

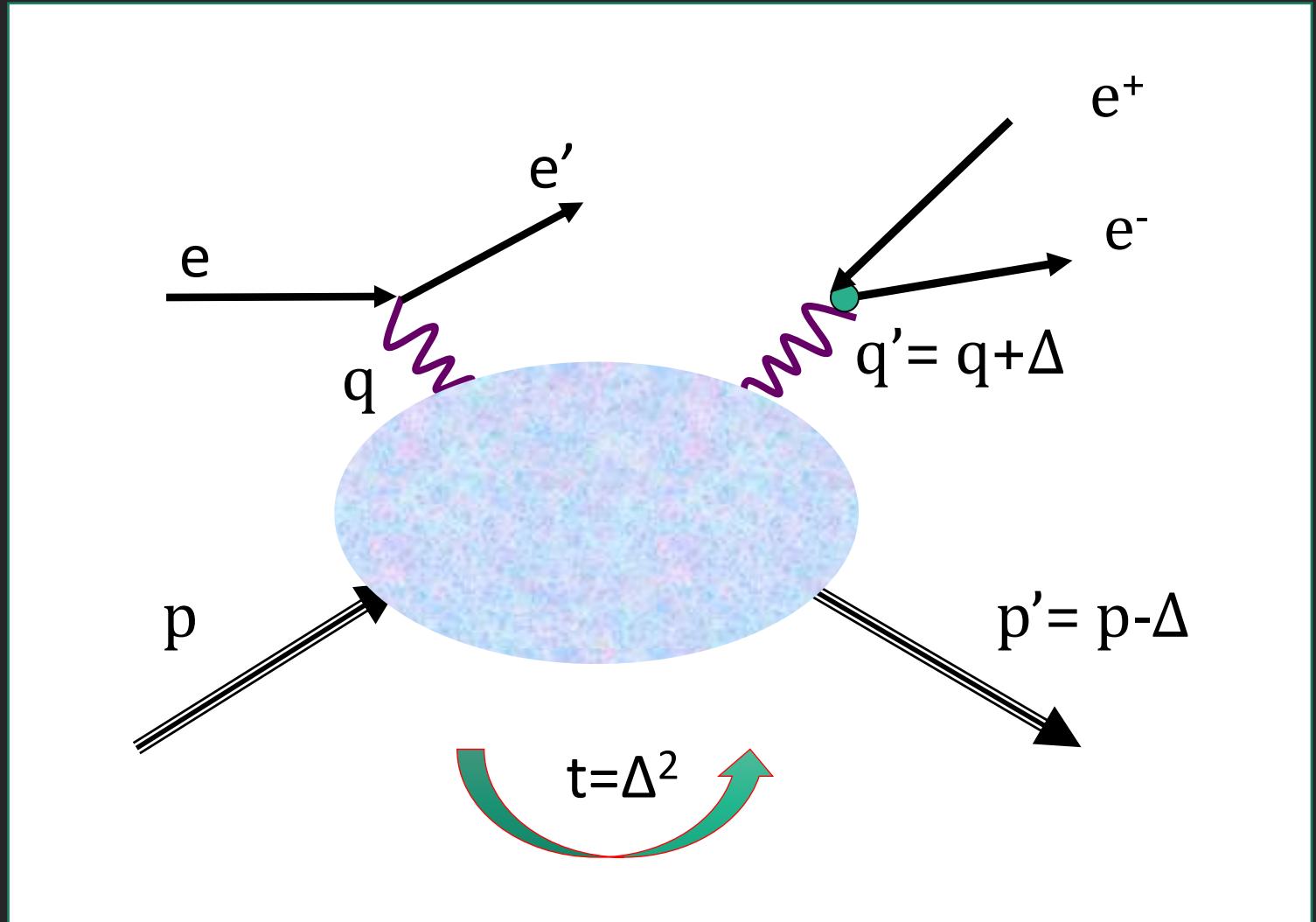
Theory of DVCS in Nuclei

CORE Meeting
June 30, 2020

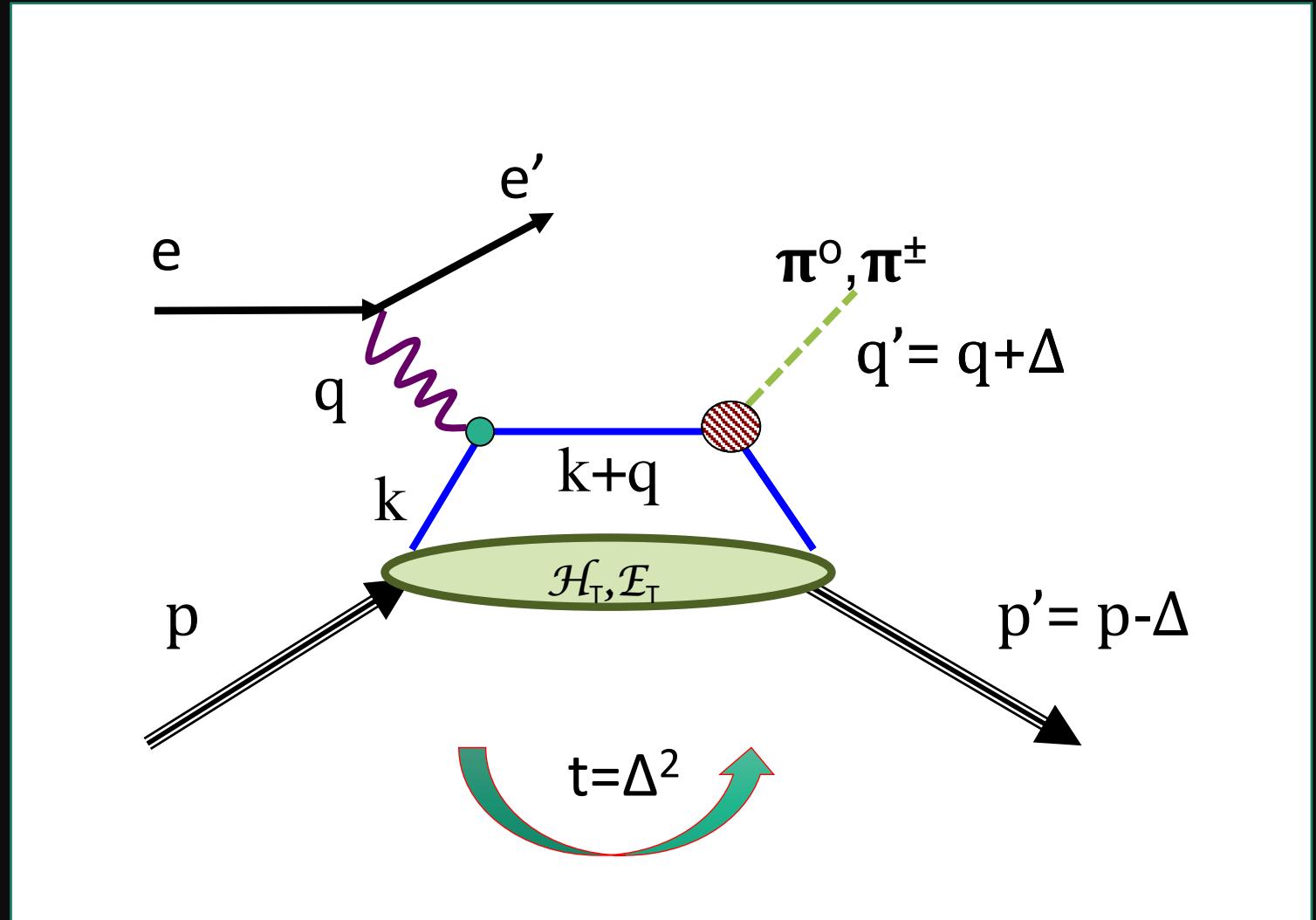
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Deeply Virtual Exclusive Processes from proton

- Deeply Virtual Compton Scattering
- Timelike Compton Scattering
- DDVCS



- Deeply Virtual Meson Production
- Exclusive Drell Yan



Nuclei with Spin 0: Why is DVCS on ${}^4\text{He}$ interesting?

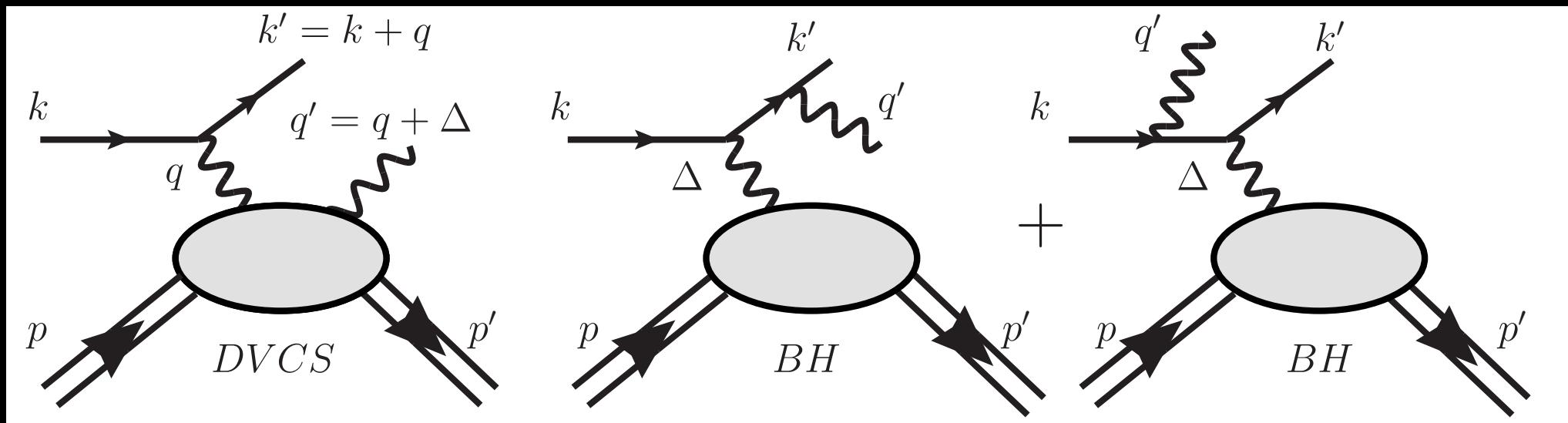
1. A simpler cross section structure for GPD studies

1. Kinematic twists (work in progress)
2. Dynamic twist

$H_3^{||}$ corresponds to a configuration with a longitudinally polarized quark in an unpolarized target associated with the quark's spin-orbit contribution

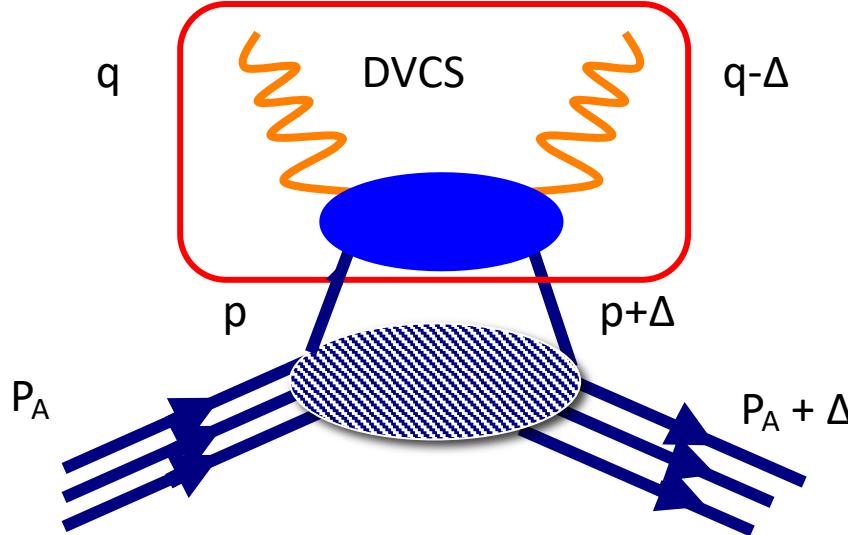
- 1 twist-two GPD, H
 - 2 twist-three GPDs, H_3 and $H_3^{||}$
-
1. Key to understanding the origin of EMC effect/shadowing/anti-shadowing

Structure of cross section

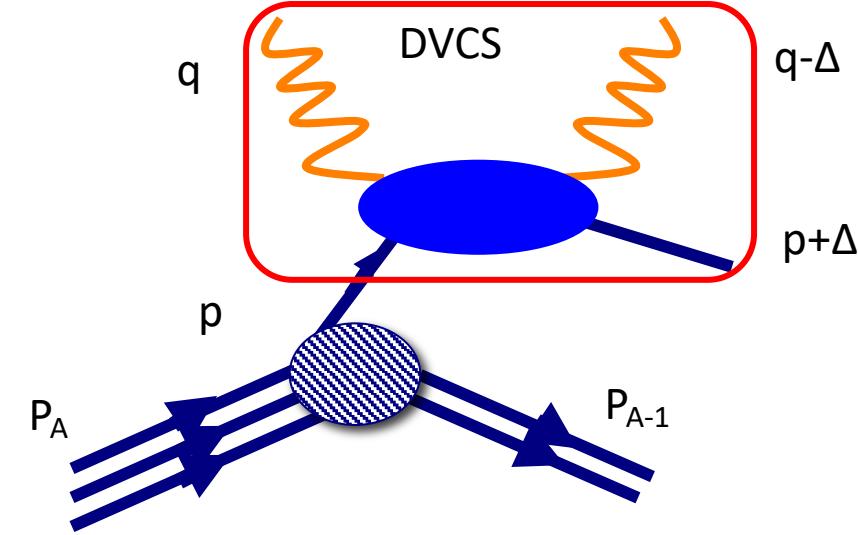


Deeply virtual processes from nuclei: Coherent vs. Incoherent
(SL and S. Taneja, [hep-ph/0505123](#), [hep-ph/0504027](#) PRC, 2006)

Coherent



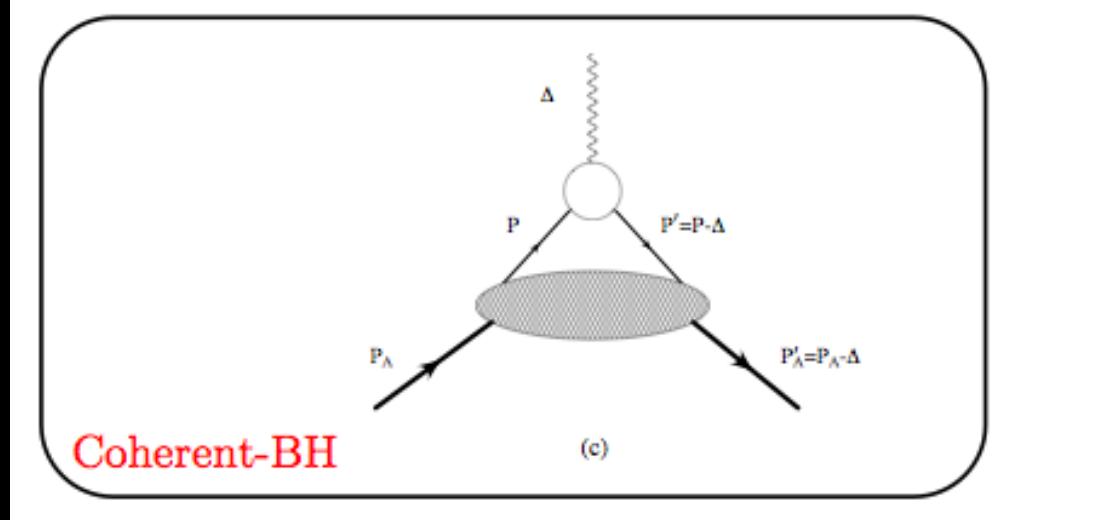
Incoherent



Off-forward nucleon spectral function

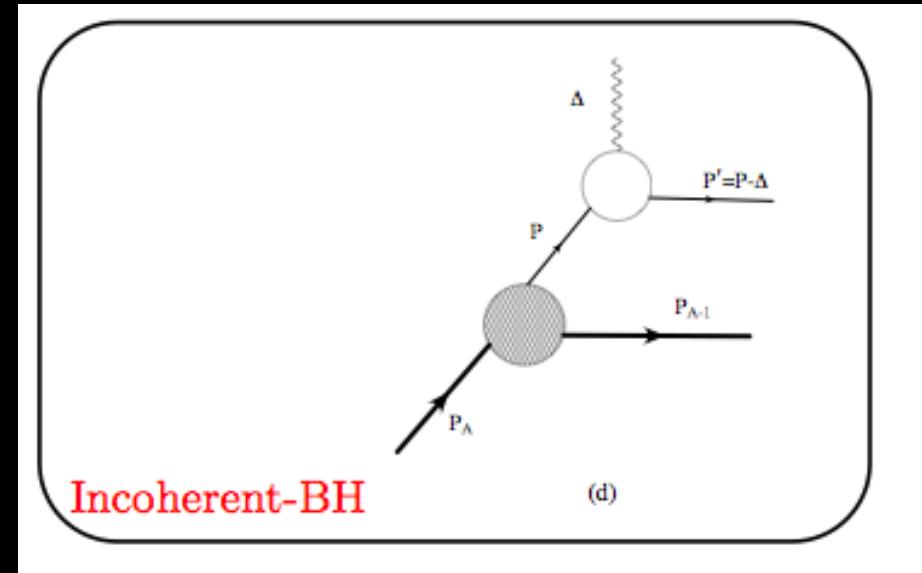
“Standard forward” nucleon spectral function

Coherent/Incoherent contributions also in Bethe Heitler processes from nuclei



Coherent-BH

(c)



Incoherent-BH

(d)

Non-forward spectral function

⇒ Interference Term for Coherent DVCS & BH

$$\mathcal{I}_{coh}(\zeta, t) = \mathcal{K} H^A(\zeta, t) \times Z^2 F^A(t)$$

$$H^A(\zeta, t) = \int \frac{d^2 P_\perp dY}{2(2\pi)^3} \mathcal{N} \rho^A(Y, P^2; \zeta, t) H^N \left(\frac{\zeta}{Y}, \frac{\zeta}{Y}, t; P^2; \right)$$

↑ off-forward EMC-effect ↑

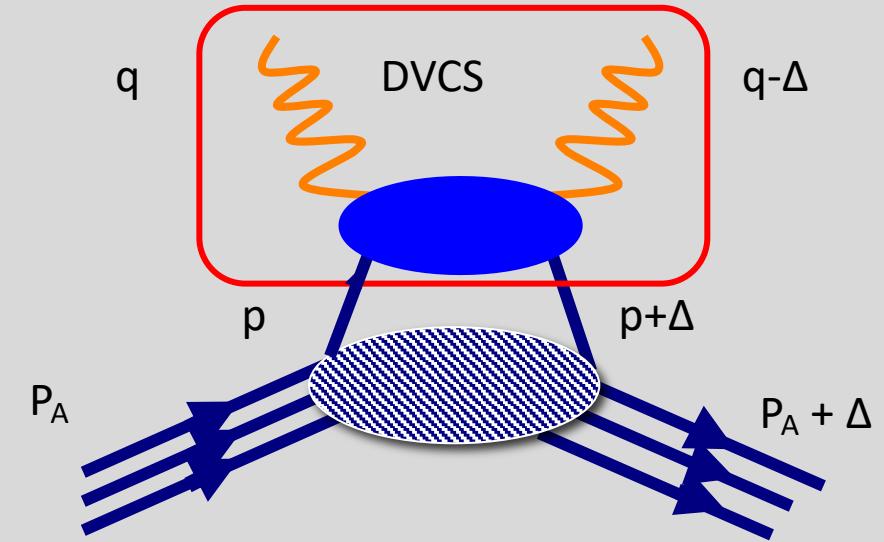
Forward spectral function

⇒ Interference Term for Incoherent DVCS & BH

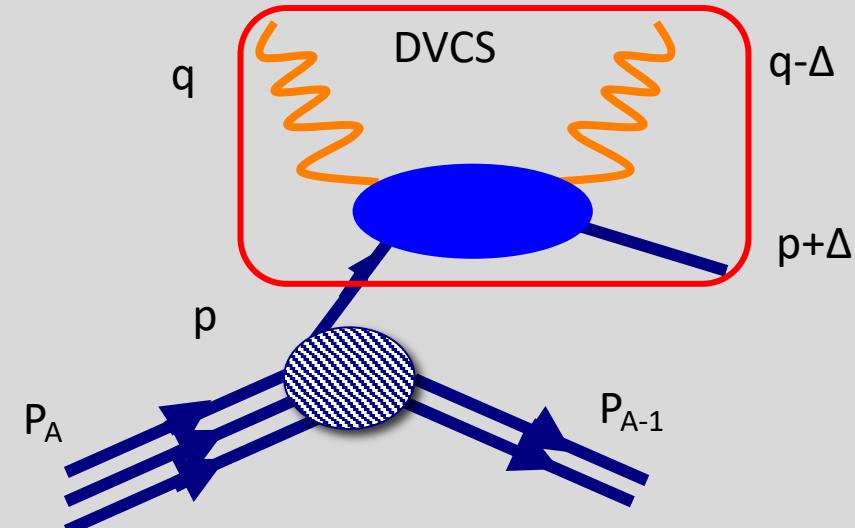
$$\mathcal{I}_{inc}(\zeta, t) = \mathcal{K} H_0^A(\zeta, t) \times Z F_1^N(t)$$

$$H_0^A(\zeta, t) = \int \frac{d^2 P_\perp dY}{2(2\pi)^3} \mathcal{N} \rho_0^A(Y, P^2) H^N \left(\frac{\zeta}{Y}, \frac{\zeta}{Y}, t; P^2 \right)$$

Coherent



Incoherent



DVCS in spin $\frac{1}{2}$ compared to spin 0

$$\frac{d^5\sigma_{DVCS}}{dx_B j dQ^2 d|t| d\phi d\phi_S} = \text{twist two GPDs}$$

$$= \text{twist three GPDs}$$

+

+

+

+

+

+

+

$$\frac{\alpha^3}{16\pi^2(s - M^2)^2\sqrt{1 + \gamma^2}} |T_{DVCS}|^2$$

$$\frac{\Gamma}{Q^2(1 - \epsilon)} \left\{ \begin{array}{l} F_{UU,T} + \epsilon F_{UU,L} + \epsilon \cos 2\phi F_{UU}^{\cos 2\phi} \\ \sqrt{\epsilon(\epsilon + 1)} \left[\cos \phi F_{UU}^{\cos \phi} + \sin \phi F_{UU}^{\sin \phi} \right] \\ \lambda_e i) \sqrt{2\epsilon(1 - \epsilon)} \sin \phi F_{LU}^{\sin \phi} \end{array} \right.$$

H_3

$H_3^{||}$

$$\begin{aligned} & S_L \left[F_{UL} + \sqrt{\epsilon(\epsilon + 1)} \sin \phi F_{UL}^{\sin \phi} + \epsilon \sin 2\phi F_{UL}^{\sin 2\phi} \right] \\ & \lambda_e \left[\sqrt{1 - \epsilon^2} F_{LL} + 2 \lambda_e \sqrt{\epsilon(1 - \epsilon)} \cos \phi F_{LL}^{\cos \phi} \right] \end{aligned}$$

$$| S_T | \left[\sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,I}^{\sin(\phi - \phi_S)} \right) \right]$$

$$\epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)}$$

$$+ \sqrt{2\epsilon(1 + \epsilon)} \left(\sin \phi_S F_{UT}^{\sin \phi_S} + \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \right)$$

$$\begin{aligned} & \lambda_e S_L \left[\sqrt{1 - \epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} + \sqrt{2\epsilon(1 - \epsilon)} \cos \phi_S F_{LT}^{\cos \phi_S} \right. \\ & \quad \left. + \sqrt{2\epsilon(1 - \epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \right] \end{aligned}$$

$$\begin{aligned} \frac{d^5\sigma_{\mathcal{I}}}{dx_{Bj}dQ^2d|t|d\phi d\phi_S} &= e_l \Gamma (T_{BH}^* T_{DVCS} + T_{DVCS}^* T_{BH}) \\ &= e_l \frac{\Gamma}{Q^2 |t|} \left\{ F_{UU}^{\mathcal{I}} + (2h)F_{LU}^{\mathcal{I}} + (2\Lambda)F_{UL}^{\mathcal{I}} + (2h)(2\Lambda)F_{LL}^{\mathcal{I}} + (2\Lambda_T)F_{UT}^{\mathcal{I}} + (2h)(2\Lambda_T)F_{LT}^{\mathcal{I}} \right\} \end{aligned}$$

BH-DVCS interference

Unpolarized

$$F_{UU}^{\mathcal{I}} = F_{UU}^{\mathcal{I},tw2} + \frac{K}{\sqrt{Q^2}} F_{UU}^{\mathcal{I},tw3}$$

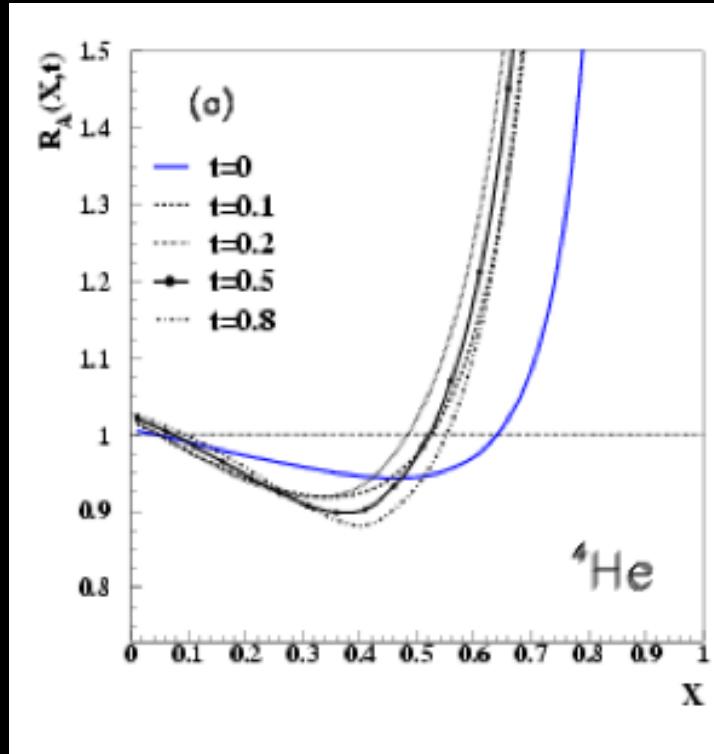
This formalism supersedes “harmonics”-based formalism

B. Kriesten, A.Meyer, SL, et al. arXiv:1903.05742

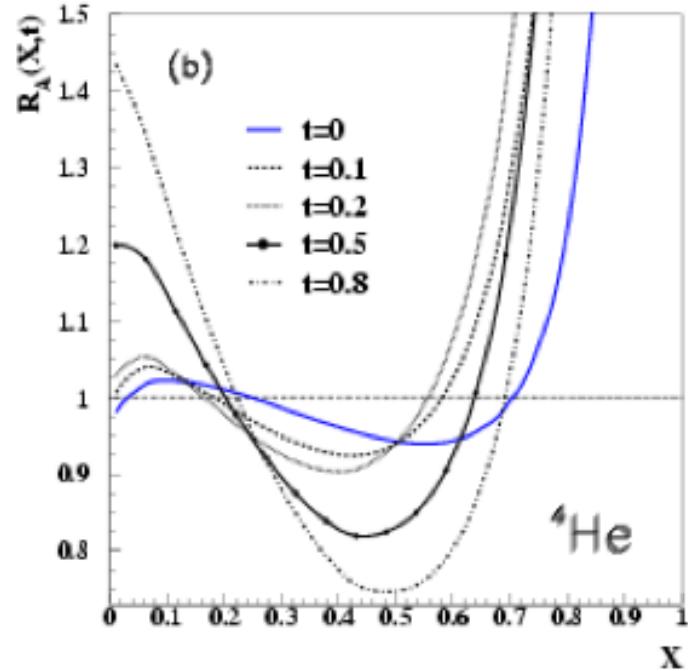
$$\begin{aligned} c_0^T &= -8(2-y) \frac{t}{Q^2} F_A \operatorname{Re}\{\mathcal{H}_A\} \\ &\quad \times \left\{ (2-x_A)(1-y) - (1-x_A)(2-y)^2 \left(1 - \frac{t_{min}}{Q^2}\right) \right\}, \\ c_1^T &= 8K(2y - y^2 - 2)F_A \operatorname{Re}\{\mathcal{H}_A\}, \\ s_1^T &= 8Ky(2-y)F_A \operatorname{Im}\{\mathcal{H}_A\}. \end{aligned}$$

Some predictions

No off-shell effects



Off-shell effects



Off shell effects grow with t because k_T becomes more important in this region

Conclusions and Outlook

- Many future possibilities for exploring the nature of nuclear modifications of the proton structure thanks to the accessibility of deeply virtual exclusive reactions
- Re-interactions/twist three is important: transverse d.o.f. connection between k_T and b
- Future Jlab and EIC Exclusive experiments using nuclei will provide the much needed laboratory to study QCD in coordinate space: vast phenomenology scenario
- BUT we need to start using the correct formalism