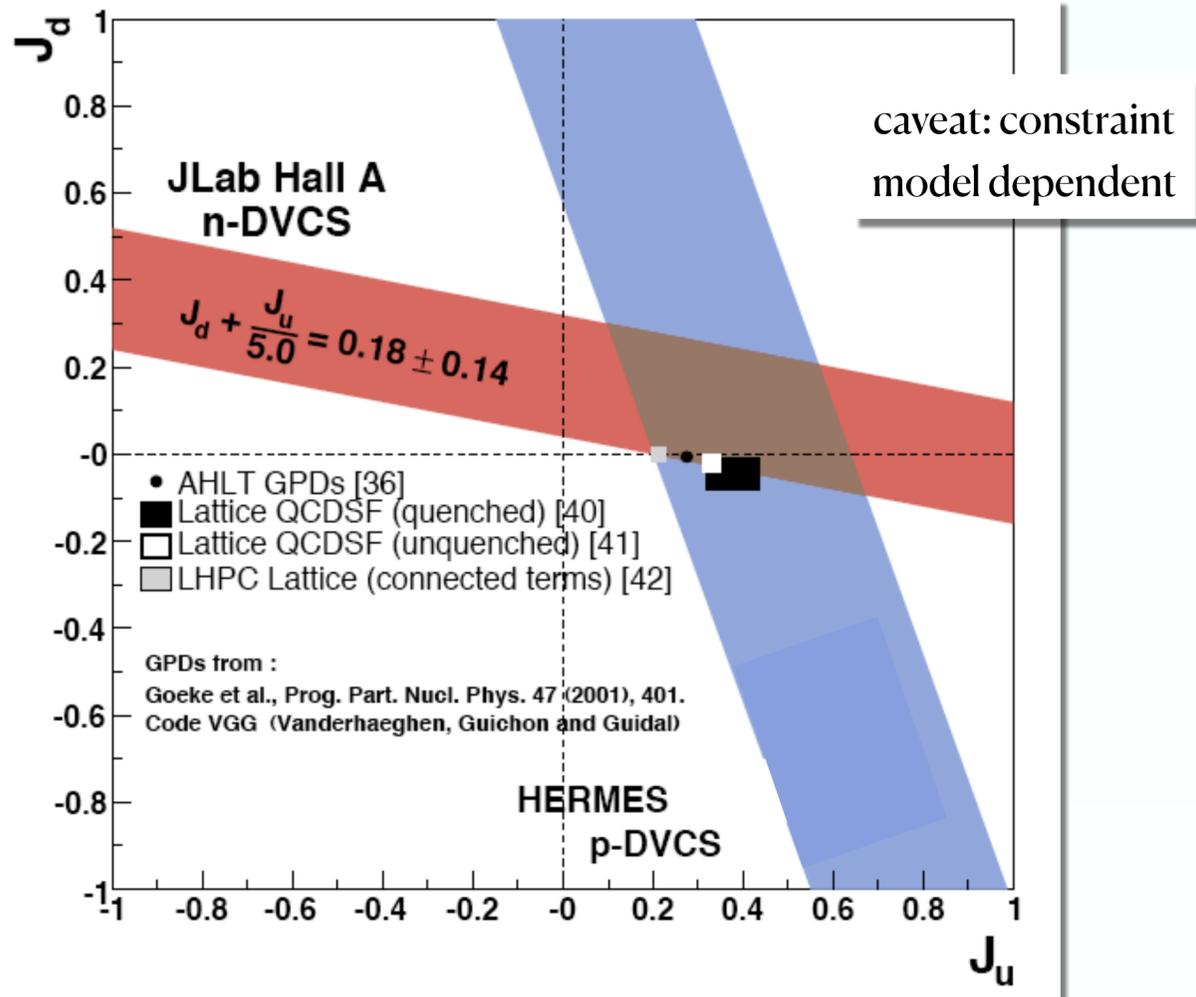
GPD E linked to orbital angular momentum

Measurements sensitive to GPD E allow (in principle) to access the total angular momentum of partons, J_q .

Ji sum rule for the nucleon:

[Ji, PRL 78 (1997) 610]

$$J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H^q(x, \xi, t) + E^q(x, \xi, t)]$$



[Hall A PRL 99, 242501 (2007)]

caveat: constraint model dependent

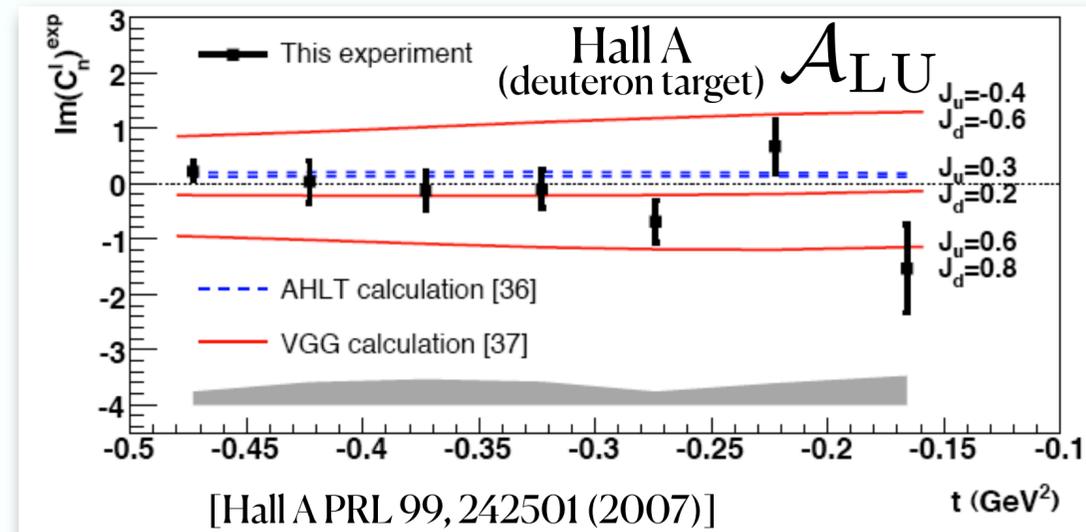
HERMES and Hall-A DVCS asymmetries

(A) HERMES: $ep^\uparrow \rightarrow ep\gamma$:

$H-E$ (transversely polarized proton target) A_{UT}

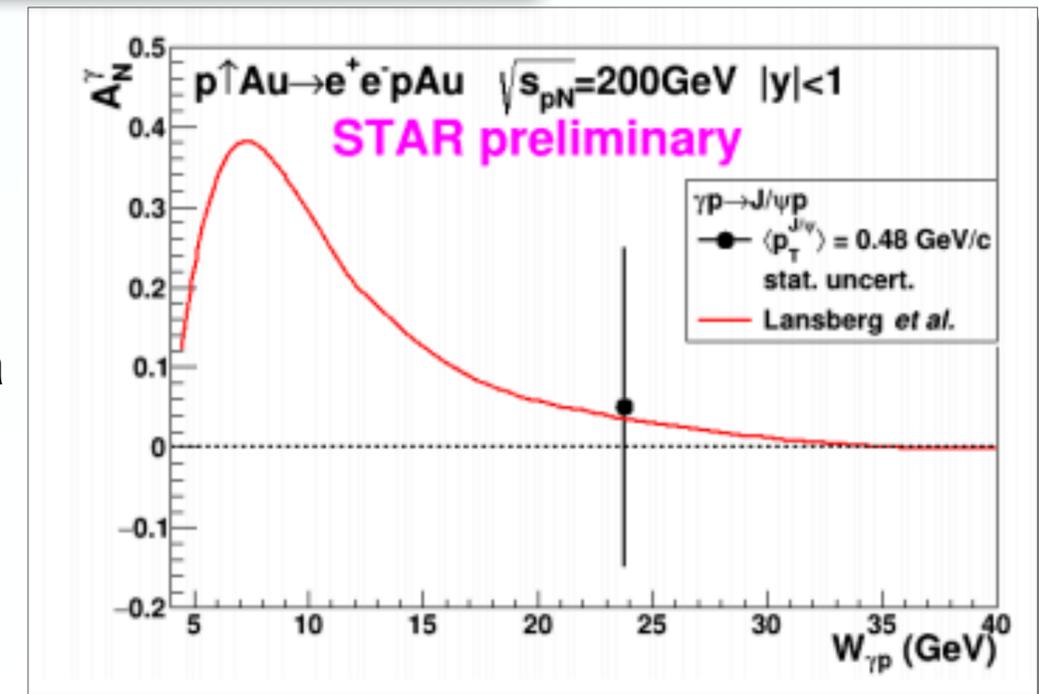
(B) Hall A: $e^-n \rightarrow e^-n\gamma$:

E dominant for the neutron A_{LU}



STAR exclusive J/Psi asymmetry

- More data from JLab12 and RHIC (STAR) to come.
- STAR: exclusive J/Psi production in ultra-peripheral $p^\uparrow p$ collisions (UPC) \rightarrow gluon GPD E



Vector meson production and decay

