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# Transverse momentum imaging

## Lecture 5

Hampton University Graduate School (e-HUGS) 2021

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# Recap++ of lectures 3,4

# Quark TMD PDFs (spin 1/2)

		quark pol.		
		U	L	T
nucleon pol.	U	$f_1$		$h_1^\perp$
	L		$g_{1L}$	$h_{1L}^\perp$
	T	$f_{1T}^\perp$	$g_{1T}$	$h_1, h_{1T}^\perp$

At leading twist: 8 TMD PDFs  
(similar classification for gluons)

The **symmetries of QCD** play a crucial role in this classification

- **Black**: time-reversal even AND collinear
- **Blue**: time-reversal even
- **Red**: time-reversal odd (*process dependence*)

Quark inside spin 1/2 hadron

# Quark TMD PDFs (spin 1/2)

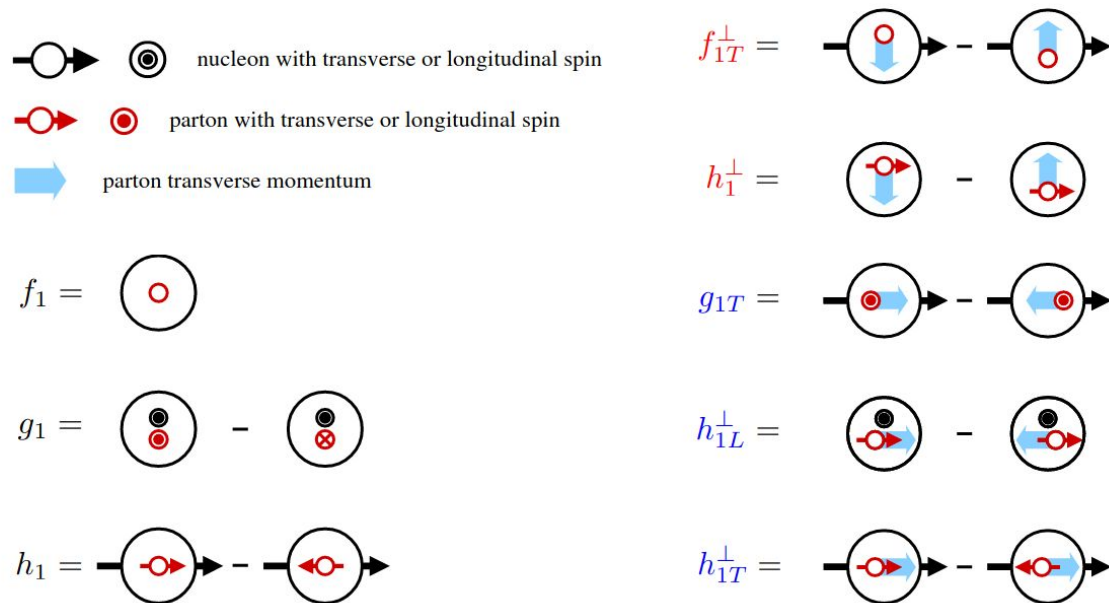
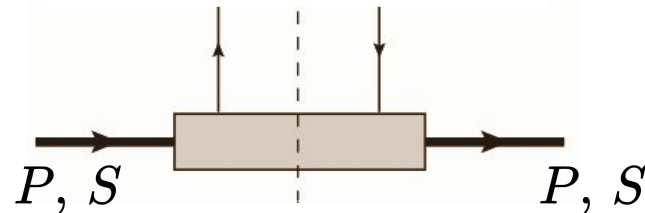


Figure 3.5: Probabilistic interpretation of twist-2 transverse-momentum-dependent distribution functions. To avoid ambiguities, it is necessary to indicate the directions of quark's transverse momentum, target spin and quark spin, and specify that the proton is moving out of the page, or alternatively the photon is moving into the page.

# Gauge invariant quark correlator

$$\Phi_{ij}(k, P, S) = \int \frac{d^4\xi}{(2\pi)^4} e^{i k \cdot \xi} \langle PS | \bar{\psi}_j(0) U(0, \xi) \psi_i(\xi) | PS \rangle$$



**GAUGE INVARIANT!**

$$U(x) = e^{i \alpha^a(x) t^a}$$

The Wilson line “bridges” the non-locality and makes the operator gauge invariant

$$U(0, \xi) \rightarrow U(0) U(0, \xi) U^\dagger(\xi)$$

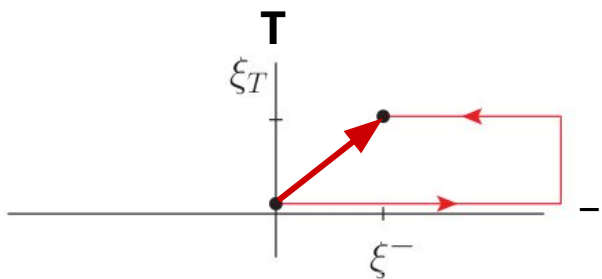
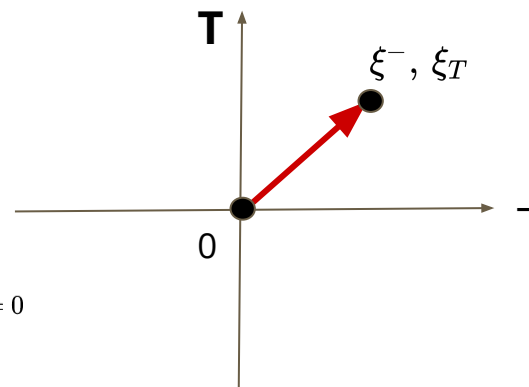
$$\bar{\psi}_j(0) U(0, \xi) \psi_i(\xi) \rightarrow \bar{\psi}_j(0) U^\dagger(0) U(0) U(0, \xi) U^\dagger(\xi) U(\xi) \psi_i(\xi) = \bar{\psi}_j(0) U(0, \xi) \psi_i(\xi)$$



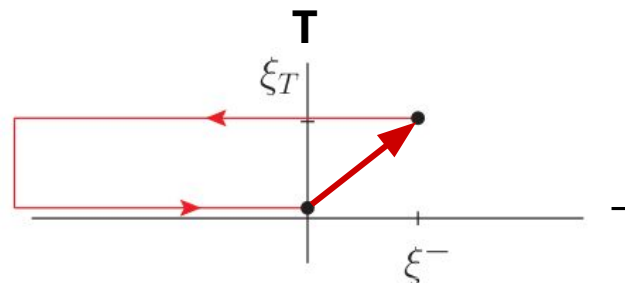
Eventually the correlator and the (TMD) PDFs **depend on the gauge link and its path** in spacetime

# Gauge links for TMD PDFs

$$\begin{aligned} \Phi_{ij}^{[U]}(x, \mathbf{p}_T, S) &= \int dp^+ dp^- \delta(p^+ - xP^+) \Phi^{[U]}(p, P, S) = \\ &= \int \frac{d\xi^- d^2\xi_T}{2\pi} e^{ip \cdot \xi} \langle PS | \bar{\psi}_j(0) U(0, \xi) \psi_i(\xi) | PS \rangle_{\xi^+ = 0} \end{aligned}$$



$U^{[+]}$  Future pointing (SIDIS)



$U^{[-]}$  Past pointing (Drell-Yan)

# Process dependence

The interplay between **time reversal** and **gauge symmetry** generates **relations** between the two configurations:

$$f_1^{a \ [+]}(x, k_T^2) = f_1^{a \ [-]}(x, k_T^2)$$

$$f_{1T}^{a \perp \ [+]}(x, k_T^2) = -f_{1T}^{a \perp \ [-]}(x, k_T^2)$$

T-even distribution

striking consequence  
of the symmetries of QCD

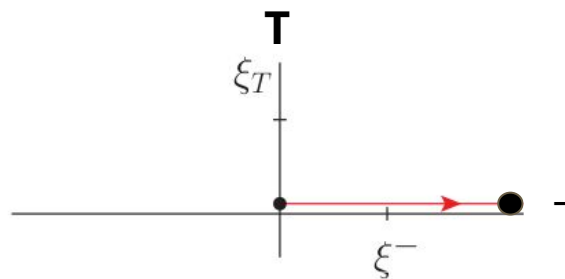
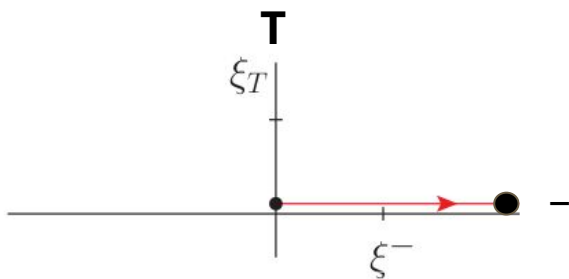
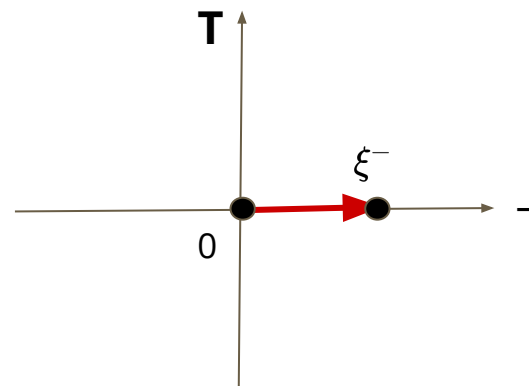
T-odd distribution



Sign-change relation for the Sivers function : not yet confirmed experimentally

# Gauge links for collinear PDFs (simpler)

$$\begin{aligned}\Phi_{ij}^{[U]}(x, S) &= \int dk^+ dk^- d^2\mathbf{k}_T \delta(k^+ - xP^+) \Phi^{[U]}(k, P, S) = \\ &= \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle PS | \bar{\psi}_j(0) U(0, \xi) \psi_i(\xi) | PS \rangle_{\xi^+ = \xi_T = 0}\end{aligned}$$



In the collinear limit the two gauge links reduce to the same object



# Factorized cross section

$$pp \longrightarrow \gamma / Z \longrightarrow l\bar{l} + X$$

$$\frac{d\sigma}{dq_T} \sim \mathcal{H} f_1(x_a, k_{T_a}, Q, Q^2) f_1(x_b, k_{T_b}, Q, Q^2) \delta^{(2)}(q_T - k_{T_a} - k_{T_b})$$

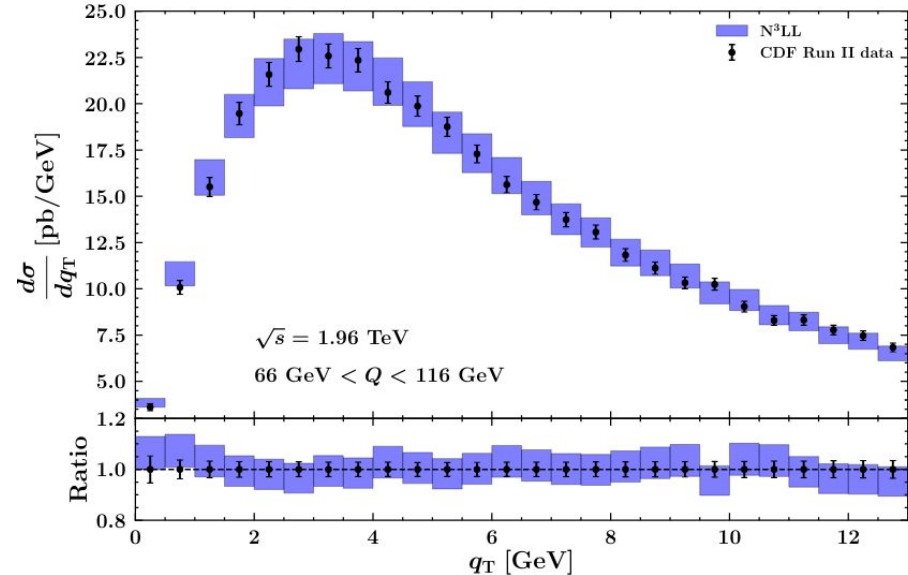
Renormalized  
TMDs

[TMD region,  $q_T \ll Q$ ]

+  $\mathcal{O}(q_T/Q) + \mathcal{O}(\Lambda/Q)$  [large  $q_T$  and low  $Q$  corrections]

Description of data:  
essential an  
approach with  
predictive power

Factorization  $\rightarrow$  renormalization  
(evolution) of TMDs

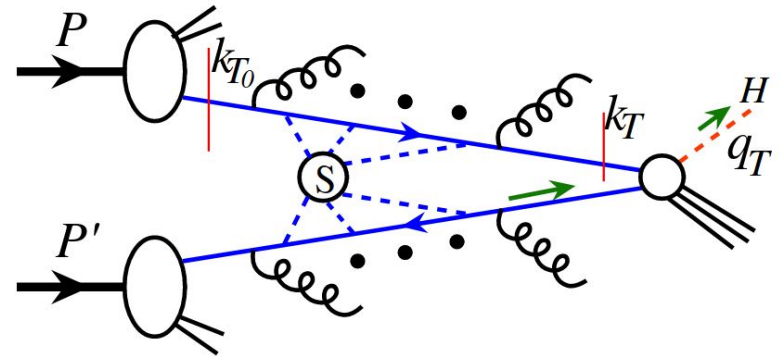


# Physical intuition

$$pp \longrightarrow \gamma / Z \longrightarrow l\bar{l} + X$$

$$\frac{d\sigma}{dq_T} \sim \mathcal{H} f_1(x_a, k_{T_a}, Q, Q^2) f_1(x_b, k_{T_b}, Q, Q^2) \delta^{(2)}(q_T - k_{T_a} - k_{T_b})$$

- The TMDs reproduce the structure of the **IR poles** in the cross section (same non-perturbative physics)
- The **observed transverse momentum** is accounted for by the transverse momenta of **quarks**
- The quark transverse momentum has **radiative** (perturbative) and **intrinsic** (non-perturbative) components
- Renormalization = **evolution** equations tell us how to distinguish between the two



# Evolved TMD distribution

$$F_a(x, b_T^2; \mu, \zeta) = F_a(x, b_T^2; \mu_0, \zeta_0) \quad \rightarrow \text{TMD distribution at initial scales}$$

$$\times \exp \left[ \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \gamma_F \left( \alpha_s(\mu'), \frac{\zeta}{\mu'^2} \right) \right] \quad \rightarrow \text{evolution in } \mu$$

$$\times \left( \frac{\zeta}{\zeta_0} \right)^{-D(b_T \mu_0, \alpha_s(\mu_0))} \quad \rightarrow \text{evolution in } \zeta$$

→ this solution is valid **at low  $b_T$**

See e.g. <https://inspirehep.net/literature/1785810> for more details (but also JCC book, etc.)

# Non perturbative components

$$F_a(x, b_T^2; \mu, \zeta) = F_a(x, b_T^2; \mu_0, \zeta_0) \quad \rightarrow \text{TMD distribution at initial scales}$$

$$\times \exp \left[ \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \gamma_F \left( \alpha_s(\mu'), \frac{\zeta}{\mu'^2} \right) \right] \quad \rightarrow \text{evolution in } \mu$$

$$\times \left( \frac{\zeta}{\zeta_0} \right)^{-D(b_T \mu_0, \alpha_s(\mu_0))} \boxed{+ g_K(b_T; \lambda)} \quad \rightarrow \text{evolution in } \zeta$$

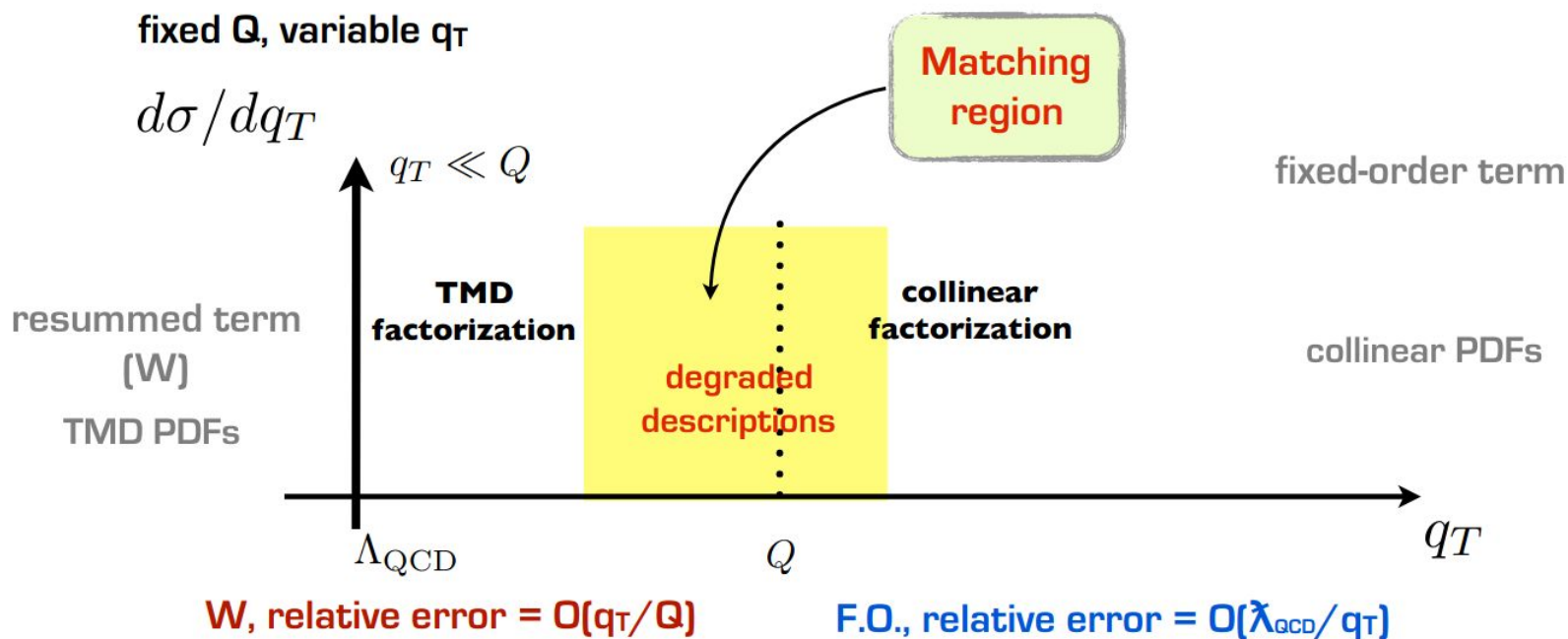
**Non-pert. correction to evolution (large  $b_T$ )**

$$F_a(x, b_T^2; \mu_0, \zeta_0) = \sum_b C_{a/b}(x, b_T^2, \mu_0, \zeta_0) \otimes f_b(x, \mu_0) \boxed{F_{NP}(b_T; \lambda)}$$

**Intrinsic contribution (large  $b_T$ )**

See e.g. <https://inspirehep.net/literature/1785810> for more details (but also JCC book, etc.)

# Matching TMD and collinear factorization

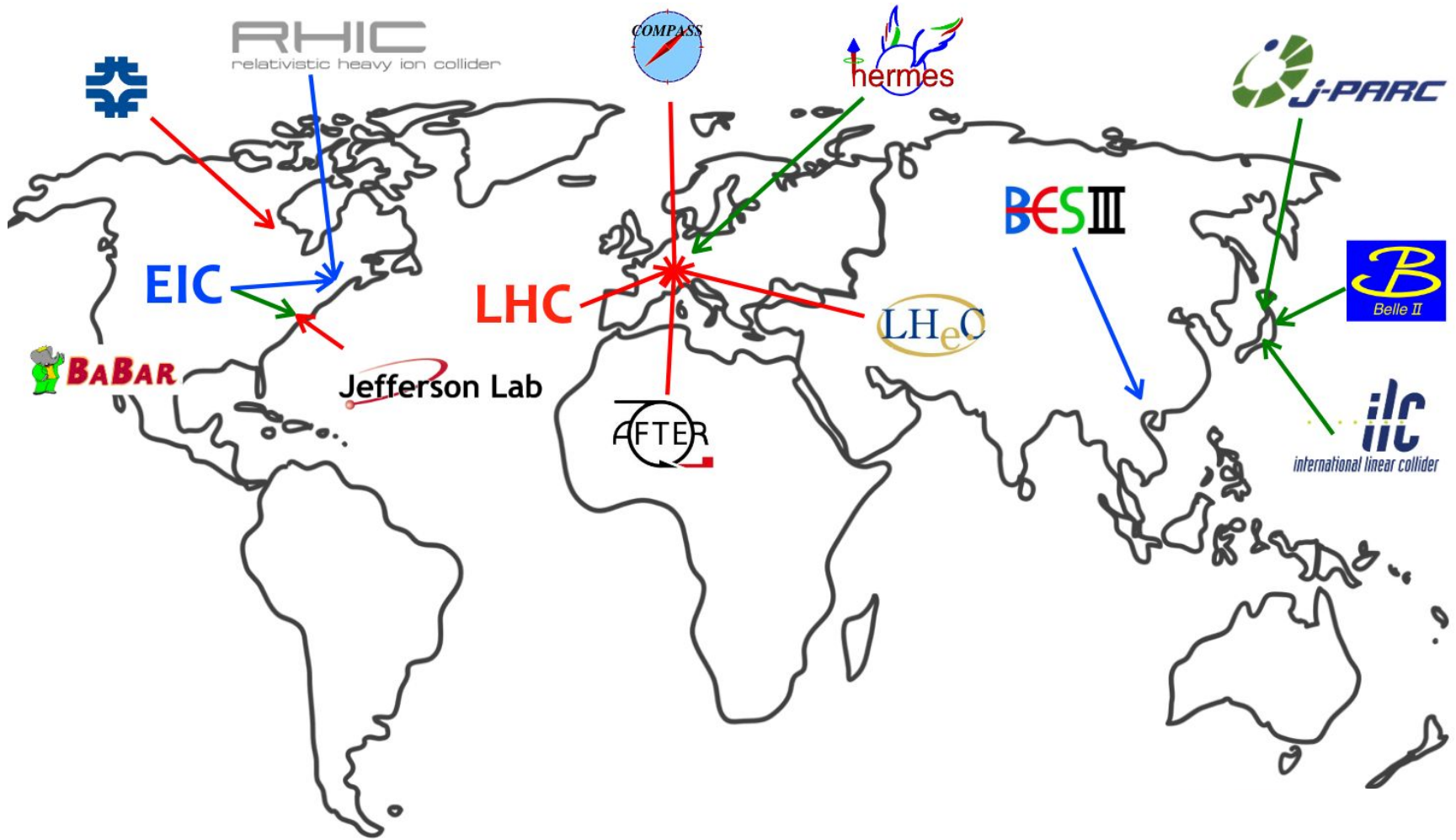


# Plan of these lectures

1. DIS and partons
2. From DIS to SIDIS
3. Symmetries and universality
4. Factorization, evolution, matching
5. Phenomenology

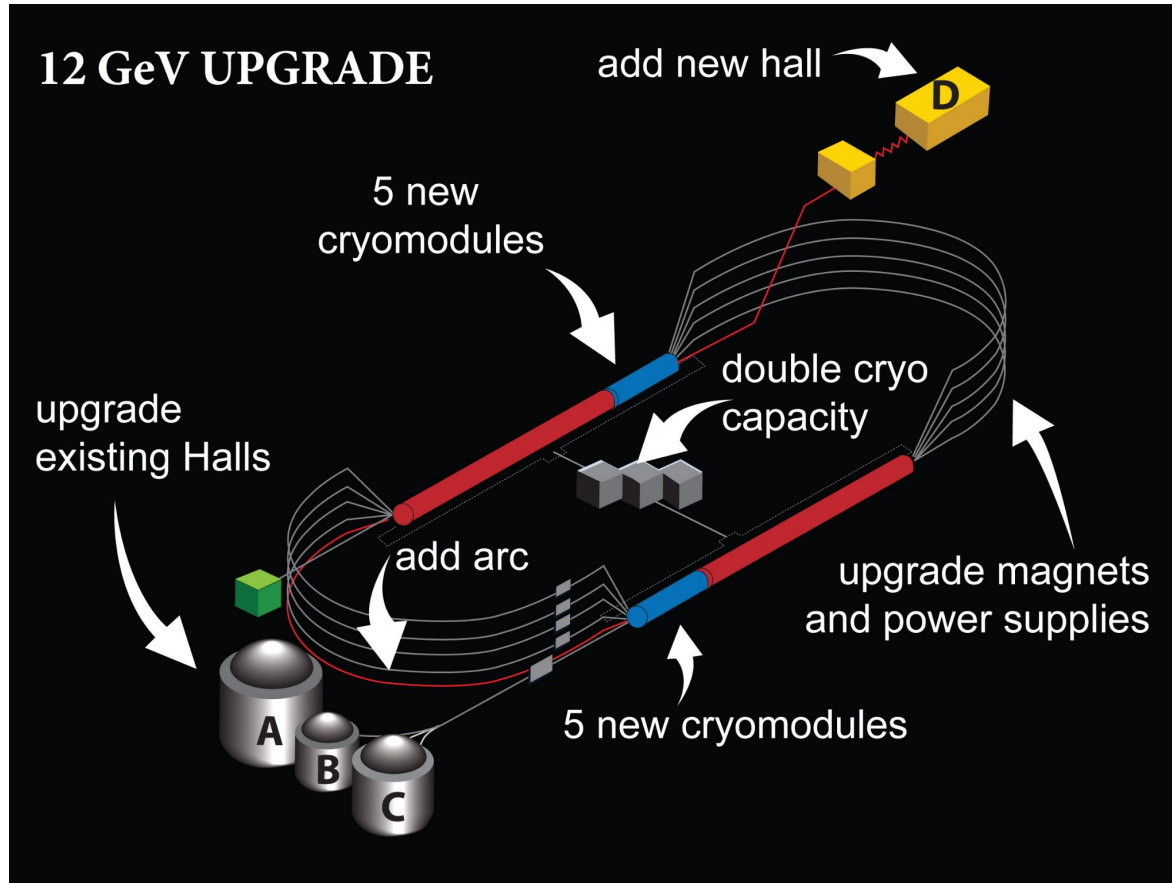
# 5.1 Experiments

(a short selection)





# CEBAF at Jefferson Lab



CEBAF:

Continuous Electron  
Beam  
Accelerator Facility

Built in 1984,  
recently completed a  
major upgrade  
from 6 GeV to 12 GeV  
+ one new hall



- Hall A & C: hadron structure, high luminosity
- Hall B: hadron structure,  $4\pi$  coverage
- Hall D: hadron spectroscopy

# The Electron-Ion Collider

- 00 home
- 01 about
- 02 goals
- 03 design
- 04 benefits
- 05 status
- 06 news

about

benefits

goals

status

design

news



The Electron-Ion Collider is a proposed machine for delving deeper than ever before into the building blocks of matter, so that we may better understand the matter within us and its role in the universe around us..

<https://www.jlab.org/eic>





## Precision 3D imaging of protons and nuclei

An Electron-Ion Collider will take three-dimensional precision snapshots of the internal structure of protons and atomic nuclei.

00 home

01 about

02 goals

03 design

04 benefits

05 status

06 news



## Solving the Mystery of Proton Spin

An EIC would reveal how the teeming quarks and gluons inside the proton combine their spins to generate the proton's overall spin.



## Search for Saturation

A unique form of matter, the color glass condensate, may be produced for study for the first time by an EIC, providing deeper insight into gluons and their interactions.

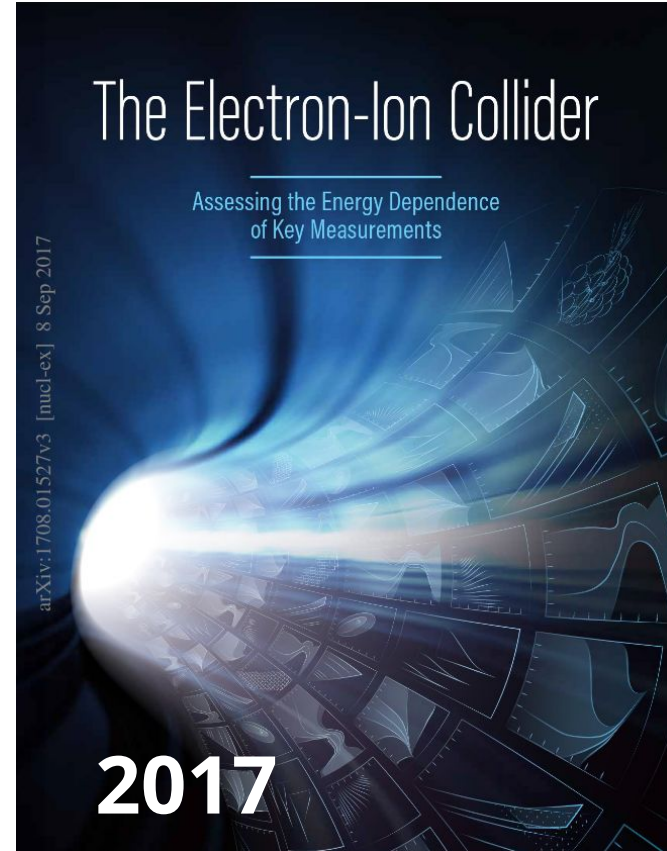
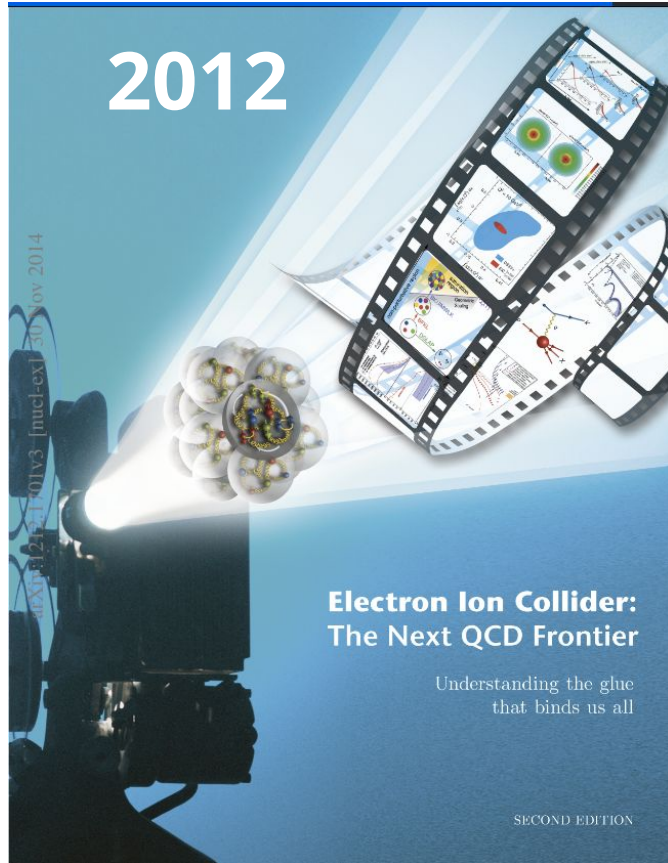


## Quark and Gluon Confinement

Experiments at an EIC would cast fresh light on the mystery of why quarks or gluons can never be observed in isolation but must remain confined within protons and nuclei.

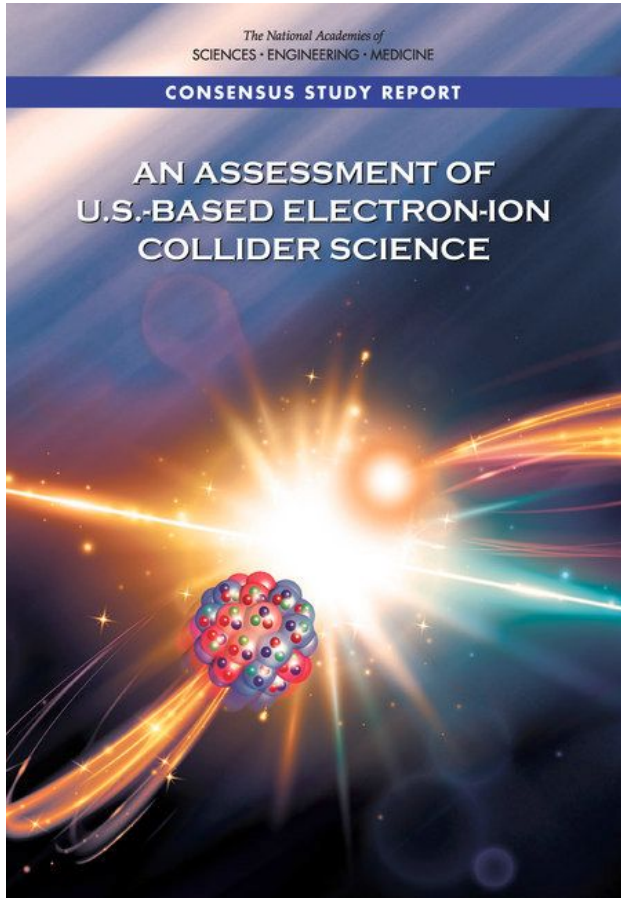
# The Electron-Ion Collider (EIC)

<https://www.bnl.gov/eic/>



# The Electron-Ion Collider (EIC)

<https://www.bnl.gov/eic/>



**2018**

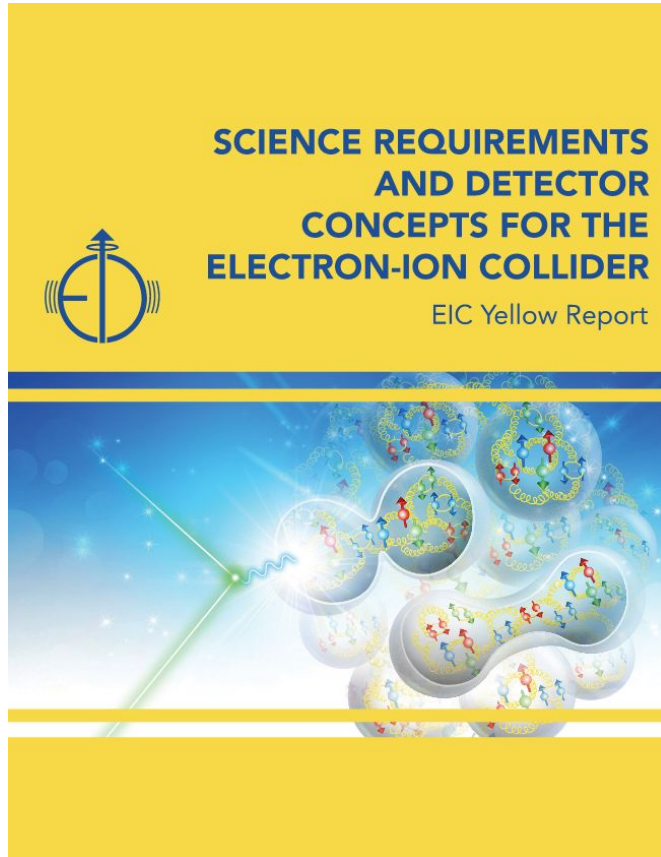
**2019:** DoE critical decision 0

**2020:** site selection (BNL)

**>2030 (?):** operations



# The Electron-Ion Collider (EIC)



**2021**

The "EIC Yellow Report"  
A community effort to line out the science requirements and detector concepts for the EIC

More details from the EIC Users Group website:  
<http://eicug.org/>

# A fixed-target program at the LHC



Contents lists available at [ScienceDirect](#)

## Physics Reports

journal homepage: [www.elsevier.com/locate/physrep](http://www.elsevier.com/locate/physrep)



## A fixed-target programme at the LHC: Physics case and projected performances for heavy-ion, hadron, spin and astroparticle studies



C. Hadjidakis<sup>1,a</sup>, D. Kikoła<sup>2,a</sup>, J.P. Lansberg<sup>1,\*a</sup>, L. Massacrier<sup>1,a</sup>,  
M.G. Echevarria<sup>3,4,b</sup>, A. Kusina<sup>5,b</sup>, I. Schienbein<sup>6,b</sup>, J. Seixas<sup>7,8,9,b</sup>, H.S. Shao<sup>10,b</sup>,  
A. Signori<sup>11,3,12,b</sup>, B. Trzeciak<sup>13,14,b</sup>, S.J. Brodsky<sup>15</sup>, G. Cavoto<sup>16</sup>, C. Da Silva<sup>17</sup>,  
F. Donato<sup>18</sup>, E.G. Ferreira<sup>19,20</sup>, I. Hřivnáčová<sup>1</sup>, A. Klein<sup>17</sup>, A. Kurepin<sup>21</sup>,  
C. Lorcé<sup>22</sup>, F. Lyonnet<sup>23</sup>, Y. Makdisi<sup>24</sup>, S. Porteboeuf Houssais<sup>25</sup>, C. Quintans<sup>8</sup>,  
A. Rakotozafindrabe<sup>26</sup>, P. Robbe<sup>1</sup>, W. Scandale<sup>27</sup>, N. Topilskaya<sup>21</sup>, A. Uras<sup>28</sup>,  
J. Wagner<sup>29</sup>, N. Yamanaka<sup>1,32,30,31</sup>, Z. Yang<sup>33</sup>, A. Zelenski<sup>24</sup>

<https://doi.org/10.1016/j.physrep.2021.01.002>



## 5.2 Collinear PDFs

# PDFs: what do we know

See <https://inspirehep.net/literature/1801417>

2020 PDFLATTICE REPORT

5

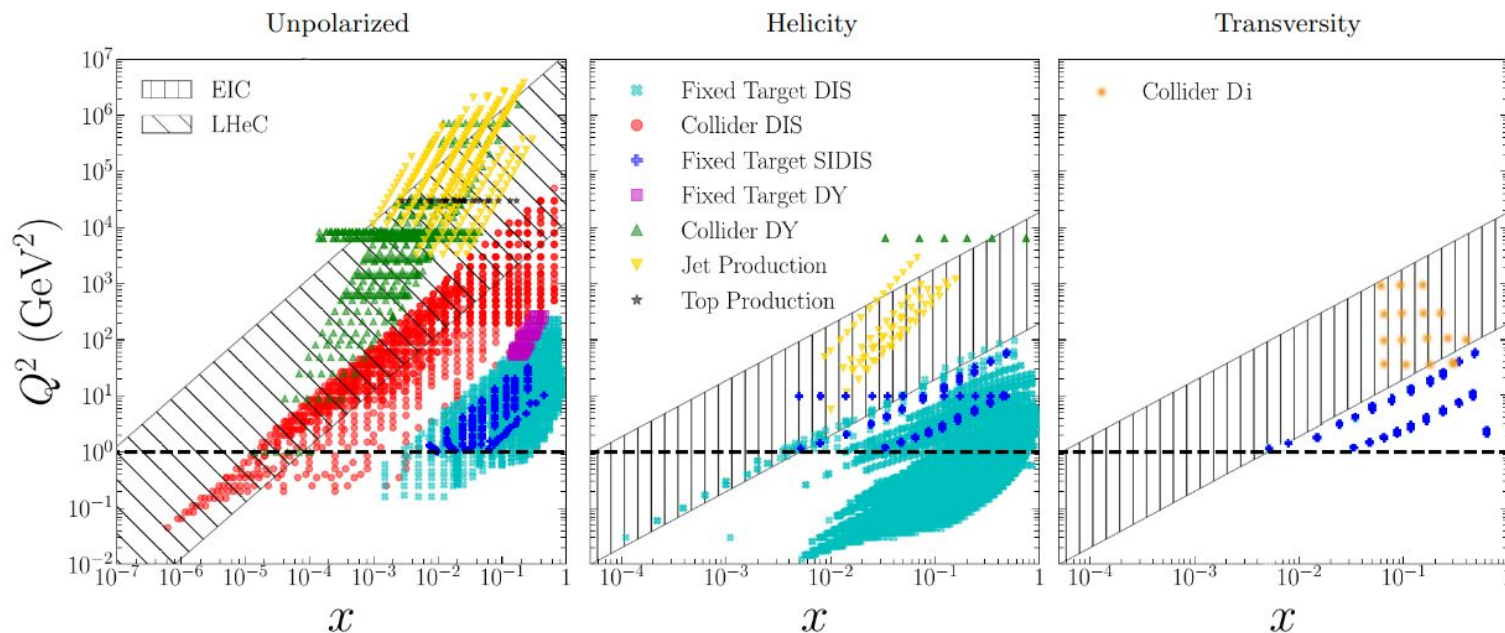


FIG. 1 The kinematic coverage in the  $(x, Q^2)$  plane of the hadronic cross-section data for the processes commonly included in global QCD analyses of collinear unpolarized, helicity, and transversity PDFs. The extended kinematic ranges attained by the LHeC and the EIC are also displayed. See Fig. 1 of Ref. (Ethier and Nocera, 2020) for unpolarized nuclear PDFs.

# PDFs: what do we know

See <https://inspirehep.net/literature/1801417>

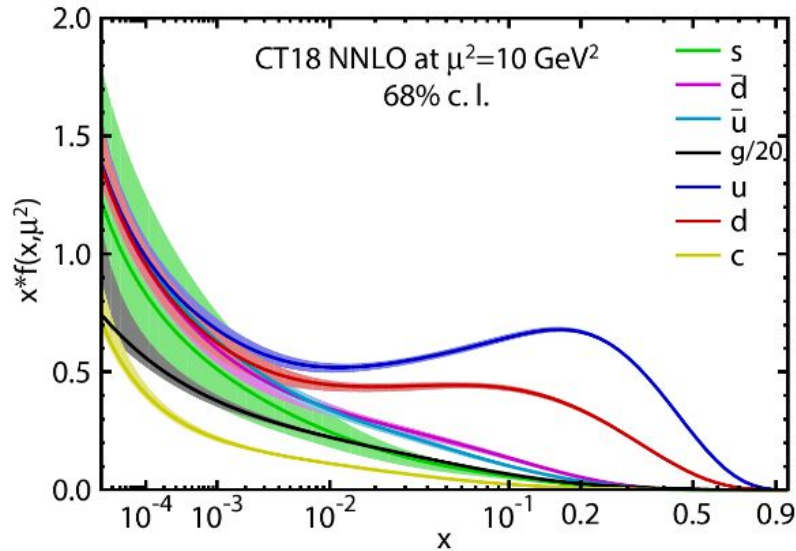


FIG. 2 The CT18 PDFs at  $\mu^2 = 10 \text{ GeV}^2$  for the  $xu$ ,  $x\bar{u}$ ,  $xd$ ,  $x\bar{d}$ ,  $xs = x\bar{s}$ , and  $xg$  PDFs. Error bands correspond to the 68% confidence level. Figure from (Kovářik *et al.*, 2019).

Many extractions available  
See e.g. the LHAPDF library  
<https://lhapdf.hepforge.org/>

(and W. Melnitchouk's talk)

# PDFs: what do we know

See <https://inspirehep.net/literature/1801417>

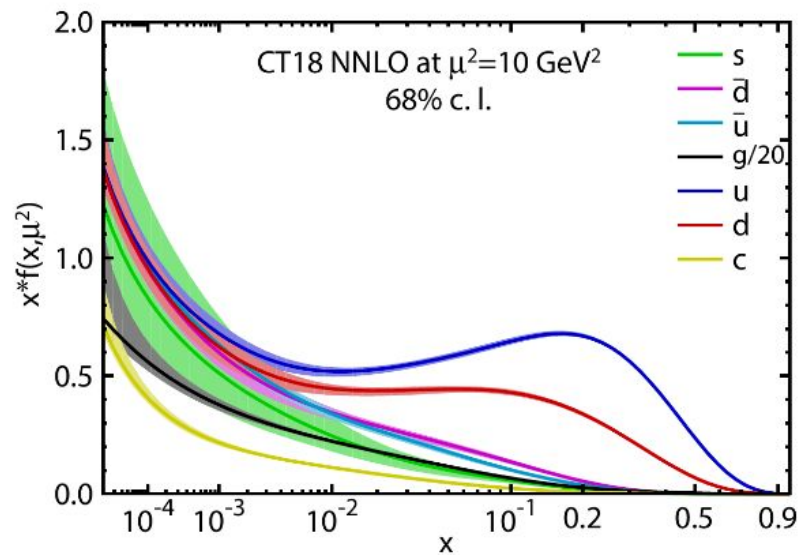


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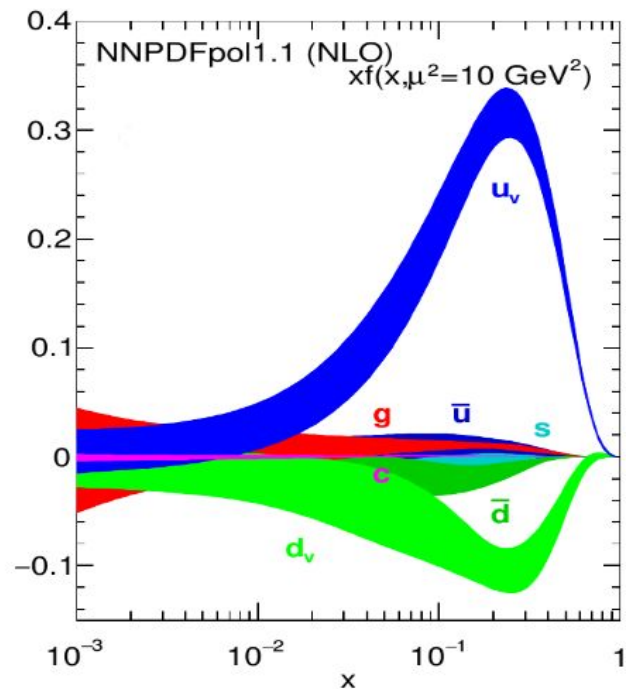


FIG. 5 The helicity PDFs from the NNPDFPOL1.1 parton set at  $\mu^2 = 10 \text{ GeV}^2$ . Figure from (Tanabashi *et al.*, 2018).

# PDFs: what do we know

See <https://inspirehep.net/literature/1801417>

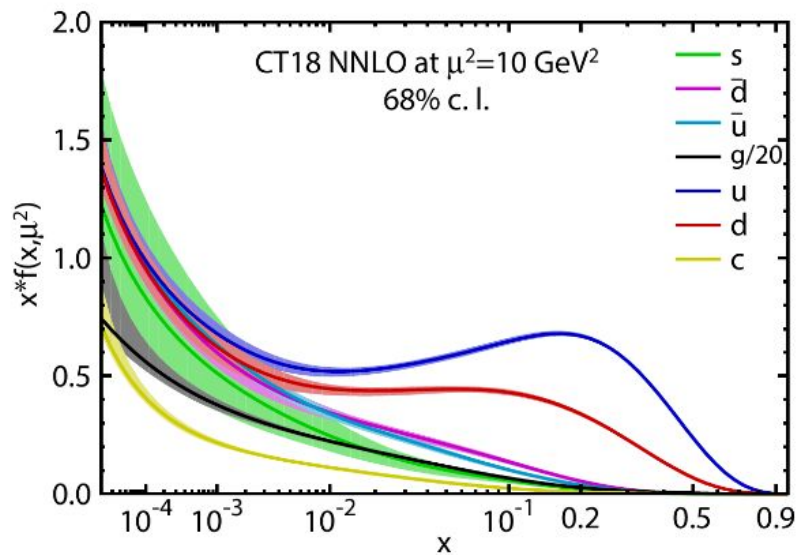


FIG. 2 The CT18 PDFs at  $\mu^2 = 10 \text{ GeV}^2$  for the  $xu$ ,  $x\bar{u}$ ,  $xd$ ,  $x\bar{d}$ ,  $xs = x\bar{s}$ , and  $xg$  PDFs. Error bands correspond to the 68% confidence level. Figure from (Kovářík *et al.*, 2019).

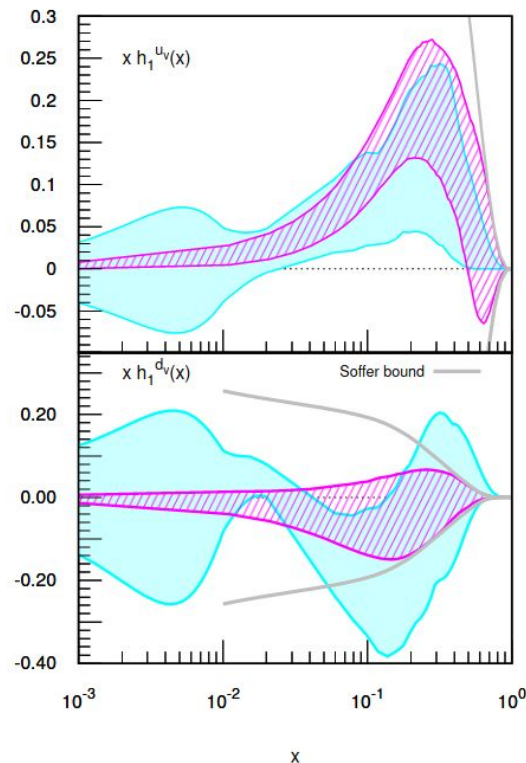
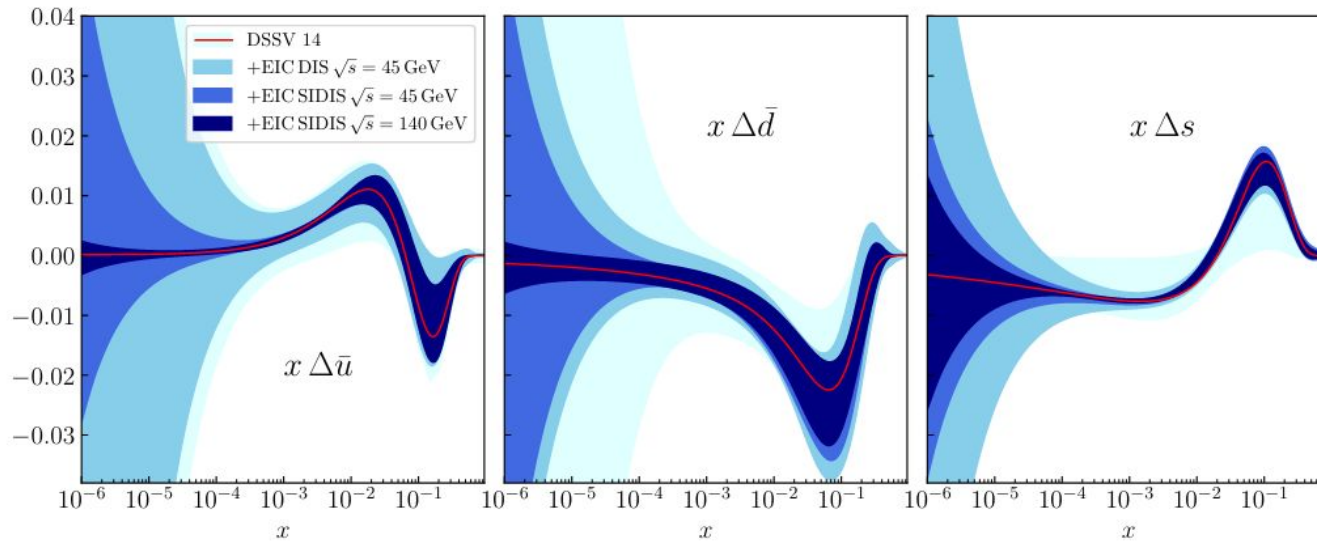


FIG. 6 The transversity  $x h_1(x)$  at 90% CL. Upper (lower) plot for valence up (down) component. Gray lines represent the Softer bound. Darker (pink) band for the PV18 global fit of (Radici and Bacchetta, 2018) at  $Q^2 = 2.4 \text{ GeV}^2$ . Lighter (cyan) band for the MEX19 constrained analysis of (Benel *et al.*, 2020) at the average scale of the data.

# Impact studies

See <https://inspirehep.net/literature/1851258>

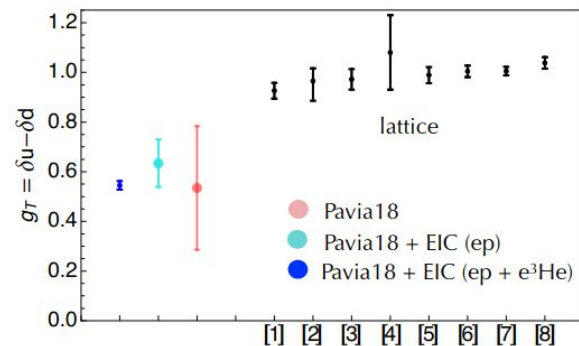
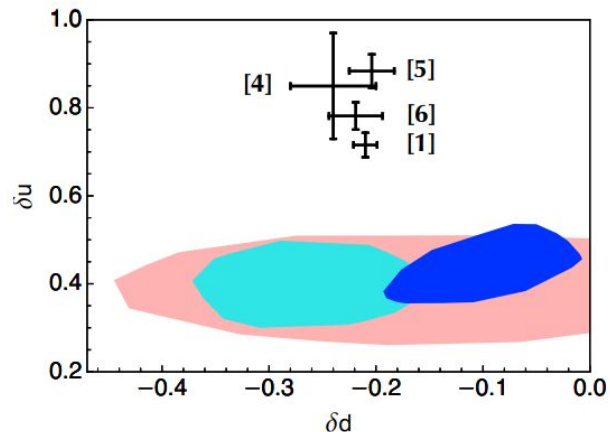
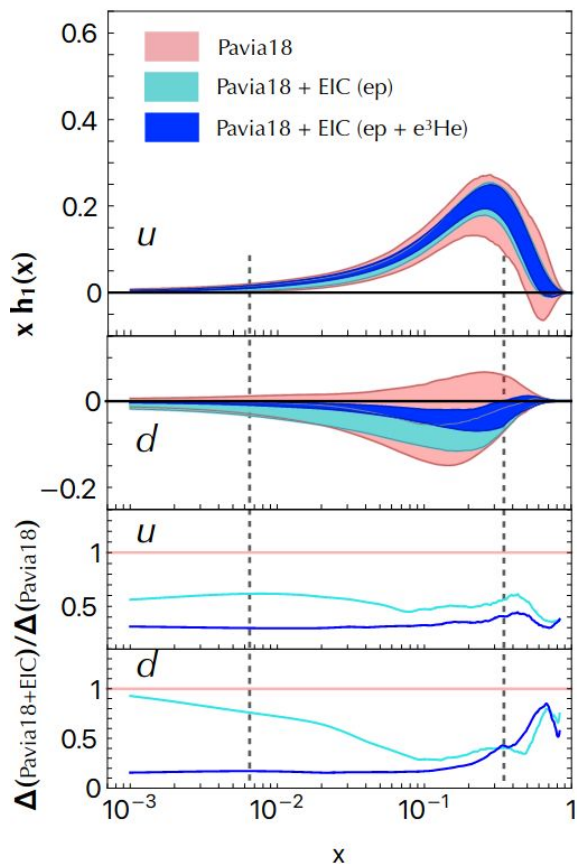


**Figure 7.19:** Impact of SIDIS measurements at the EIC on the sea quark helicities  $x\Delta\bar{u}$ ,  $x\Delta\bar{d}$  and  $x\Delta s$  as a function of  $x$  at  $Q^2 = 10 \text{ GeV}^2$ .



# Impact studies

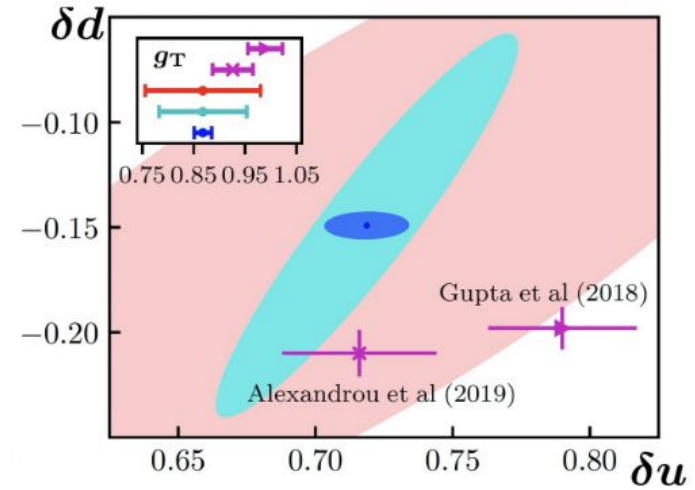
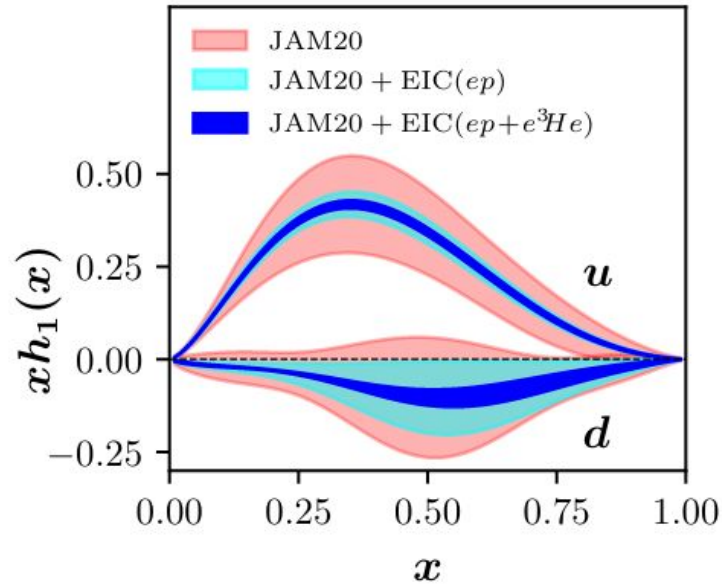
See <https://inspirehep.net/literature/1851258>



Transversity PDF and tensor charge

# Impact studies

See <https://inspirehep.net/literature/1851258>

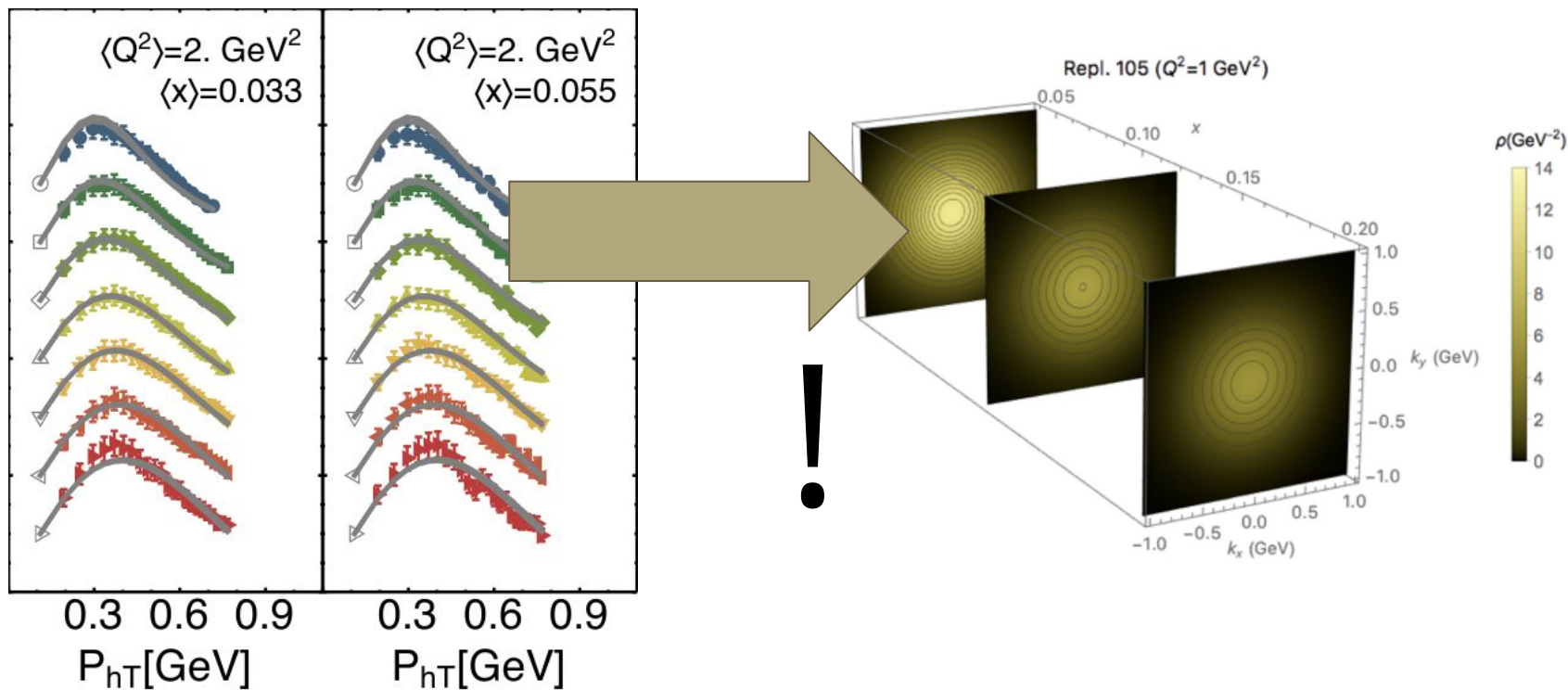


Transversity PDF and tensor charge



## 5.3 Unpolarized TMDs

# Transverse momentum imaging



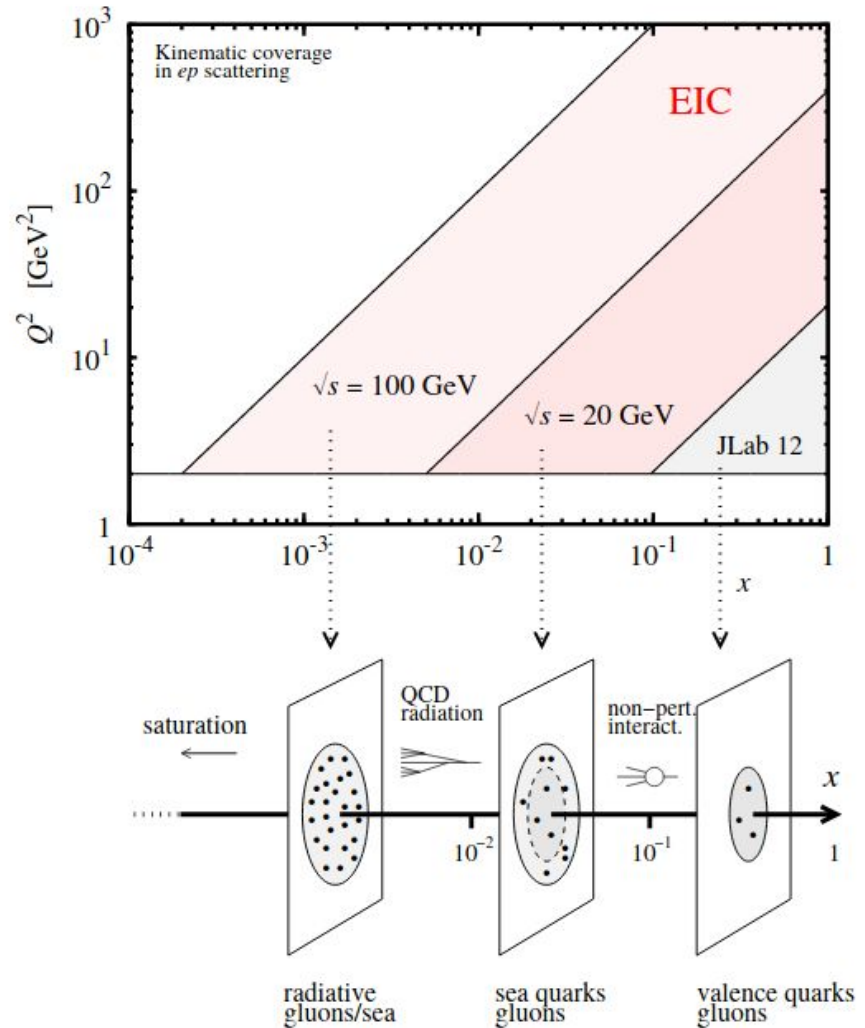
# Imaging from SIDIS

Importance of complementary experiments

from JLab 12 GeV, Hermes, Compass  
to the EIC

**zooming** into hadron structure

Credit picture: C. Weiss



# Non perturbative components

$$F_a(x, b_T^2; \mu, \zeta) = F_a(x, b_T^2; \mu_0, \zeta_0) \quad \rightarrow \text{TMD distribution at initial scales}$$

$$\times \exp \left[ \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \gamma_F \left( \alpha_s(\mu'), \frac{\zeta}{\mu'^2} \right) \right] \quad \rightarrow \text{evolution in } \mu$$

Calculable in pQCD

$$\times \left( \frac{\zeta}{\zeta_0} \right)^{-D(b_T \mu_0, \alpha_s(\mu_0)) + g_K(b_T; \lambda)} \quad \rightarrow \text{evolution in } \zeta$$

**Non-pert. corrections (large  $b_T$ )**

$$F_a(x, b_T^2; \mu_0, \zeta_0) = \sum_b C_{a/b}(x, b_T^2, \mu_0, \zeta_0) \otimes \underline{f_b(x, \mu_0)} F_{NP}(b_T; \lambda)$$

**Prior knowledge assumed (?)**

See e.g. <https://inspirehep.net/literature/1785810> for more details (but also JCC book, etc.)

# Separating small and large $b_T$

One needs to “separate” the small (perturbative)  $b_T$  region from the large (non-perturbative)  $b_T$  region:

$$\alpha_s(\mu = \mu_b \sim 1/b_T) \longrightarrow b_T < b_{max}$$

Avoid the Landau pole of QCD

$$\int_{\mu_b \sim 1/b}^Q \gamma_F, \mu_b < Q \longrightarrow b_T > b_{min}$$

Otherwise gluon “absorption” instead of “emission”

# Separating small and large $b_T$

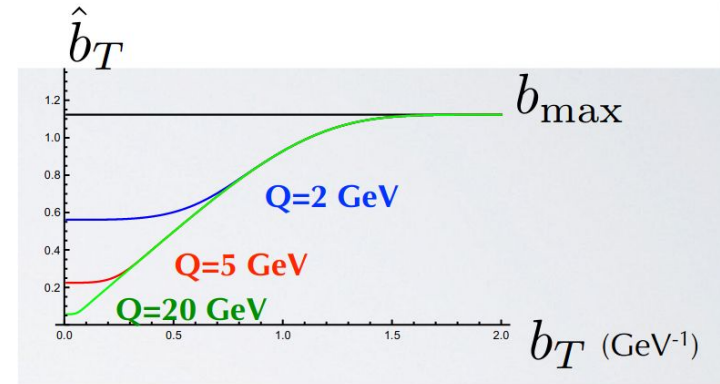
One needs to “separate” the small (perturbative)  $b_T$  region from the large (non-perturbative)  $b_T$  region:

$$\hat{b}(b_T; b_{\min}, b_{\max}) = b_{\max} \left( \frac{1 - e^{-b_T^4/b_{\max}^4}}{1 - e^{-b_T^4/b_{\min}^4}} \right) \begin{array}{l} \nearrow b_{\max}, \quad b_T \rightarrow +\infty \\ \searrow b_{\min}, \quad b_T \rightarrow 0 \end{array}$$

$$b_{\max} = 2e^{-\gamma E}$$

$$b_{\min} = 2e^{-\gamma E} / Q$$

These choices guarantee that for  $Q=1$  GeV the TMD coincides with the NP model



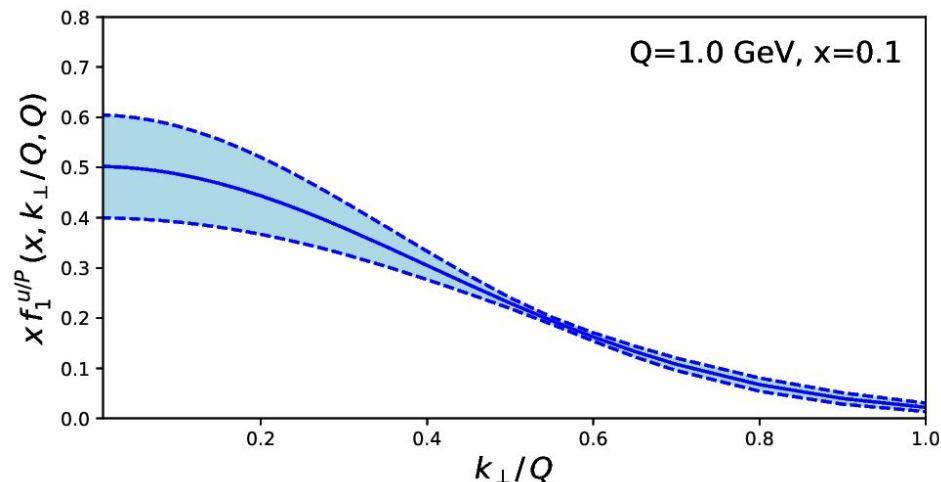
For more details see <https://inspirehep.net/literature/1520011>

# Recent fits of unpolarized TMDs

	Framework	HERMES	COMPASS	DY	Z production	N of points	$\chi^2/N_{\text{points}}$
Pavia 2017 <a href="#">arXiv:1703.10157</a>	NLL	✓	✓	✓	✓	8059	1.55
SV 2017 <a href="#">arXiv:1706.01473</a>	NNLL'	✗	✗	✓	✓	309	1.23
BSV 2019 <a href="#">arXiv:1902.08474</a>	NNLL'	✗	✗	✓	✓	457	1.17
SV 2019 <a href="#">arXiv:1912.06532</a>	NNLL'	✓	✓	✓	✓	1039	1.06
Pavia 2019 <a href="#">arXiv:1912.07550</a>	N <sup>3</sup> LL	✗	✗	✓	✓	353	1.02

# Global fit of unpolarized TMDs (PV17)

Imaging from **SIDIS** data (Hermes and Compass) **and Drell-Yan** data (Fermilab, low energy and Z)

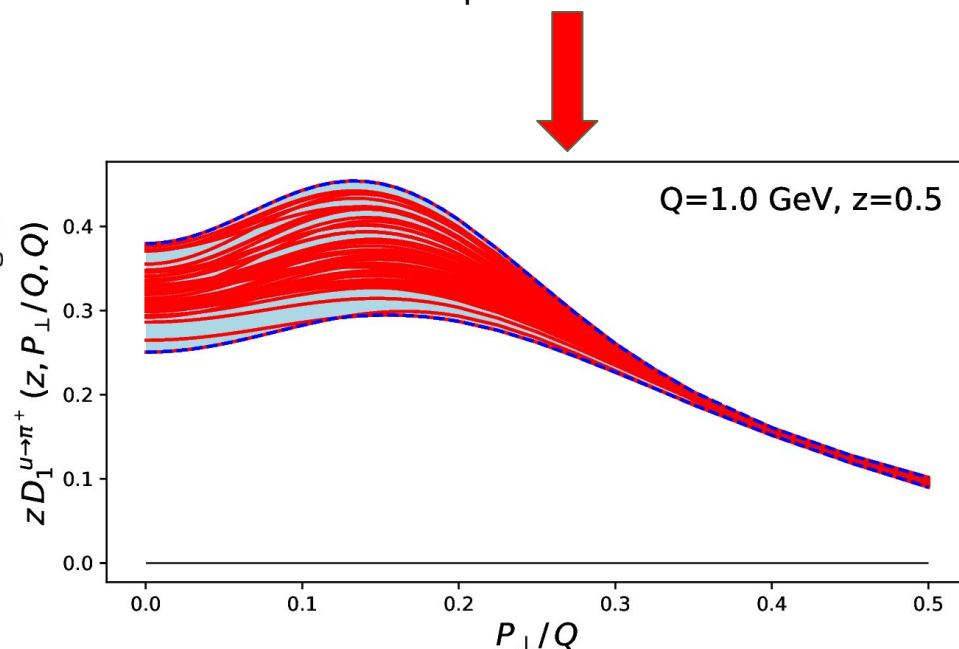


Combining SIDIS and Drell-Yan:  
Possibility to disentangle  
hadron structure and formation

See <https://inspirehep.net/literature/1520011>

← Unpolarized TMD PDF

Unpolarized TMD FF





# Some open questions

A non-exhaustive *personal* list of open questions:

- deepen our understanding of **sea** quarks
- **flavor structure** of TMDs
- experimental confirmation of **sign change** relation
- **gluon** observables and **spin-1** effects
- what can **hadronization** teach us about **confinement**?
- interplay between **nuclear/hadron** and **high-energy** physics
- ..

# Predictive power

Small  $bT$  → **perturbative (radiation)** contributions to TMD PDF

Large  $bT$  → **non-perturbative (intrinsic)** contributions to TMD PDF

Exercise:

In which kinematic regions is the TMD PDF dominated by small / large  $bT$  contributions ?

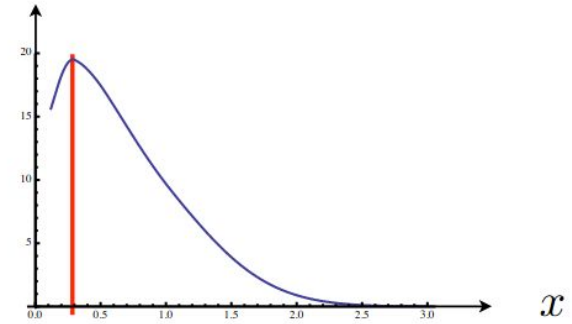
Hint: think about the shape of the TMD PDF in  $bT$  space and where it peaks

See e.g. <https://inspirehep.net/literature/1785810> for more details (but also JCC book, etc.)

# Saddle-point approximation

Function  $f$  with a maximum " $x_0$ " in  $(a,b)$

$$I(x_0, A) = \int_a^b dx e^{Af(x)} = e^{Af(x_0)} \sqrt{\frac{2\pi}{A(-f''(x_0))}} \left(1 + \mathcal{O}\left(\frac{1}{A}\right)\right)$$



One can apply this approximation to the TMD PDF and plot the position of the peak in  $b_T$  as a function of  $x$  and  $Q$ :

$$\frac{d}{db_T} \left\{ \ln \left[ b_T^2 F_a(x, b_T^2; Q, Q^2) \right] \right\}_{b_T=b_T^{sp}} = 0$$

The TMD PDF is dominated by perturbative contributions at large  $Q$  and small  $x$

See e.g. <https://inspirehep.net/literature/1785810> for more details (but also JCC book, etc.)

# 5.4 Polarized TMDs

# Spin asymmetries

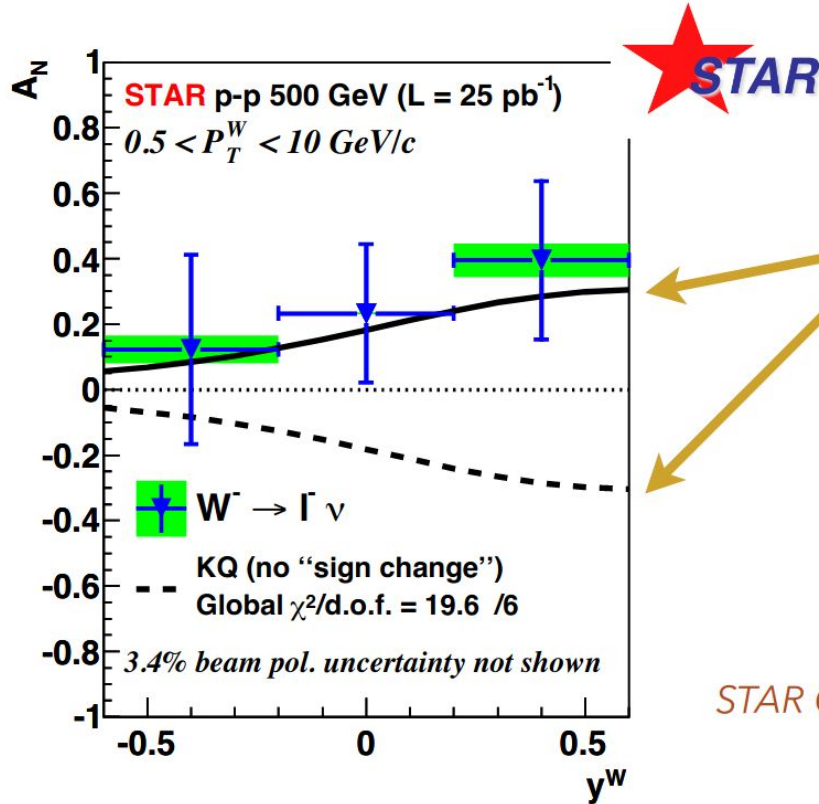
$$A_{UT} \sim \frac{d\sigma(\uparrow) - d\sigma(\downarrow)}{d\sigma(\uparrow) + d\sigma(\downarrow)} \quad \longrightarrow \quad \text{Polarized structure functions / unpolarized one}$$

$$\downarrow$$
$$f_1 \otimes D_1$$

Asymmetries in general have the benefits to :

- Single out specific structure functions by using Fourier analysis
- Reduce the effect of systematic uncertainties common to denominator and numerator (e.g. acceptance effects)
- Knowledge of the unpolarized cross section (denominator) is required in order to study the numerator

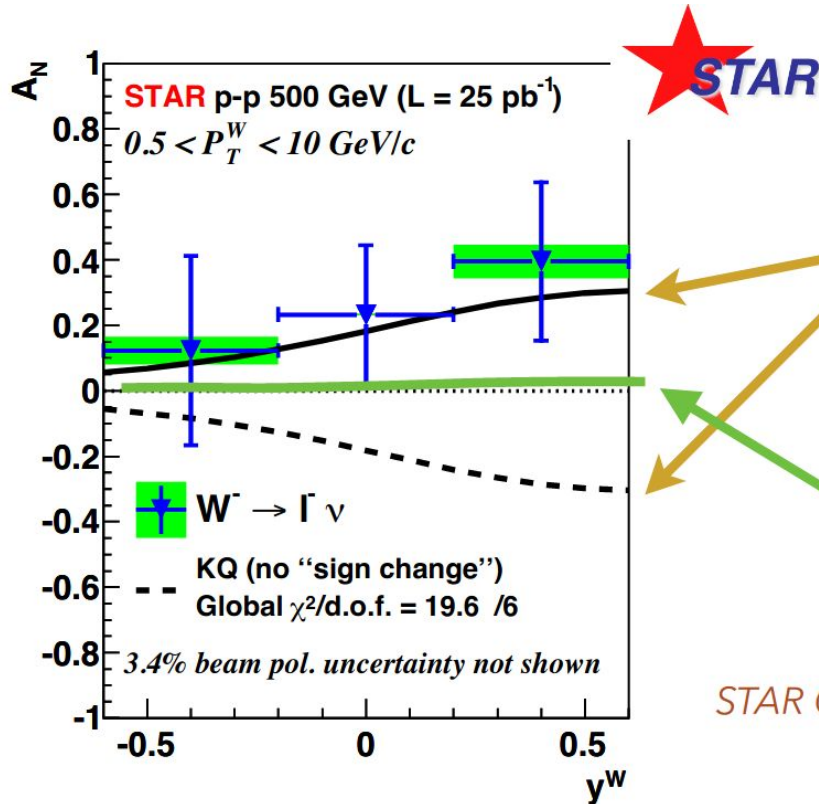
# Sign-change for Sivers function



first evidence  
of sign change?

STAR Collab. [arXiv:1511.06003](https://arxiv.org/abs/1511.06003)

# Sign-change for Sivers function



first evidence  
of sign change?

prediction with TMD  
evolution equations

STAR Collab. [arXiv:1511.06003](https://arxiv.org/abs/1511.06003)

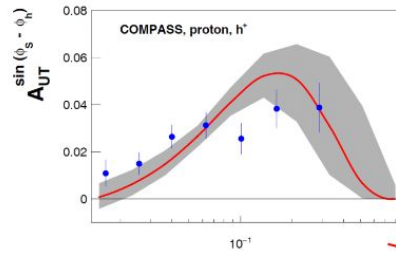
# Sign-change for Sivers function

## Sivers asymmetry in Semi-Inclusive DIS



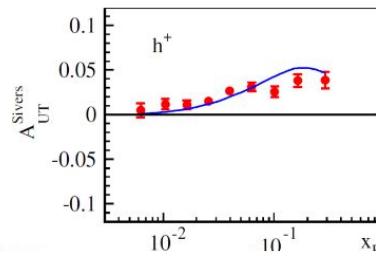
**DGLAP (2016)**

M. Anselmino et al., [arXiv:1612.06413](https://arxiv.org/abs/1612.06413)



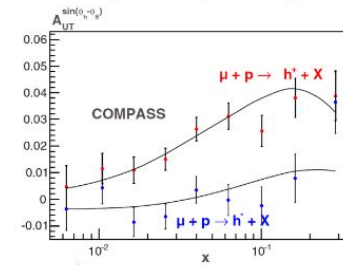
**TMD-1 (2014)**

M. G. Echevarria et al. [PRD89,074013](https://arxiv.org/abs/1407.0740)



**TMD-2 (2013)**

P. Sun, F. Yuan, [PRD88, 114012](https://arxiv.org/abs/1307.1140)





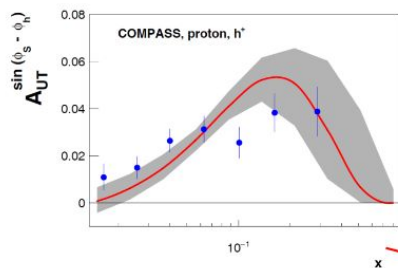
# Sign-change for Sivers function

Sivers asymmetry in Semi-Inclusive DIS



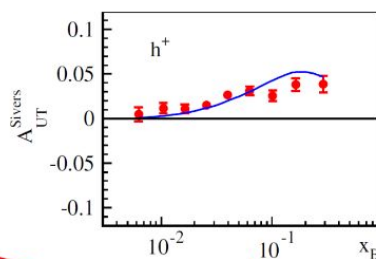
**DGLAP** (2016)

M. Anselmino et al., [arXiv:1612.06413](https://arxiv.org/abs/1612.06413)



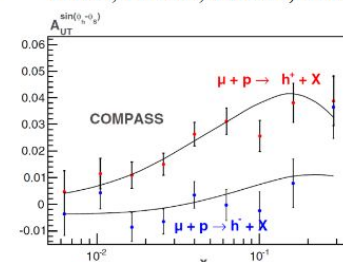
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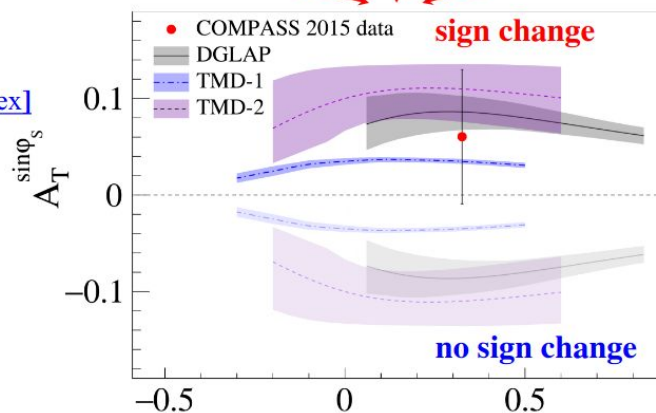


**New! 03 April 2017**

**COMPASS**

[CERN-EP-2017-059](https://arxiv.org/abs/1704.00488)

[arXiv:1704.00488\[hep-ex\]](https://arxiv.org/abs/1704.00488)

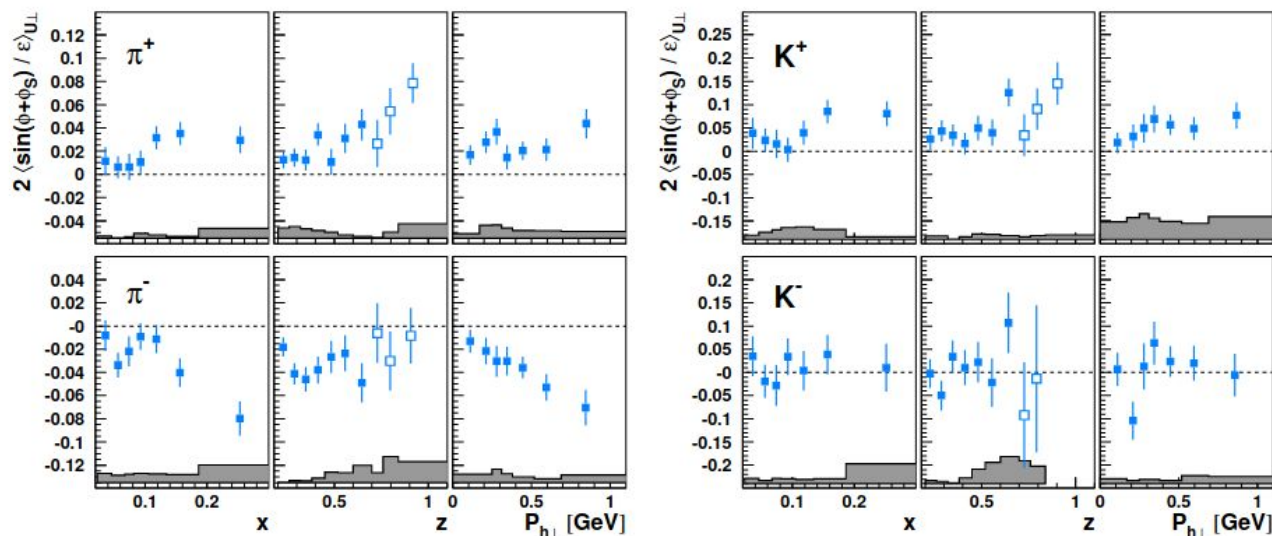


Sivers asymmetry in  
Drell-Yan

# Other TMDs from SIDIS

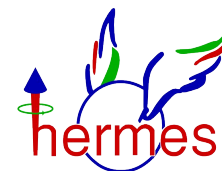
see <https://inspirehep.net/literature/1806922>

Transversity TMD PDF & Collins TMD FF



$$\frac{h_1 \otimes H_1^{\perp}}{f_1 \otimes D_1}$$

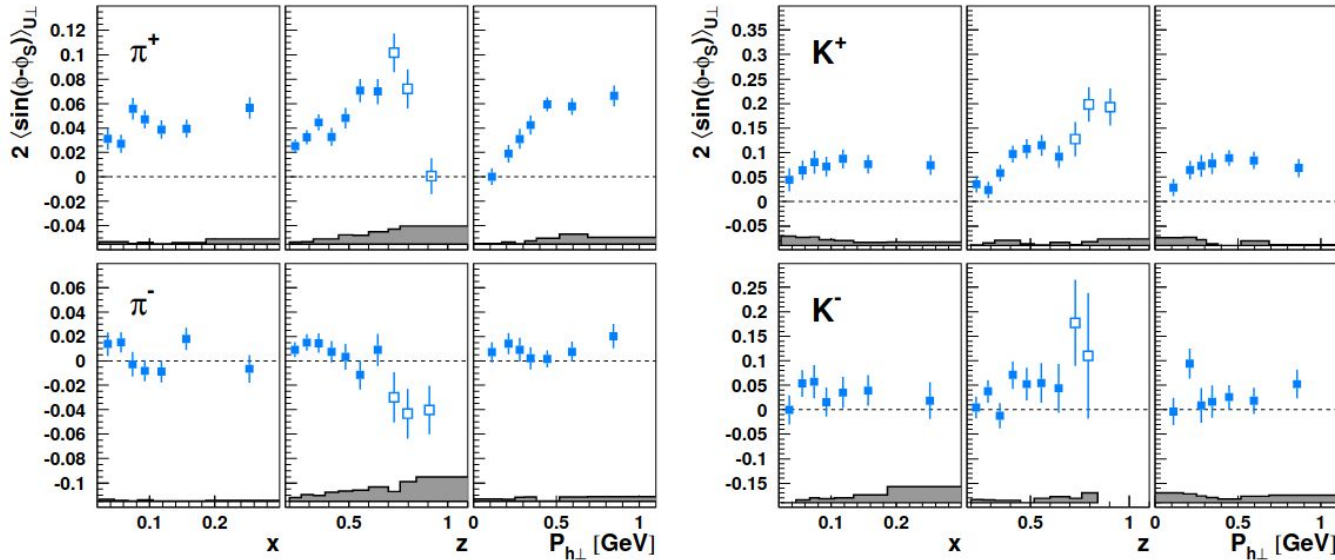
**Figure 8.** Collins SFA for charged mesons (left: pions; right: kaons) presented either in bins of  $x$ ,  $z$ , or  $P_{h\perp}$ . Data at large values of  $z$ , marked by open points in the  $z$  projection, are not included in the other projections. Systematic uncertainties are given as bands, not including the additional scale uncertainty of 7.3% due to the precision of the target-polarization determination.



# Other TMDs from SIDIS

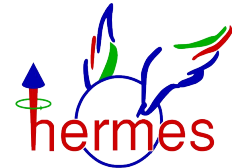
see <https://inspirehep.net/literature/1806922>

Sivers TMD PDF



$$\frac{f_{1T}^\perp \otimes D_1}{f_1 \otimes D_1}$$

**Figure 12.** Sivers SFA for charged mesons (left: pions; right: kaons) presented either in bins of  $x$ ,  $z$ , or  $P_{h\perp}$ . Data at large values of  $z$ , marked by open points in the  $z$  projection, are not included in the other projections. Systematic uncertainties are given as bands, not including the additional scale uncertainty of 7.3% due to the precision of the target-polarization determination.

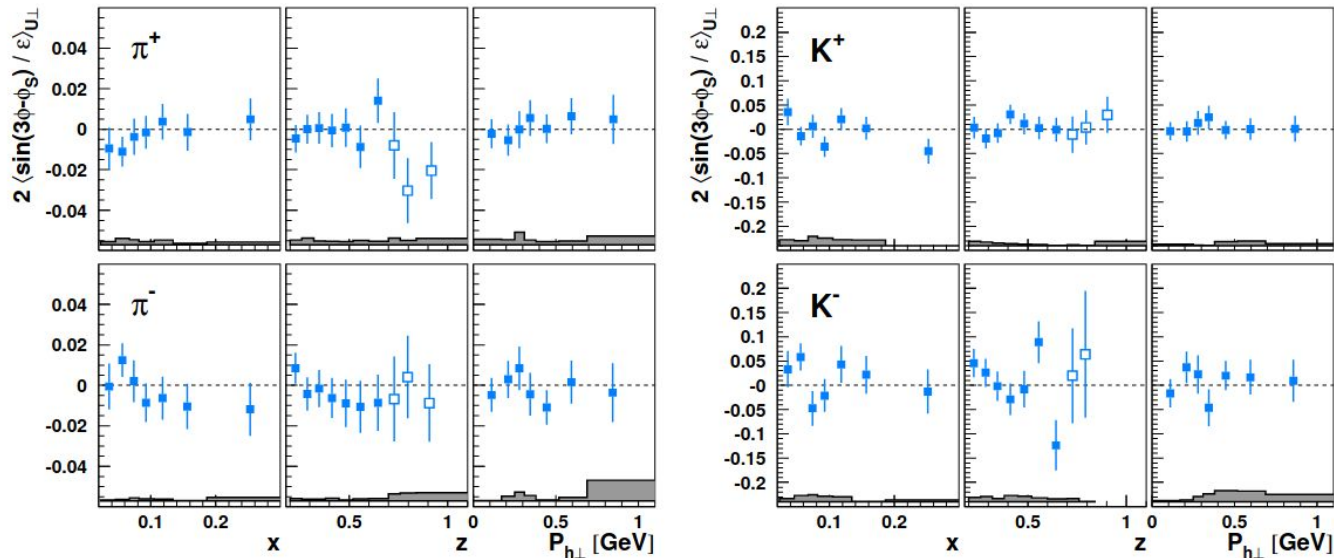


# Other TMDs from SIDIS

see <https://inspirehep.net/literature/1806922>

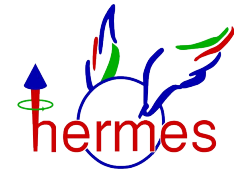
Pretzelosity TMD PDF & Collins TMD FF

(vanishing signal ..?)



$$\frac{h_{1T}^{\perp} \otimes H_1^{\perp}}{f_1 \otimes D_1}$$

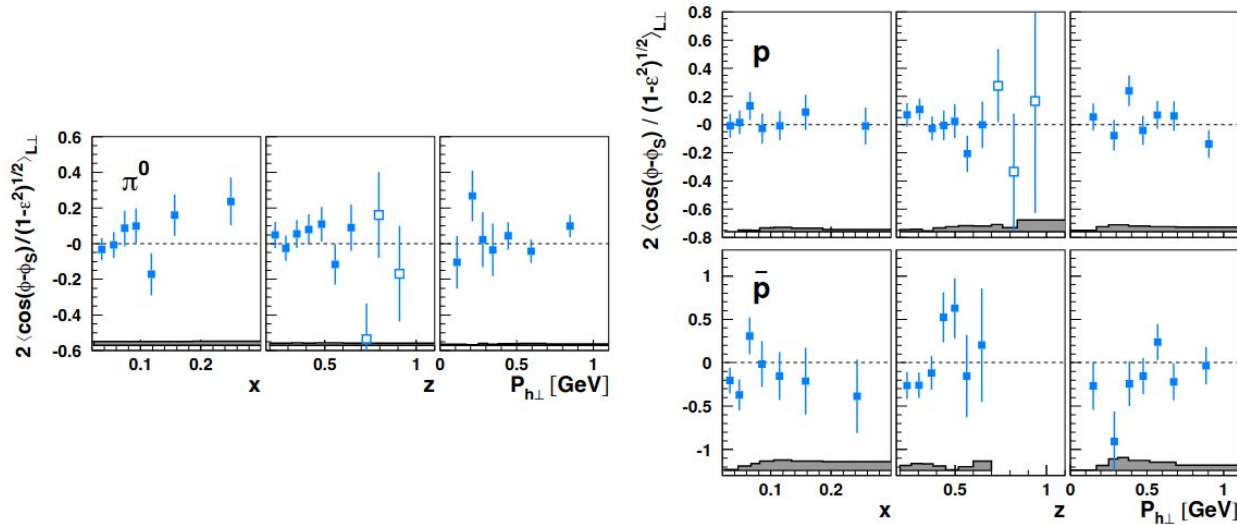
**Figure 19.** Pretzelosity SFA for charged mesons (left: pions; right: kaons) presented either in bins of  $x$ ,  $z$ , or  $P_{h\perp}$ . Data at large values of  $z$ , marked by open points in the  $z$  projection, are not included in the other projections. Systematic uncertainties are given as bands, not including the additional scale uncertainty of 7.3% due to the precision of the target-polarization determination.



# Other TMDs from SIDIS

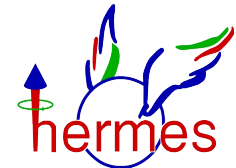
see <https://inspirehep.net/literature/1806922>

Worm-gear TMD PDF



**Figure 22.** The  $2\langle\cos(\phi - \phi_S)/\sqrt{1 - \epsilon^2}\rangle_{L\perp}^h$  amplitudes for  $\pi^0$  (left), protons, and antiprotons (right) presented either in bins of  $x$ ,  $z$ , or  $P_{h\perp}$ . Data at large values of  $z$ , marked by open points in the  $z$  projection, are not included in the other projections (no such high- $z$  points are available for antiprotons due to a lack of precision). Systematic uncertainties are given as bands, not including the additional scale uncertainty of 8.0% due to the precision in the determination of the target and beam polarizations.

$$\frac{g_{1T} \otimes D_1}{f_1 \otimes D_1}$$





# Other TMDs from SIDIS

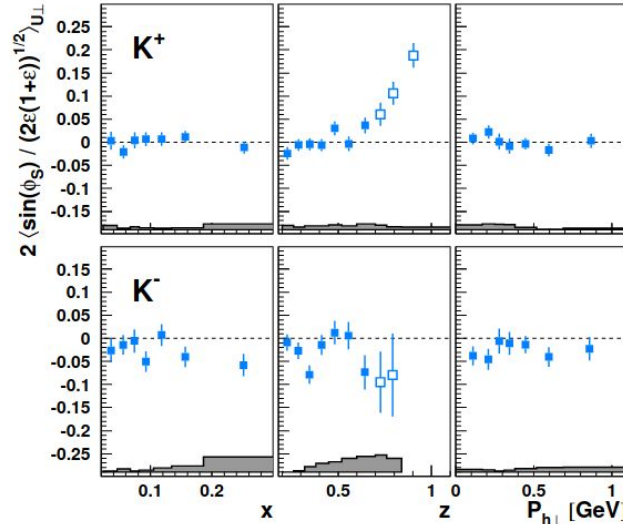
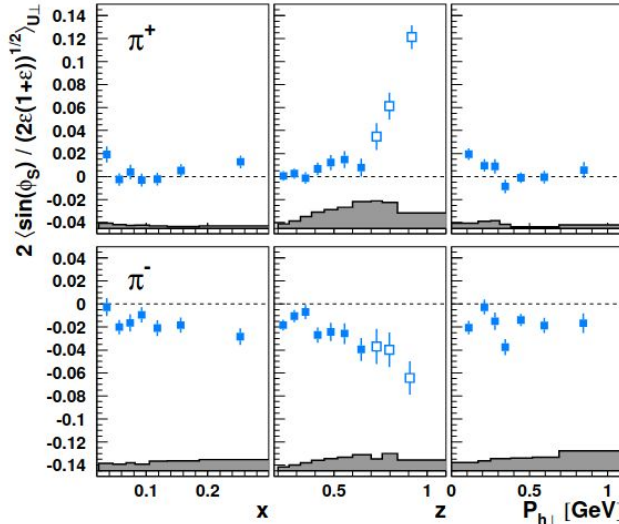
see <https://inspirehep.net/literature/1806922>

**Twist 3** TMD PDFs and TMD FFs

$$\frac{f_T \otimes D_1}{f_1 \otimes D_1}$$

$$\frac{h_1 \otimes \tilde{H}}{f_1 \otimes D_1}$$

$$\frac{h_T \otimes H_1^\perp}{f_1 \otimes D_1}$$



$$\frac{g_{1T} \otimes \tilde{G}^\perp}{f_1 \otimes D_1}$$

$$\frac{h_T^\perp \otimes H_1^\perp}{f_1 \otimes D_1}$$

$$\frac{f_{1T}^\perp \otimes \tilde{D}^\perp}{f_1 \otimes D_1}$$

**Figure 25.** The  $2\langle \sin(\phi_S) / \sqrt{2\epsilon(1+\epsilon)} \rangle_{U\perp}^h$  amplitudes for charged mesons (left: pions; right: kaons) presented either in bins of  $x$ ,  $z$ , or  $P_{h\perp}$ . Data at large values of  $z$ , marked by open points in the  $z$  projection, are not included in the other projections. Systematic uncertainties are given as bands, not including the additional scale uncertainty of 7.3% due to the precision of the target-polarization determination.

