Study of TMDs in DIS single-jet production



- Lorenzo Rossi
- In collaboration with: Alessandro Bacchetta, Yiannis Makris
 - University of Pavia & INFN
 - EIC User Group Early Career Workshop 2022



Istituto Nazionale di Fisica Nucleare



Jet-TMDFF

Phenomenological Estimates



Jet-TMDFF

Phenomenological Estimates











 $\bigstar \quad p_T \ll Q^2 = -q^2$





 $\bigstar \quad p_T \ll Q^2 = -q^2$







 $\bigstar \quad p_T \ll Q^2 = -q^2$





Usual TMDPDFs and TMDFFs



 $\bigstar \quad p_T \ll Q^2 = -q^2$









$$F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q;\mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x) \right]$$

 $x, b_T, Q; \mu, \zeta_2)$



$$F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q;\mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x) \right]$$

 $x, b_T, Q; \mu, \zeta_2) J_{i \rightarrow \text{jet}}(Q, R; \mu_j)$



$$F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q;\mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x) \right]$$

 $x, b_T, Q; \mu, \zeta_2) J_{i \rightarrow \text{jet}}(Q, R; \mu_j) \mathcal{S}_J(b_T, R; \mu_{Sj}, \zeta_R)$



$$F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q; \mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x) \right]$$

where $\mathscr{B}_0[fD] = \frac{x}{2\pi} \int_0^\infty f_0^{\infty}$

 $db_T b_T J_0(b_T p_T) f(x, b_T; Q^2) D(z, b_T^2; Q^2).$



е

 $F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q;\mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x,b_T,Q;\mu,\zeta_2) J_{i \to \text{jet}}(Q,R;\mu_j) \, \mathscr{S}_J(b_T,R;\mu_{Sj},\zeta_R) \right]$



е



 $F_{UU,T}^{\text{jet}} = \sum_{i:} H_{ij}(Q;\mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x,b_T,Q;\mu,\zeta_2) J_{i \to \text{jet}}(Q,R;\mu_j) \, \mathscr{S}_J(b_T,R;\mu_{Sj},\zeta_R) \right]$

е



е



$$F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q; \mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x, b_T) \right]$$

 $T, Q; \mu, \zeta_2) J_{i \to jet}(Q, R; \mu_j) \mathcal{S}_J(b_T, R; \mu_{Sj}, \zeta_R)$

е



$$F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q; \mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x, b_T) \right]$$

 $T, Q; \mu, \zeta_2) J_{i \to \text{jet}}(Q, R; \mu_j) \mathcal{S}_J(b_T, R; \mu_{Sj}, \zeta_R)$

е



$$F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q; \mu_Q) \mathscr{B}_0 \left[f_{1j}^{\text{uns.}}(x, b_T) \right]$$

 $T, Q; \mu, \zeta_2) J_{i \to \text{jet}}(Q, R; \mu_j) \mathcal{S}_J(b_T, R; \mu_{Sj}, \zeta_R)$

е



Jet-TMDFF

Phenomenological Estimates

It's useful to introduce a new function called Jet-TMDFF:

 $D_{i \to \text{jet}}(b_T, Q, R; \mu, \zeta_1) = \frac{J_{i \to \text{jet}}(Q, R; \mu_J) \mathcal{S}_J(b_T, R; \mu_R, \zeta_R)}{\sqrt{S(b; \mu_S, \zeta_S)}}$

It's useful to introduce a new function called Jet-TMDFF:

 $D_{i \to \text{jet}}(b_T, Q, R; \mu, \zeta_1) = \frac{J_{i \to \text{jet}}(Q, R; \mu_J) \mathcal{S}_J(b_T, R; \mu_R, \zeta_R)}{\sqrt{S(b; \mu_S, \zeta_S)}}$

$$F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q;\mu_Q) \mathscr{B}_0 \Big[f_{1j}^{\text{uns.}}(b_T, x) \Big]$$

 $x, Q; \mu, \zeta_2) \sqrt{S(b; \mu_S, \zeta_S)} D_{i \rightarrow \text{jet}}(b_T, Q, R; \mu, \zeta_1)$

It's useful to introduce a new function called Jet-TMDFF:

$$D_{i \to \text{jet}}(b_T, Q, R; \mu, \zeta_1) =$$

$$\Rightarrow \qquad F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q; \mu_Q) \mathscr{B}_0 \Big[f_{1j}^{\text{uns.}}(b_T, x, f_{1j}) \Big]$$
$$\Rightarrow \qquad F_{UU,T}^{\text{jet}} = \sum_{ij} H_{ij}(Q; \mu_Q) \mathscr{B}_0 \Big[f_{1j}^{\text{sub}} \Big]$$

 $J_{i \to \text{jet}}(Q, R; \mu_J) \mathcal{S}_J(b_T, R; \mu_R, \zeta_R)$ $\sqrt{S(b; \mu_S, \zeta_S)}$

 $Q; \mu, \zeta_2) \sqrt{S(b; \mu_S, \zeta_S)} D_{i \rightarrow \text{jet}}(b_T, Q, R; \mu, \zeta_1)$

 ${}^{\mathrm{lb}}(b_T, x, Q; \mu, \zeta_2) D_{i \to \mathrm{jet}}(b_T, Q, R; \mu, \zeta_1)$

Definition:





Diagrams contributing at next-to-leading order in the perturbative calculation of the jet Soft Function. There are other eight mirror diagrams, which are not shown in this figure.

е



е

е



е

e

Ρ

NLO expression



e



NLO expression

$$\mathcal{S}_{J}(b) = 1 + \frac{\alpha_{s}C_{F}}{\pi} \left[\frac{1}{2\epsilon^{2}} - \frac{1}{\epsilon\alpha} - \frac{1}{\epsilon} \ln \frac{1}{\mu} - \ln \left(\frac{\mu^{2}}{\mu b^{2}}\right) \ln \left(\frac{\nu}{\mu \tan(\frac{R}{2})}\right) \right] + c$$

 $\mathcal{O}(\alpha^2) + \mathcal{O}(\epsilon^2) \, .$

e



NLO expression

$$\mathcal{S}_{J}(b) = 1 + \frac{\alpha_{s}C_{F}}{\pi} \left[\frac{1}{2\varepsilon^{2}} - \frac{1}{\varepsilon\alpha} - \frac{1}{\varepsilon} \ln \frac{1}{\mu} - \ln \frac{1}{\omega} \ln \frac{1}{\omega} - \ln \left(\frac{\mu^{2}}{\mu b^{2}}\right) \ln \left(\frac{\nu}{\mu \tan(\frac{R}{2})}\right) \right] + \delta dt$$

 $\mathcal{O}(\alpha^2) + \mathcal{O}(\epsilon^2) \, .$
Let's remind the definition of the Jet-TMDFF:

 $D_{i \to \text{jet}}(b_T, Q, R; \mu, \zeta_1) = \frac{J_{i \to \text{jet}}(Q, R; \mu_J) \mathcal{S}_J(b_T, R; \mu_R, \zeta_R)}{\sqrt{S(b; \mu_S, \zeta_S)}}$

Let's remind the definition of the Jet-TMDFF:

Next-to-Leading Order Expression:

$$\begin{split} D_{i \to jet} &= 1 + \frac{\alpha_s(\mu_J)}{2\pi} \Big[\frac{3C_F}{2} \ln \Big(\frac{\mu_J^2}{Q^2 \tan^2 \frac{R}{2}} \Big) + \frac{C_F}{2} \ln^2 \Big(\frac{\mu_J^2}{Q^2 \tan^2 \frac{R}{2}} \Big) + d_J^{q, \text{alg}} \Big] + \frac{\alpha_s(\mu_R)C_F}{\pi} \Big) + \\ &- \frac{\pi^2}{6} - \ln^2 \Big(\frac{\mu_R^2}{\mu_b^2} \Big) - 4 \ln \Big(\frac{\mu_R^2}{\mu_b^2} \Big) \ln \Big(\frac{\nu_R}{\mu_R \tan \frac{R}{2}} \Big) \Big] - \frac{\alpha_s(\mu_S)C_F}{4\pi} \Big[- \ln^2 \Big(\frac{\mu_S^2}{\mu_b^2} \Big) + \\ &+ 4 \ln \Big(\frac{\mu_S^2}{\mu_b^2} \Big) \ln \Big(\frac{\mu_0}{\nu_S} \Big) - \frac{\pi^2}{6} \Big] \,. \end{split}$$

 $D_{i \to \text{jet}}(b_T, Q, R; \mu, \zeta_1) = \frac{J_{i \to \text{jet}}(Q, R; \mu_J) \mathcal{S}_J(b_T, R; \mu_R, \zeta_R)}{\sqrt{1-2}}$

$$\sqrt{S(b;\mu_S,\zeta_S)}$$

10

Jet-TNDFF





Choice of initial scales:



$\bigstar \quad \mu_R = \mu_S = \mu_0 \qquad \bigstar \quad \nu_S = \sqrt{\zeta_S} = \sqrt{\zeta_0} = \mu_0 \qquad \bigstar \quad \nu_R = \sqrt{\zeta_R} = \sqrt{\zeta_0} \tan(R/2) = \mu_0 \tan(R/2)$





$$\mu_R = \mu_S = \mu_0 \qquad \bigstar \qquad \nu_S = \sqrt{\zeta_S} = \sqrt{\zeta_0} = \mu_0 \qquad \bigstar \qquad \nu_R = \sqrt{\zeta_R} = \sqrt{\zeta_0} \tan(R/2) = \mu_0 \tan(R/2)$$

$$D_{i \to \text{jet}} = 1 + \frac{\alpha_s(\mu_J)}{2\pi} \Big[\frac{3C_F}{2} \ln\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + \frac{C_F}{2} \ln^2\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + d_J^{q,alg} \Big] + -4 \frac{\alpha_s(\mu_0)C_F}{\pi} \ln\Big(\frac{\mu_0^2}{\mu_b^2}\Big) \ln\Big(\frac{\nu_R}{\nu_S \tan(R/2)}\Big)$$



$d\sigma_{\text{jet}} \sim H_{ij}(Q;\mu_H) \otimes D_{i \to \text{jet}}(b,Q,R;\mu_1,\zeta_1) \times f_{j \leftarrow P}^{sub.}(x,b,Q;\mu_2,\zeta_2)$

 $d\sigma_{\text{jet}} \sim H_{ij}(Q;\mu_H) \otimes D_{i \to \text{jet}}(b,Q,R;\mu_1,\zeta_1) \times f_{j \leftarrow P}^{sub.}(x,b,Q;\mu_2,\zeta_2)$ $\gamma_{\mu}^{D_{i \to jet}} + \gamma_{\mu}^{H} + \gamma_{\mu}^{f} = 0 \qquad \qquad \gamma_{\nu}^{D_{i \to jet}} + \gamma_{\nu}^{f} = 0$

 $d\sigma_{\text{jet}} \sim H_{ij}(Q;\mu_H) \otimes D_{i \to \text{jet}}(b,Q,R;\mu_1,\zeta_1) \times f_{j \leftarrow P}^{sub.}(x,b,Q;\mu_2,\zeta_2)$ $\gamma_{\mu}^{D_{i \to jet}} + \gamma_{\mu}^{H} + \gamma_{\mu}^{f} = 0$

$$\gamma_{\nu}^{D_{i \to \text{jet}}} + \gamma_{\nu}^{f} = 0$$

 $d\sigma_{\text{sidis}} \sim H_{ii}(Q;\mu_H) \otimes D_{i \to h}^{sub.}(b,Q,z;\mu_1,\zeta_1) \times f_{i \leftarrow P}^{sub.}(x,b,Q;\mu_2,\zeta_2)$

 $\checkmark \qquad \gamma_{\mu}^{D_{i \to jet}} + \gamma_{\mu}^{H} + \gamma_{\mu}^{f} = 0$

 $d\sigma_{\text{sidis}} \sim H_{ij}(Q;\mu_H) \otimes D_{i \to h}^{sub.}(h)$ $\gamma_{\mu}^{D_{i \to h}} + \gamma_{\mu}^{H} + \gamma_{\mu}^{f} = 0$

 $d\sigma_{\text{jet}} \sim H_{ij}(Q;\mu_H) \otimes D_{i \to \text{jet}}(b,Q,R;\mu_1,\zeta_1) \times f_{j \leftarrow P}^{sub.}(x,b,Q;\mu_2,\zeta_2)$

$$\gamma_{\nu}^{D_{i \to \text{jet}}} + \gamma_{\nu}^{f} = 0$$

$$(b, Q, z; \mu_1, \zeta_1) \times f^{sub.}_{j \leftarrow P}(x, b, Q; \mu_2, \zeta_2)$$

$$\gamma_{\nu}^{D_{i \to h}} + \gamma_{\nu}^{f} = 0$$

 $\gamma_{\mu}^{D_{i \to jet}} + \gamma_{\mu}^{H} + \gamma_{\mu}^{f} = 0$

 $\gamma_{\mu}^{D_{i \to h}} + \gamma_{\mu}^{H} + \gamma_{\mu}^{f} = 0$

 $d\sigma_{\text{jet}} \sim H_{ij}(Q;\mu_H) \otimes D_{i \to \text{jet}}(b,Q,R;\mu_1,\zeta_1) \times f_{j \leftarrow P}^{sub.}(x,b,Q;\mu_2,\zeta_2)$

$$\gamma_{\nu}^{D_{i \to \text{jet}}} + \gamma_{\nu}^{f} = 0$$

 $d\sigma_{\text{sidis}} \sim H_{ij}(Q;\mu_H) \otimes D^{sub.}_{i \to h}(b,Q,z;\mu_1,\zeta_1) \times f^{sub.}_{j \leftarrow P}(x,b,Q;\mu_2,\zeta_2)$

$$\gamma_{\nu}^{D_{i\to h}} + \gamma_{\nu}^{f} = 0$$

 $\gamma_{\mu}^{D_{i \to jet}} + \gamma_{\mu}^{H} + \gamma_{\mu}^{f} = 0$

 $\gamma^{D_{i \to h}}_{\mu} + \gamma^{H}_{\mu} + \gamma^{f}_{\mu} = 0$

 $d\sigma_{\text{jet}} \sim H_{ij}(Q;\mu_H) \otimes D_{i \to \text{jet}}(b,Q,R;\mu_1,\zeta_1) \times f_{j \leftarrow P}^{sub}(x,b,Q;\mu_2,\zeta_2)$

$$\gamma_{\nu}^{D_{i \to \text{jet}}} + \gamma_{\nu}^{f} = 0$$

 $d\sigma_{\text{sidis}} \sim H_{ij}(Q;\mu_H) \otimes D^{sub.}_{i \to h}(b,Q,z;\mu_1,\zeta_1) \times f^{sub.}_{j \leftarrow P}(x,b,Q;\mu_2,\zeta_2)$

$$\gamma_{\nu}^{D_{i\to h}} + \gamma_{\nu}^{f} = 0$$

 $\gamma_{\mu}^{D_{i \to \text{jet}}} + \gamma_{\mu}^{H} + \gamma_{\mu}^{f} = 0$

 $\gamma^{D_{i \to h}}_{\mu} + \gamma^{H}_{\mu} + \gamma^{f}_{\mu} = 0$

 $d\sigma_{\text{jet}} \sim H_{ij}(Q;\mu_H) \otimes D_{i \to \text{jet}}(b,Q,R;\mu_1,\zeta_1) \times f_{j \leftarrow P}^{sub}(x,b,Q;\mu_2,\zeta_2)$

$$\gamma_{\nu}^{D_{i \to \text{jet}}} + \gamma_{\nu}^{f} = 0$$

 $d\sigma_{\text{sidis}} \sim H_{ij}(Q;\mu_H) \otimes D^{sub.}_{i \to h}(b,Q,z;\mu_1,\zeta_1) \times f^{sub.}_{j \leftarrow P}(x,b,Q;\mu_2,\zeta_2)$

$$\gamma_{\nu}^{D_{i\to h}} + \gamma_{\nu}^{f} = 0$$

The Jet-TMDFF has the same evolution of TMD-FF.

$$D_{i \to \text{jet}} = 1 + \frac{\alpha_s(\mu_J)}{2\pi} \Big[\frac{3C_F}{2} \ln\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + \frac{C_F}{2} \ln^2\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + d_J^{q,alg} \Big] + \frac{4\alpha_s(\mu_0)C_F}{\pi} \ln\Big(\frac{\mu_0^2}{\mu_b^2}\Big) \ln\Big(\frac{\nu_R}{\nu_S \tan(R/2)}\Big)$$

$$D_{i \to \text{jet}} = 1 + \frac{\alpha_s(\mu_J)}{2\pi} \Big[\frac{3C_F}{2} \ln\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + \frac{C_F}{2} \ln^2\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + d_J^{q,alg} \Big] + \frac{4\alpha_s(\mu_0)C_F}{\pi} \ln\Big(\frac{\mu_0^2}{\mu_b^2}\Big) \ln\Big(\frac{\nu_R}{\nu_S \tan(R/2)}\Big)$$

$$D_{i \to jet} = 1 + \frac{\alpha_s(\mu_J)}{2\pi} \Big[\frac{3C_F}{2} \ln\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + \frac{C_F}{2} \ln^2\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + d_J^{q,alg} \Big] + \frac{4\alpha_s(\mu_0)C_F}{\pi} \ln\Big(\frac{\mu_0^2}{\mu_b^2}\Big) \ln\Big(\frac{\nu_R}{\nu_S \tan(R/2)}\Big)$$

$$D_{i \to \text{jet}} = 1 + \frac{\alpha_s(\mu_J)}{2\pi} \Big[\frac{3C_F}{2} \ln\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + \frac{C_F}{2} \ln^2\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + d_J^{q,alg} \Big] + \frac{4\alpha_s(\mu_0)C_F}{\pi} \ln\Big(\frac{\mu_0^2}{\mu_b^2}\Big) \ln\Big(\frac{\nu_R}{\nu_S \tan(R/2)}\Big)$$

We have two different scales due to the Jet function and the Soft Functions.

$$D_{i \to \text{jet}} = 1 + \frac{\alpha_s(\mu_J)}{2\pi} \Big[\frac{3C_F}{2} \ln\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + \frac{C_F}{2} \ln^2\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + d_J^{q,alg} \Big] + \frac{4\alpha_s(\mu_0)C_F}{\pi} \ln\Big(\frac{\mu_0^2}{\mu_b^2}\Big) \ln\Big(\frac{\nu_R}{\nu_S \tan(R/2)}\Big)$$

We have two different scales due to the Jet function and the Soft Functions.

We have to do a sub-evolution of Jet Function:

$$D_{i \to jet} = 1 + \frac{\alpha_s(\mu_J)}{2\pi} \Big[\frac{3C_F}{2} \ln\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + \frac{C_F}{2} \ln^2\Big(\frac{\mu_J^2}{Q^2 \tan^2(R/2)}\Big) + d_J^{q,alg} \Big] + \frac{4\alpha_s(\mu_0)C_F}{\pi} \ln\Big(\frac{\mu_0^2}{\mu_b^2}\Big) \ln\Big(\frac{\nu_R}{\nu_S \tan(R/2)}\Big)$$

We have two different scales due to the Jet function and the Soft Functions.

We have to do a sub-evolution of Jet Function:

$$U_J[\mu_0 \leftarrow \mu_J] = \exp\left\{\int_{\mu_J}^{\mu_0} \frac{d\mu}{\mu} \gamma_J(\mu)\right\}.$$



Jet-Sidis Process

Jet-TMDFF

Phenomenological Estimates

(1) $G(R, b_T) = e^{-\frac{g_3(R)b_T^2}{4}}$

(1) $G(R, b_T) = e^{-\frac{g_3(R)b_T^2}{4}}$

(2) $S_{\tilde{K}}(b_T) = e^{g_{\tilde{K}} \ln Q^2} = e^{-\frac{g_2 b_T^2}{4} \ln \zeta_R}$

Collins-Soper kernel

(1) $G(R, b_T) = e^{-\frac{g_3(R)b_T^2}{4}}$

(2) $S_{\tilde{K}}(b_T) = e^{g_{\tilde{K}} \ln Q^2} = e^{-\frac{g_2 b_T^2}{4} \ln \zeta_R}$



To parametrize the non-perturbative effects in the TMD-PDFs, we use 200 sets of values of the free parameters obtained in the so-called MAP22 extraction. arXiv:2206.07598

(1) $G(R, b_T) = e^{-\frac{g_3(R)b_T^2}{4}}$

(2) $S_{\tilde{K}}(b_T) = e^{g_{\tilde{K}} \ln Q^2} = e^{-\frac{g_2 b_T^2}{4} \ln \zeta_R}$

Fixed by the MAP22' proton replicas.



We made some phenomenological estimates in two different situations at NNLL accuracy:



We made some phenomenological estimates in two different situations at NNLL accuracy:



By fixing $g_3 = 0.35$ GeV² and by changing the set of parametrization of the







We made some phenomenological estimates in two different situations at NNLL accuracy:







We made some phenomenological estimates in two different situations at NNLL accuracy:

We expect NP effect relatively high!




Phenomenological Estimates



Phenomenological Estimates





Obtained Results:

- New form of the Structure Functions
- Computation of the NLO expressions of \mathcal{S}_J
- Computation of the unpolarized cross section at NNLL accuracy
- Phenomenological Estimates



Obtained Results:

- New form of the Structure Functions
- Computation of the NLO expressions of \mathcal{S}_{I}
- Computation of the unpolarized cross section at NNLL accuracy
- Phenomenological Estimates

- **Possible outlooks:**
- Extension to N^3LL accuracy
- Predictions for the Sivers polarized Structure Function $F_{IITT}^{\sin(\phi_h \phi_s)}$ UT,T
 - Reanalyzing of data from HERA

