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# Transverse momentum imaging

## Lecture 2

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# Plan of these lectures

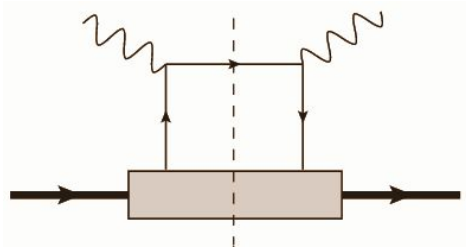
1. DIS and partons
2. From DIS to SIDIS
3. Symmetries and universality
4. Factorization, evolution, matching
5. Phenomenology

## 2. From DIS to SIDIS

# Where is transverse momentum?

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S) \quad \text{INCLUSIVE DIS} \rightarrow \text{differential in } x_B$$

We need a process with an **“experimental handle” on transverse momentum**, for example **Semi Inclusive DIS**



quark-antiquark

$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(x, S) \gamma^\mu \gamma^+ \gamma^\nu]$$

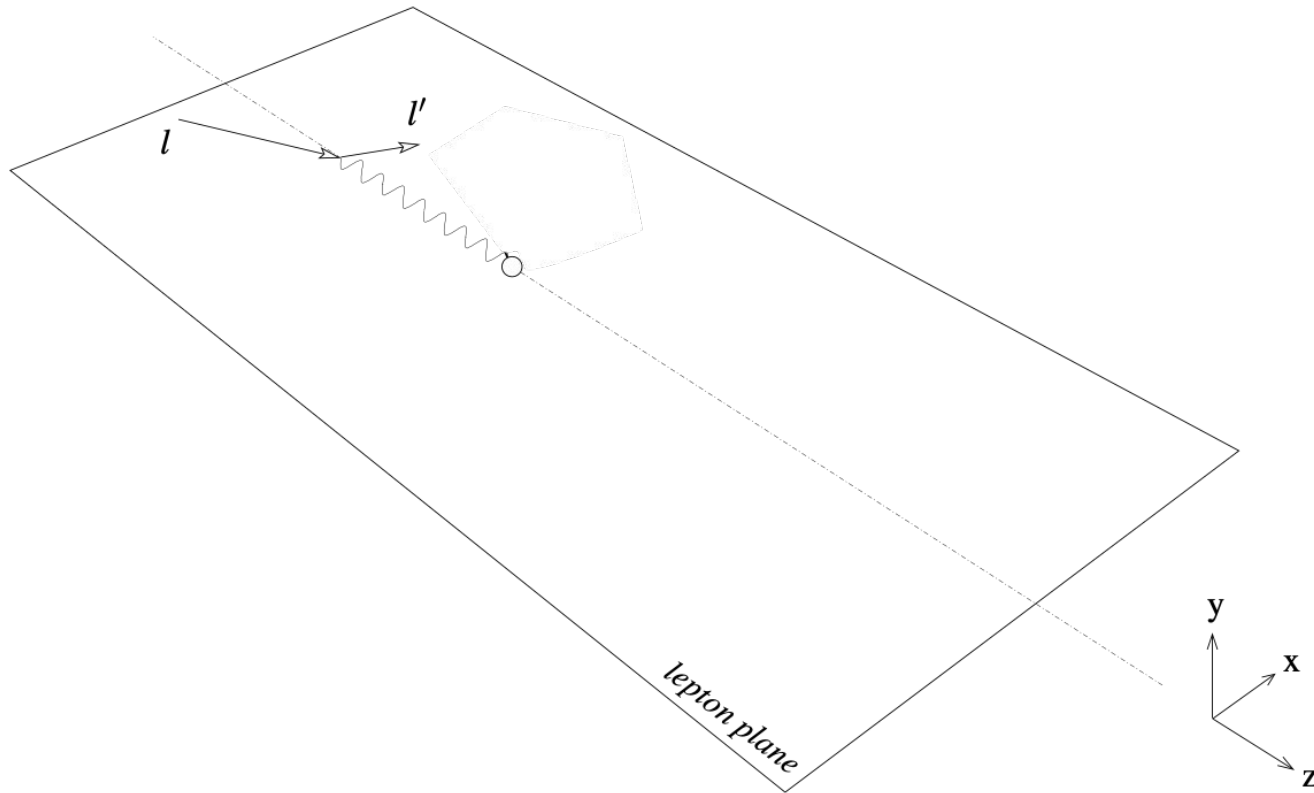
$\phi(x, S)$ : “collinear” quark correlator

$x_B \simeq x \equiv k^+ / P^+ \rightarrow$  measure collinear parton dynamics

The quark transverse momentum is integrated out in DIS

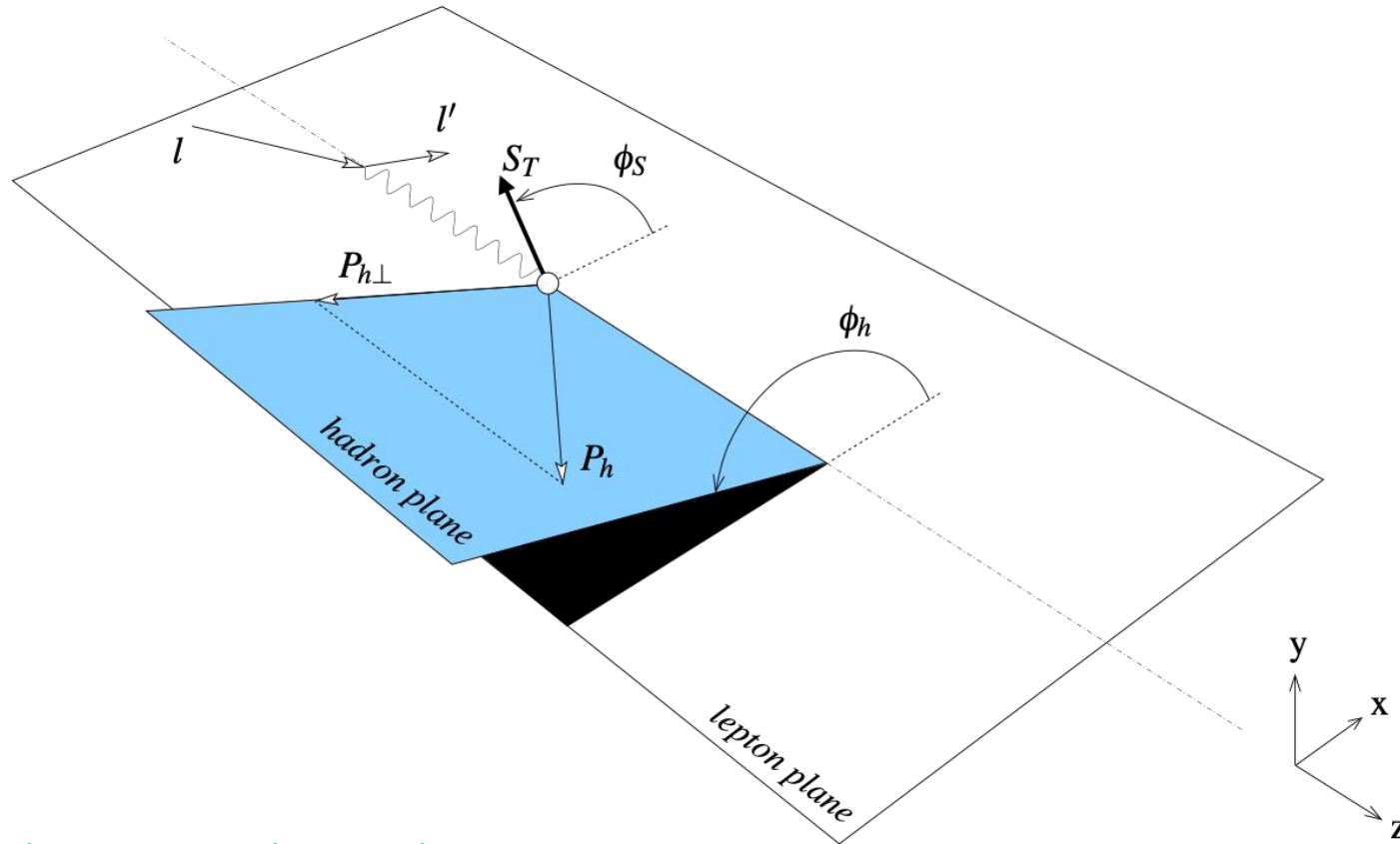
# Deep-inelastic scattering

$$l(\ell) + N(P) \rightarrow l'(\ell') + X(P_X)$$



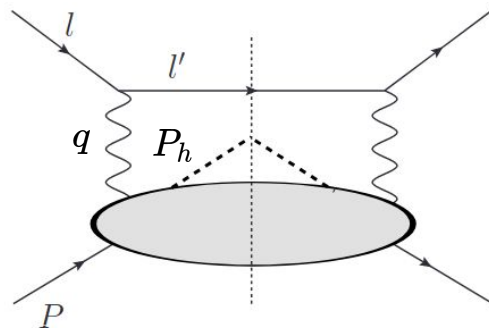
# Semi-Inclusive DIS

$$\ell(l) + N(P) \rightarrow \ell(l') + h(P_h) + X,$$



<https://inspirehep.net/literature/732275>

# Cross section DIS vs SIDIS



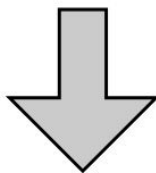
$$x_B = \frac{Q^2}{2P \cdot q}$$

$$z_h = \frac{P \cdot P_h}{P \cdot q}$$

“Handle” on collinear parton dynamics

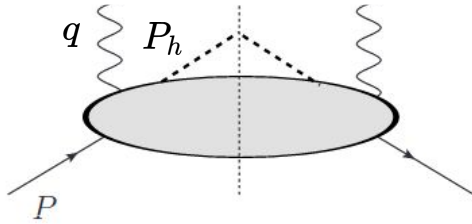
$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S) \quad \text{DIS}$$

“Handle” on transverse parton dynamics



$$\frac{d^6\sigma}{dx_B dy dz_h d\phi_S d\phi_h dP_{h\perp}^2} = \frac{\alpha^2 y}{2z_h Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S, P_h) \quad \text{SIDIS}$$

# Hadronic tensor (unpolarized)



Compared to DIS, there are five structure functions instead of two for unpolarized target

They depend on two extra variables

$$2MW^{\mu\nu}(q, P, S) = \frac{2z_h}{x_B} \left[ -g_{\perp}^{\mu\nu} F_{UU,T}(x_B, z_h, P_{h\perp}^2, Q^2) + \hat{t}^{\mu}\hat{t}^{\nu} F_{UU,L}(x_B, z_h, P_{h\perp}^2, Q^2) \right. \\ \left. + \left( \hat{t}^{\mu}\hat{h}^{\nu} + \hat{t}^{\nu}\hat{h}^{\mu} \right) F_{UU}^{\cos\phi_h}(x_B, z_h, P_{h\perp}^2, Q^2) + \left( \hat{h}^{\mu}\hat{h}^{\nu} + g_{\perp}^{\mu\nu} \right) F_{UU}^{\cos 2\phi_h}(x_B, z_h, P_{h\perp}^2, Q^2) \right. \\ \left. - i \left( \hat{t}^{\mu}\hat{h}^{\nu} - \hat{t}^{\nu}\hat{h}^{\mu} \right) F_{LU}^{\sin\phi_h}(x_B, z_h, P_{h\perp}^2, Q^2) \right],$$

$$\hat{h} = \frac{P_{h\perp}}{|P_{h\perp}|}$$



# SIDIS cross section (unpolarized)

$$\frac{d\sigma}{dx dy d\phi_S dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{x y Q^2} \frac{y^2}{2(1-\varepsilon)} \left\{ \begin{aligned} & F_{UU,T}(x, z, P_{h\perp}^2, Q^2) \\ & + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} \\ & + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \end{aligned} \right\}$$

5 structure functions for unpolarized target

# SIDIS cross section (polarized nucleon - spin 1/2)

$$\begin{aligned}
 & \frac{d\sigma}{dx dy d\phi_S dz d\phi_h dP_{h\perp}^2} \\
 &= \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} \right. \\
 &+ \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} + S_L \left[ \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 &+ S_L \lambda_e \left[ \sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 &+ S_T \left[ \sin(\phi_h - \phi_S) \left( F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} \right. \\
 &+ \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} \\
 &+ \left. \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] + S_T \lambda_e \left[ \sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} \right. \\
 &+ \left. \left. \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\}
 \end{aligned}$$

18 structure functions  
for polarized nucleon target

# SIDIS cross section (polarized deuteron - spin 1)

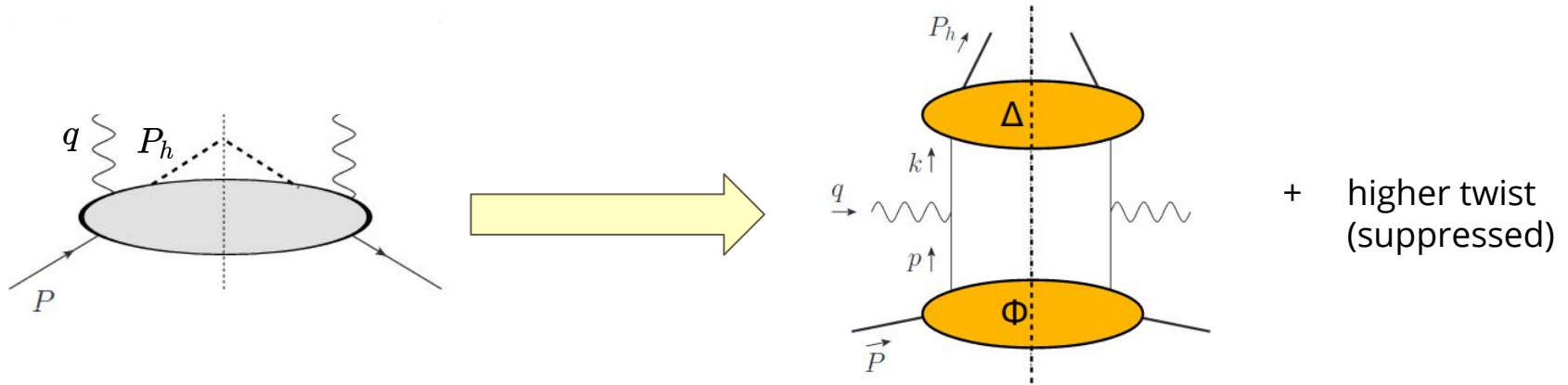
?

Again, up to now no partons ...

How do quarks and gluons emerge in this description?

# Partonic interpretation

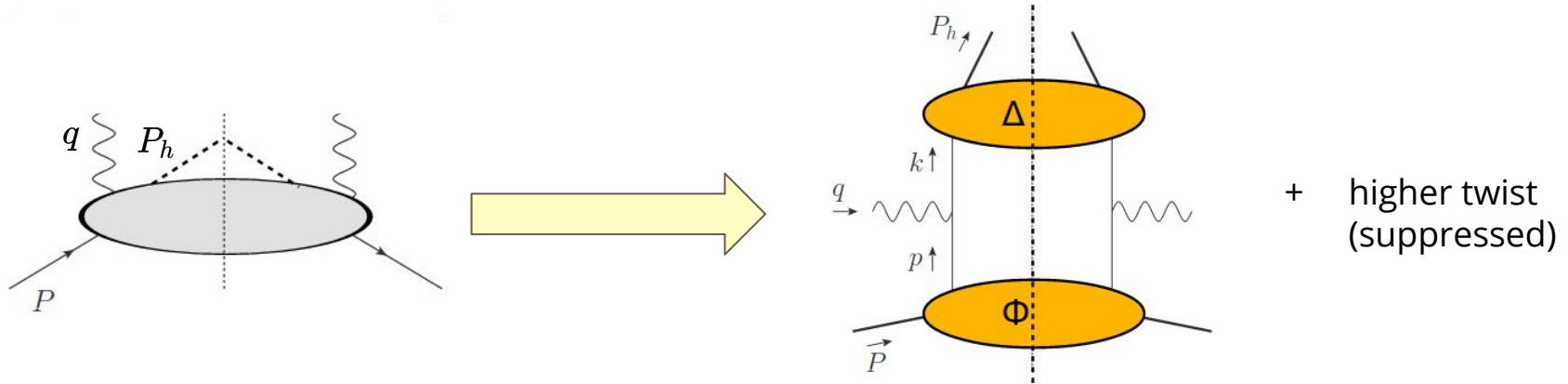
$$2 M W_{\mu\nu}(q, P, S, P_h) = \sum_X \int \frac{d^3 P_X}{2E_X} \delta^4(P + q - P_X) \langle PS | J_\mu^\dagger(0) | P_h P_X \rangle \langle P_h P_X | J_\nu(0) | PS \rangle$$



The presence of an identified hadron does not allow us to use the commutator form  
 → **OPE not applicable**

Use "**diagrammatic approach**" → use quark correlation functions for hadron structure and formation : it corresponds to the result in **TMD factorization** (when there is one)

# Partonic interpretation

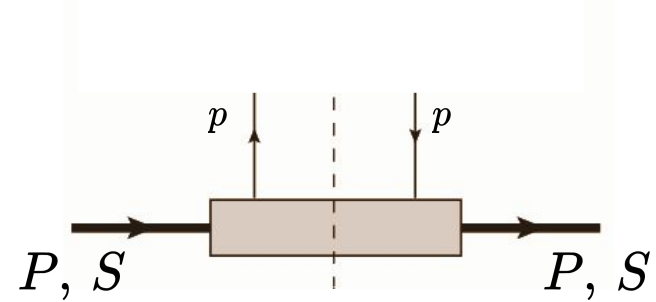


$$2MW^{\mu\nu}(q, P, S, P_h) = \frac{2z_h}{x_B} \mathcal{C} \left[ \text{Tr}(\Phi(x_B, \mathbf{p}_T, S) \gamma^\mu \Delta(z_h, \mathbf{k}_T) \gamma^\nu) \right]$$

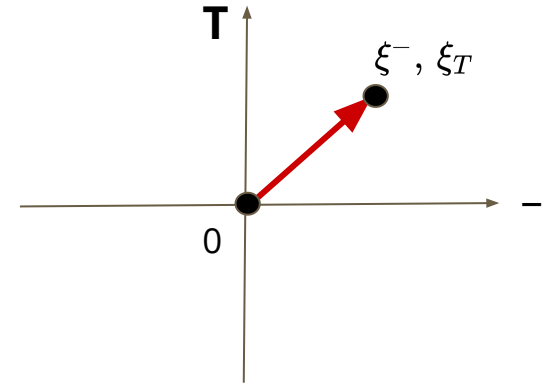
$$\mathcal{C}[wfD] = \sum x e_a^2 \int d^2 \mathbf{p}_T d^2 \mathbf{k}_T \delta^{(2)}(\mathbf{p}_T - \mathbf{k}_T - \mathbf{P}_{h\perp}/z) w(\mathbf{p}_T, \mathbf{k}_T) f^a(x, p_T^2) D^a(z, k_T^2)$$

# Quark (distribution) correlator

$$\Phi_{ij}(p, P, S) = \int \frac{d^4\xi}{(2\pi)^4} e^{i p \cdot \xi} \langle PS | \bar{\psi}_j(0) \psi_i(\xi) | PS \rangle$$

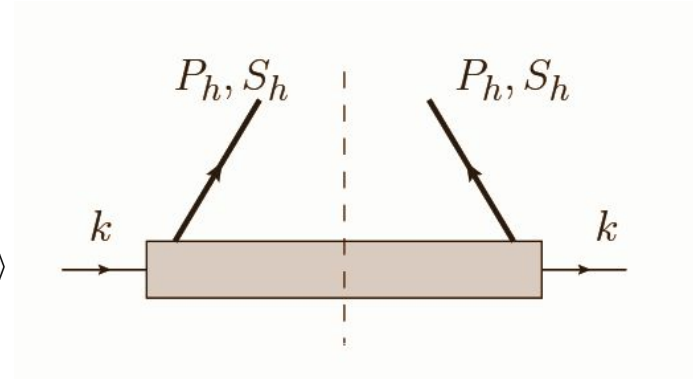


$$\begin{aligned} \Phi_{ij}(x, \mathbf{p}_T, S) &= \int dp^+ dp^- \delta(p^+ - xP^+) \Phi(p, P, S) = \\ &= \int \frac{d\xi^- d^2\xi_T}{2\pi} e^{i p \cdot \xi} \langle PS | \bar{\psi}_j(0) \psi_i(\xi) | PS \rangle_{\xi^+ = 0} \end{aligned}$$

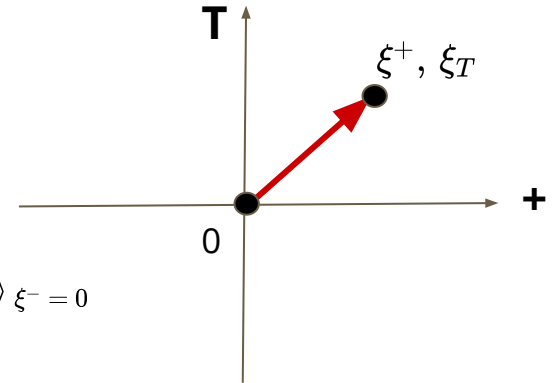


# Quark fragmentation correlator

$$\Delta_{ij}^h(k, P_h, S_h) = \sum_X \int \frac{d^4\xi}{(2\pi)^4} e^{i k \cdot \xi} \langle 0 | \psi_i(\xi) | X P_h S_h \rangle \langle X P_h S_h | \bar{\psi}_j(0) | 0 \rangle$$



$$\begin{aligned} \Delta_{ij}(z, \mathbf{k}_T, S_h) &= \int dk^- dk^+ \delta\left(k^- - \frac{1}{z} P_h^-\right) \Delta_{ij}(k, P_h, S_h) = \\ &= \sum_X \int \frac{d\xi^+ d^2\xi_T}{2\pi} e^{i k \cdot \xi} \langle 0 | \bar{\psi}_i(\xi) | X P_h S_h \rangle \langle X P_h S_h | \bar{\psi}_j(0) | 0 \rangle_{\xi^- = 0} \end{aligned}$$



# Quark TMD distribution functions (spin 1/2)

		quark pol.		
		U	L	T
nucleon pol.	U	$f_1$		$h_1^\perp$
	L		$g_{1L}$	$h_{1L}^\perp$
	T	$f_{1T}^\perp$	$g_{1T}$	$h_1, h_{1T}^\perp$

At leading twist: 8 TMD PDFs

(similar classification for gluons)

The **symmetries of QCD** play a crucial role in this classification

- **Black**: time-reversal even AND collinear
- **Blue**: time-reversal even
- **Red**: time-reversal odd (*process dependence*)

Quark inside spin 1/2 hadron



# Quark TMD distribution functions (spin 1)

Quarks	$\gamma^+$	$\gamma^+ \gamma^5$	$i\sigma^{i+} \gamma^5$
U	$f_1$		$h_1^\perp$
L		$g_1$	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}$	$h_1, h_{1T}^\perp$
LL	$f_{1LL}$		$h_{1LL}^\perp$
LT	$f_{1LT}$	$g_{1LT}$	$h_{1LT}, h_{1LT}^\perp$
TT	$f_{1TT}$	$g_{1TT}$	$h_{1TT}, h_{1TT}^\perp$

At leading twist: **18 (!)** TMD PDFs  
(similar classification for gluons)

The **symmetries of QCD** play a crucial role in this classification

Quark inside spin 1 hadron

# Quark TMD fragmentation functions

		quark pol.		
		U	L	T
hadron pol.	U	$D_1$		$H_1^\perp$
	L		$G_{1L}$	$H_{1L}^\perp$
	T	$D_{1T}^\perp$	$G_{1T}$	$H_1, H_{1T}^\perp$

At leading twist:  
8 TMD FFs and  
3 collinear FFs (diagonal)

The **symmetries of QCD** play  
a crucial role in this classification

# Structure functions and TMDs

$$F_{UU,T} = \mathcal{C}[f_1 D_1],$$

$$F_{UU}^{\cos 2\phi_h} = \mathcal{C}\left[-\frac{2(\hat{\mathbf{h}} \cdot \mathbf{k}_T)(\hat{\mathbf{h}} \cdot \mathbf{p}_T) - \mathbf{k}_T \cdot \mathbf{p}_T}{MM_h} h_1^\perp H_1^\perp\right]$$

$$F_{UT,T}^{\sin(\phi_h - \phi_s)} = \mathcal{C}\left[-\frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} f_{1T}^\perp D_1\right],$$

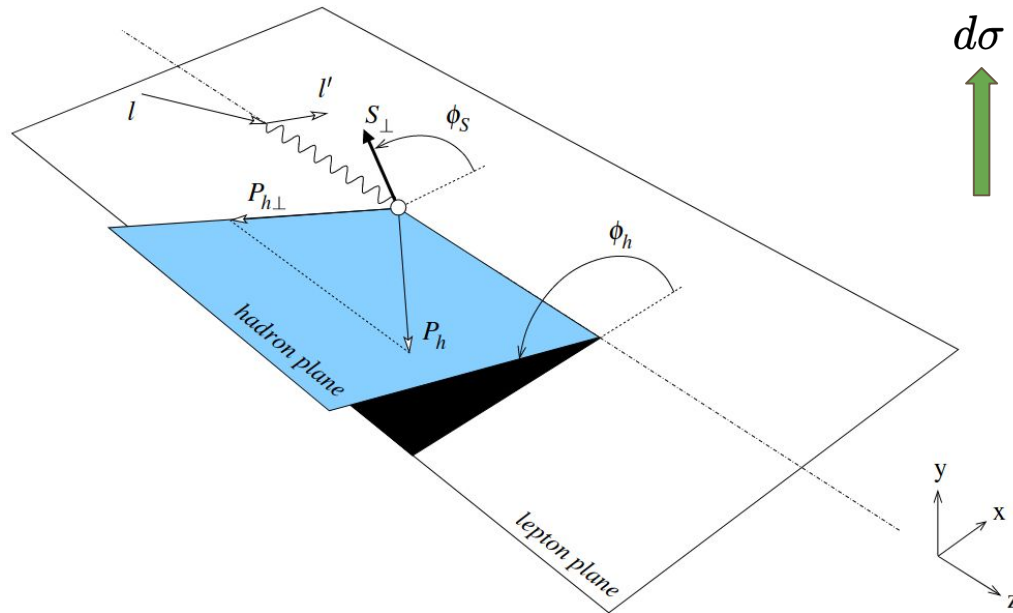
Etc. ...

The **LHS** of these equations can be **measured** and the **RHS** is expressed in terms of **partonic quantities (TMDs)**

→ **transverse momentum imaging**

$$\mathcal{C}[w f D] = \sum_x e_a^2 \int d^2 \mathbf{p}_T d^2 \mathbf{k}_T \delta^{(2)}(\mathbf{p}_T - \mathbf{k}_T - \mathbf{P}_{h\perp}/z) w(\mathbf{p}_T, \mathbf{k}_T) f^a(x, p_T^2) D^a(z, k_T^2)$$

# Transverse momentum imaging



Calculable  
in perturbation theory

$$d\sigma \sim \mathcal{H}$$

$$\text{TMD PDF} \otimes \text{TMD FF}$$

The partonic "maps", to be  
extracted from data

$$\ell N \rightarrow \ell' h X$$



JLab, EIC, ...