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# Transverse momentum imaging

## Lecture 1

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# Introduction

# The core of matter

galaxies, stars, ...

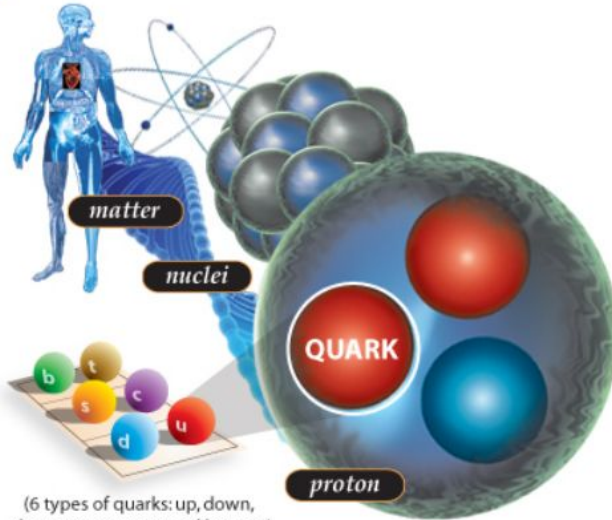
human body

molecules

DNA helix

atoms

nuclei



(6 types of quarks: up, down, charm, strange, top and bottom)

quark flavor

composite  
elementary

hadrons  
(protons and the like)

elementary particle physics

credit picture : Ohio State Univ.

Protons and neutrons: > 99 % of the visible mass in our world

# Quantum Chromodynamics (QCD)

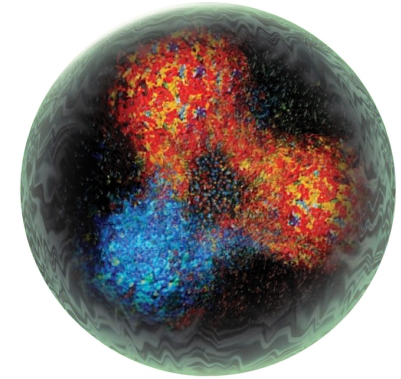
**quarks** and **gluons** (partons) are the elementary degrees of freedom in QCD, but they manifest only in bound states (**hadrons**)

$$\mathcal{L}_{QCD} = -\frac{1}{4} F^{a,\mu\nu} F_{a,\mu\nu} + \bar{\psi}(i\not{D} - m)\psi$$

$$D_{\mu} = \partial_{\mu} - igT^a A_{\mu}^a$$

gluon field

quark field



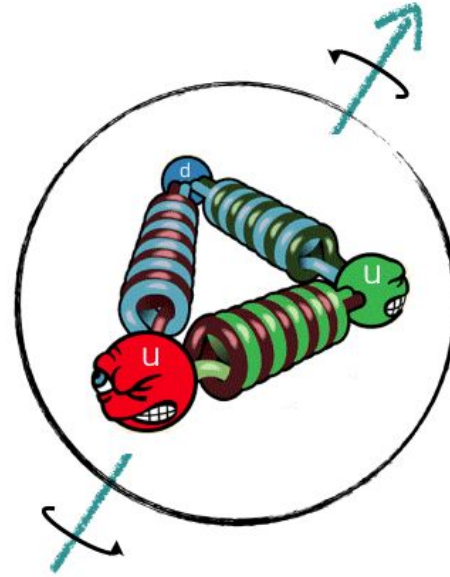
Can we understand the properties of hadrons in terms of quarks and gluons?

# Global properties

Can we understand the

**mass, spin, size  
of hadrons**

in terms of  
quarks and gluons?



# Confinement

Can we understand

**hadron formation  
and  
confinement**

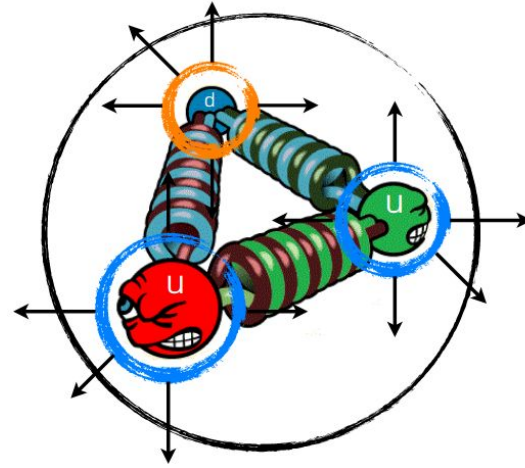
in terms of  
quarks and gluons?



# Internal structure

Can we understand the  
**structure of hadrons**

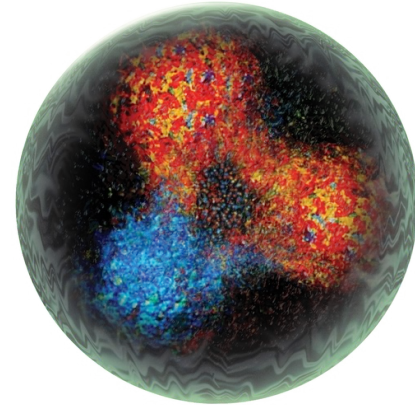
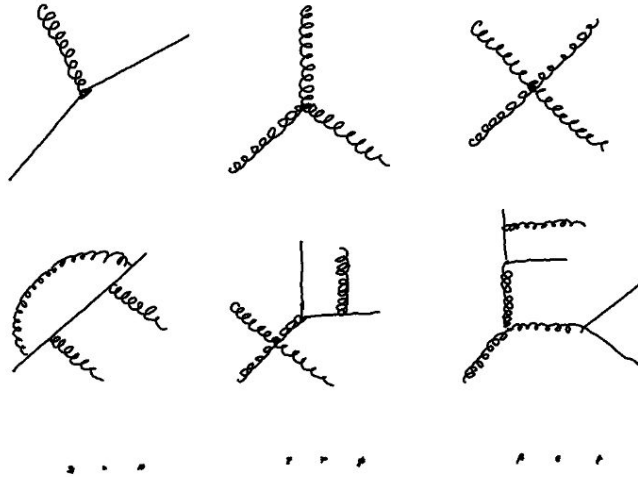
in terms of  
quarks and gluons?



# How should we “use” QCD ?

Expansion of observable in powers of the coupling constant  $\alpha$ :

$$\mathcal{O}(Q) \sim \mathcal{O}^{(0)} + \alpha_s^1(Q) \mathcal{O}^{(1)} + \alpha_s^2(Q) \mathcal{O}^{(2)} + \alpha_s^3(Q) \mathcal{O}^{(3)} \dots = ??$$



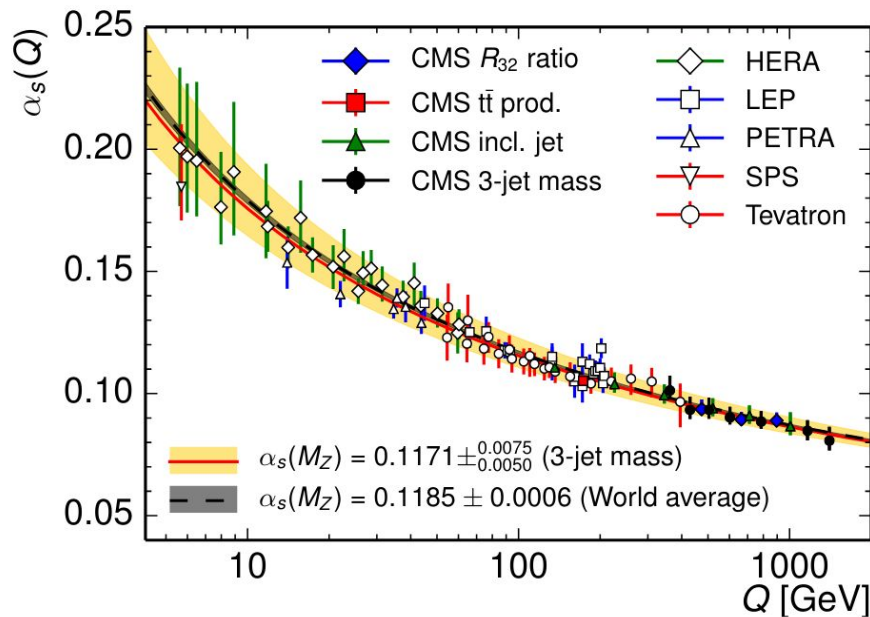
$$Q \sim M_N \sim 1 \text{ GeV}$$



# How should we “use” QCD ?

Expansion of observable in powers of the coupling constant  $\alpha$ :

$$\mathcal{O}(Q) \sim \mathcal{O}^{(0)} + \alpha_s^1(Q) \mathcal{O}^{(1)} + \alpha_s^2(Q) \mathcal{O}^{(2)} + \alpha_s^3(Q) \mathcal{O}^{(3)} \dots = ??$$



High energy  $\rightarrow$  convergence  
 $\rightarrow$  perturbative QCD

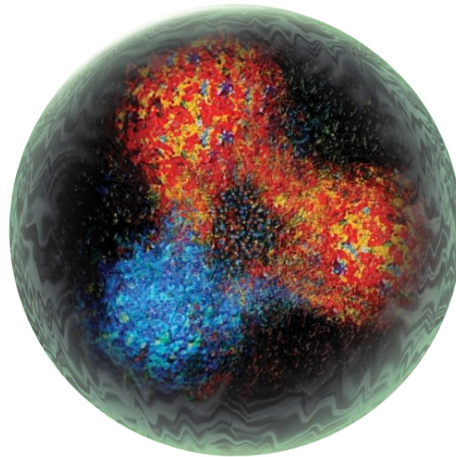
Low energy (hadronic scales)  
 $\rightarrow$  non-perturbative QCD

need alternative techniques

# Hadronic physics

Two macro areas to investigate:

1. **Hadron structure**: “hadrons  $\rightarrow$  partons” transition



# Hadronic physics

Two macro areas to investigate:

1. **Hadron structure**: “hadrons  $\rightarrow$  partons” transition
2. **Hadron formation**: “partons  $\rightarrow$  hadrons” transition  
(*hadronization*)



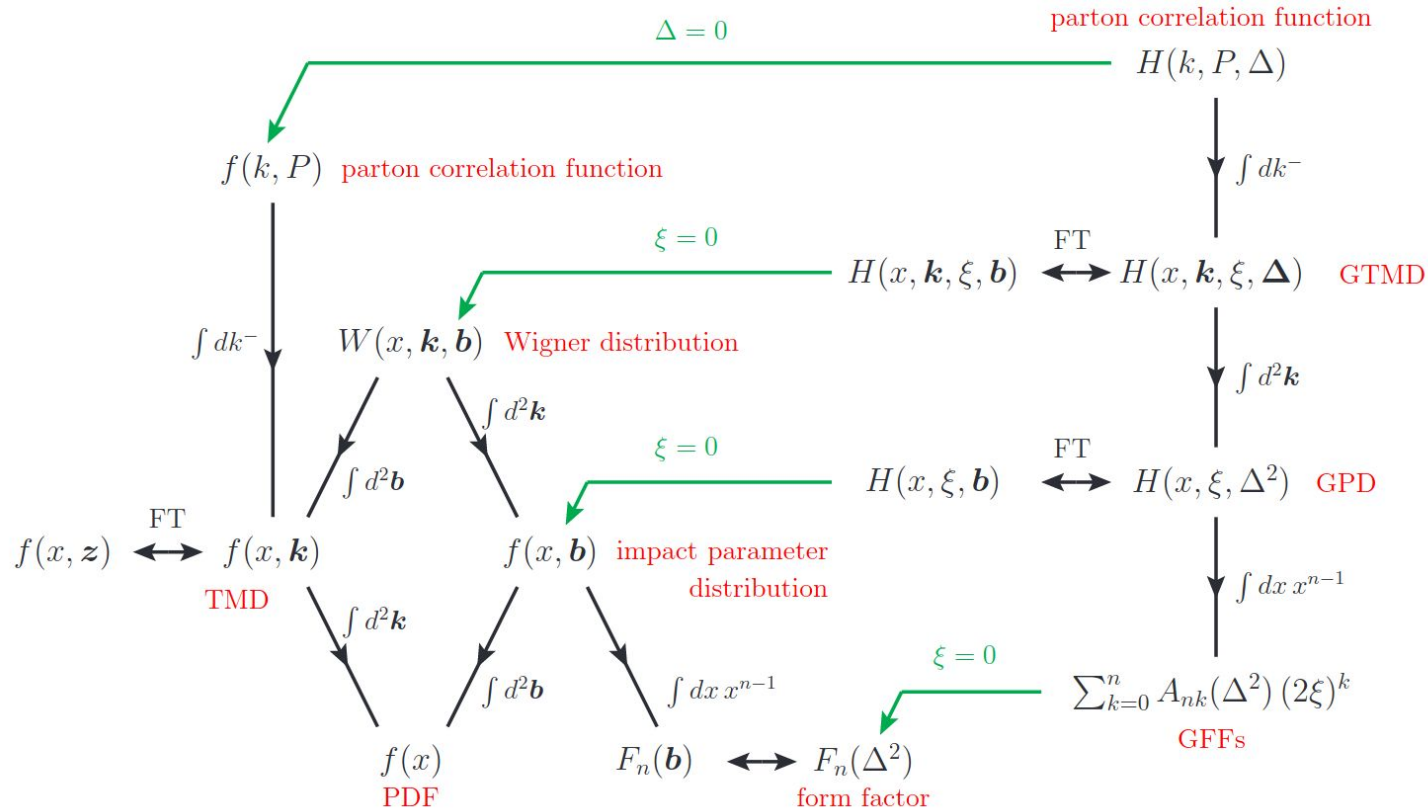
# Hadronic physics

The motivations are:

1. **Conceptual** : understand confinement, dynamical breaking of chiral symmetry ...  
*make sense of the world!*
2. **Practical** : improve our understanding of scattering experiments ...  
and *unlock >99% of the energy stored in matter!*



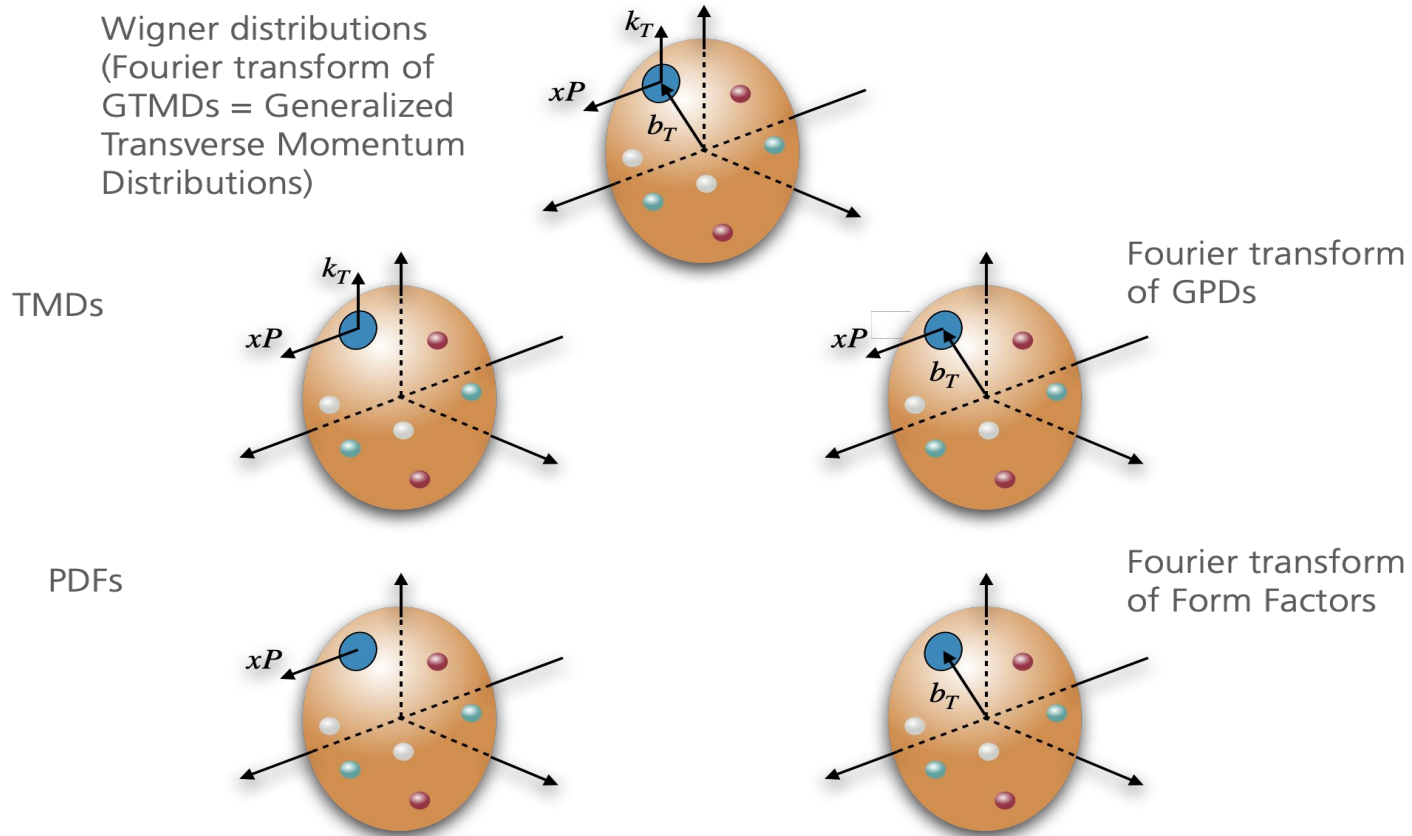
# The hadron structure landscape



Credit picture: M. Diehl - [arXiv 1512.01328]

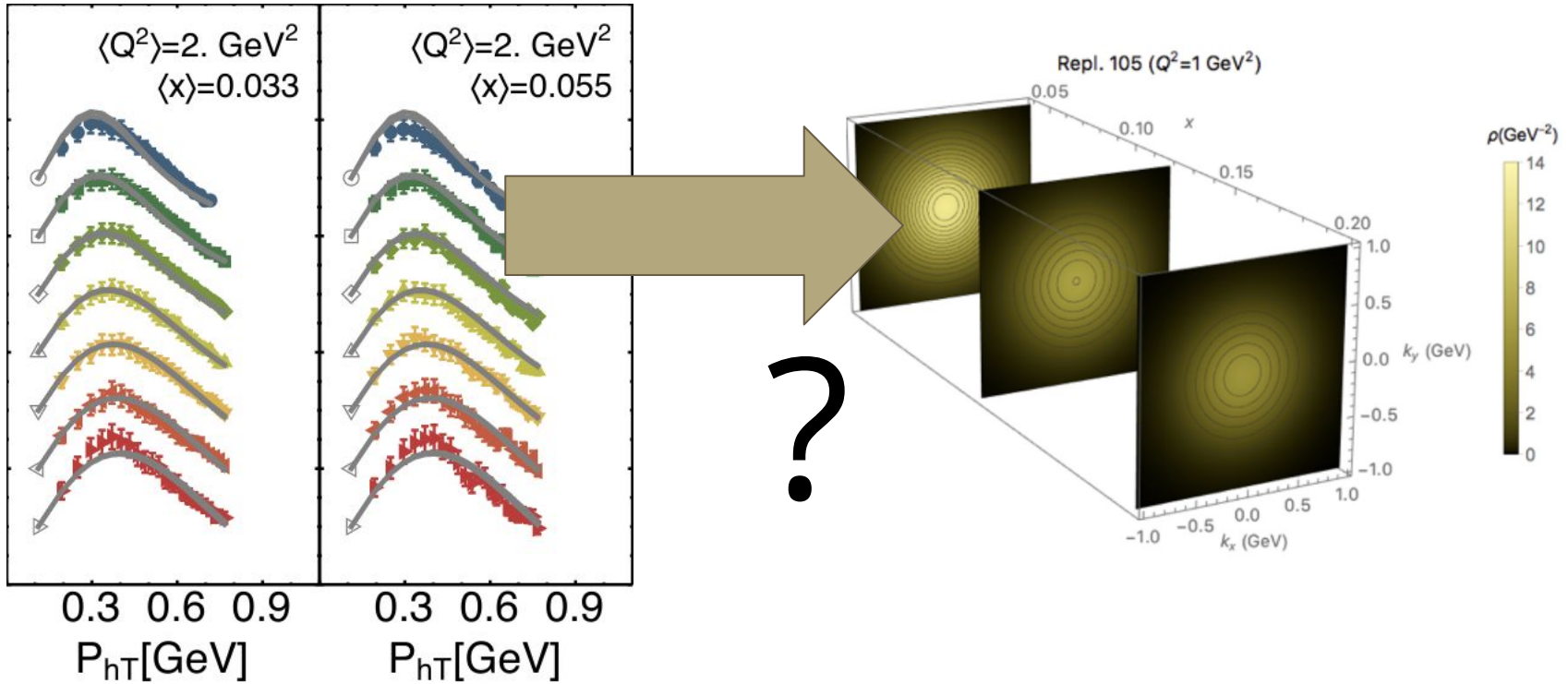
# The hadron structure landscape

Wigner distributions  
(Fourier transform of  
GTMDs = Generalized  
Transverse Momentum  
Distributions)



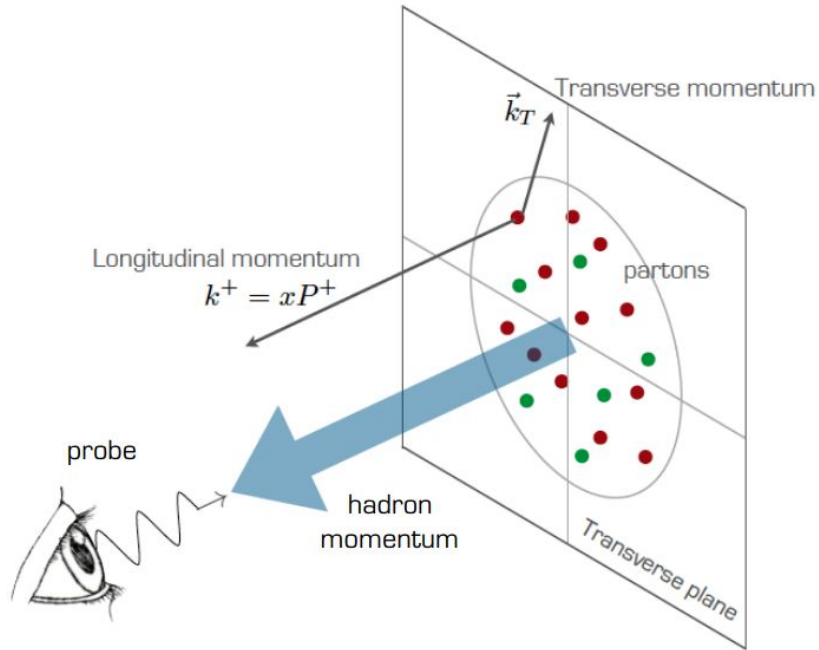
see, e.g., C. Lorcé, B. Pasquini, M. Vanderhaeghen, JHEP 1105 (11)

# Transverse momentum imaging



# Parton distribution functions (PDFs)

“Maps” of hadron *structure* in momentum space



$$f_1(x)$$

1D structure  
in momentum space

$$f_1(x, k_T^2)$$

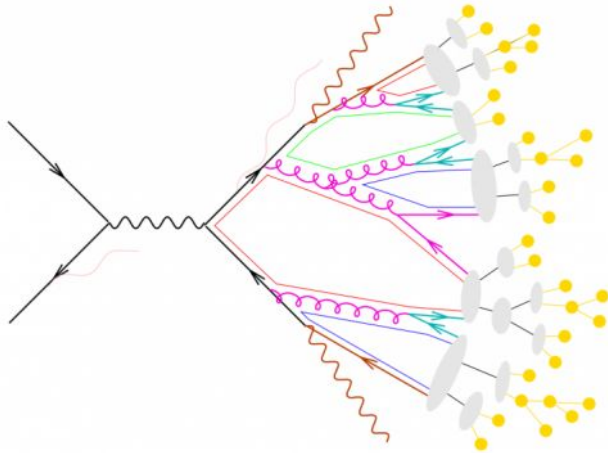
3D structure  
in momentum space

Credit picture: A. Bacchetta



# Fragmentation functions (FFs)

“Maps” of hadron *formation* in momentum space



$D_1^h(z)$       single-hadron collinear FF

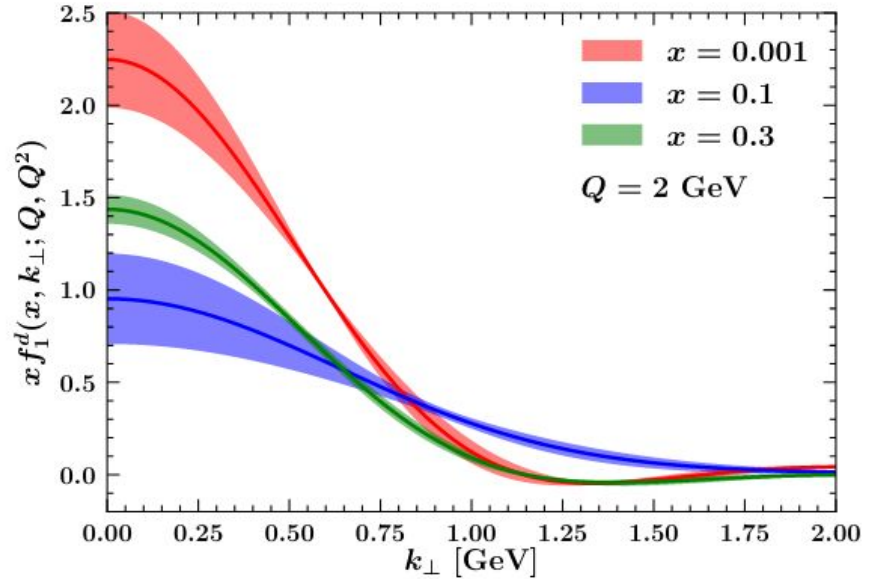
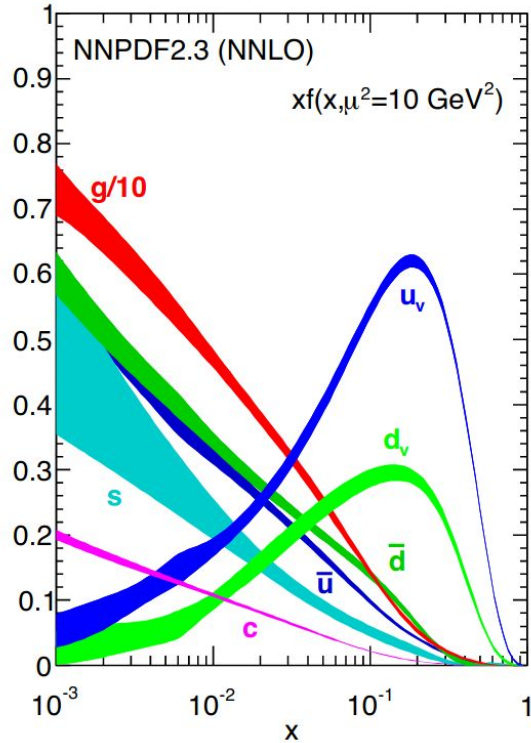
$D_1^h(z, P_T^2)$       single-hadron TMD FF

$D_1^{h_1 h_2}(z, \zeta)$       di-hadron FF

$J(s)$       inclusive jet FF

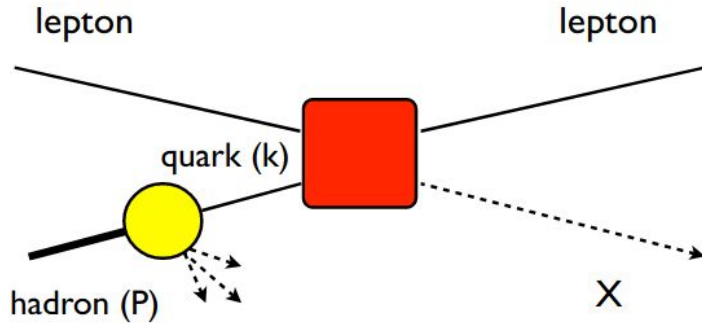
$\mathcal{G}^h(s, z)$       in-jet FF

# collinear & TMD PDFs



arXiv 1912.07550 (PV19 extraction)

# Operator definition (PDFs)

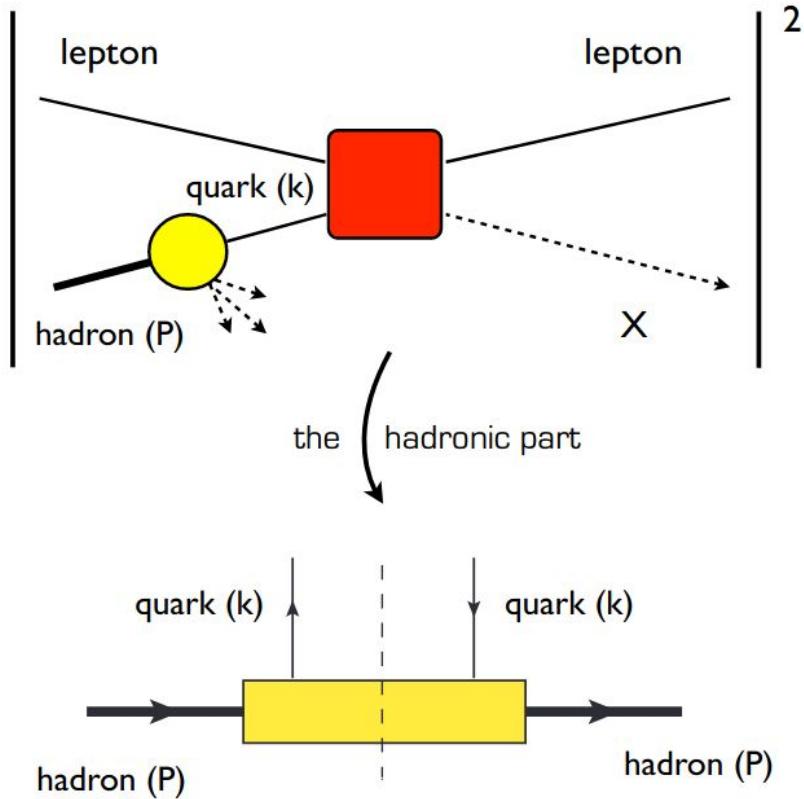


Scattering process with hadron in initial state :  
(e.g. Deep Inelastic Scattering - DIS)

need a "hadron  $\rightarrow$  parton" transition

(Parton Distribution Function)

# Operator definition (PDFs)



PDFs defined as traces of  $\Phi$  :

$$F^{[U]}(x, k_T^2) \sim \text{Tr} [\Phi \Gamma] , \quad \Gamma = \gamma^+ , \dots$$

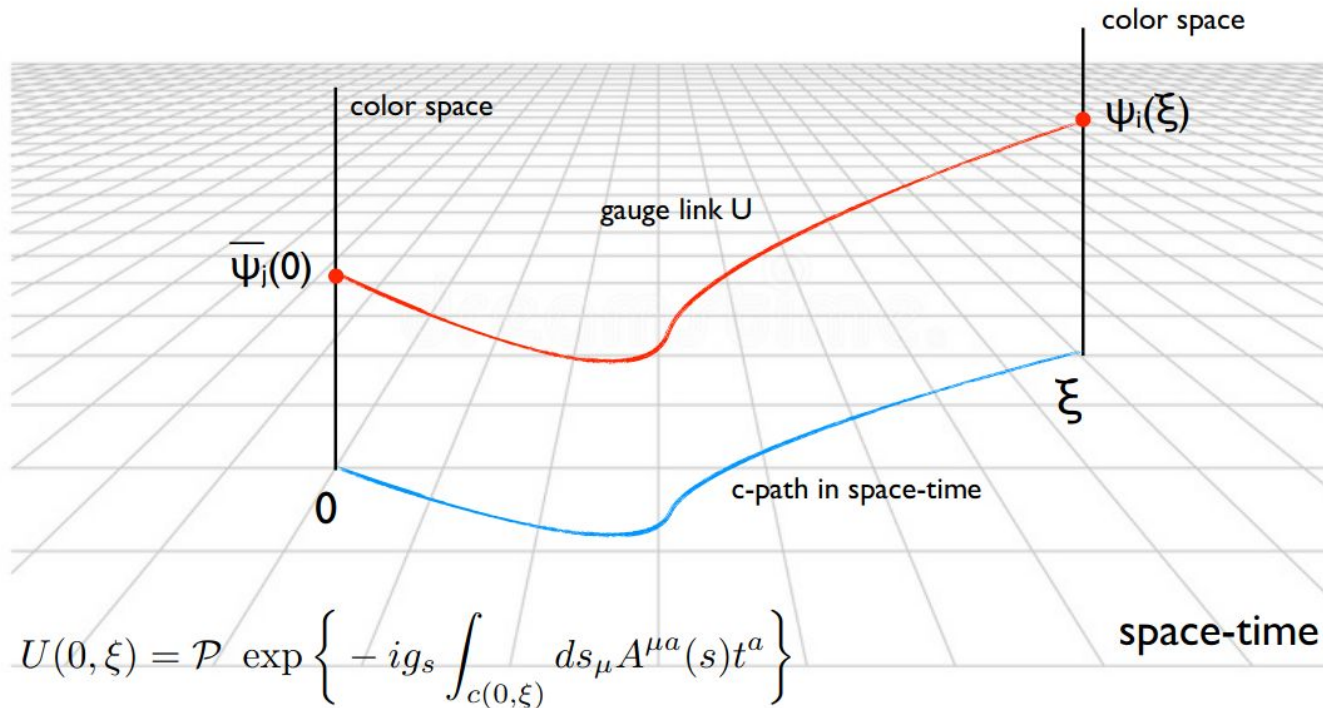
(**8 functions** that depend on parton kinematics and **gauge link U**)

Hadronic part described as a **universal** "quark-quark correlation function" in space-time

$$\Phi_{ij}(k, P) = \text{F.T.} \langle P | \bar{\psi}_j(0) U \psi_i(\xi) | P \rangle$$

# Geometric structure

$$\Phi(k, P) = \text{F.T.} \langle P | \bar{\psi}_j(0) U \psi_i(\xi) | P \rangle \longrightarrow f_1^a [U](x, k_T^2) \not{P} + \dots$$



A selection of  
useful references

# The HUGS pedagogical page

This is a list of references in preparation for and in support of the HUGS program. Further specific references will be suggested by the speakers. You are also welcome to browse the similarly aimed CTEQ pedagogical page, and to send us your comments and suggestions ([hugs@jlab.org](mailto:hugs@jlab.org)).

## General textbooks

- Donnelly, Formaggio, Holstein, Milner, Surov - *Foundations of Nuclear and Particle Physics* (2017)
  - Short, focused chapters covering practically all past, present, and near future HUGS topics!
- Povh, Rith, Scholz, Zetsche, Rodejohann - *Particles and Nuclei* (2015)
  - Good introductory level text
- Griffiths - *Introduction to Elementary Particles* (2008)
  - Another good introductory level text, more focused on the elementary particle aspects
- Halzen, Martin - *Quarks and leptons* (2008)
  - More advanced, treats QCD in some detail

<https://www.jlab.org/education/hugs/references>

## (Perturbative) QCD

- W. K. Tung, *Perturbative QCD and the parton structure of the nucleon*
- K. Kovarik, P. M. Nadolski, D. E. Soper, *Hadron structure in high-energy collisions*
- B. Poetter, *Calculational Techniques in Perturbative QCD: The Drell-Yan Process*
- Textbooks:
  - J. Collins - *Foundations of Perturbative QCD* (2011)
  - Kovchegov, Levin - *Quantum Chromodynamics at High Energy* (2012)
- Check also:
  - The "[Suggested QCD literature](#)" list of references by T. Rogers

## 3D Structure of Nucleons

- Introductory:
  - A. Bacchetta, *Transverse Momentum Distributions* (a.k.a. "Trento lectures", 2012)
  - P. Mulders, *Transverse-momentum distributions and beyond: setting up the nucleon tomography*, lectures at the Galileo Galilei Institute (2015)
  - M. Diehl, *Introduction to GPDs and TMDs*, Eur.Phys.J. A52 (2016) 149
- P. Mulders, *Transverse momentum dependence in structure functions in hard scattering processes*
- M. Diehl, *Lectures on GPDs*, Varenna (ITA), 2011
- M. Diehl, *Generalized Parton Distributions*, Phys.Rept. 388 (2003) 41



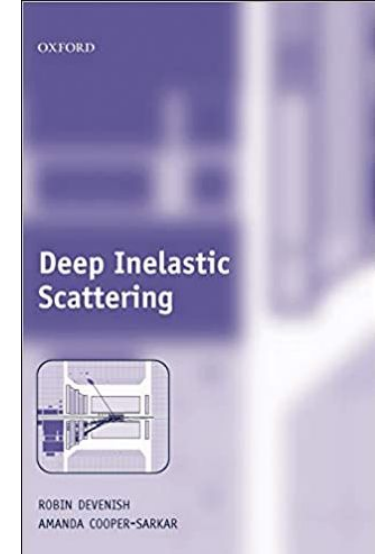
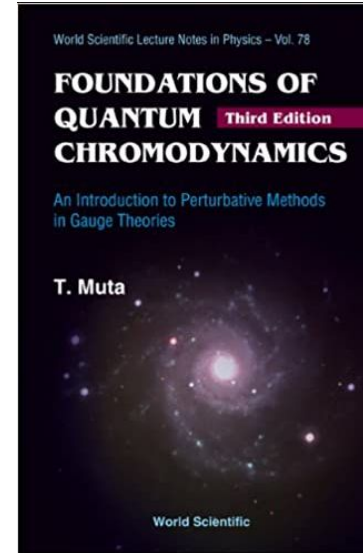
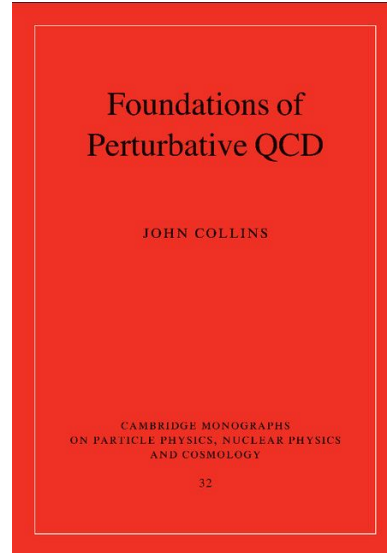
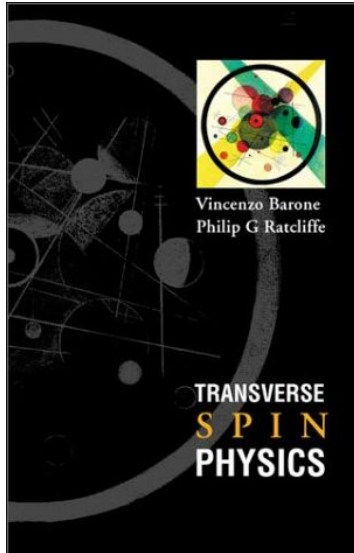
# Lecture notes

- Barone - Cabeo lecture notes: [https://www.fe.infn.it/cabeo\\_school/2010/cabeo\\_school\\_2010.pdf](https://www.fe.infn.it/cabeo_school/2010/cabeo_school_2010.pdf)
- Bacchetta - Trento lecture notes: [https://www2.pv.infn.it/~bacchett/teaching/Bacchetta\\_Trento2012.pdf](https://www2.pv.infn.it/~bacchett/teaching/Bacchetta_Trento2012.pdf)
- Jaffe - Erice lecture notes: <https://arxiv.org/pdf/hep-ph/9602236.pdf>
- Mulders - GGI lecture notes: <http://www.nat.vu.nl/~mulders/tmdreview-vs3.pdf>
- ...



# Books

- Barone, Ratcliffe: *Transverse Spin Physics*
- Collins: *Foundations of perturbative QCD*
- Devenish, Cooper-Sarkar: *Deep Inelastic Scattering*
- Muta: *Foundations of Quantum Chromodynamics*
- ...

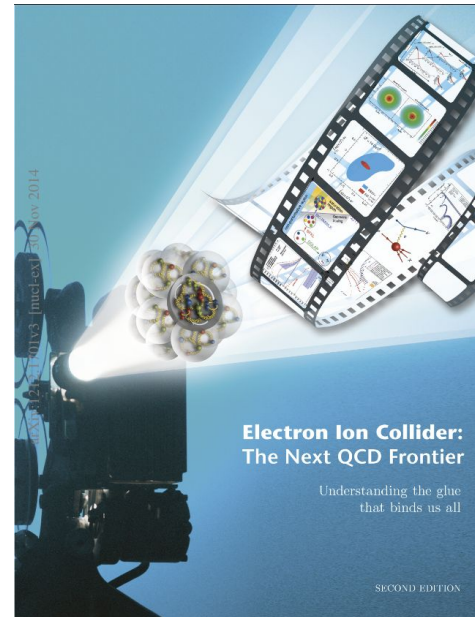


# Papers

- EPJ-A topical issue: *The 3D structure of the nucleon*  
[https://link.springer.com/journal/10050/topicalCollection/AC\\_628286e999d9a60c9a780398df15f93d](https://link.springer.com/journal/10050/topicalCollection/AC_628286e999d9a60c9a780398df15f93d)
- Diehl: *Introduction to GPDs and TMDs*  
<https://inspirehep.net/literature/1408303>
- Bacchetta et al.: *Single spin asymmetries: the Trento conventions*  
<https://inspirehep.net/literature/660999>
- Collins: *Light cone variables, rapidity and all that*  
<https://inspirehep.net/literature/443368>
- Metz-Vossen: *Parton fragmentation functions*  
<https://inspirehep.net/literature/1475000>
- Scimemi: *A short review on recent developments in TMD factorization and implementation*  
<https://inspirehep.net/literature/1716549>
- ...

# Experimental overviews

- Dudek et al.: *Physics opportunities with the 12 GeV upgrade at Jefferson Lab*  
<https://inspirehep.net/literature/1125972>
- Accardi et al.: *Electron Ion Collider: The next QCD Frontier - understanding the glue that binds us all*  
<https://inspirehep.net/literature/1206324>
- ...

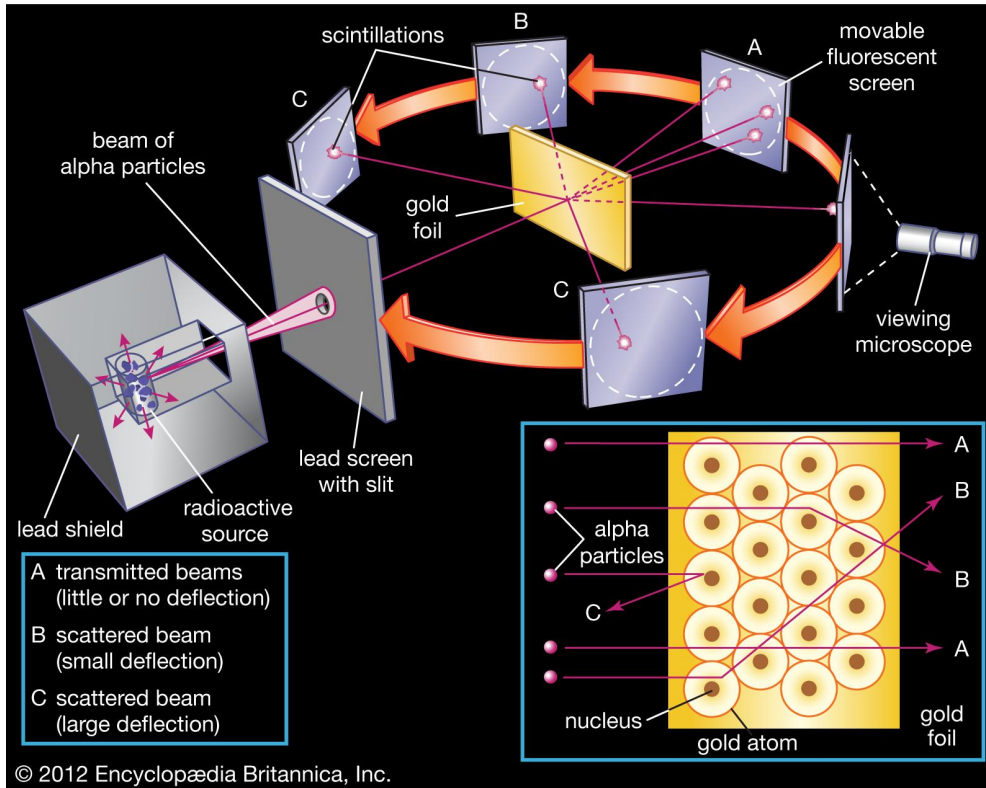


# Plan of these lectures

1. DIS and partons
2. From DIS to SIDIS
3. Symmetries and universality
4. Factorization, evolution, matching
5. Phenomenology

# 1. DIS & partons

# Geiger / Marsden / Rutherford experiment



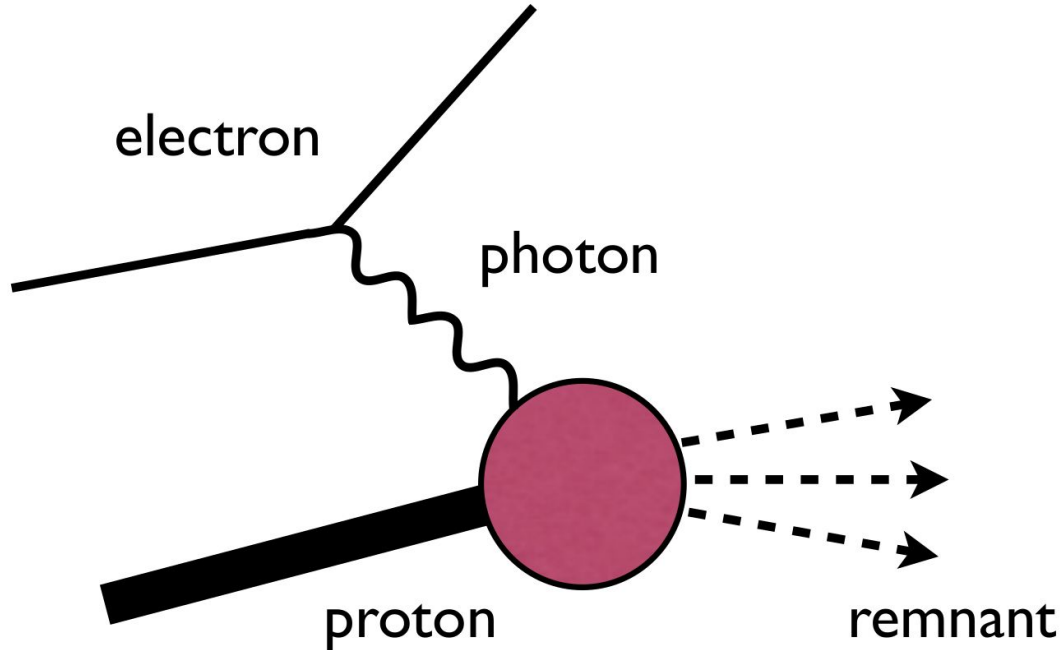
~1910

Scattering of alpha particles on gold:

discovery of the atomic nucleus

# Deep-inelastic scattering

$$l(\ell) + N(P) \rightarrow l'(\ell') + X(P_X)$$



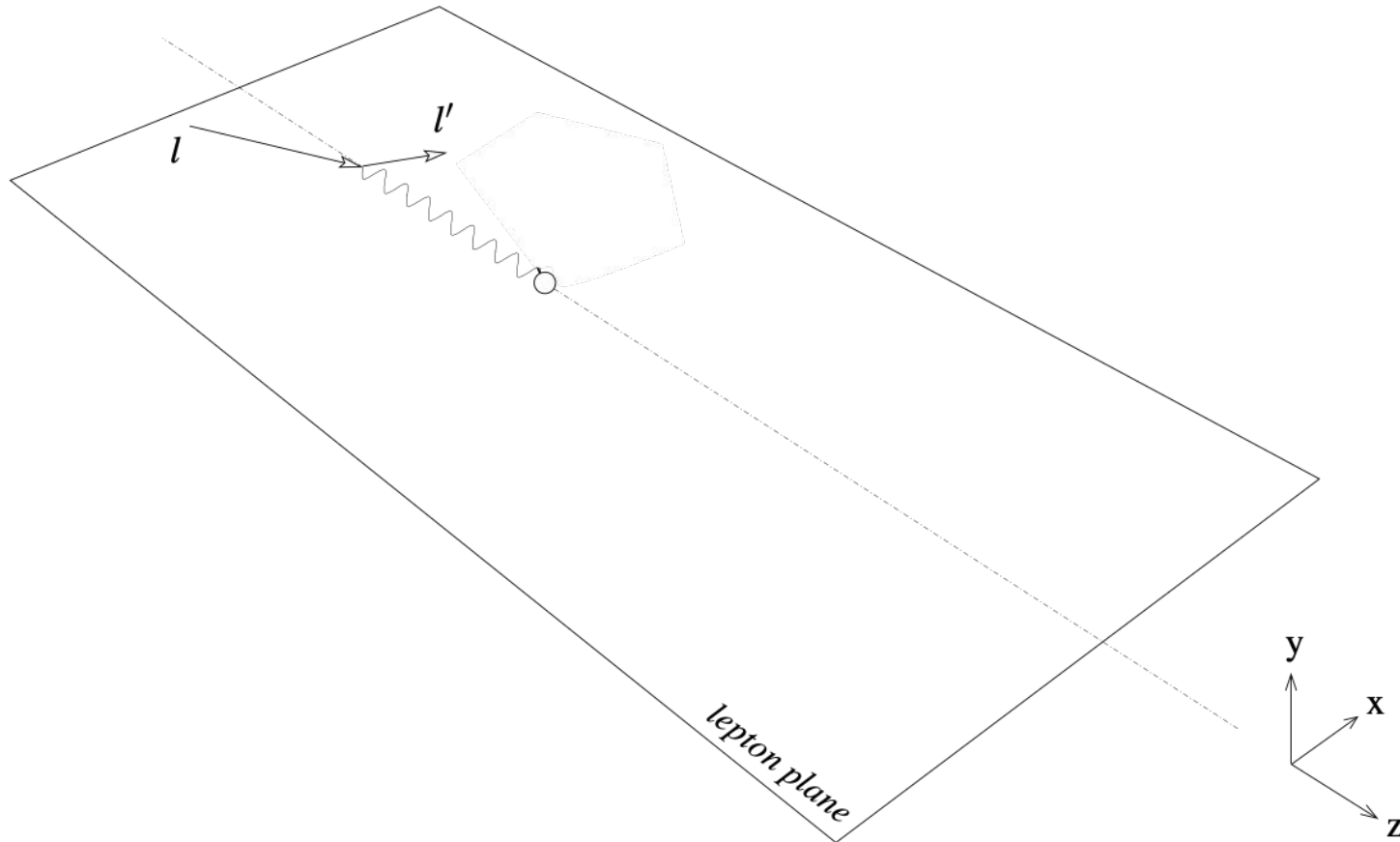
MIT-SLAC experiments ('60-'70)

Scattering of electrons off protons to test hadrons' substructure:

First evidence of free point-like spin- $\frac{1}{2}$  constituents (partons) inside the proton

# Deep-inelastic scattering

$$l(\ell) + N(P) \rightarrow l'(\ell') + X(P_X)$$





# Light cone variables

Choice of a basis :

$$\{n_+, n_-\}$$

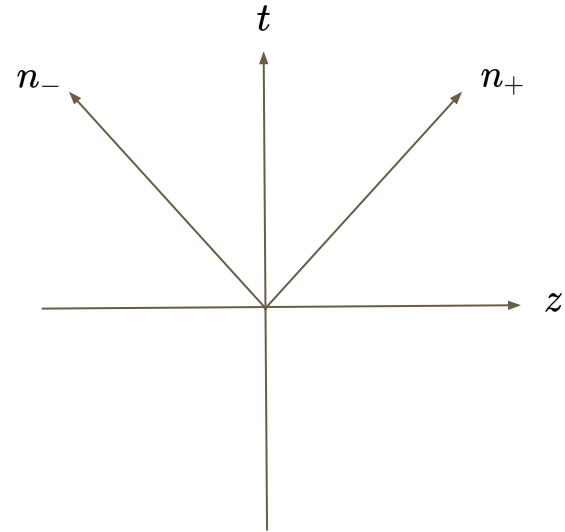
$$n_+^2 = n_-^2 = 0$$

$$n_+ \cdot n_- = 1$$

Projectors on the transverse space:

$$g_T^{\mu\nu} = g^{\mu\nu} - n_+^\mu n_-^\nu - n_-^\mu n_+^\nu$$

$$\epsilon_T^{\mu\nu} = \epsilon^{\mu\nu\rho\sigma} n_-^\rho n_+^\sigma$$



$$V^\mu = (V^0, V^1, V^2, V^3) = [V^+, V^-, \mathbf{V}_T]$$

$$V^\pm = \frac{1}{\sqrt{2}}(V^0 \pm V^3), \quad \mathbf{V}_T = (V^1, V^2)$$

$$V^2 = 2V^+V^- - |\mathbf{V}_T|^2$$

# Kinematics

With a nucleon target, we have four “external” vectors at our disposal: “spin”,  $P$ ,  $\ell$ ,  $\ell'$

We can build the following invariants

$$s = (P + \ell)^2$$

$$W^2 = (P + q)^2$$

$$Q^2 = -q^2 = -(\ell - \ell')^2$$

... and variables :

$$x_B = \frac{Q^2}{2P \cdot q}$$

$$y = \frac{P \cdot q}{P \cdot \ell}$$

Deep-inelastic regime (*Bjorken limit*):

$$Q^2, P \cdot q \rightarrow +\infty \quad (\gg M^2 \text{ in practice})$$

$x_B$  fixed

# Spin 1/2

The spin is described by means of a density operator (matrix, standard Quantum Mechanics)

Spin  $J \rightarrow$  (up to) rank  $2J$  tensor (e.g.: spin  $\frac{1}{2} \rightarrow$  rank 1 tensor = spin vector)

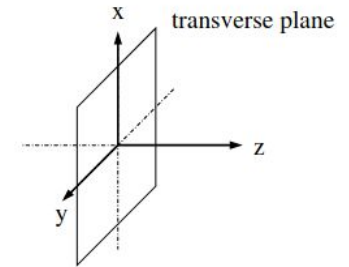
Spin 1/2

$$\rho = \frac{1}{2} (1 + S^i \sigma^i)$$

$\rightarrow$  identity operator 1 and Pauli matrices

$\rightarrow$  spin 3-vector  $S \quad S^i = (S_T^x, S_T^y, S_L)$

(z chosen as longitudinal direction)



Covariant spin vector

$$S^\mu = (0, S^i)$$

# Spin 1

see <https://inspirehep.net/literature/530045> for more details

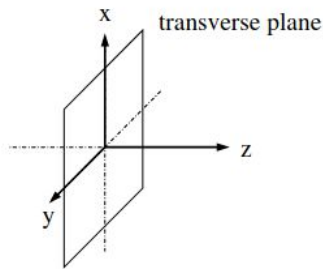
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Spin 1

$$\rho = \frac{1}{3} \left( 1 + \frac{3}{2} S^i \Sigma^i + 3 T^{ij} \Sigma^{ij} \right) \rightarrow \text{spin 3-vector "S" and spin tensor "T"}$$

$\rightarrow$  identity 1, 3D Pauli matrices  $\Sigma$  and their generalization to rank-2



$$\mathbf{S} = (S_T^x, S_T^y, S_L),$$

$$\mathbf{T} = \frac{1}{2} \begin{pmatrix} -\frac{2}{3}S_{LL} + S_{TT}^{xx} & S_{TT}^{xy} & S_{LT}^x \\ S_{TT}^{xy} & -\frac{2}{3}S_{LL} - S_{TT}^{xx} & S_{LT}^y \\ S_{LT}^x & S_{LT}^y & \frac{4}{3}S_{LL} \end{pmatrix}$$

# Spin - take home message

Spin 1/2

$$\rho = \frac{1}{2} (1 + S^i \sigma^i) \quad \rightarrow \text{spin 3-vector "S"}$$

Spin 1

$$\rho = \frac{1}{3} (1 + \frac{3}{2} S^i \Sigma^i + 3 T^{ij} \Sigma^{ij}) \quad \rightarrow \text{spin 3-vector "S" and spin tensor "T"}$$

A polarized deuteron (spin 1) has more "degrees of freedom" compared to a polarized nucleon (spin 1/2). This leads to a richer spin structure:

Additional structure functions and partonic distributions (see also E. Long's lectures)

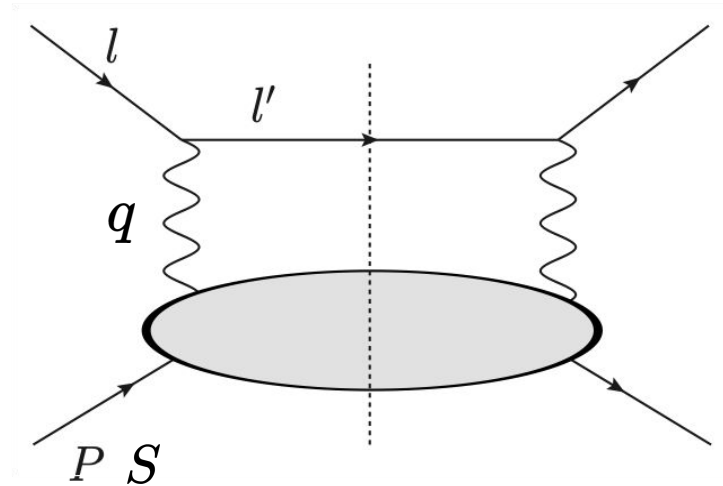
# Cross section

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S)$$

CUT DIAGRAM notation for one-photon exchange approximation.

Represents the product of two Feynman amplitudes ( → cross section), one at the left and one at the right of the “cut”

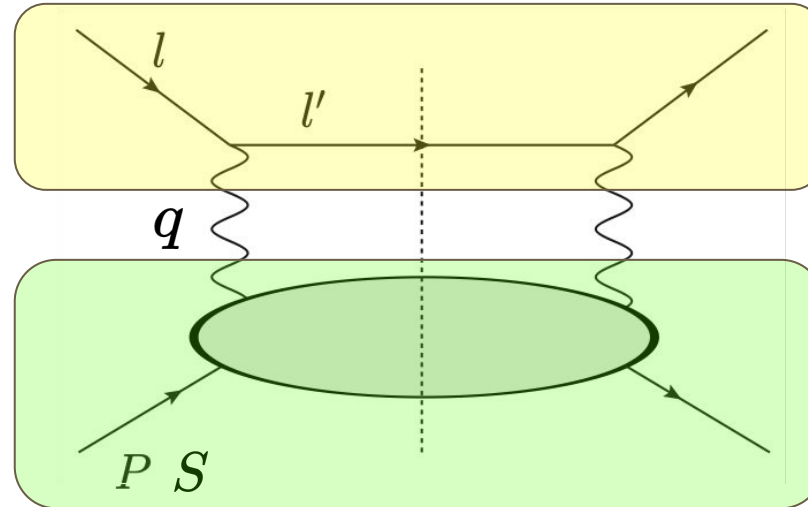
The dashed line represents the “cut”: particles that go to the final state (on-shell)



# Cross section

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S)$$

Leptonic tensor - QED  
(completely calculable)



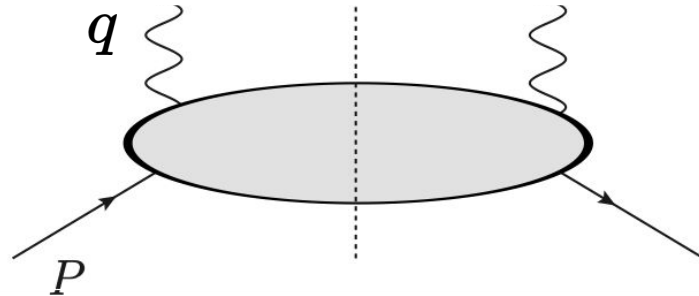
Hadronic tensor - QCD  
(NOT completely calculable)

# The hadronic tensor (unpolarized)

$$2 M W_{\mu\nu}(q, P) = \sum_X \int \frac{d^3 P_X}{2E_X} \delta^4(P + q - P_X) \langle P | J_\mu^\dagger(0) | P_X \rangle \langle P_X | J_\nu(0) | P \rangle$$

The scattering electron “feels” the electromagnetic current  $J$  in the target

$$J_\mu(\xi) = : \bar{\psi}(\xi) Q \gamma_\mu \psi(\xi) : \quad (\text{in case of weak interaction (W,Z) the current is different})$$



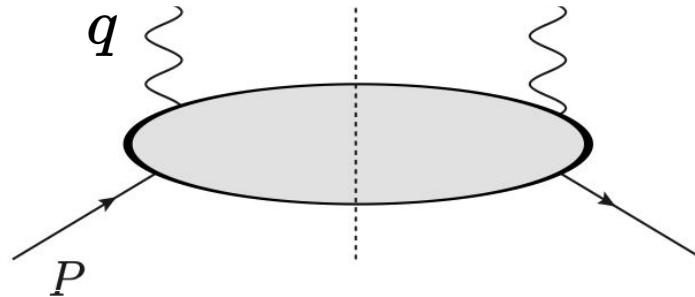


# The hadronic tensor (unpolarized)

We can parametrize  $W$  using the available vectors and the symmetries of the theory:

$$2 M W_{\mu\nu}(q, P) = \left[ A g_{\mu\nu} + B \frac{q_\mu q_\nu}{q^2} + C \frac{P_\mu P_\nu}{M^2} + D \frac{P_\mu q_\nu + P_\nu q_\mu}{M^2} \right]$$

Conditions: **parity** invariance, **time-reversal** invariance, **gauge** invariance and *hermiticity*



Weak interaction: no parity  $\rightarrow$  additional terms!  
Spin  $\rightarrow$  additional terms!

# Structure functions

$$\begin{aligned}
 M W^{\mu\nu}(q, P) &= \left( \frac{q^\mu q^\nu}{q^2} - g^{\mu\nu} \right) F_1(x_B, Q^2) + \frac{1}{P \cdot q} \left( P^\mu - \frac{P \cdot q}{q^2} q^\mu \right) \left( P^\nu - \frac{P \cdot q}{q^2} q^\nu \right) F_2(x_B, Q^2) = \\
 &= -g_{\perp}^{\mu\nu} F_{UU,T}(x_B, Q^2) + \hat{t}^\mu \hat{t}^\nu F_{UU,L}(x_B, Q^2)
 \end{aligned}$$

F1, F2: "standard" unpolarized DIS structure functions

FT, FL: structure functions in the {t, z} basis (direct connection with the photon polarization)

Orthogonal and normalized basis

$$z^\mu = -\hat{q}^\mu \equiv q^\mu / Q$$

$$\hat{t}^\mu = \frac{2x_B}{Q\sqrt{1+\gamma^2}} \left( P^\mu - \frac{P \cdot q}{q^2} q^\mu \right)$$

Transverse projectors

$$g_{\perp}^{\mu\nu} = g^{\mu\nu} + \hat{q}^\mu \hat{q}^\nu - \hat{t}^\mu \hat{t}^\nu$$

$$\epsilon_{\perp}^{\mu\nu} = \epsilon^{\mu\nu\rho\sigma} \hat{t}_\rho \hat{q}_\sigma$$

# Polarized case - spin 1/2

$$W^{\mu\nu}(q, P, S) \sim -g_{\perp}^{\mu\nu} F_{UU,T} + \hat{t}^{\mu} \hat{t}^{\nu} F_{UU,L} \\ + i S_L \epsilon_{\perp}^{\mu\nu} F_{LL} + i \left( \hat{t}^{\mu} \epsilon_{\perp}^{\nu\rho} - \hat{t}^{\nu} \epsilon_{\perp}^{\mu\rho} \right) S_{\rho} F_{LT}^{\cos \phi}$$

Two additional structure functions for the nucleon:

**longitudinal** and **transverse** target polarization → related to “standard” g1 and g2 functions

Transverse beam polarization is proportional to electron mass and thus suppressed

# Polarized case - spin 1

$$\begin{aligned} W^{\mu\nu}(q, P, S) \sim & -g_{\perp}^{\mu\nu} F_{UU,T} + \hat{t}^{\mu} \hat{t}^{\nu} F_{UU,L} \\ & + i S_L \epsilon_{\perp}^{\mu\nu} F_{LL} + i \left( \hat{t}^{\mu} \epsilon_{\perp}^{\nu\rho} - \hat{t}^{\nu} \epsilon_{\perp}^{\mu\rho} \right) S_{\rho} F_{LT}^{\cos\phi} \\ & + \{b_1, b_2, b_3, b_4\} \longrightarrow T^{\mu\nu} \text{ tensor polarized terms} \end{aligned}$$

Two additional structure functions for the nucleon:

**longitudinal** and **transverse** target polarization  $\rightarrow$  related to “standard” g1 and g2 functions

Transverse beam polarization is proportional to electron mass and thus suppressed

For a deuteron there are four additional structures associated with the **tensor polarization!**

# Cross section (polarized nucleon)

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S)$$

$$\frac{d\sigma}{dx_B dy d\phi_S} = \frac{2\alpha^2}{x_B y Q^2} \left\{ \left(1 - y + \frac{y^2}{2}\right) F_{UU,T} + (1 - y) F_{UU,L} + S_L \lambda_e y \left(1 - \frac{y}{2}\right) F_{LL} \right. \\ \left. + |S_T| \lambda_e y \sqrt{1 - y} \cos \phi_S F_{LT}^{\cos \phi_s} \right\} \quad \text{F...: functions of } x, Q$$

Up to now no partons ...

How do quarks and gluons emerge in this description?

For a summary see e.g. <https://inspirehep.net/literature/732275>

# Light cone dominance

$$2 M W_{\mu\nu}(q, P, S) = \sum_X \int \frac{d^3 P_X}{2E_X} \delta^4(P + q - P_X) \langle PS | J_\mu^\dagger(0) | P_X \rangle \langle P_X | J_\nu(0) | PS \rangle$$

By using :

- properties of delta
- completeness on the intermediate states and
- translating the argument of the current

$$2MW_{\mu\nu}(q, P, S) = \frac{1}{2\pi} \int d^4\xi e^{iq \cdot \xi} \langle PS | [J_\mu^\dagger(\xi), J_\nu(0)] | PS \rangle$$

- Causality implies:  $\xi^2 > 0$
- Riemann-Lebesgue lemma implies W is zero unless:  $\xi^+ \rightarrow 0$

Thus , the W is dominated by what happens at  $\xi^2 \simeq 0$



We can use the OPE!

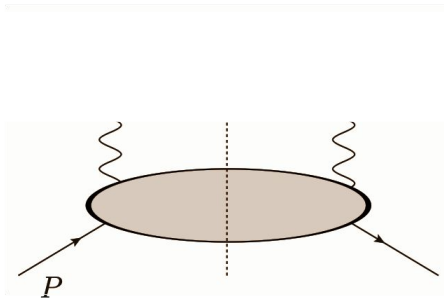
# Operator Product Expansion (OPE)

See Muta's book for details

$$A(x) B(y) = \sum_n C_n(x-y) O_n(x), \quad \text{when } |x-y| \text{ is small}$$

↑ operators      ↑ Coefficients (lower n : more dominant)

Thanks to the light cone dominance, we can apply this to the hadronic tensor:

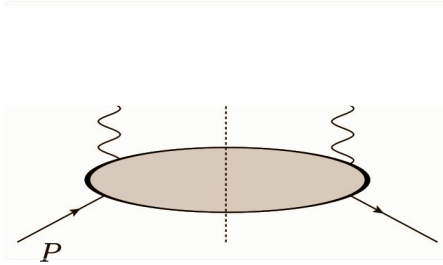


$$2MW_{\mu\nu}(q, P, S) = \frac{1}{2\pi} \int d^4\xi e^{iq \cdot \xi} \langle PS | [J_\mu^\dagger(\xi), J_\nu(0)] | PS \rangle$$

$$J_\mu(\xi) = : \bar{\psi}(\xi) Q \gamma_\mu \psi(\xi) :$$

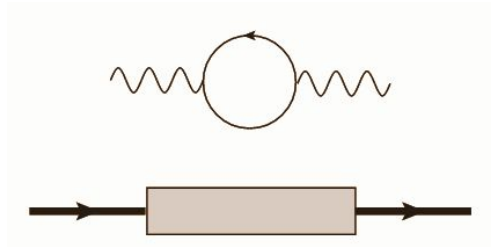
# Operator Product Expansion (OPE)

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$$2MW_{\mu\nu}(q, P, S) = \frac{1}{2\pi} \int d^4\xi e^{iq \cdot \xi} \langle PS | [J_\mu^\dagger(x), J_\nu(0)] | PS \rangle$$

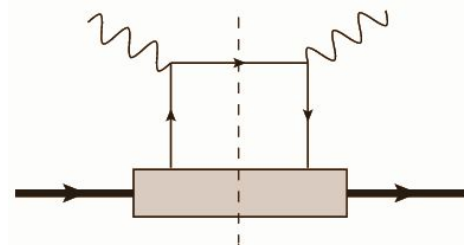
$$J_\mu(\xi) = : \bar{\psi}(\xi) Q \gamma_\mu \psi(\xi) :$$



Disconnected

irrelevant

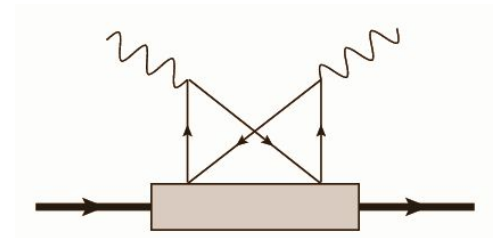
+



quark-antiquark

dominant

+

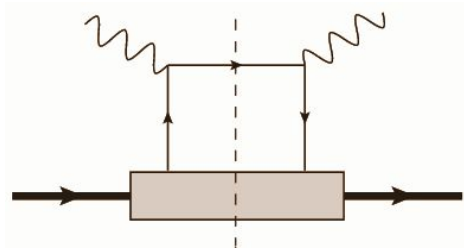
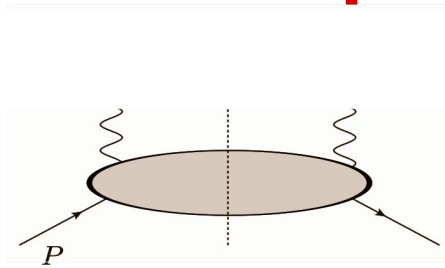


higher-twist

suppressed



# Partonic interpretation



quark-antiquark

$$2MW_{\mu\nu}(q, P, S) = \frac{1}{2\pi} \int d^4\xi e^{iq \cdot \xi} \langle PS | [J_\mu^\dagger(x), J_\nu(0)] | PS \rangle$$

$$J_\mu(\xi) = : \bar{\psi}(\xi) Q \gamma_\mu \psi(\xi) :$$

$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(x, S) \gamma^\mu \gamma^+ \gamma^\nu]$$

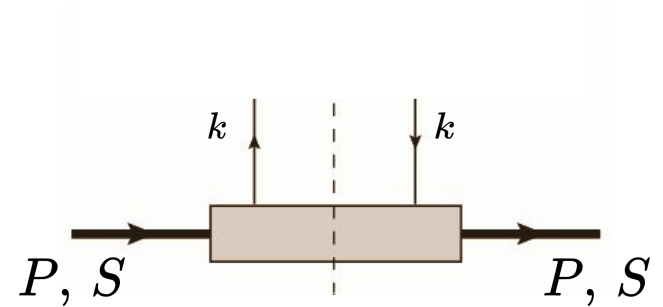
$\phi(x, S)$ : "collinear" quark correlator

$x_B \simeq x \equiv k^+ / P^+ \rightarrow$  measure collinear parton dynamics

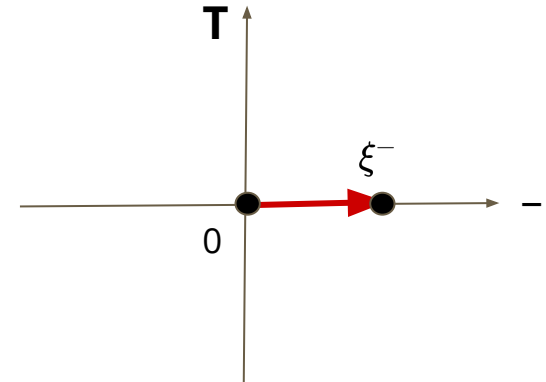
The quark transverse momentum is integrated out in DIS

# Quark correlator

$$\Phi_{ij}(k, P, S) = \int \frac{d^4\xi}{(2\pi)^4} e^{i k \cdot \xi} \langle PS | \bar{\psi}_j(0) \psi_i(\xi) | PS \rangle$$



$$\begin{aligned} \Phi_{ij}(x, S) &= \int dk^+ dk^- d^2\mathbf{k}_T \delta(k^+ - xP^+) \Phi(k, P, S) = \\ &= \int \frac{d\xi^-}{2\pi} e^{i k \cdot \xi} \langle PS | \bar{\psi}_j(0) \psi_i(\xi) | PS \rangle_{\xi^+ = \xi_T = 0} \end{aligned}$$

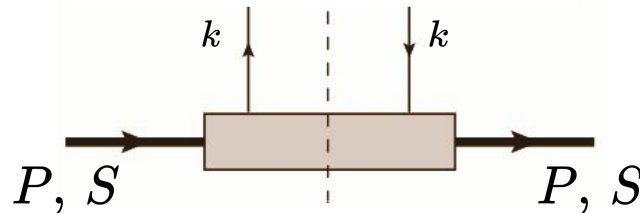


Parton physics on the light cone



# Collinear parton distribution functions

$\Phi_{ij}(k, P, S)$  : non-perturbative hadron structure matrix



$$\Phi(x, S, T) = \frac{1}{2} \boxed{f_1(x)} \not{h}_+$$

→ unpolarized PDF

$$\frac{1}{2} \boxed{g_1(x)} S_L \gamma_5 \not{h}_+$$

→ longitudinally polarized PDF  
(helicity)

$$\frac{1}{2} \boxed{h_1(x)} i\sigma_{\mu\nu} \gamma_5 n_+^\mu S_T^\nu +$$

→ transversely polarized PDF  
(transversity)

$$\frac{1}{2} \boxed{f_{1LL}(x)} S_{LL} \not{h}_+$$

→ Tensor polarized PDF

“Leading twist”  
approximation

Ok, but ...

what about transverse momentum?