



Craig Roberts ... <u>http://inp.nju.edu.cn/</u>





The Nature of Things

- In looking at the known Universe, one could be awed by the many complex things it contains.
- Even the Earth itself is complicated enough to generate questions in the minds of we observers
- > Basic amongst them are those which focus on our own very existence and composition.
- Here, too, there are many levels to be explored, running right down to the nuclei at the core of every atom and molecule
- > Even deeper, to the neutrons and protons (nucleons) that constitute those nuclei.
- Faced with all this, physicists (scientists) nevertheless assume that a few simple mathematical rules should be sufficient to provide a complete explanation of everything we can now perceive, and which might become perceptible in future.
- This may be correct
- Or it might be hubris



L'ANIMAL-MACHINE

L'animal-machine est une hypothèse éthologique selon laquelle les animaux sont des machines. Comme les machines, les animaux seraient des assemblages de pièces et rouages, dénués de conscience ou de pensée. D'un point de vue religieux, l'application du mécanisme à la vie revient à nier l'âme

des bêtes qui périssent donc entièrement au moment de leur mort. Poussée à l'extrême, notamment par Nicolas Malebranche, cette conception implique que leurs cris et gémissements ne peuvent être que le reflet de dysfonctionnements dans les « rouages » plutôt que l'expression d'une souffrance. Même si cette vision du problème est complètement décalée par rapport à la vision moderne, elle peine à être délogée par des conceptions plus en adéquation avec les avancées scientifiques récentes.



In the nineteenth century, Descartes was revered for his mechanistic physiology and theory that animal bodies are machines (that is, are constituted by material mechanisms, governed by the laws of matter alone).

Reductionist Perspective

Holism = Emergentism

- One might define emergent phenomena as those features of Nature which don't readily admit explanation solely in terms of known or conjectured rules.
- The concept is at least as old as Aristotle (384-322 BC), who argued that a compound item can have (emergent) properties in the whole which are not explicable merely through the independent actions of the item's constituent parts
- His view is often represented by the statement: "The whole is more than the sum of its parts".
- In this sense, emergence has its origins in the Greek "sunergos": "together" plus "working" = origin of our concept of synergy, viz. things working together more effectively than could be anticipated from their independent actions in isolation.
- This perspective is typically contrasted with that described as reductionism; namely, the view that everything in Nature can ultimately be viewed as no more complex, in principle, than,

e.g. a (very good) watch, which is clearly a complex object; but, equally clearly, not more than the sum of its parts.





Standard Model of Particle Physics

- Standard Model (SM) offers a description of all known fundamental physics except for gravity, and gravity is something that has no discernible effect when particles are studied a few at a time.
- Since LHC's discovery of the Higgs in 2012, the Higgs Boson has been promoted to the Centre of Things
- Standard Model has

17 particles and 19 parameters, most of which relate to the Higgs and all of which must be determined through comparison with experiment

SM supposedly describes most powerful forces in Nature; yet, somewhat unsatisfactory



Standard Model of Particle Physics

- Strong Interactions in the Standard Model are supposed to be described by quantum chromodynamics (QCD)
- Only two of the parameters are intrinsic to QCD
 - Higgs enters through current-quark masses
- > One of them $-\theta_{QCD}$ appears to be zero (exactly or almost) ... know this because nucleon EDM is unmeasurably small
- Just one parameter remains to be fixed
- ➢ Perhaps science has a chance of <u>understanding</u> QCD ∈ SM?



Emergent Phenomena ... in the Standard Model(?)

Existence of our Universe depends critically on, *inter alia*, the following empirical facts:

- Proton is massive
 - *i.e.*, the mass-scale for strong interactions is vastly different to that of electromagnetism
- Proton is absolutely stable
 - Despite being a composite object constituted from three valence quarks
- Pion is unnaturally light (not massless, but lepton-like mass)
 - Despite being a strongly interacting composite object built from a valence-quark and valence antiquark



Emergence: low-level rules producing high-level phenomena, with enormous apparent complexity

Emergence of Hadron Mass

Standard Model of Particle Physics has one *known* mass-generating mechanism

- = Higgs Boson ... impacts are critical to evolution of Universe as we know it
- > However, Higgs boson is alone responsible for just ~ 1% of the visible mass in the Universe
- Proton mass budget

Only 9 MeV/939 MeV is directly from Higgs

Evidently, Nature has another very effective mechanism for producing mass:

Emergent Hadron Mass (EHM)

- ✓ Alone, it produces 94% of the proton's mass—
- Remaining 5% is generated by constructive interference between EHM and Higgs-boson

ASS re son 882 • chiral limit mass • EHM+HB feedback = HB current mass

proton mass budget



8 (78)

Emergence of Hadron Mass - Basic Questions

- > What is the origin of EHM?
- Does it lie within the Standard Model, i.e., within QCD
- What are the connections with ...
 - Gluon and quark confinement?
 - Dynamical chiral symmetry breaking (DCSB)?
 - Nambu-Goldstone modes = $\pi \& K$?
- What is the role of Higgs in modulating observable properties of hadrons?
 - Critically, without Higgs mechanism of mass generation, π and K would be indistinguishable
- What is and wherefrom mass?



FIG. 1.1. Mass budgets for A-proton, B-kaon and C-pion, drawn using a Poincaré invariant decomposition. There are crucial differences. The proton's mass is large in the chiral limit, *i.e.* even in the absence of Higgs couplings into QCD. This nonzero chiral-limit component is an expression of emergent hadronic mass (EHM) in the SM. Conversely and yet still owing to EHM via its dynamical chiral symmetry breaking (DCSB) corollary, the kaon and pion are massless in the chiral limit – they are the SM's Nambu-Goldstone modes [24–27]. (See Eq. (2.22) below.) (Units MeV, separation at $\zeta = 2$ GeV, produced using information from Refs. [8, 21–23].)





All mass is interaction.

— Richard P. Feynman —

In QCD, so is the absence of mass



It's not just about MASS



EIC Early Career Researcher Workshop 2022 July 24-25

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Emergent Hadron Mass (EHM)

- > In Nature, mass and length⁻¹ are effectively the same thing
- Thus, asking for the origin of 99% of visible mass in the Universe is possibly/probably the same as asking what is the source of the proton's size

Confinement scale!

- Confinement is far more than the statement that Nature contains only colour-singlet combinations of gluons and quarks
 - Those combinations have fm-scale sizes this is crucial
 - It wouldn't be confinement if the scales were Å size

Confinement ... only guarantee of proton stability

- If one can't measure them, what are the gluons and quarks in the QCD Lagrangian?
- Are they anything more than a theoretical artifice; useful things for calculations in perturbation theory, but practically irrelevant when resolving detectable hadrons?
- What degrees-of-freedom should be used to compute and *understand* hadron properties?



Quantum Chromodynamics

$$L = \frac{1}{4} G^a_{\mu\nu}(x) G^a_{\mu\nu}(x) + \bar{\psi} \left[\gamma \cdot \partial_x + m + ig \, \frac{\lambda^a}{2} \gamma \cdot A^a(x) \right] \psi(x)$$
$$G^a_{\mu\nu}(x) = \partial_\mu A^a_\nu(x) - \partial_\nu A^a_\mu(x) - f^{abc} A^b_\mu(x) A^c_\nu(x)$$

One-line Lagrangian – expressed in terms of gluon and quark partons
 Which are NOT the degrees-of-freedom measured in detectors
 Questions

- What are the asymptotic detectable degrees-of-freedom?
- How are they built from the Lagrangian degrees-of-freedom?
- Is QCD really the theory of strong interactions?
- > Is QCD really a theory? \Rightarrow Implications far beyond Standard Model



GENESIS



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Pinch Technique: Theory and Applications Daniele Binosi & Joannis Papavassiliou Phys. Rept. 479 (2009) 1-152





- $\blacktriangleright \text{ Running gluon mass}$ $d(k^2) = \frac{\alpha(\zeta)}{k^2 + m_g^2(k^2;\zeta)}$ $\alpha_s(0) = 2.77 \approx 0.9\pi, \ m_g^2(0) = (0.46 \,\text{GeV})^2$
- Gluons are *cannibals* a particle class whose members become massive by eating each other!

Recent reviews ...

Emergent Hadron Mass in Strong Dynamics Daniele Binosi, e-Print: 2203.00942 [hep-ph] Few Body Syst. 63 (2022) 2, 42

Emergence of mass in the gauge sector of QCD Joannis Papavassiliou, arXiv:2207.04977 [hep-ph] To appear in Chin. Phys. C.



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IR Behaviour of QCD



Gluon mass scale

- > QCD permits a gluon mass scale to emerge the character of QCD is sufficient
 - However, the emergence of a gluon mass scale is not necessary
- > The magnitude of the gluon mass-scale is not contained within the SM
 - On the other hand, we exist; hence, it can be measured
- > QCD is a renormalizable quantum field theory
 - Perturbation theory is a valid tool for high-energy phenomena
 - Can fix the scale by something perturbative, *e.g.* the value of QCD's running coupling at the mass of the τ lepton
 - Once that number is fixed, then nonperturbative tools can predict the value of the gluon mass function on the entire domain of spacelike momenta



Modern Understanding Grew Slowly from *Ancient* Origins

More than 40 years ago

Dynamical mass generation in continuum quantum chromodynamics, J.M. Cornwall, Phys. Rev. D **26** (1981) 1453 ... ~ 1050 citations



➤ Owing to strong self-interactions, gluon partons ⇒ gluon quasiparticles, described by a mass function that is large at infrared momenta



Truly mass from nothing An interacting theory, written in terms of massless gluon fields, produces dressed gluon fields that are characterised by a mass function that is large at infrared momenta



- ✓ QCD fact
- Continuum theory and lattice simulations agree

✓ Empirical verification?



Today: Gluon Mass Function

Gluon mass function characterised by renormalisation group invariant IR mass

 $\widehat{m}_0 = 0.43(1) \text{ GeV} \approx \frac{1}{2}m_{\text{proton}}$

- > The value is a prediction in the sense that it is $\frac{1}{2}m_{\text{proton}}$
- Follows that long-wavelength gluons are screened from interactions
- Such screening eliminates the Landau pole
- Solves the Gribov ambiguity

Locating the Gribov horizon

Fei Gao (Peking U. and Peking U., SKLNPT), Si-Xue Qin (Chongqing U.), Craig D. Roberts (Argonne, PHY), Jose Rodriguez-Quintero (Huelva U.) (Jun 14, 2017) Published in: *Phys.Rev.D* 97 (2018) 3, 034010 • e-Print: 1706.04681 [hep-ph]

🖹 pdf 🛛 🤗 DOI 🛛 🖃 cite



J. Papavassiliou

2014 total downloads



https://www.youtube.com/watch?v=xlkvz2V8ZRw



Effective charge from lattice QCD, Zhu-Fang Cui, Jin-Li Zhang et al.,

NJU-INP 014/19, arXiv:1912.08232 [hep-ph], Chin. Phys. C 44 (2020)

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More info



OCD's Running Coupling

"Asymptotic freedom comes with a flip side"

- "Imaginatively called infrared slavery"
- Based on perturbative analyses, appearance of a Landau pole at infrared momenta, became common to assume coupling runs to infinity as k² runs to zero
- 2.0 1.5 Perturbative-QCD running coupling extrapolated onto infrared domain $k^2 < 1 \text{ GeV}^2$ 0.5 0.0 k^2 / GeV^2
- Many people interpreted/still interpret this as being synonymous with confinement
- *Is the notion correct?*





(78)

EUROPEAN CENTRE FOR THEORETICAL STUDIES IN NUCLEAR PHYSICS AND RELATED AREAS AND Physics (DSEMP2014) Trento, Italy, September 22-26, 2014

Process independent effective charge = running coupling

- Modern theory enables unique QCD analogue of "Gell-Mann – Low"
 - running charge to be rigorously defined and calculated
- Analysis of QCD's gauge sector yields a parameter-free prediction
- > N.B. Qualitative change in $\hat{\alpha}_{Pl}(k)$ at $k \approx \frac{1}{2} m_p$
- No Landau Pole
 - "Infrared Slavery" picture is not correct
- Below k ~ m̂₀, interactions become scale independent, just as they were in the Lagrangian; so, QCD becomes practically conformal again

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The QCD Running Coupling,

A. Deur, S. J. Brodsky and G. F. de Teramond, Prog. Part. Nucl. Phys. 90 (2016) 1-74

Process independent strong running coupling

Daniele Binosi et al., arXiv:1612.04835 [nucl-th], Phys. Rev. D 96 (2017) 054026/1-7

Effective charge from lattice QCD, Zhu-Fang Cui et al., NJU-INP 014/19, <u>arXiv:1912.08232</u> [hep-ph], <u>Chin. Phys. C **44** (2020) 083102/1-10</u>

Perturbation Theory and the Onset of Nonperturbative Physics



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Chiral limit = essentially nonperturbative

EHM Basics

> Absent Higgs boson couplings, the Lagrangian of QCD is scale invariant

≻ Yet ...

- Massless gluons become massive
- A momentum-dependent scale-expressing charge is produced
- Massless quarks become massive
- EHM is expressed in
 - EVERY strong interaction observable
- Challenge to Theory =

Elucidate all observable consequences of these phenomena and highlight the paths to measuring them

Challenge to Experiment =

Test the theory predictions so that

the boundaries of the Standard Model can finally be drawn





QCD Fact Pion (Nambu-Goldstone modes) and mass

- → Higgs boson couplings \rightarrow 0
- Pion exists and is massless
- Pion Bethe-Salpeter amplitude

EHM demands equivalence between one-body mass and two-body correlation strength in Nature's most fundamental Nambu-Goldstone bosons



Pion wave function

quark mass function

Thus, enigmatically, properties of the nearly massless pion are the cleanest expression of EHM in the Standard Model !







$$\langle p(P)|T_{\mu\nu}|p(P)\rangle = -P_{\mu}P_{\nu}$$

$$\langle p_0(P)|T_{\mu\mu}|p_0(P)\rangle = -P^2 = m_{p_0}^2 = \langle p_0(P)|\Theta_0|p_0(P)\rangle$$

Prima Facie, the QCD Lagrangian CANNOT support a nonzero value of this expectation value.

QCD's scale anomaly



G. Krein (Sao Paulo, IFT), T.C. Peixoto (Sao Paulo, IFT and Unlisted, BR) (Nov 23, 2020) Published in: *Few Body Syst.* 61 (2020) 4, 49 • e-Print: 2011.11615 [hep-ph]

^{49 • e-Print: 2011.11615 [hep-ph]} EHM Measurements How does the mass of the nucleon arise?

- One suggestion = QCD scale anomaly
- Measurement = femtoscopic low-momentum correlation measurements
- > QCD multipole expansion ... J/ψ and/or Y chromopolarisability relates forward scattering amplitude to a key matrix element in the origin of the nucleon mass problem

viz. the average chromoelectric gluon distribution in the nucleon:

$$\langle N | \left[(g\boldsymbol{E}^{a})^{2} - (g\boldsymbol{B}^{a})^{2} \right] | N \rangle = -\frac{1}{2} \langle N | g^{2} G^{a}_{\mu\nu} G^{a\mu\nu} | N \rangle = \frac{16\pi^{2}}{9} m_{N} \leq \langle (g\boldsymbol{E})^{2} \rangle$$

Once considered viable ...

electromagnetic vector-meson production reaction ($V = J/\psi, \Upsilon$) $e = e + p \rightarrow e' + V + p$

to access purely hadronic process

$$V + p \rightarrow V + p$$



 $\gamma^*(Q)$

 A_{V_p}



Vector-meson production & Vector meson dominance



Vector-meson production and vector meson dominance, Yin-Zhen Xu (徐胤禛), Si-Yang Chen (陈斯阳), Zhao-Qian Yao (姚照千), D. Binosi, Z.-F. Cui (崔著钫) and C. D. Roberts, <u>NJU-INP 044/21</u>, <u>arXiv:2107.03488 [hep-ph]</u>, Eur. Phys. J. C 81 (2021) 895/1-11

Trace Anomaly via Heavy-meson photo-production?

> Is there a relation between the electromagnetic vector-meson production reaction

 $e + p \rightarrow e' + V + p$ and the purely hadronic process $V + p \rightarrow V + p$?

- > Answer depends on the fidelity of single-pole VMD assumption
- If this VMD assumption is false ...
 - Then all existing cross-section estimates and conclusions based on VMD must be reviewed
- > For EIC this means ... will need to rethink access to and understanding of
 - hidden-charm pentaquark production
 - origin of the proton mass via trace anomaly



Photon Vacuum Polarisation

- Consider the photon vacuum polarisation. Statement of simple fact: $Q^2 \Pi(Q^2) \stackrel{Q^2+m_V^2 \simeq 0}{=} \bar{e}_V^2 \frac{2f_V^2 \bar{m}_V^2}{Q^2 + m_V^2}$ decay constant
- > Thus, on $Q^2 + m_V^2 \simeq 0$ the timelike photon is indistinguishable from the vector-meson.
- This is merely the statement that e+ e- collisions with a tuned centre-of-mass energy can be used to produce vector-mesons.
- ► <u>However</u>: VMD assumption asserts that the photon is also indistinguishable from the vector-meson on $Q^2 \simeq 0$: $Q^2 \Pi(Q^2) = \frac{2\bar{e}_V^2 f_V^2}{m_V^2}$
 - Very different statement
- Plainly false because photon is massless & associated Ward-Green-Takahashi identity demands:

$$Q^2 \Pi(Q^2) \Big|_{Q^2 \simeq 0} \equiv 0 \quad \neq \frac{2\bar{e}_V^2 f_V^2}{m_V^2}$$



Photon-quark vertex

- Can one save VMD by looking at photon-quark vertex ?
- \succ This statement is always true: $\epsilon^{\lambda} \cdot \Gamma^{\gamma}(k;Q) \stackrel{Q^2 + m_V^2 \simeq 0}{=} \frac{2f_V m_V}{Q^2 + m_V^2} \Gamma_V^{\lambda}(k;Q)$
- \succ Are there any conditions under which $\Gamma_{\nu}^{\gamma}(k;Q)|_{Q^2 \simeq 0}$ has a link to an on-shell vector-meson that may be approximated by the VMD current-field identity?

➢ No.

Example. VMD validity would require

this ratio to be unity:

Leading term in vector meson bound-state amplitude

$$R_V(k^2; Q^2 = 0) := \frac{2f_V}{m_V} \frac{F_1^0(k^2; -m_V^2)}{G_1^0(k^2; Q^2 = 0)}$$

Leading term in photon-quark vertex





 $\gamma^*(Q)$

 A_{V_n}

Might be workable for light mesons but plainly false for heavy mesons Craia Roberts: cdroberts@nju.edu.cn 411 "On Mass and Matter



34 (78) Vector-meson production and vector meson dominance, Yin-Zhen Xu (徐胤禛), Si-Yang Chen (陈斯阳), Zhao-Qian Yao (姚照千), D. Binosi, Z.-F. Cui (崔著钫) and C. D. Roberts, <u>NJU-INP 044/21</u>, <u>arXiv:2107.03488 [hep-ph]</u>, Eur. Phys. J. C 81 (2021) 895/1-11

EHM Measurements EHM Measurements How does the mass of the nucleon arise?

- > Many other tests failed by VMD assumption for heavy vector-mesons:
- See also ...
 - Deciphering the mechanism of near-threshold J/ψ photoproduction, M.-L. Du, V.
 Baru, F.-K. Guo et al., Eur. Phys. J. C 80, 1053 (2020)
 - Perturbative QCD analysis of near threshold heavy quarkonium photoproduction at large momentum transfer, Peng Sun, Xuan-Bo Tong, Feng Yuan, Phys. Lett. B 822 (2021) 136655
- There is no objective, model-independent means by which to connect $e + p \rightarrow e' + V + p$ with $V + p \rightarrow V + p$
- Hence, vector meson photoproduction does not provide a path to the QCD trace anomaly
- Is there a viable alternative?





Era of Meson Targets

A elastic form factors $\gamma(q)$ $\gamma(q)$ $\gamma($



- JLab & EIC & EicC
 - High luminosity electron (+ ion) beams
 - Access to meson targets via the Sullivan Process,
 - i.e., a baryon's "meson cloud"
- AMBER @ CERN SPS
 - High-intensity beams of pions

($\geq 10^7$ pions/sec in Phase-1 = approved) and kaons (5 × 10⁶ kaons/sec Phase-2 = proposal

being prepared)

- Drell-Yan, J/ψ production, prompt photon production
 - ... from proton and nuclear targets


π&KDAs&Form Factors

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Wave Functions of Nambu Goldstone Bosons

- Physics Goals:
 - Pion and kaon distribution amplitudes (DAs $\varphi_{\pi,\kappa}$)
 - Nearest thing in quantum field theory to a Schrödinger wave function
 - Consequently, fundamental to understanding π and K structure.
- Scientific Context:
 - For 40 years, the x-dependence of the pion's dominant DA has been controversial.
 - Modern theory \Rightarrow EHM is expressed in the *x*-dependence of pion and kaon DAs
 - Pion DA is a direct measure of the dressed-quark running mass in the chiral limit.
 - Kaon DA is asymmetric around the midpoint of its domain of support (0<x<1)
 - Degree of asymmetry is signature of constructive interference between EHM and HB mass-generating mechanisms

Insights into the Emergence of Mass from Studies of Pion and Kaon Structure, Craig D. Roberts, David G. Richards, Tanja Horn and Lei Chang, NJU-INP 034/21, <u>arXiv: 2102.01765 [hep-ph]</u>, Prog. Part. Nucl. Phys. **120** (2021) 103883/1-65



Meson leading-twist DAs

- Continuum results exist & IQCD results arriving
 - Common feature = broadening
 - Origin = EHM
- > NO differences between $\pi \& K$ if EHM is all there is
 - Differences arise from Higgs-modulation of EHM mechanism
 - "Contrasting π & K properties reveals Higgs wave on EHM ocean"



FIG. 10. Fit of the $P_z = 4\frac{2\pi}{L}$ pion (left) and kaon (right) data to the analytical form in Bjorken-*x* space, compared with previous calculations (with only central values shown). Although we do not impose the symmetric condition m = n, both results for the pion and kaon are symmetric around x = 1/2 within error. Craig Roberts: cdroberts@nju.edu.cn 411 "On Mass and Matter"

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- Kaon DA vs pion DA
 - almost as broad
 - peak shifted to x=0.4(5)
 - $-\langle \xi^2 \rangle = 0.24(1), \langle \xi \rangle = 0.035(5)$
- ERBL evolution logarithmic
 - Broadening & skewing persist to <u>very</u> large resolving scales – beyond LHC

Hard Exclusive Processes e.g, meson form factors

Normalisation of large Q² behaviour of meson form factors is largely determined by EHM, expressed in

✓ f_M = order parameter for DCSB (EHM)

✓ $\varphi_M(x)$ = meson wave function

Parton distribution functions (PDFs) and parton distribution amplitudes (PDAs) are connected through the hadron's light-front wave function

$$\begin{split} \varphi(x) &\sim \int d^2 k_\perp \psi(x, k_\perp^2) \,, \\ q(x) &\sim \int d^2 k_\perp |\psi(x, k_\perp^2)|^2 \end{split}$$

Hadron's LFWF is the key.

- ✓ EHM is expressed in every hadron LFWF
- ✓ LFWF correlates all observables

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The impetus for measuring $F_M(Q^2)$ at large momentum transfers is a need to understand and validate a strict prediction of QCD [32–34], viz. $\exists Q_0 \gg \Lambda_{\text{QCD}}$ such that

$$Q^2 F_M(Q^2) \stackrel{Q^2 > Q_0^2}{\approx} 16\pi \alpha_s(Q^2) f_M^2 w_M^2(Q^2), \quad (1)$$

where α_s is the one-loop strong running coupling, $\Lambda_{\rm QCD} \approx 0.3 \,\text{GeV}, f_{\pi} = 0.092 \,\text{GeV}, f_K = 0.110 \,\text{GeV} [1];$ and $w_M^2 = e_{\bar{q}} w_{\bar{q}}^2(Q^2) + e_u w_u^2(Q^2),$

$$\boldsymbol{w}_f = \frac{1}{3} \int_0^1 dx \, \boldsymbol{g}_f(x) \, \varphi_M(x; Q^2) \,, \qquad (2)$$

$$g_u = \frac{1}{x}, \ g_{q^-} = \frac{1}{1-x}$$

- High-energy nuclear and particle physics have long focused on finding evidence for power-law scaling in experimental data.
- Certainly, an important first step; BUT it should be remembered that QCD is not found in scaling laws.
- Instead, since quantum field theory requires deviations from strict scaling, then QCD is found in the existence and character of scaling violations.



(78)

Wave Functions of Nambu Goldstone Bosons

- Scientific Context:
 - Today, continuum and lattice are delivering consistent predictions for $\varphi_{\pi,\kappa}(x)$
 - Can signature features of interference between EHM and HB, which are manifest in low-order Mellin moments of the DAs, be probed experimentally?
- Measurement

А

Elastic electromagnetic form factors

Breakaway from scaling at $Q^2 \approx 6 \ GeV^2$ Thereafter, Q^2 dependence

is largely described by hard-scattering formula with realistic DA Reveal QCD scaling violations for 1st time in hard elastic scattering





Hard Exclusive Processes

- Considering hard exclusive processes & focusing on meson electromagnetic form factors, modern theory predicts that
 - scaling violations become apparent to experiment on $Q^2 > 10 {\rm ~GeV^2}$
 - magnitude of any given form factor on a sizeable domain above $Q^2 = 10 \text{ GeV}^2$ determined by the physics of emergent hadronic mass (EHM).
- Experiments focused in this area are of the <u>greatest</u> importance.

- ✓ Location of maximum points to beginning of pQCD evolution = scaling violations
- Height of ratio measures relative strength of emergent- and Higgsgenerated mass for s- and u-quarks





Parton Distribution Functions



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NG Boson Distribution Functions

- Physics Goals:
 - Precise data that can be used to determine
 Pion and Kaon Distribution Functions valence, sea and glue
 - Provide the first complete charts of the internal structure of Nature's most fundamental Nambu-Goldstone bosons.
- > Today:
 - Existing pion data are more than 40-years-old
 - That data only covers the valence-quark domain
 - A forty-year controversy, with doubts persisting over whether the data agree with QCD predictions or challenge the truth of QCD
 - Regarding the kaon, worldwide, only 8 points of data exist



Insights into the Emergence of Mass from Studies of Pion and Kaon Structure, Craig D. Roberts, David G. Richards, Tanja Horn and Lei Chang, NJU-INP 034/21, <u>arXiv: 2102.01765 [hep-ph]</u>, Prog. Part. Nucl. Phys. **120** (2021) 103883/1-65

NG Boson Distribution Functions

- Physics Goals:
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 Pion and Kaon Distribution Functions valence, sea and glue



 Provide the first complete charts of the internal structure of Nature's most fundamental Nambu-Goldstone bosons.

> Future:

- JLab, EIC, EicC
 - ⇒ pion and kaon elastic electromagnetic form factors ... reveal and quantify scaling
 - violations in hard exclusive processes ... hard prediction of QCD, never seen
 - \Rightarrow pion and kaon valence quark distribution functions at large x_B
- AMBER
 - \rightarrow precision data to chart π and K structure: DFs of valence, sea and glue.
 - → Glue is particularly important ... because controversial, yet prominent theory predicts that pions contain (almost) zero glue.

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Controversy over pion valence DF

- QCD-improvement of parton model leads to the following statement:
 - At any $\zeta > \zeta_{\rm H}$ for which experiment can be interpreted through parton distributions, then

$$x\simeq 1 \Rightarrow q^{\pi}(x;\zeta) \propto (1-x)^{\beta=2+\gamma}, \gamma>0$$

Consequence

- Any analysis of DY, DIS, etc. experiment which returns β <2 conflicts with QCD.
- Observations
 - All existing internally-consistent calculations preserve connection between large-k² behaviour of interaction and large-x behaviour of DF: $J=0 \dots (1/k^2)^n \Leftrightarrow (1-x)^{2n}$
- > No existing calculation with n=1 produces anything other than $(1-x)^2$
- Internally-consistent calculation that preserve RG properties of QCD,
 - then 2 \rightarrow 2+ γ , γ >0, at any factorisation-valid scale

> Controversy:

Contemporary data fits yield (1-x)^{1+γ}

- Is QCD wrong?
- Is existing large-x data a true measure of $q^{\pi}(x; \zeta)$?
- Are fitter-favoured analysis methods correct/complete?

π valence-quark distributions 23 Years of Theory Evolution \rightarrow 2022

- Developments in continuum-QCD enabled 1st parameter-free predictions of valence, glue and sea distributions within the pion
 - Reveal that $u^{\pi}(x; \zeta)$ is <u>hardened</u> by EHM
- Novel lattice-QCD algorithms beginning to yield results for pointwise behaviour of $u^{\pi}(x; \zeta)$
- Agreement between
 new continuum prediction for u^π(x; ζ) [Ding:2019lwe]
 and recent lattice-QCD result [Sufian:2019bol]
- Real strides toward understanding pion structure.
- Standard Model prediction: stronger than ever before
- After 30 years new era dawning in which the ultimate experimental checks can be made: M2 beam-line @ CERN ... JLab12 ... EIC ... EicC

 $\beta^{\text{contm}}(\zeta_5) = 2.66(12)$ $\beta^{\text{lattice}}(\zeta_5) = 2.45(58)$



π DFs ... Modern Theory Predictions vs Traditional Phenomenological Fits to Data

> Valence:

- momentum fraction similar
- Phenomenological Fits ... profile much harder
 & inconsistent with QCD prediction
- ➤ Glue:
 - Qualitative similarities on $x \ge 0.05$, but marked quantitative disagreement, especially on complementary domain
 - Both continuum prediction and fit are very different from early phenomenology
 - Should be tested in new experiments that are directly sensitive to the pion's gluon content.
 - E.g., prompt photon & J/Ψ production
- Sea:
 - Prediction and fit disagree on entire x-domain
 - If pion's gluon content is considered uncertain, then fair to describe sea-quark distribution as empirically unknown
 - Motivation for the collection and analysis of DY data with π^{\pm} beams on isoscalar targets





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48 (78)

π DFs ... Modern Theory Predictions vs Traditional Phenomenological Fits to Data





49 (78)

Breaking news for glue in π : Continuum (Eur. Phys. J. C 80 (2020) 1064/1-20) & Lattice Predictions (arXiv: 2104.06372)

Two distinct methods for tackling QCD Agree quantitatively on $g^{\pi}(x)$

- Phenomenological analyses exhibit qualitatively different behaviour
- Highlights need for new data and improved phenomenology in order to turn that data into a real test of QCD and our understanding of Nambu-Goldstone modes.
- > AMBER @ CERN can provide the necessary precise data.



Regarding the distribution of glue in the pion, Lei Chang (常雷) and Craig D Roberts, <u>e-Print: 2106.08451 [hep-ph]</u>, <u>Chin. Phys. Lett.</u> **38** (8) (2021) 081101/1-6</u> - Editors' Suggestion

Status: Kaon

- ➤ Little empirical information available on K DFs ⇒ no recent phenom. inferences.
 - Valence-quark distributions: results from models and a single, recent IQCD study
 - Kaon's glue and sea distributions: <u>no results</u>
- > One piece of available experimental information:

 $\frac{u_K(x)}{u_\pi(x)}$

- Continuum prediction for ratio is consistent with data.
- But, given large errors, this ratio is very forgiving of even large differences between various calculations of the individual DFs used to produce the ratio.
 - New, precise data critical if this ratio to be used as path to understanding the Standard Model's Nambu-Goldstone modes;
 - Results for $u_{\pi}(x;\zeta_5)$, $u_{\kappa}(x;\zeta_5)$ separately = better.

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Programmes at JLab, AMBER Proposals at EIC and EicC

Tackling the kaon structure function at EicC Gang Xie et al. e-Print: 2109.08483 [hep-ph]. Chin. Phys. C 46 (2022) 6, 064107

Entirety of Human history

– only 8 kaon data in existenceCERN ... Phys. Lett. B 93 (1980) 354-356



51 *(78)*

Status: Kaon

- ➤ Little empirical information available on K DFs ⇒ no recent phenom. inferences.
 - Valence-quark distributions: results from models and a single, recent IQCD study
 - Kaon's glue and sea distributions: <u>no predictions</u> ... until now
- Glue and Sea Predictions:
 - DFs similar to those in the pion
 - Detailed comparison requires use of mass-dependent splitting functions.
 - Development underway ... Preliminary conclusions:
 - i. Light-front momentum fraction carried by s-quarks in the kaon increases by \sim 5%;
 - ii. Compensated by a commensurate decrease in fractions carried by glue (-1%) and sea (-2%).

Kaon parton distributions, Z.-F. Cui et al., arXiv:2006.14075 [hep-ph], <u>Eur. Phys. J. C 80 (2020) 1064/1-20</u>



Proton & Deep Inelastic Scattering

- The ratio of neutron and proton structure functions at large x is keen discriminator between competing pictures of proton structure
- > Example:
 - Only scalar diquark in the proton (no axial-vector): $\lim_{x \to 1} \frac{F_2^n(x)}{F_2^p(x)} = \frac{1}{4}$
 - No correlations in the proton wave function (SU(4) spin-flavour) $\lim_{x \to 1} \frac{F_2^n(x)}{F_2^p(x)} = \frac{2}{3}$
- Experiments have been trying to deliver reliable data on this ratio for fifty years!
- MARATHON a more-than ten-year effort, using a tritium target at JLab, has delivered precise results

D. Abrams, et al., Measurement of the Nucleon Fn2/Fp2 Structure Function Ratio by the Jefferson Lab MARATHON Tritium/Helium-3 Deep Inelastic Scattering Experiment – arXiv:2104.05850 [hep-ex], Phys. Rev. Lett. **128** (2022) 132003

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FIG. 2: The F_2^n/F_2^p ratio plotted versus the Bjorken x from the JLab MARATHON experiment. Also shown are JLab Hall B BoNuS data [56], and a band based on the fit of the SLAC data as provided in Ref. [46], for the MARATHON kinematics $[Q^2 = 14 \cdot x \text{ (GeV}/c)^2]$ (see text). All three experimental data sets include statistical, point to point systematic, and normalization uncertainties.



Valence Quark Ratio in the Proton Zhu-Fang Cui, Fei Gao, Daniele Binosi, Lei Chang, Craig D. Roberts, and Sebastian M. Schmidt Chin. Phys. Lett. 2022, 39 (4): 041401 . DOI: 10.1088/0256-307X/39/4/041401 Abstract HTML The PDF (571KB)

- New mathematical method for interpolation and extrapolation of data
 - based on continued-fraction representation of functions, augmented by statistical sampling
- Delivers model-independent prediction for all ratios
 - No reference to models or physics theories
- Provides benchmark against which all pictures of nucleon structure can be measured
- Probability that scalar diquark only models of nucleon might be consistent with available data is 1/7,000,000

by statistical sampling





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Valence quark ratio in the proton, Zhu-Fang Cui (崔著钫), Fei Gao (高飞), Daniele Binosi, Lei Chang (常雷), C. D. Roberts and S. M. Schmidt, NJU-INP 049/21, e-print: <u>2108.11493 [hep-ph]</u>, Chin. Phys. Lett. *Express* **39** (04) (2022) 041401/1-5: <u>Express Letter</u>

- New mathematical method for interpolation and extrapolation of data
 - based on continued-fraction representation of functions, augmented by statistical sampling
- Delivers model-independent prediction for all ratios
 - No reference to models or physics theories
- Provides benchmark against which all pictures of nucleon structure can be measured in future.
- Probability that scalar diquark only models of nucleon might be consistent with available data is 1/7,000,000

MARATHON EXPERIMENT - Schlessinger point method



FIG. 3. $\lim_{x\to 1} F_2^n(x)/F_2^p(x)$. MARATHON-based SPM prediction compared with results inferred from: nuclear DIS [35]; Dyson-Schwinger equation analyses (DSE) [57, 58]; quark counting (helicity conservation) [59]; and a phenomenological fit (CJ15) [53]. The vertical red line marks the Nachtmann lower-limit, Eq. (3); and row 3 is the average in Eq. (8).

- Today, notwithstanding the enormous expense of time and effort, much must still be learnt before proton and pion structure may be considered understood in terms of DFs
- Most simply, what are the differences, if any, between the distributions of partons within the proton and the pion?
- The question of similarity/difference between proton and pion DFs has particular resonance today as science seeks to explain EHM
- How are obvious macroscopic differences between protons and pions expressed in the structural features of these two bound-states?



Figure 1: Left panel-A. In terms of QCD's Lagrangian quanta, the proton, p, contains two valence up (u) quarks and one valence down (d) quark; and also infinitely many gluons and sea quarks, drawn here as "springs" and closed loops, respectively. The neutron, as the proton's isospin partner, is defined by one u and two d valence quarks. *Right panel*-B. The pion, π^+ , contains one valence u-quark, one valence \bar{d} -quark, and, akin to the proton, infinitely many gluons and sea quarks. (In terms of valence quarks, $\pi^- \sim d\bar{u}$ and $\pi^0 \sim u\bar{u} - d\bar{d}$.)



Proton and pion DFs - QCD predictions

 \succ Valence-quark domain: there is a scale $\zeta_H < m_p$ at which \cdot

$$\begin{bmatrix} d^p(x;\zeta_{\mathcal{H}}), u^p(x;\zeta_{\mathcal{H}}) \stackrel{x\simeq 1}{\propto} (1-x)^3 \\ \bar{d}^{\pi}(x;\zeta_{\mathcal{H}}), u^{\pi}(x;\zeta_{\mathcal{H}}) \stackrel{x\simeq 1}{\propto} (1-x)^2 \end{bmatrix}$$

. 3

$$\succ \zeta > m_p$$
: val. $\propto (1-x)^{\beta_{p,\pi}}$, $\beta_p = 3 + \gamma_p$, $\beta_\pi = 2 + \gamma_\pi$

- Gluon DFs:
$$\beta_{p,\pi}^{\text{glue}} \ge \beta_{p,\pi}^{\text{val}} + 1$$

- Sea DFs: $\beta_{p,\pi}^{\text{sea}} \ge \beta_{p,\pi}^{\text{val}} + 2$

- \checkmark These are simple consequence of DGLAP equations.
- ✓ Notably, argument can be reversed:

if large-x glue or sea DF exponent is smaller than that of valence DF at any given scale, then it is smaller at all lower scales.

 \checkmark DF with lowest exponent defines the valence degree-of-freedom.

Proton and pion distribution functions in counterpoint Continuum Schwinger function methods

Proton via Faddeev equation



Pion via Bethe-Salpeter equation



Resolving the Bethe-Salpeter kernel, Si-Xue Qin and Craig D Roberts, NJU-INP 026/20, <u>arXiv:2009.13637 [hep-ph]</u>, <u>Express Letter</u> in <u>Chin.</u> Phys. Lett. **38** (7) (2021) 071201/1-6 - Editors' Suggestion





Figure 1: Left panel – A. In terms of QCD's Lagrangian quanta, the proton, p, contains two valence up (u) quarks and one valence down (d) quark; and also infinitely many gluons and sea quarks, drawn here as "springs" and closed loops, respectively. The neutron, as the proton's isospin partner, is defined by one u and two d valence quarks. Right panel – B. The pion, π^+ , contains one valence u-quark, one valence \bar{d} -quark, and, akin to the proton, infinitely many gluons and sea quarks. (In terms of valence quarks, $\pi^- \sim d\bar{u}$ and $\pi^0 \sim u\bar{u} - d\bar{d}$.)

Proton and pion distribution functions in counterpoint, Ya Lu (陆 $\overline{\mathbb{W}}$) et al., NJU-INP 056/22, e-Print: 2203.00753 [hep-ph], Phys. Lett. B 830 (2022) 137130

> Valence-quark domain: there is a scale $\zeta_H < m_p$ at which

$$\succ \zeta > m_p$$
: val. $\propto (1-x)^{\beta_{p,\pi}}$, $\beta_p = 3 + \gamma_p$, $\beta_\pi = 2 + \gamma_\pi$

- Gluon DFs: $\beta_{p,\pi}^{\text{glue}} \ge \beta_{p,\pi}^{\text{val}} + 1$ Sea DFs: $\beta_{p,\pi}^{\text{sea}} \ge \beta_{p,\pi}^{\text{val}} + 2$
- > Further, no simultaneous global fits to proton and pion data have ever been performed
- Largely because pion data are scarce
- Existing approaches are unlikely to yield definitive answers because practitioners typically ignore QCD constraints

$$\begin{bmatrix} d^p(x;\zeta_{\mathcal{H}}), u^p(x;\zeta_{\mathcal{H}}) \stackrel{x\simeq 1}{\propto} (1-x)^3 \\ \bar{d}^{\pi}(x;\zeta_{\mathcal{H}}), u^{\pi}(x;\zeta_{\mathcal{H}}) \stackrel{x\simeq 1}{\propto} (1-x)^2 \end{bmatrix}$$

- \checkmark These are simple consequence of DGLAP equations.
- \checkmark Notably, argument can be reversed:

if large-x glue or sea DF exponent is smaller than that of valence DF at any given scale, then it is smaller at all lower scales.

- \checkmark DF with lowest exponent defines the valence degree-of-freedom.
- ✓ Modern global analyses of proton DIS and related data encompass fits with role of glue and valenceguarks reversed!
- ✓ Proton has valence glue but no valence quarks!



Proton and pion distribution functions in counterpoint, Ya Lu (陆 $\overline{\mathbb{W}}$) et al., NJU-INP 056/22, e-Print: 2203.00753 [hep-ph], Phys. Lett. B 830 (2022) 137130

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- > Further, no simultaneous global fits to proton and pion data have ever been performed
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$$\begin{bmatrix} d^p(x;\zeta_{\mathcal{H}}), u^p(x;\zeta_{\mathcal{H}}) \stackrel{x\simeq 1}{\propto} (1-x)^3 \\ \bar{d}^{\pi}(x;\zeta_{\mathcal{H}}), u^{\pi}(x;\zeta_{\mathcal{H}}) \stackrel{x\simeq 1}{\propto} (1-x)^2 \end{bmatrix}$$

- \checkmark These are simple consequence of DGLAP equations.
- ✓ CT18: large-x power of glue distribution at the scale ζ = mass_{charm} is (almost) identical to that of valencequarks. With this behavior, the proton has valencegluon degrees of freedom at all scales. That would make the proton a hybrid baryon, which it is not.
- \checkmark CT18Z: large-x power of glue distribution is $a_2=1.87$, whereas that on the valence quarks is $a_2=3.15$, *i.e.*, at ζ = mass_{charm} valence-quarks are subleading degrees-of-freedom. Instead, gluons dominate on what is typically called the valence-quark domain.



Proton and pion distribution functions in counterpoint, Ya Lu (陆亚) et al., NJU-INP 056/22, e-Print: 2203.00753 [hep-ph], Phys. Lett. B 830 (2022) 137130

Symmetry-preserving analyses using continuum Schwinger function methods (CSMs) deliver hadron scale DFs that agree with QCD constraints

$$\langle x \rangle_{u_p}^{\zeta_{\mathcal{H}}} = 0.687 , \ \langle x \rangle_{d_p}^{\zeta_{\mathcal{H}}} = 0.313 , \ \langle x \rangle_{u_{\pi}}^{\zeta_{\mathcal{H}}} = 0.5$$

- > Valence-quark degrees-of-freedom carry all hadron's momentum at ζ_H
- > Diquark correlations in proton, induced by EHM

 $\Rightarrow u_V(x) \neq 2d_V(x)$

- Proton and pion valence-quark DFs have markedly different behaviour
 - $u^{\pi}(x; \zeta_H)$ is Nature's most dilated DF
 - "Obvious" because of $(1 x)^2$ vs. $(1 x)^3$ behaviour and preservation of this unit difference under evolution
 - Also "hidden" = strong EHM-induced broadening

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Proton valence-quark DFs: Continuum cf. Lattice

- Owing to difficulties in handling so-called disconnected contributions, the calculation of individual valence DFs using lattice-regularised QCD (IQCD) is problematic
- IQCD results are typically only available for isovector distributions, from which disconnected contributions vanish in the continuum limit.
- Comparison of isovector distributions

 $u^p(x;\zeta_3) - d^p(x;\zeta_3)$

Completely different approaches; yet good agreement, especially since refinements of both calculations may be anticipated.



- ✓ <u>Continuum</u>: Proton and pion distribution functions in counterpoint, Ya Lu (陆亚) et al., NJU-INP 056/22, e-Print: 2203.00753 [hep-ph]
- ✓ <u>Lattice</u>: Nucleon Isovector Unpolarized Parton Distribution in the Physical-Continuum Limit, H.-W. Lin et al., arXiv:2011.14971 [hep-lat]



Proton and pion distribution functions in counterpoint - glue and sea

Predicted CSM glue-in-pion DF is confirmed by recent IQCD calculation

[*Regarding the distribution of glue in the pion,* Lei Chang (常雷) and Craig D Roberts, e-Print: 2106.08451 [hep-ph], Chin. Phys. Lett. 38 (8) (2021) 081101/1-6]

- Glue-in-π DF possess significantly more support on the valence domain than the glue-in-p DF
- Sea-in-π DF possess significantly more support on the valence --domain than sea-in-p DFs.
- s and c sea DFs are commensurate in size with those of the lightquark sea DFs
- For s-and c-quarks, too, the pion DFs possess significantly greater support on the valence domain than the kindred proton DFs.
- These outcomes are measurable expressions of EHM



х



Asymmetry of antimatter in the proton

- > Pauli blocking: gluon splitting produces $d + \overline{d}$ in preference to $u + \overline{u}$
- Comparison with SeaQuest data

[J. Dove, et al., *The asymmetry of antimatter in the proton*, Nature 590 (7847) (2021) 561–565.]

Gottfried sum rule

$$\int_{0.004}^{0.8} dx \left[\bar{d}(x;\zeta_3) - \bar{u}(x;\zeta_3) \right] = 0.116(12)$$

Most recent result from global fits [CT18]:
 0.110(80)





Neutron/Proton structure function ratio

- Ratio 1⁺/0⁺ diquarks in proton wave function is measure of EHM
- Structure function ratio is clear window onto $d_V(x)/u_V(x)$

$$\frac{F_2^n(x;\zeta)}{F_2^p(x;\zeta)} = \frac{\mathcal{U}(x;\zeta) + 4\mathcal{D}(x;\zeta) + \Sigma(x;\zeta)}{4\mathcal{U}(x;\zeta) + \mathcal{D}(x;\zeta) + \Sigma(x;\zeta)}$$
$$U(x;\zeta) = u(x;\zeta) + \bar{u}(x;\zeta), \ D(x;\zeta) = d(x;\zeta) + \bar{d}(x;\zeta)$$
$$\Sigma(x;\zeta) = s(x;\zeta) + \bar{s}(x;\zeta) + c(x;\zeta) + \bar{c}(x;\zeta)$$

Comparison with MARATHON data

[D. Abrams, *et al.*, Measurement of Nucleon F_2^n/F_2^p Structure Function Ratio by the Jefferson Lab MARATHON Tritium/Helium-3 Deep Inelastic Scattering Experiment – arXiv:2104.05850 [hep-ex], Phys. Rev. Lett. (2022) *in press*]

Agreement with modern data on entire x-domain – parameter free prediction

Walence quark ratio in the proton, Zhu-Fang Cui, (崔著钫), Fei Gao (高飞), Daniele Binosi, Lei Chang (常雷), Craig D. Roberts and Sebastian M. Schmidt, <u>NJU-INP 049/21</u>, e-print: <u>2108.11493</u>
[hep-ph], Chin. Phys. Lett. Express **39** (04) (2022) 041401/1-5: <u>Express Letter</u>

- CSM prediction = presence of axialvector diquark correlation in the proton
- ✓ Responsible for ≈ ^{0.} 40% of proton charge





Probability that scalar diquark only models of nucleon might be consistent with available data is 1/7,000,000

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- CSMs have delivered 1st unified body of predictions for all proton and pion DFs valence, glue, and four-flavour-separated sea.
- Within mesons & baryons that share familial flavour structure, light-front momentum fractions carried by identifiable, distinct parton classes are identical at any scale.

Table 2: Low-order Mellin moments, $\langle x^m \rangle_{p_H}^{\zeta_3}$, of the DFs drawn in Figs. 2B – 3, measured in %. As an illustration of the numerical accuracy of our evolution procedure, we note that $\langle x \rangle_{\zeta_{\pi}}^{\zeta_3}$ and $\langle x \rangle_{\zeta_{p}}^{\zeta_3}$ differ by only 0.3%. Uncertainties associated with $\zeta_{\mathcal{H}} \rightarrow \zeta_{\mathcal{H}}(1 \pm 0.05)$ are shown. To simplify comparisons with phenomenological fits to relevant data, results for $\langle x^m \rangle_{p_H}^{\zeta_2}$, $\zeta_2 = 2$ GeV, are also listed. The m = 1, 2, 3 moments of the proton isovector distribution, [u - d], are: $\zeta_2 - 17.9(8)\%$, 5.1(3)%, 1.8(2)%; and $\zeta_3 - 16.6(7)\%$, 4.5(3)%, 1.6(1)%

pion	u^{π}	\bar{d}^{π}	${\mathscr G}^\pi$	\mathcal{S}^u_π	$\mathcal{S}^{ar{d}}_{\pi}$	\mathcal{S}^s_π	\mathcal{S}^c_π
$\langle x \rangle^{\zeta_2}$	24.0(1.1)	24.0(1.1)	41.0(1.2)	3.3(3)	3.3(3)	2.65(22)	1.33(5)
$\langle x^2 \rangle^{\zeta_2}$	9.5(7)	9.5(7)	3.7(1)	0.27(1)	0.27(1)	0.21(1)	0.092(2)
$\langle x^3 \rangle^{\zeta_2}$	4.7(4)	4.7(4)	0.92(6)	0.057(1)	0.057(1)	0.044(0)	0.018(1)
$\langle x \rangle^{\zeta_3}$	22.1(1.0)	22.1(1.0)	42.9(1.0)	3.7(3)	3.7(3)	3.0(2)	1.83(6)
$\langle x^2 \rangle^{\zeta_3}$	8.4(6)	8.4(6)	3.5(1)	0.27(1)	0.27(1)	0.22(1)	0.120(3)
$\langle x^3 \rangle^{\zeta_3}$	4.0(3)	4.0(3)	0.82(5)	0.056(0)	0.056(0)	0.044(0)	0.022(1)
proton	u^p	d^{p}	${\cal J}^p$	\mathcal{S}_p^u	\mathcal{S}_p^d	\mathcal{S}_p^s	\mathcal{S}_p^c
$\langle x \rangle^{\zeta_2}$	32.9(1.4)	15.0(0.7)	40.9(1.1)	2.9(2)	3.7(3)	2.64(22)	1.32(5)
$\langle x^2 \rangle^{\zeta_2}$	8.7(6)	3.6(2)	2.4(1)	0.14(1)	0.21(1)	0.13(0)	0.059(2)
$\langle x^3 \rangle^{\zeta_2}$	2.9(3)	1.1(1)	0.39(2)	0.019(0)	0.030(1)	0.019(0)	0.008(0)
$\langle x \rangle^{\zeta_3}$	30.4(1.3)	13.8(0.6)	42.8(1.0)	3.3(3)	4.1(3)	3.0(2)	1.82(6)
$\langle x^2 \rangle^{\zeta_3}$	7.7(5)	3.2(2)	2.2(1)	0.15(1)	0.21(1)	0.14(0)	0.075(2)
$\langle x^3 \rangle^{\zeta_3}$	2.5(2)	0.9(1)	0.35(2)	0.019(0)	0.028(0)	0.019(0)	0.010(1)



- CSMs have delivered 1st ever unified body of predictions for all proton and pion DFs valence, glue, and four-flavour-separated sea.
- Within mesons & baryons that share familial flavour structure, light-front momentum fractions carried by identifiable, distinct parton classes are identical at any scale.
- On the other hand, *smoking gun for EHM*, *x*-dependence of DFs is strongly hadron dependent
 - At any resolving scale, ζ , those in the pion are the hardest (most dilated).
- > All CSM DFs comply with QCD constraints on endpoint (low- and high-x) scaling behaviour.
- However, existing global fits ignore QCD constraints, so:
 - Fail to deliver realistic DFs, even from abundant proton data
 - Meson data almost nonexistent and controversial results from fits
- Only after imposing QCD constraints on future phenomenological data fits will it be possible to draw reliable pictures of hadron structure.
- Especially important for attempts to expose and understand differences between Nambu-Goldstone bosons and seemingly less complex hadrons.



Many, Many Other Expressions of EHM

➤ EHM ⇒ formation of strong nonpointlike quark+quark (diquark) correlations within baryons

- All baryons, including those with one or more heavy quarks
- > Nucleons possess 0^+ isoscalar & 1^+ isovector correlations
 - Marathon data

 \Rightarrow Probability that scalar diquark only models might be consistent is $\frac{1}{7.000.000}$

- > Nucleon resonances contain more correlations ... 0^- isoscalar, 1^- isoscalar & 1^- isovector
- $\succ \Delta$ -baryon resonances ... 1^- isoscalar & 1^- isovector
- Nucleon-elastic & nucleon-to-resonance transition form factors can test these and other structural predictions
- > On to ... GPDs, TMDs, TDAs ... gluon+quark correlations within baryons, etc.

Progress demands Synergy between Experiment + Phenomenology + Theory





EHM Measurements Meson Structure



- i. Letter of Intent: <u>A New QCD facility at the M2 beam line of the CERN SPS.</u> This document covers all ideas for future experiments as of January 2019.
- ii. Proposal for Phase-1: <u>COMPASS++/AMBER: Proposal for Measurements at the M2 beam line of the CERN SPS Phase-1: 2022-2024</u>. This document covers the three phase-1 experiments (start in 2022).
- iii. Pion and Kaon Structure at the Electron-Ion Collider, Arlene C. Aguilar et al., <u>arXiv:1907.08218</u> [nucl-ex], Eur. Phys. J. A 55 (2019) 190/1-15
- *iv.* Strong QCD from Hadron Structure Experiments, S. J. Brodsky et al., <u>arXiv:2006.06802 [hep-ph]</u>, <u>Int. J. Mod. Phys. E 29</u> (2020) 08, 2030006/1-122
- v. Selected Science Opportunities for the EicC, Xurong Chen, Feng-Kun Guo, Craig D. Roberts and Rong Wang, <u>arXiv:2008.00102 [hep-ph]</u>, Few Body Syst. 61 (2020) 4, 43/1-37. Invited contribution to the Special Issue: "New Trends in Hadron Physics: a Few-Body Perspective"
- vi. Insights into the Emergence of Mass from Studies of Pion and Kaon Structure, Craig D. Roberts, David G. Richards, Tanja Horn and Lei Chang, <u>arXiv: 2102.01765 [hep-ph]</u>, Prog. Part. Nucl. Phys. 120 (2021) 103883/1-65
- vii. Electron-Ion Collider in China, D. P. Anderle et al., arXiv:2102.09222 [nucl-ex], Front. Phys. (Beijing) 16 (2021) 6, 64701
- viii. Revealing the structure of light pseudoscalar mesons at the Electron-Ion Collider, John Arrington et al., <u>arXiv:</u> 2102.11788 [nucl-ex], J. Phys. G 48 (2021) 7, 075106

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AMBER @ CERN SPS: Excerpts from the Science Roadmap

Desired Quantities	Scientific Grounds	Experiments
Pion and kaon distribution amplitudes $(DAs) - \varphi(x)$ {Nearest thing in quantum field theory to a Schrödinger wave function}	 Theory predicts EHM is expressed in the shape of DAs 40-year controversy over x-dependence of φ Higgs-boson modulation of EHM measured by <i>skewing</i> of K DA cf. π DA Today, finally, continuum and lattice theory delivering consistent predictions for x-dependence Signature features of EHM and HB-interference manifest in low-order DA moments 	One 20-year-old FNAL experiment on π DA yielded controversial results – in some ways, internally inconsistent Low-order moments of DAs are accessible via π & K beam–induced production of 3-meson final states AMBER@CERN with high-intensity π & K beams stands tall in the world with the capacity to place novel constraints on NG-mode DAs
Pion and kaon distribution functions (DFs) {provide detailed structural information}	 How does NG-mode structure differ from that of proton? After all, mass profiles are completely different 30-year controversy: x-dependence of π valence-quark DF & Only 10 data points in existence for kaon ZERO empirical information on sea and glue in π & K Gluon content? Controversial yet prominent theory predictions that π contains no glue. What are NG modes? 	Existing experiments are more than 40- years-old. Data only on valence-quark domain. AMBER@CERN anticipates high- luminosity, high-energy π & K beams for DY & J/ ψ -production & prompt- γ production \Rightarrow especially capable of providing precision data to answer ALL these questions



AMBER @ CERN SPS: Excerpts from the Science Roadmap

Desired Quantities	Scientific Grounds	Experiments
Pion and kaon distribution amplitudes $(DAs) - \varphi(x)$ {Nearest thing in quantum field theory to a Schrödinger wave function} Sullivan process \Rightarrow Meson form factors at large Q ²	 Theory predicts EHM is expressed in the shape of DAs 40-year controversy over x-dependence of φ Higgs-boson modulation of EHM measured by <i>skewing</i> of K DA cf. π DA Today, finally, continuum and lattice theory delivering consistent predictions for x-dependence Signature features of EHM and HB-interference manifest in low-order DA moments 	One 20-year-old FNAL experiment on π DA yielded controversial results – in some ways, internally inconsistent Low-order moments of DAs are accessible via π & K beam–induced production of 3-meson final states AMBER@CERN with high-intensity π & K beams stands tall in the world with the capacity to place novel constraints on NG-mode DAs
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Revealing the structure of light pseudoscalar mesons at the Electron-Ion Collider,

Revealing the structure of light pseudoscalar mesons of the Lieu on 101, control, John Arrington et al., arXiv: 2102.11788 [nucl-ex], J. Phys. G 48 (2021) 7, 075106 Selected Science Questions for EIC

Science Question	Key Measurement[1]	Key Requirements[2]
What are the quark and gluon energy contributions to the pion mass?	Pion structure function data over a range of x and Q^2 .	 Need to uniquely determine e + p → e' + X + n (low -t) CM energy range ~10-100 GeV Charged- and neutral currents desirable
Is the pion full or empty of gluons as viewed at large Q^2 ?	Pion structure function data at large Q^2 .	 CM energy ~100 GeV Inclusive and open-charm detection
What are the quark and gluon energy contributions to the kaon mass?	Kaon structure function data over a range of x and Q^2 .	 Need to uniquely determine Λ, Σ⁰: e + p → e' + X + Λ/Σ⁰ (low -t) CM energy range ~10-100 GeV
Are there more or less gluons in kaons than in pions as viewed at large Q^2 ?	Kaon structure function data at large Q^2 .	 CM energy ~100 GeV Inclusive and open-charm detection
Can we get quantitative guidance on the emergent pion mass mechanism?	Pion form factor data for $Q^2 = 10{\text{-}}40 \ (\text{GeV/c})^2$.	 Need to uniquely determine exclusive process e + p → e' + π⁺ + n (low -t) e-p and e-d at similar energies CM energy ~10-75 GeV
What is the size and range of interference between emergent-mass and the Higgs-mass mechanism?	Kaon form factor data for $Q^2 = 10{\text{-}}20 \;(\text{GeV/c})^2$.	 Need to uniquely determine exclusive process e + p → e' + K⁺ + Λ (low -t) L/T separation at CM energy ~10-20 GeV e-p Λ/Σ⁰ ratios at CM energy ~10-50 GeV
What is the difference between the impacts of emergent- and Higgs-mass mechanisms on light-quark behaviour?	Behaviour of (valence) up quarks in pion and kaon at large x	 CM energy ~20 GeV (lowest CM energy to access large-x region) Higher CM energy for range in Q² desirable
What is the relationship between dynamically chiral symmetry breaking and confinement?	Transverse-momentum dependent Fragmentation Functions of Meson quarks into pions and kaons fragment	 Collider kinematics desirable (as compared to fixed-target kinematics) CM energy range ~20-140 GeV
nia Paharts: cdrabarts@niu adu cn. 111 "On Mass and Mattar"	functions	


Revealing the structure of light pseudoscalar mesons at the Electron-Ion Collider, John Arrington *et al.*, <u>arXiv: 2102.11788 [nucl-ex]</u>, <u>J. Phys. G 48 (2021) 7, 075106</u>

Selected Science Questions for EIC

More speculative observables		
What is the trace anomaly contribution to the pion mass?	Elastic J/ψ production at low W off the pion.	 Need to uniquely determine exclusive process e + p → e' + π⁺ + J/Ψ + n (low -t) High luminosity (10³⁴⁺) CM energy ~70 GeV
Can we obtain tomographic snapshots of the pion in the transverse plane? What is the pressure distribution in a pion?	Measurement of DVCS off pion target as defined with Sullivan process	 Need to uniquely determine exclusive process e + p → e' + π⁺ + γ + n (low -t) High luminosity (10³⁴⁺) CM energy ~10-100 GeV
Are transverse momentum distributions universal in pions and protons?	Hadron multiplicities in SIDIS off a pion target as defined with Sullivan process	 Need to uniquely determine scattered off pion: e + p → e + h + X + n (low -t) High luminosity (10³⁴⁺) e-p and e-d at similar energies desirable CM energy ~10-100 GeV

Revealing the structure of light pseudoscalar mesons at the Electron-Ion Collider, John Arrington *et al.*, <u>arXiv: 2102.11788 [nucl-ex]</u>, <u>J. Phys. G 48 (2021) 7, 075106</u>

Selected Science Questions for EIC



Revealing pion and kaon structure via generalised parton distributions, <u>Khépani</u> <u>Raya</u> et al. <u>e-Print: 2109.11686 [hep-ph]</u>, Chin. Phys. C **46** (01) (2022) 013107/1-22

- The analyses predict that mass-squared gravitational form factors are stiffer than electromagnetic form factors
- ✓ Reveal that K pressure profiles are tighter than profiles
- ✓ Both π and K mesons sustain near-core pressures at magnitudes similar to that expected at the core of neutron stars

Emergent Hadron Mass



- > QCD is unique amongst known fundamental theories of natural phenomena
 - The degrees-of-freedom used to express the scale-free Lagrangian are not directly observable
 - Massless gauge bosons become massive, with no "human" interference
 - Gluon mass ensures a stable, infrared completion of the theory through the appearance of a running coupling that saturates at infrared momenta, being everywhere finite
 - Massless fermions become massive, producing
 - Massive baryons and simultaneously Massless mesons
- > These emergent features of QCD are expressed in every strong interaction observable
- They can also be revealed via
 - EHM interference with Nature's other known source of mass = Higgs
- We are capable of building facilities that can validate these concepts, proving QCD to be the 1st well-defined four-dimensional quantum field theory ever contemplated
- > This may open doors that lead far beyond the Standard Model



Jhis is the Fascination and Vzgency of Physics

- Is there a Lagrangian for Nature?
- If so, then the natural mass/length scales for all things are contained within
- Much more than that ...
- > If there is \mathcal{L}_{Nature} , then Logic, as expressed in Mathematics, is innate to Nature
- > Then, Mathematics is not something we invented to describe Nature
 - It is a secret of Nature that we discovered
- > So that Nature is Mathematics and Mathematics describes everything
 - Including us.
- > This may be true
 - If so, then there <u>are</u> answers to the questions we have long asked
 - And answers to questions we have not yet thought to ask



The Nature of Things

Mature

Jhis is the Fascination and Vzgency of Physics

- Is there a Lagrangian for Nature?
- > If so, then the natural mass/length scales for all things are contained within
- > Much more than that ...
- If there is There are theories of many things,
- Then, Mat – It is a s But is there a theory of everything?
- > So that Nature is Mathematics and Mathematics describes everything
 - Including us.
- This may be true
 - If so, then there <u>are</u> answers to the questions we have long asked
 - And answers to questions we have not yet thought to ask



The Nature of Things

LNature



There are theories of many things But is there a theory of everything



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Quantum Chromodynamics

- Definition: "Well-defined" quantum field theory
 - finite number of renormalization constants, $\{Z_i, i = 1, ..., N\}$
 - $\{Z_i, i = 1, ..., N\}$ can all be computed and remain bounded real numbers as any regularization scale is removed
 - Quantum electrodynamics fails this test owing to Landau pole in ultraviolet
 - Weak interactions are essentially perturbative because addition of Higgs scalar introduces enormous infrared scale that suppresses all nonperturbative effects But Higgs boson mass itself is quadratically divergent ⇒ not renormalizable
- > Nonperturbative QCD has come a long way in the last few years.
- Beginning to look increasingly likely that QCD is humanity's first well-defined fourdimensional quantum field theory
- Stands alone as an internally consistent theory
 - a predictive tool as written, with nothing needing to be added





Hadron Radii

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Nambu Goldstone Bosons: How Big are They?

- Physics Measurements:
 - Pion and kaon charge radii
- Scientific Context:
 - Hadron masses and radii are a direct expression of the links between EHM and confinement
 - Example: within Nambu-Goldstone boson domain of pseudoscalar-meson mass trajectory

 $f_{0-} r_{0-} \approx constant$ product of meson leptonic decay constant (f_{0-}) and charge radius (r_{0-}) is practically constant

- f_{0-} is an order parameter EHM. Equally, so is r_{π}
- Precise measurements of π & K radii will reveal compositeness (confinement) scale for Nambu-Goldstone bosons





Size of Nambu Goldstone Bosons

- > Data Existing:
 - 70+ years after their discovery, knowledge of pion and kaon radii is very limited.
 - Pion charge radius is poorly known.
 - m_{π} measured to one-part in a million;
 - r_{π} only constrained to 1% & even this is debatable
 - The most precise data (charged-pion scattering at CERN) are more than 30-years-old and do not reach to very low momentum transfers
 - The most recent data (Selex) yield a result with errors more than 10-times larger than CERN _
 - New theory methods enable model-independent extraction of radii from precise low-Q² data
 - New method \Rightarrow r_{π}: 2.4 σ tension with PDG 2021
 - No information beyond least-squares fit in K data
 - E.g. AMBER, with high-intensity π & K beams, capable of delivering precise data that are required
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This data is all we have ... It cannot reasonably be used to claim a result for the kaon charge radius



- Kaon radius is practically unknown
 - 15 imprecise scattering points, on Q² > 0.015 GeV², analysed using a simple monopole fit, which is today known to be a poor approach
 - Recall that > 60-years of *ep* scattering measurements, analysed this way, led to proton radius extractions that are likely wrong by an amount in excess of 7σ.

Enable 1st objective extraction of kaon size

Pion charge radius from

pion+electron

elastic

scattering

data, Zhu-

Fang Cui (崔著

钫), Daniele

Binosi, C.D.

INP 047/21.

Roberts, S.M.

Schmidt, NIU-

arXiv:2108.04 948 [hep-ph],

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