SMEFT projections of neutral-current PVDIS asymmetries at the EIC

Kağan Şimşek

Northwestern University

in collaboration with

Radja Boughezal, Alexander Emmert, Tyler Kutz, Sonny Mantry,
Michael Nycz, Frank Petriello, Daniel Wiegand, and Xiaochao Zheng

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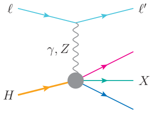
- We study NC DIS cross-section asymmetries at EIC.
- BSM effects are parametrized in SMEFT framework.
- Higher-dimensional operators are built of existing SM particles with Wilson coefficients as effective couplings at UV scale Λ:

$$\mathscr{L}_{\text{SMEFT}} = \mathscr{L}_{\text{SM}} + \sum_{n>4} \frac{1}{\Lambda^{n-4}} \sum_k C_k^{(n)} O_k^{(n)}$$

- All new physics is assumed to be heavier than all SM states and accessible collider energy.
- We focus on semi-leptonic 4-fermion $O_k^{(n)}$ at $n = 6$.
- We find that the EIC can
  - probe complementarily and competitively to LHC DY
  - resolve blind spots observed in LHC NC DY data fits

We study the NC DIS in the process $\ell + H \to \ell' + X$, where $\ell = e^-, e^+$ and $H = p, D$:



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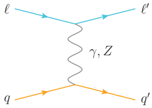
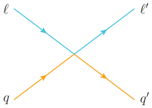
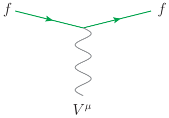


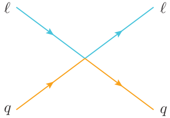


to leading order at parton level

We parameterize the vertex factors in terms of vector and axial couplings:



$$i\gamma_\mu g_1^{(fV)} + i\gamma_\mu \gamma_5 g_5^{(fV)},$$



$$\begin{aligned} & i[\gamma_\mu][\gamma^\mu] g_{11}^{(\ell q)} \\ & + i[\gamma_\mu][\gamma^\mu \gamma_5] g_{15}^{(\ell q)} \\ & + i[\gamma_\mu \gamma_5][\gamma^\mu] g_{51}^{(\ell q)} \\ & + i[\gamma_\mu \gamma_5][\gamma^\mu \gamma_5] g_{55}^{(\ell q)} \end{aligned}$$

We don't consider Yukawa or dipole interactions because they are suppressed by fermion masses, which we assume to vanish.

SMEFT operators shift the usual vector and axial SM couplings in a gauge-invariant way: e.g.

$$g_1^{(fZ)} = g_V^f + \mathcal{O}(C_k), \quad g_5^{(fZ)} = g_A^f + \mathcal{O}(C_k)$$

Operators that contribute to the ffV and $\ell\ell qq$ vertices at dimension 6 are (Grzadkowski *et al.* [1008.4884]):

ffV	$\ell\ell qq$
$O_{\varphi\ell}^{(1)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{\ell}\gamma^\mu \ell)$	$O_{\ell q}^{(1)} = (\bar{\ell}\gamma_\mu \ell)(\bar{q}\gamma^\mu q)$
$O_{\varphi\ell}^{(3)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \tau^I \varphi)(\bar{\ell}\gamma^\mu \tau^I \ell)$	$O_{\ell q}^{(3)} = (\bar{\ell}\gamma_\mu \tau^I \ell)(\bar{q}\gamma^\mu \tau^I q)$
$O_{\varphi e} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$	$O_{eu} = (\bar{e}\gamma_\mu e)(\bar{u}\gamma^\mu u)$
$O_{\varphi q}^{(1)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$	$O_{ed} = (\bar{e}\gamma_\mu e)(\bar{d}\gamma^\mu d)$
$O_{\varphi q}^{(3)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \tau^I \varphi)(\bar{q}\gamma^\mu \tau^I q)$	$O_{\ell u} = (\bar{\ell}\gamma_\mu \ell)(\bar{u}\gamma^\mu u)$
$O_{\varphi u} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$	$O_{\ell d} = (\bar{\ell}\gamma_\mu \ell)(\bar{d}\gamma^\mu d)$
$O_{\varphi d} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$	$O_{qe} = (\bar{q}\gamma_\mu q)(\bar{e}\gamma^\mu e)$

There is one more:

$O_{\varphi WB} = (\varphi^\dagger \tau^I \varphi) W^I_{\mu\nu} B^{\mu\nu}$ $\Rightarrow$ causes kinetic mixing of W^3 and B

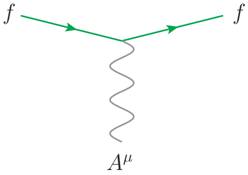
$\Rightarrow$ universally shifts the ffV vertices after diagonalization that gives physical photon and Z boson states

The *ffV* operators are already strongly bounded by Z and W pole observables (Dawson & Giardino [1909.02000]):

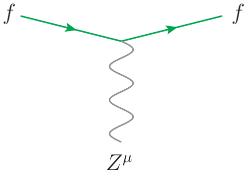
ffV	C_k	95% CL, $\Lambda = 1$ TeV
$O_{\varphi\ell}^{(1)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{\ell}\gamma^\mu \ell)$	$C_{\varphi\ell}^{(1)}$	$[-0.043, 0.012]$
$O_{\varphi\ell}^{(3)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \tau^I \varphi)(\bar{\ell}\gamma^\mu \tau^I \ell)$	$C_{\varphi\ell}^{(3)}$	$[-0.012, 0.0029]$
$O_{\varphi e} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$	$C_{\varphi e}$	$[-0.013, 0.0094]$
$O_{\varphi q}^{(1)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$	$C_{\varphi q}^{(1)}$	$[-0.027, 0.043]$
$O_{\varphi q}^{(3)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \tau^I \varphi)(\bar{q}\gamma^\mu \tau^I q)$	$C_{\varphi q}^{(3)}$	$[-0.011, 0.014]$
$O_{\varphi u} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$	$C_{\varphi u}$	$[-0.072, 0.091]$
$O_{\varphi d} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$	$C_{\varphi d}$	$[-0.16, 0.060]$
$O_{\varphi WB} = (\varphi^\dagger \tau^I \varphi) W_{\mu\nu}^I B^{\mu\nu}$	$C_{\varphi WB}$	$[-0.0088, 0.0013]$

Thus, we restrict our attention only to the operators contributing to the $\ell\ell qq$ vertex, which leaves us with seven Wilson coefficients of interest: C_{eu}, C_{ed}, $C_{\ell q}^{(1)}$, $C_{\ell q}^{(3)}$, $C_{\ell u}$, $C_{\ell d}$, and C_{qe}.

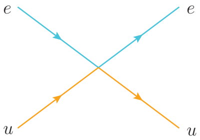
Since we consider contributions only to the $\ell\ell qq$ interaction, we assume the usual SM ffV vertices in our analysis:



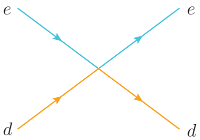
$$g_1^{(fA)} = -eQ_f$$
$$g_5^{(fA)} = 0$$



$$g_1^{(fZ)} = g_V^f$$
$$g_5^{(fZ)} = g_A^f$$



$$
\begin{aligned}
g_{11}^{(eu)} &= \tfrac{1}{4}\left[C_{eu} + (C_{\ell q}^{(1)} - C_{\ell q}^{(3)}) + C_{\ell u} + C_{qe}\right] \\
g_{15}^{(eu)} &= \tfrac{1}{4}\left[C_{eu} - (C_{\ell q}^{(1)} - C_{\ell q}^{(3)}) + C_{\ell u} - C_{qe}\right] \\
g_{51}^{(eu)} &= \tfrac{1}{4}\left[C_{eu} - (C_{\ell q}^{(1)} - C_{\ell q}^{(3)}) - C_{\ell u} + C_{qe}\right] \\
g_{55}^{(eu)} &= \tfrac{1}{4}\left[C_{eu} + (C_{\ell q}^{(1)} - C_{\ell q}^{(3)}) - C_{\ell u} - C_{qe}\right]
\end{aligned}
$$



the same as for *eeuu* but with $u \to d$ and $C_{\ell q}^{(1)} - C_{\ell q}^{(3)} \to C_{\ell q}^{(1)} + C_{\ell q}^{(3)}$

Amplitude and cross section for $\ell + q \to \ell' + q'$:

$$\mathscr{M} = \mathscr{M}_\gamma + \mathscr{M}_Z + \mathscr{M}_\times \Rightarrow \mathrm{d}\sigma^{\lambda_\ell \lambda_q} = \frac{\mathrm{d}^2\sigma}{\mathrm{d}x\,\mathrm{d}Q^2} = \frac{1}{16\pi x^2 s^2}\,|\mathscr{M}|^2 + \mathscr{O}(C_k^2)$$

Asymmetry definitions:

- unpolarized PV asymmetries: $A_{\mathrm{PV}} = \dfrac{\mathrm{d}\sigma_\ell}{\mathrm{d}\sigma_0}$
- polarized PV asymmetries: $\Delta A_{\mathrm{PV}} = \dfrac{\mathrm{d}\sigma_H}{\mathrm{d}\sigma_0}$
- lepton-charge asymmetries: $A_{\mathrm{LC}} = \dfrac{\mathrm{d}\sigma_0(e^+H) - \mathrm{d}\sigma_0(e^-H)}{\mathrm{d}\sigma_0(e^+H) + \mathrm{d}\sigma_0(e^-H)}$

where

$$\mathrm{d}\sigma_0 = \frac{1}{4}\sum_q f_{q/H}[\mathrm{d}\sigma^{++} + \mathrm{d}\sigma^{+-} + \mathrm{d}\sigma^{-+} + \mathrm{d}\sigma^{--}] : \text{unpol. } \ell + \text{unpol. } H$$

$$\mathrm{d}\sigma_\ell = \frac{1}{4}\sum_q f_{q/H}[\mathrm{d}\sigma^{++} + \mathrm{d}\sigma^{+-} - \mathrm{d}\sigma^{-+} - \mathrm{d}\sigma^{--}] : \text{pol. } \ell + \text{unpol. } H$$

$$\mathrm{d}\sigma_H = \frac{1}{4}\sum_q \Delta f_{q/H}[\mathrm{d}\sigma^{++} - \mathrm{d}\sigma^{+-} + \mathrm{d}\sigma^{-+} - \mathrm{d}\sigma^{--}] : \text{unpol. } \ell + \text{pol. } H$$

Data sets, shown with beam energies and nominal annual luminosities:

D1	5 GeV $\times$ 41 GeV eD, 4.4 fb^{-1}
D2	5 GeV $\times$ 100 GeV eD, 36.8 fb^{-1}
D3	10 GeV $\times$ 100 GeV eD, 44.8 fb^{-1}
D4	10 GeV $\times$ 137 GeV eD, 100 fb^{-1}
D5	18 GeV $\times$ 137 GeV eD, 15.4 fb^{-1}
P1	5 GeV $\times$ 41 GeV ep, 4.4 fb^{-1}
P2	5 GeV $\times$ 100 GeV ep, 36.8 fb^{-1}
P3	10 GeV $\times$ 100 GeV ep, 44.8 fb^{-1}
P4	10 GeV $\times$ 275 GeV ep, 100 fb^{-1}
P5	18 GeV $\times$ 275 GeV ep, 15.4 fb^{-1}
P6	18 GeV $\times$ 275 GeV ep, 100 fb^{-1}

P6: Yellow Report reference setting [2103.05419]

Data set labels:
D, P: unpolarized PV asymmetry
ΔD, ΔP: polarized PV asymmetry
LD, LP: lepton-charge asymmetry

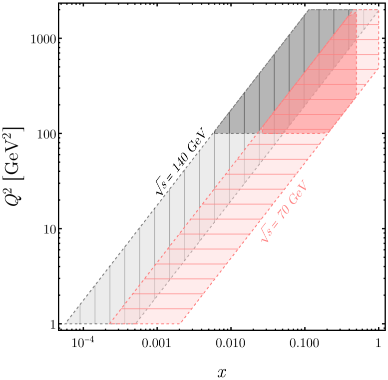
Cuts on projected data:

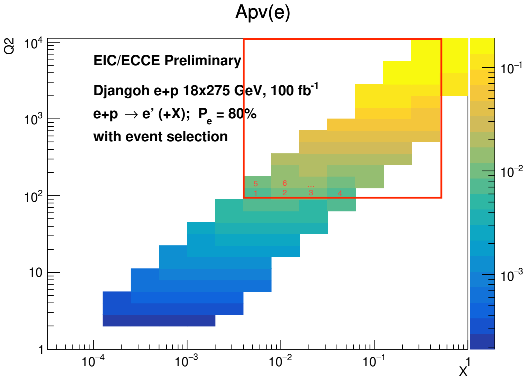
$Q > 1$ GeV	to avoid nonperturbative QCD
$y > 0.1$	to avoid bin migration and unfolding uncertainty
$y < 0.9$	to avoid high photoproduction background due to final-state hadron
$\|\eta\| < 3.5$	to restrict events in main acceptance of ECCE detector
$E' > 2$ GeV	to have high-purity e^- samples

Additional cuts in SMEFT analysis:

$x < 0.5$	to avoid *large* uncertainties from nonperturbative QCD and nuclear dynamics
$Q > 10$ GeV	

Kinematic region of the data sets ($\sqrt{s} = 70\text{-}140$ GeV, $0.1 \le y \le 0.9$):



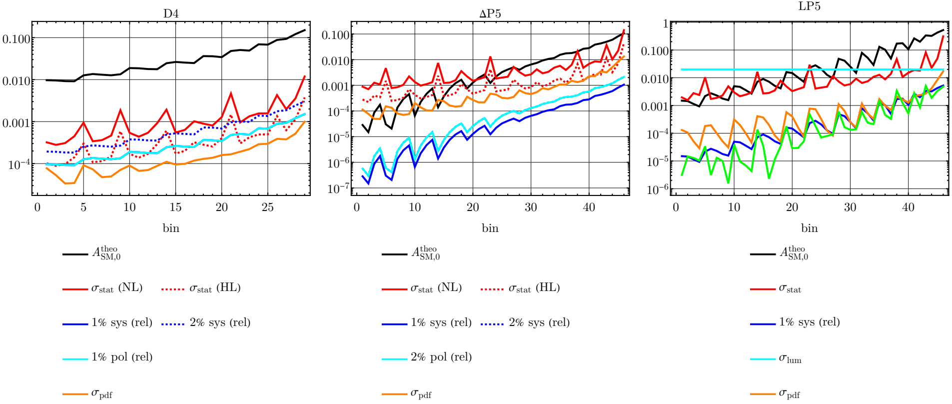


The shaded region on the left and the red box on the right indicate the kinematic region and *good* bins used in our SMEFT analysis, respectively.

Anticipated uncertainty components:

Error type	A_{PV} (D, P)	ΔA_{PV} (ΔD, ΔP)	A_{LC} (LD, LP)
statistical (NL)	$\sigma_{\mathrm{stat}} = \frac{1}{P_\ell \sqrt{N}}$	$\frac{P_\ell}{P_H}\sigma_{\mathrm{stat}}$	$\sqrt{10} P_\ell \sigma_{\mathrm{stat}}$
statistical (HL)	$\frac{1}{\sqrt{10}}\sigma_{\mathrm{stat}}$	$\frac{1}{\sqrt{10}}\frac{P_\ell}{P_H}\sigma_{\mathrm{stat}}$	✗
uncorrelated systematic	1% rel.	1% rel.	1% rel.
fully correlated beam polarization	1% rel.	2% rel.	✗
fully correlated luminosity	✗	✗	2% abs.
uncorrelated QED NLO	✗	✗	$5\% \times (A_{\mathrm{LC}}^{\mathrm{NLO}} - A_{\mathrm{LC}}^{\mathrm{Born}})$
fully correlated PDF	✓	✓	✓

PDF sets used: `NNPDF3.1 NLO` and `NNPDFpol1.1`



- Bins on the horizontal axes are the *good* x and Q^2 bins.
- Stat error dominates in PV asymmetries in NL case.
- Systematic and beam-polarization errors become comparable to stat error in HL case.
- Luminosity error dominates in LC asymmetries.
- Stat error competes with luminosity error at high-x high-Q^2 bins.
- PDF errors are the least dominant in unpolarized PV asymmetries but become significant in the polarized case.

Pseudodata generation:

$$A^{(e)}_{\text{pseudo},b} = A_{\text{SM},b} + r^{(e)}_b \sigma^{\text{unc}}_b + r'^{(e)} \sigma^{\text{cor}}_b$$

$$b \in \text{Range}(N_{\text{bin}}), \quad e \in \text{Range}(N_{\text{exp}}), \quad N_{\text{exp}} = 10^3, \quad r^{(e)}_b, r'^{(e)} \sim \mathcal{N}(0,1)$$

$$\sigma^{\text{unc}}_b = \sigma_{\text{stat},b} \oplus \sigma_{\text{sys},b} \qquad \sigma^{\text{unc}}_b = \sigma_{\text{stat},b} \oplus \sigma_{\text{sys},b} \oplus \sigma_{\text{nlo},b}$$

$$\sigma^{\text{cor}}_b = \sigma_{\text{pol},b} \qquad \sigma^{\text{cor}}_b = \sigma_{\text{lum},b}$$

SMEFT asymmetry expressions:

$$A_{\text{SMEFT},b} = A_{\text{SM},b} + \sum_{k=1}^{N_{\text{fit}}} C_k \, \delta A_{k,b} + \mathcal{O}(C_k^2), \quad N_{\text{fit}} \in \text{Range}(7)$$

χ^2 function for each pseudoexperiment:

$$\chi^{2\,(e)} = \sum_{b,b'=1}^{N_{\text{bin}}} [A_{\text{SMEFT},b} - A^{(e)}_{\text{pseudo},b}] H_{bb'} [A_{\text{SMEFT},b'} - A^{(e)}_{\text{pseudo},b'}]$$

Polarimetry and luminosity difference can be limiting factors.

- $\Rightarrow$ use data itself to constrain these systematic effects
- $\Rightarrow$ simultaneous fits of C_k with beam polarization, P, and luminosity difference, A_{lum}, in an attempt to obtain stronger bounds for C_k

Fits of C_k with P:

$$\chi^{2\,(e)} = \sum_{b,b'=1}^{N_{\text{bin}}} \left[P A_{\text{SMEFT},b} - A^{(e)}_{\text{pseudo},b}\right] \left[H_{bb'} \Big|_{\sigma_{\text{pol}}\to 0} \right] \left[P A_{\text{SMEFT},b'} - A^{(e)}_{\text{pseudo},b'}\right] + \frac{(P-\bar{P})^2}{\delta P^2}$$

unpolarized PV asymmetries:

- $|\rho(C_k, P)| \gtrsim 0.7$
- 30-50% stronger bounds

polarized PV asymmetries:

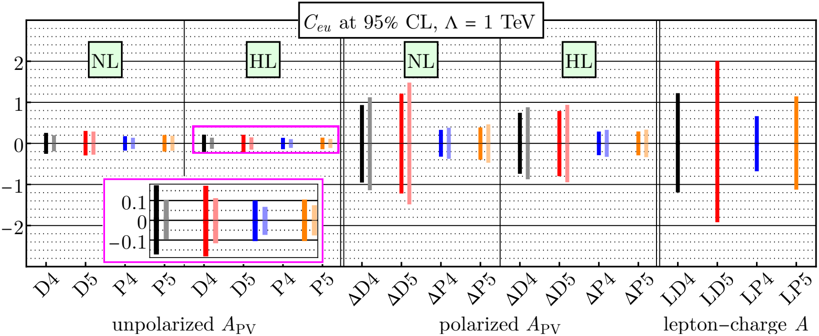
- $|\rho(C_k, P)| \lesssim 0.2$
- 15-20% weaker bounds

Improvement is more significant than worsening $\Rightarrow$ include P in fits.

Fits of C_k with A_{lum}:

$$\chi^{2\,(e)} = \sum_{b,b'=1}^{N_{\text{bin}}} \left[A_{\text{SMEFT},b} - A^{(e)}_{\text{pseudo},b} - A_{\text{lum}}\right] \left[H_{bb'} \Big|_{\sigma_{\text{lum}}\to 0} \right] \left[A_{\text{SMEFT},b'} - A^{(e)}_{\text{pseudo},b'} - A_{\text{lum}}\right]$$

Mild correlations, $|\rho(C_k, A_{\text{lum}})| \lesssim 0.4$, leading to 15-20% weaker bounds
$\Rightarrow$ do not include A_{lum} in fits.

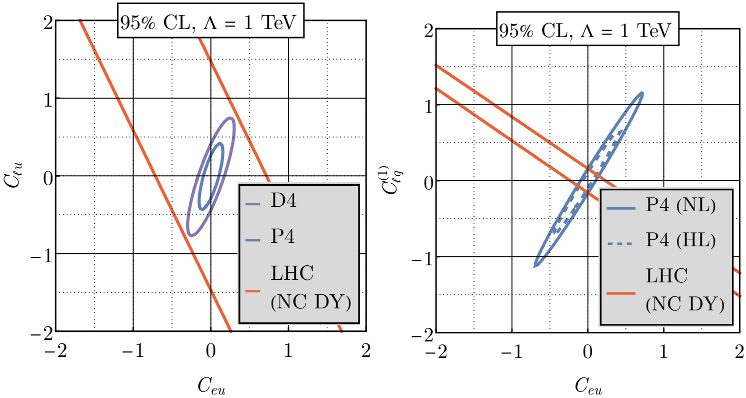


In terms of the strength of bounds:

- proton > deuteron
- high-lum. low-energy (4^{th} sets) > low-lum. high-energy (5^{th} sets)
- unpolarized PV > polarized PV > lepton-charge
- improvement: unpolarized PV > polarized PV if NL → HL

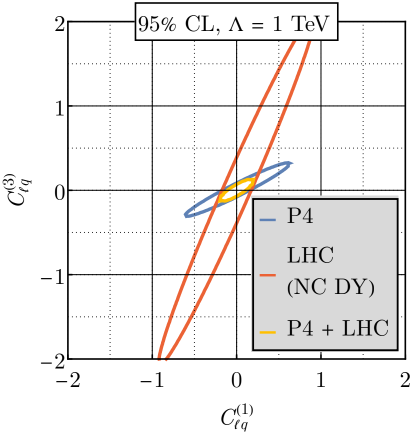
Corresponding effective UV scales: 3 TeV with NL, 4 TeV with HL

Compare the bounds from deuteron and proton data of unpolarized PV asymmetries to the 8-TeV 20-fb^{-1} LHC NC DY data (Boughezal, Petriello, & Wiegand [2004.00748, 2104.03979]):



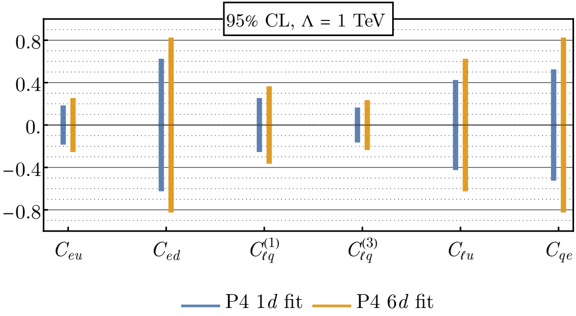
Distinct correlations: EIC fits are complementary to LHC NC DY. However, LHC fits have blind spots and exhibit flat directions, which remain even in the high-luminosity case. The EIC can resolve and constrain this parameter space strongly.

Compare proton data of unpolarized PV asymmetries to the 8-TeV 20-fb^{-1} LHC NC DY data (Boughezal, Petriello, & Wiegand [2004.00748]) when the LHC fit doesn't have a flat direction:



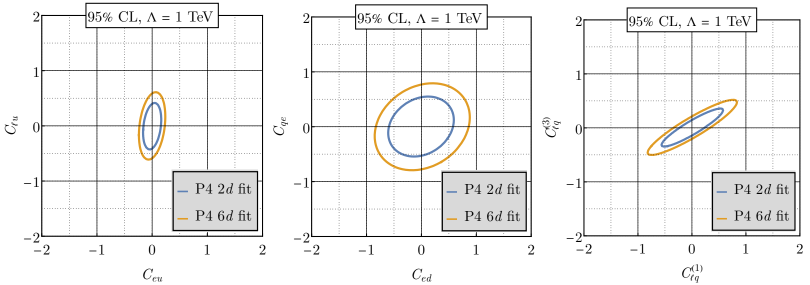
Distinct correlations again: EIC fits are complementary to LHC NC DY. Moreover, when the LHC fit gives a strong bound without showing a flat direction, the EIC can constrain the same parameter space even more strongly.

N_{fit}	N_{exp}
2	10^3
3	10^4
4	10^5
5	10^6
6	10^7
7	10^8(?!)



- Number of pseudoexperiments increases to reflect the required statistics.
- Beam polarization parameter, P, is not included here.
- Bounds become 25 to 40% weaker due to increased number of fitted parameters and correlations among them.
- Not significant worsening because correlations dominate statistical effect of increasing number of fitted parameters.

Compare the two-parameter fits of Wilson coefficients to the projections from a six-parameter fit:



- The *eeuu* vertex contains the combination $C_{\ell q}^{(1)} - C_{\ell q}^{(3)}$ and the *eedd* vertex has $C_{\ell q}^{(1)} + C_{\ell q}^{(3)}$.
- These may lead to degeneracies and flat directions in a multi-parameter fit of Wilson coefficients.
- The EIC can resolve this part of the parameter space, imposing strong bounds.

- We investigate the BSM potential of EIC in the model-independent SMEFT framework by focusing on semi-leptonic four-fermion operators at dimension 6 by giving a detailed accounting of uncertainties.
- We obtain bounds on Wilson coefficients from single-, double-, and even multiple-parameter fits by using techniques to simultaneously fit P and A_{lum} together with SMEFT parameters.

- We find that UV scales up to 3 TeV (or 4 TeV) can be probed with nominal (or $10\times$ high) annual luminosity.
- We observe that the strongest bounds come from unpolarized PV asymmetries of proton.
- EIC is shown to be complementary and competitive to LHC NC DY by
    - equally or more strongly confining the Wilson coefficients with distinct correlations;
    - resolving the degeneracies observed in the LHC data.

EIC was designed as a QCD machine and it shows strong potential for BSM physics.