HUGS: Introduction to QCD

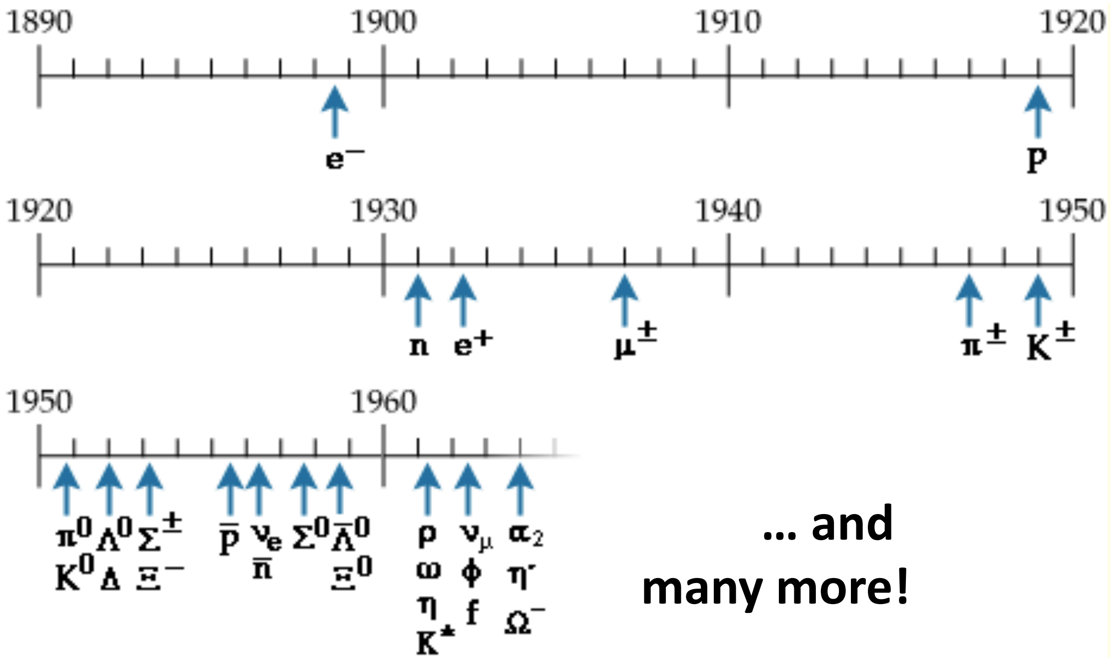
- Lec. 1: Fundamentals of QCD
- Lec. 2: Matching observed hadrons to quarks and gluons
- Lec. 3: QCD for cross sections with identified hadrons
- Lec. 4: QCD for cross sections with polarized beam(s)

HUGS 2020 was canceled due to COVID-19
HUGS 2021 will be the first virtual school

Jianwei Qiu
Theory Center
Jefferson Lab

New particles, new ideas, and new theories

- ❑ **Early proliferation of new hadrons – "particle explosion":**



… and many more!

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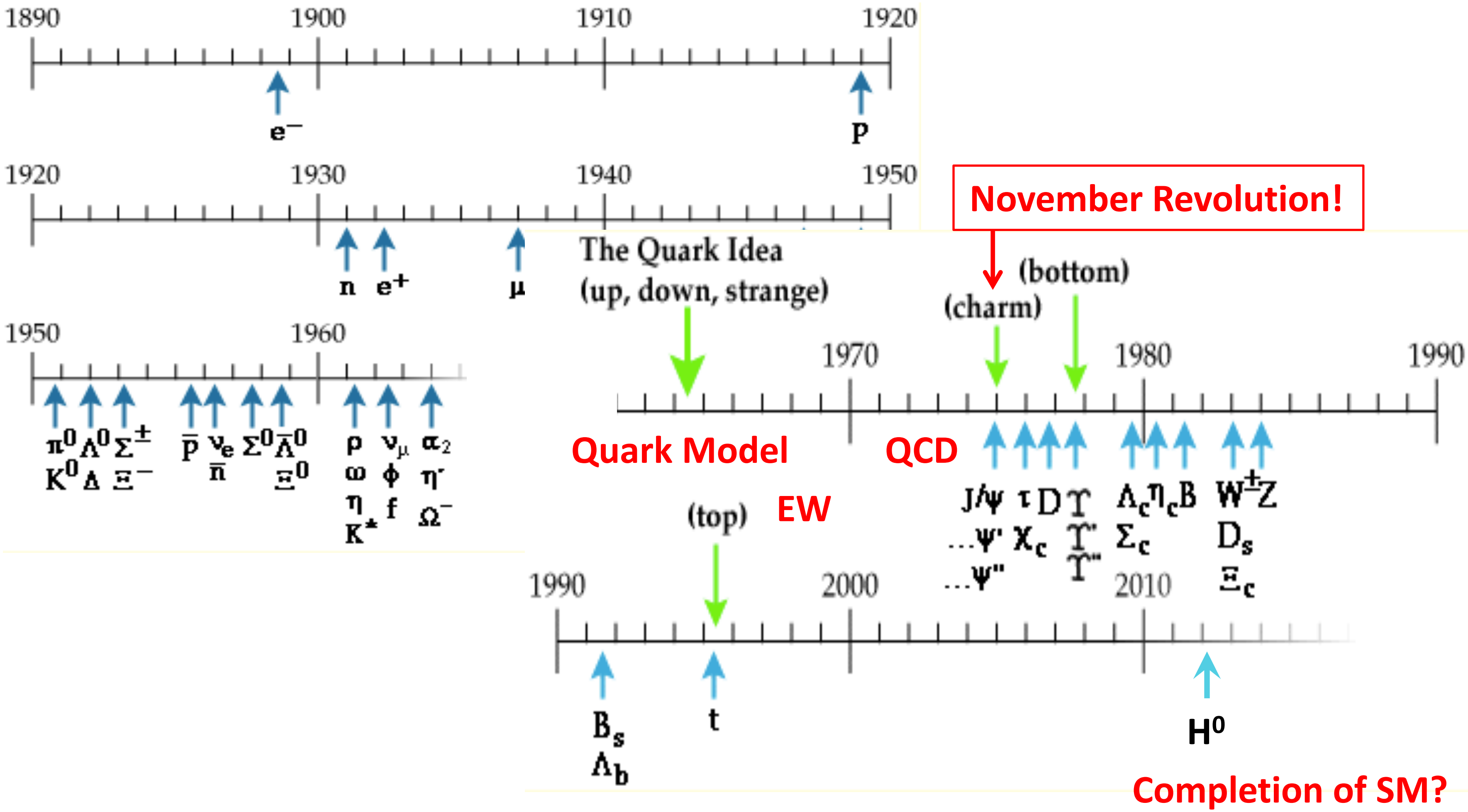

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New particles, new ideas, and new theories

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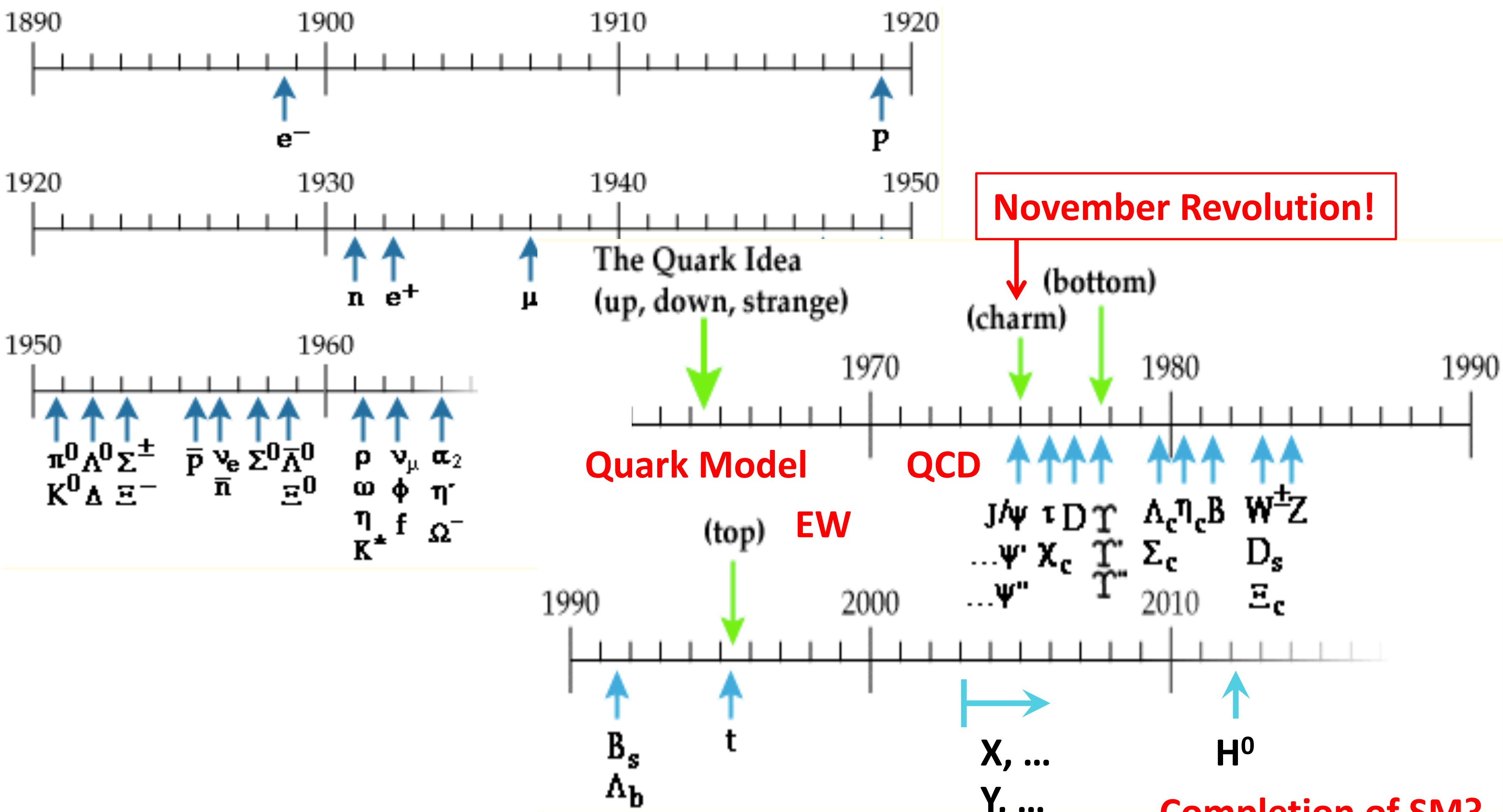

X, …

Y, …

Z, …

Pentaquark, …

How do we make sense of all of these?

Another particle explosion?

New particles, new ideas, and new theories

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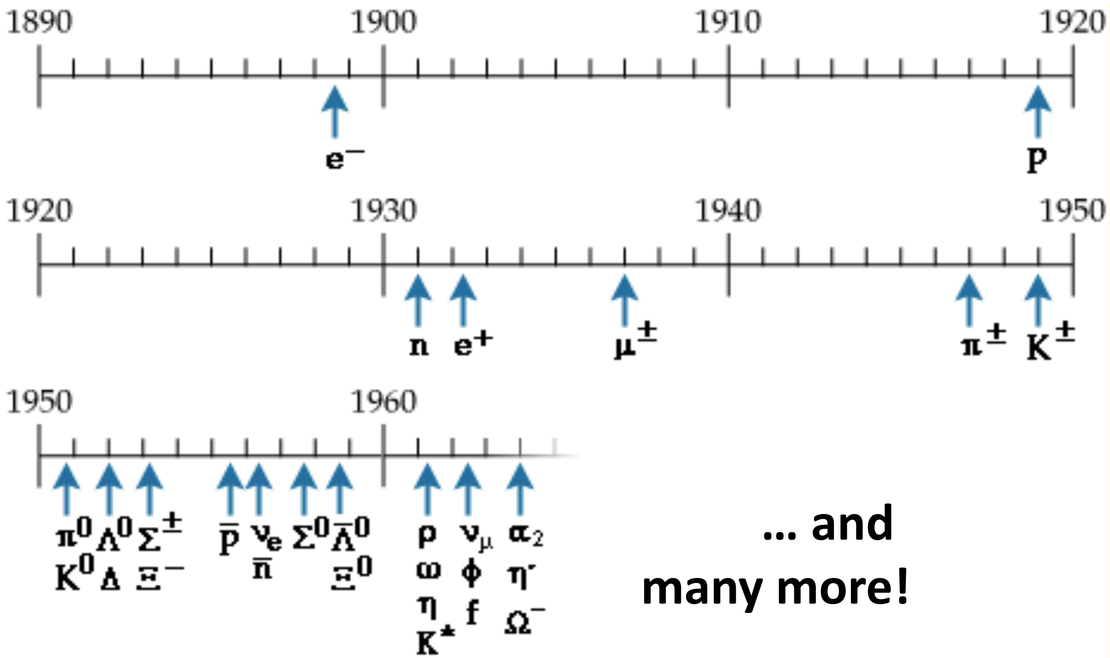

... and
many more!

❑ Nucleons has internal structure!

1933: Proton's magnetic moment

Otto Stern

Nobel Prize 1943

$$\mu_p = g_p \left(\frac{e\hbar}{2m_p}\right)$$

$$g_p = 2.792847356(23) \neq 2!$$

$$\mu_n = -1.913 \left(\frac{e\hbar}{2m_p}\right) \neq 0!$$

New particles, new ideas, and new theories

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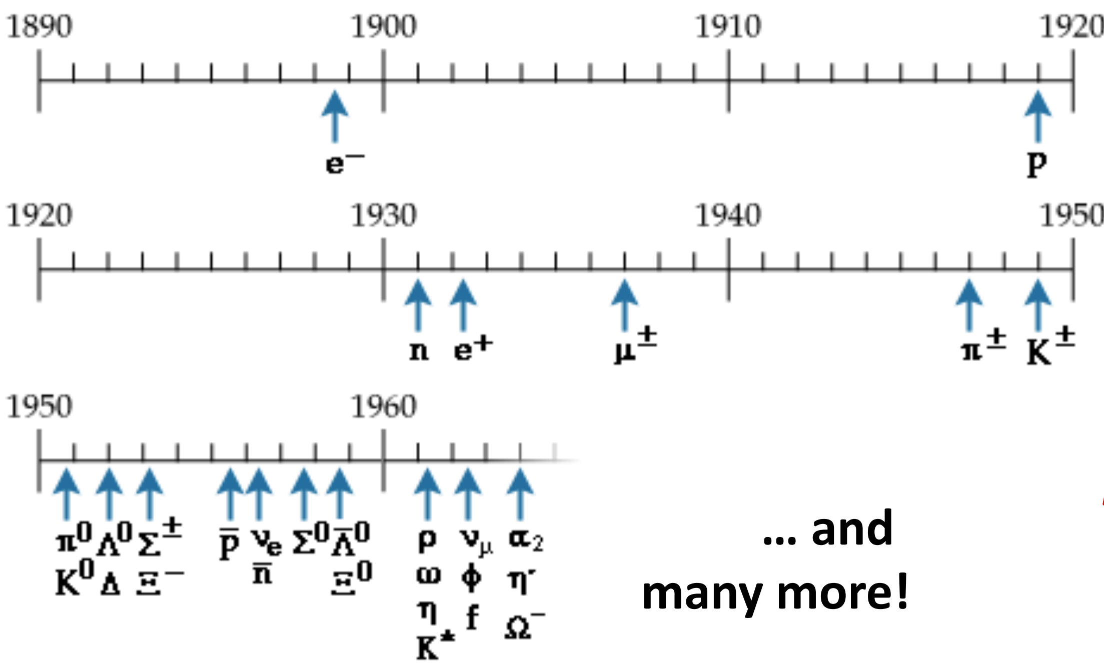
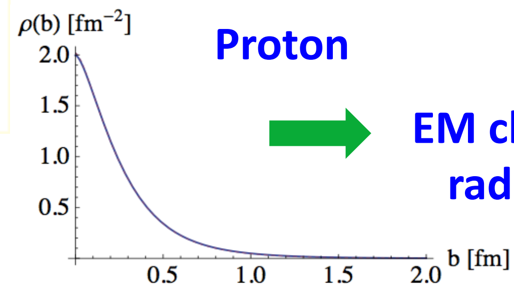


… and many more!

Electric charge distribution



Proton

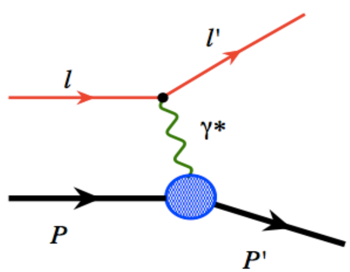
EM charge radius!

- ❑ **Nucleons has internal structure!**

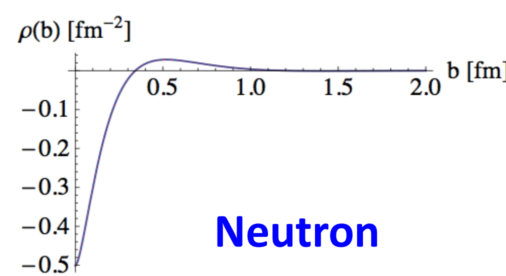
1960: Elastic e-p scattering

Robert Hofstadter

Nobel Prize 1961



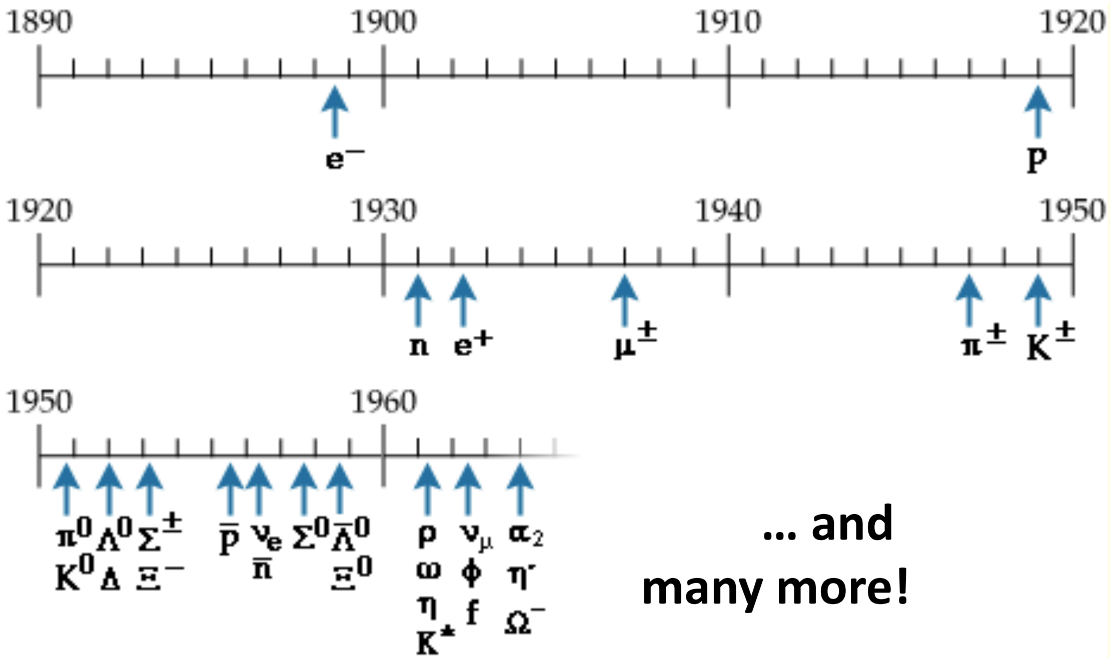
Form factors



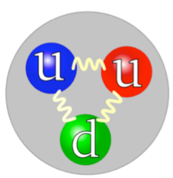
Neutron

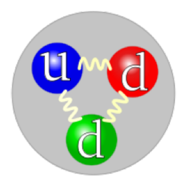
New particles, new ideas, and new theories

- Early proliferation of new hadrons – "particle explosion":



- Nucleons are made of quarks!

Proton 

Neutron 

→ Quark Model

Murray Gell-Mann
Nobel Prize, 1969

The naïve Quark Model

- **Flavor SU(3) – assumption:**

  Physical states for u, d, s, neglecting any mass difference, are represented by 3-eigenstates of the fund'l rep'n of flavor SU(3)

- **Generators for the fundamental rep'n of SU(3) – 3x3 matrices:**

  $J_i = \frac{\lambda_i}{2}$ **with** $\lambda_i, i = 1, 2, ..., 8$ **Gell-Mann matrices**

- **Good quantum numbers to label the states:**

$$J_3 = \frac{1}{2}\begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad J_8 = \frac{1}{2\sqrt{3}}\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

simultaneously diagonalized

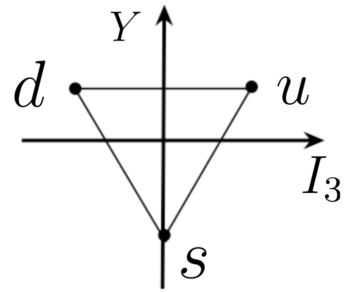
Isospin: $\hat{I}_3 \equiv J_3$ **, Hypercharge:** $\hat{Y} \equiv \frac{2}{\sqrt{3}} J_8$

- **Basis vectors – Eigenstates:** $|I_3, Y\rangle$

$$v^1 \equiv \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \Longrightarrow u = |\tfrac{1}{2}, \tfrac{1}{3}\rangle \quad v^2 \equiv \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \Longrightarrow d = |\text{-}\tfrac{1}{2}, \tfrac{1}{3}\rangle \quad v^3 \equiv \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \Longrightarrow s = |0, \text{-}\tfrac{2}{3}\rangle$$

The naïve Quark Model

❑ Quark states:



$$u = |\tfrac{1}{2}, \tfrac{1}{3}\rangle \quad d = |\text{-}\tfrac{1}{2}, \tfrac{1}{3}\rangle \quad s = |0,\text{-}\tfrac{2}{3}\rangle$$

Spin: ½

Baryon #: ***B* = ⅓**

Strangeness: ***S* = *Y* – *B*** **Electric charge:** $Q \equiv I_3 + \dfrac{Y}{2}$

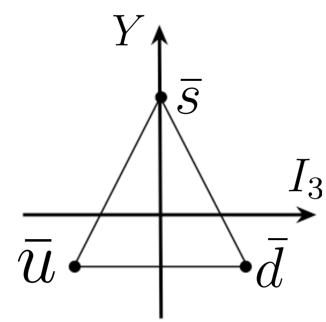
$$u \begin{cases} Q = 2/3\, e \\ s = 1/2 \\ I_3 = 1 \\ Y = 1/3 \\ B = 1/3 \\ S = 0 \end{cases} \qquad d \begin{cases} Q = -1/3\, e \\ s = 1/2 \\ I_3 = -1 \\ Y = 1/3 \\ B = 1/3 \\ S = 0 \end{cases} \qquad s \begin{cases} Q = -1/3\, e \\ s = 1/2 \\ I_3 = 0 \\ Y = -2/3 \\ B = 1/3 \\ S = -1 \end{cases}$$

❑ Antiquark states: $v_i \equiv \epsilon_{ijk}\, v^j v^k$

$$\hat{I}_3 v_1 = \epsilon_{123}[(\hat{I}_3 v^2)v^3 + v^2(\hat{I}_3 v^3)] + \epsilon_{132}[(\hat{I}_3 v^3)v^2 + v^3(\hat{I}_3 v^2)] = -\frac{1}{2}v_1$$

$$\hat{Y} v_1 = \epsilon_{123}[(\hat{Y} v^2)v^3 + v^2(\hat{Y} v^3)] + \epsilon_{132}[(\hat{Y} v^3)v^2 + v^3(\hat{Y} v^2)] = -\frac{1}{3}v_1$$

$$u \longrightarrow \bar{u} = |\text{-}\tfrac{1}{2},\text{-}\tfrac{1}{3}\rangle$$



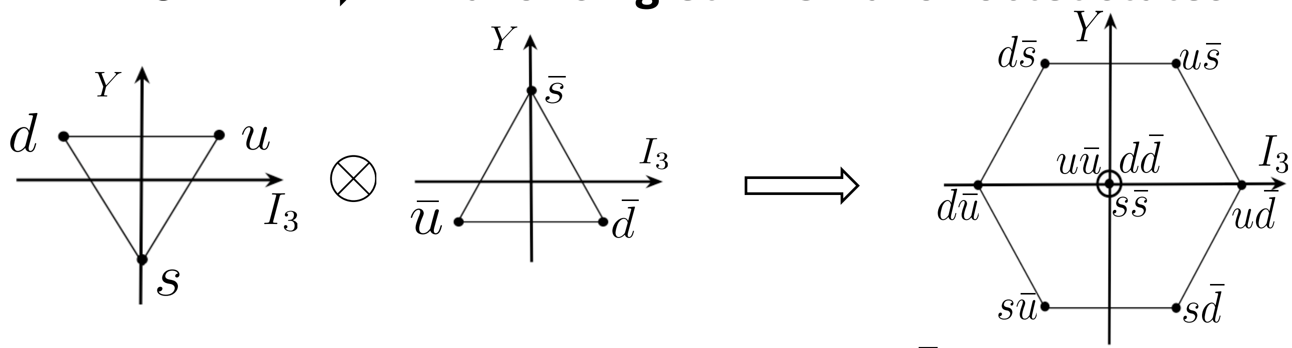
Mesons – Quark Model

Quark-antiquark $q\bar{q}$ flavor states:

❑ Group theory says:

$$q(u,d,s) = \mathbf{3}, \quad \bar{q}(\bar{u},\bar{d},\bar{s}) = \bar{\mathbf{3}}, \quad \text{of flavor SU(3)}$$

$\mathbf{3} \otimes \bar{\mathbf{3}} = \mathbf{8} \oplus \mathbf{1}$ $\Longrightarrow$ **1 flavor singlet + 8 flavor octet states**



There are three states with $I_3 = 0, Y = 0$: $u\bar{u}, d\bar{d}, s\bar{s}$

❑ Physical meson states (L=0, S=0):

✧ **Octet states:**

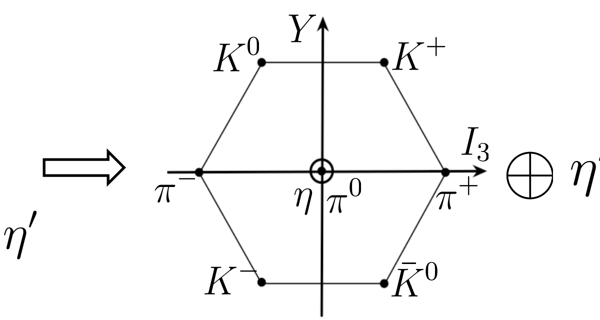
$$A = \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}) \quad \Longrightarrow \quad \pi^0$$

$$B = \frac{1}{\sqrt{6}}(u\bar{u} + d\bar{d} - 2s\bar{s}) \quad \Longrightarrow \quad \eta_8$$

✧ **Singlet states:**

$$C = \frac{1}{\sqrt{3}}(u\bar{u} + d\bar{d} + s\bar{s}) \quad \Longrightarrow \quad \eta_1$$

$\left.\begin{array}{l}\eta_8 \\ \eta_1\end{array}\right\} \eta, \eta'$

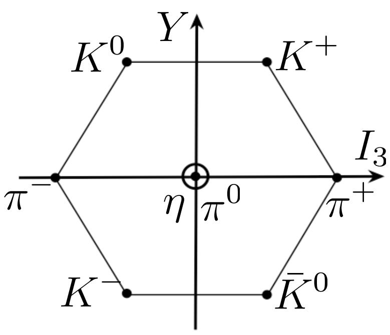


Quantum Numbers

- **Meson states:** J^{PC}
  - **Spin of $q\bar{q}$ pair:** $\vec{S} = \vec{s}_q + \vec{s}_{\bar{q}} \;\longrightarrow\; S = 0, 1$
  - **Spin of mesons:** $J = S + L$
  - **Charge conjugation:** $C = (-1)^{L+S}$

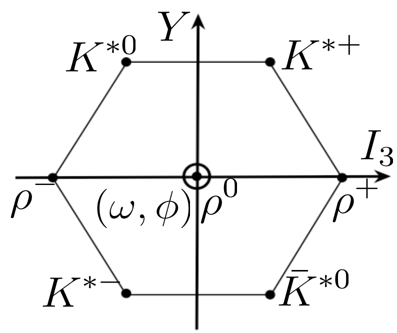
- **L=0 states:**

$J^{PC} = 0^{-+}$: ***(Y=S)***



$\oplus\ \eta'$

$J^{PC} = 1^{--}$: ***(Y=S)***



$\oplus\ \mathrm{mix}(\omega, \phi)$

Flavor singlet, spin octet

Flavor octet, spin octet

- **Color:**

No color was introduced!

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Baryons – Quark Model

3 quark qqq **states:** $B = 1$

❑ Group theory says:

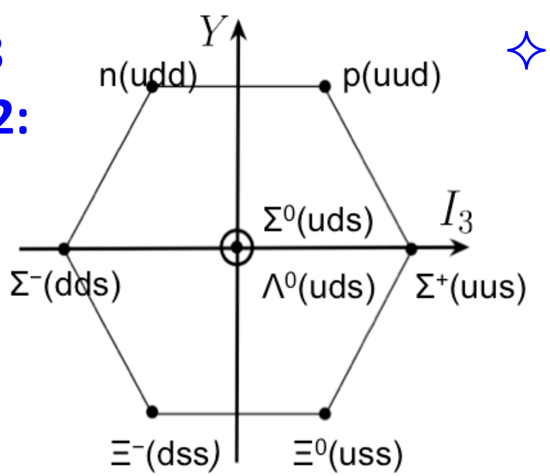
✧ **Flavor:** $\mathbf{3} \otimes \mathbf{3} \otimes \mathbf{3} = \mathbf{10}_S \oplus \mathbf{8}_{M_S} \oplus \mathbf{8}_{M_A} \oplus \mathbf{1}_A$

S: symmetric in all 3 q, M_S : symmetric in 1 and 2,
M_A : antisymmetric in 1 and 2, A : antisymmetric in all 3

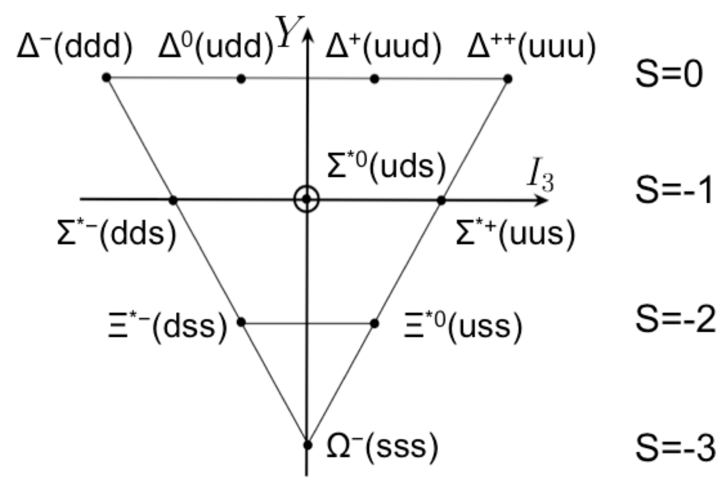
✧ **Spin:** $\mathbf{2} \otimes \mathbf{2} \otimes \mathbf{2} = \mathbf{4}_S \oplus \mathbf{2}_{M_s} \oplus \mathbf{2}_{M_A} \Longrightarrow S = \frac{3}{2}, \frac{1}{2}, \frac{1}{2}$

❑ Physical baryon states:

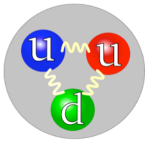
✧ **Flavor-8 Spin-1/2:**



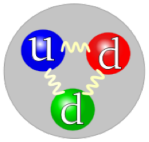
✧ **Flavor-10 Spin-3/2:**



Proton



Neutron



Δ^{++}(uuu), …

Violation of Pauli exclusive principle

Need another quantum number - color!

Color

❑ Minimum requirements:

- ✧ Quark needs to carry at least 3 different colors
- ✧ Color part of the 3-quarks' wave function needs to antisymmetric

❑ SU(3) color:

Recall: $\mathbf{3} \otimes \mathbf{3} \otimes \mathbf{3} = \mathbf{10}_S \oplus \mathbf{8}_{M_S} \oplus \mathbf{8}_{M_A} \oplus \mathbf{1}_A$

Antisymmetric color singlet state:

$$\Longrightarrow c(\text{Red, Green, Blue})$$

$$\psi_{\text{Color}}(c_1, c_2, c_3) = \frac{1}{\sqrt{6}}[\text{RGB-GRB+RBG-BRG+GBR-BGR}]$$

❑ Baryon wave function:

$$\Psi(q_1, q_2, q_3) = \psi_{\text{Space}}(x_1, x_2, x_3) \otimes \psi_{\text{Flavor}}(f_1, f_2, f_3) \otimes \psi_{\text{Spin}}(s_1, s_2, s_3) \otimes \psi_{\text{Color}}(c_1, c_2, c_3)$$

Antisymmetric Symmetric Symmetric Symmetric Antisymmetric

A complete example: Proton

- **Wave function – the state:**

$$|p\uparrow\rangle = \frac{1}{\sqrt{18}}\left[uud(\uparrow\downarrow\uparrow + \downarrow\uparrow\uparrow - 2\uparrow\uparrow\downarrow) + udu(\uparrow\uparrow\downarrow + \downarrow\uparrow\uparrow - 2\uparrow\downarrow\uparrow)\right.$$
$$\left. + duu(\uparrow\downarrow\uparrow + \uparrow\uparrow\downarrow - 2\downarrow\uparrow\uparrow)\right]$$

- **Normalization:**

$$\langle p\uparrow|p\uparrow\rangle = \frac{1}{18}[(1+1+(-2)^2) + (1+1+(-2)^2) + (1+1+(-2)^2)] = 1$$

- **Charge:**

$$\hat{Q} = \sum_{i=1}^{3}\hat{Q}_i$$

$$\langle p\uparrow|\hat{Q}|p\uparrow\rangle = \frac{1}{18}[(\tfrac{2}{3}+\tfrac{2}{3}-\tfrac{1}{3})(1+1+(-2)^2) + (\tfrac{2}{3}-\tfrac{1}{3}+\tfrac{2}{3})(1+1+(-2)^2)$$
$$+(-\tfrac{1}{3}+\tfrac{2}{3}+\tfrac{2}{3})(1+1+(-2)^2)] = 1$$

- **Spin:**

$$\hat{S} = \sum_{i=1}^{3}\hat{s}_i$$

$$\langle p\uparrow|\hat{S}|p\uparrow\rangle = \frac{1}{18}\{[(\tfrac{1}{2}-\tfrac{1}{2}+\tfrac{1}{2}) + (-\tfrac{1}{2}+\tfrac{1}{2}+\tfrac{1}{2}) + 4(\tfrac{1}{2}+\tfrac{1}{2}-\tfrac{1}{2})]$$
$$+[\tfrac{1}{2}+\tfrac{1}{2}+4\tfrac{1}{2}] + [\tfrac{1}{2}+\tfrac{1}{2}+4\tfrac{1}{2}]\} = \frac{1}{2}$$

- **Magnetic moment:**

$$\mu_p = \langle p\uparrow|\sum_{i=1}^{3}\hat{\mu}_i(\hat{\sigma}_3)_i|p\uparrow\rangle = \frac{1}{3}[4\mu_u - \mu_d]$$

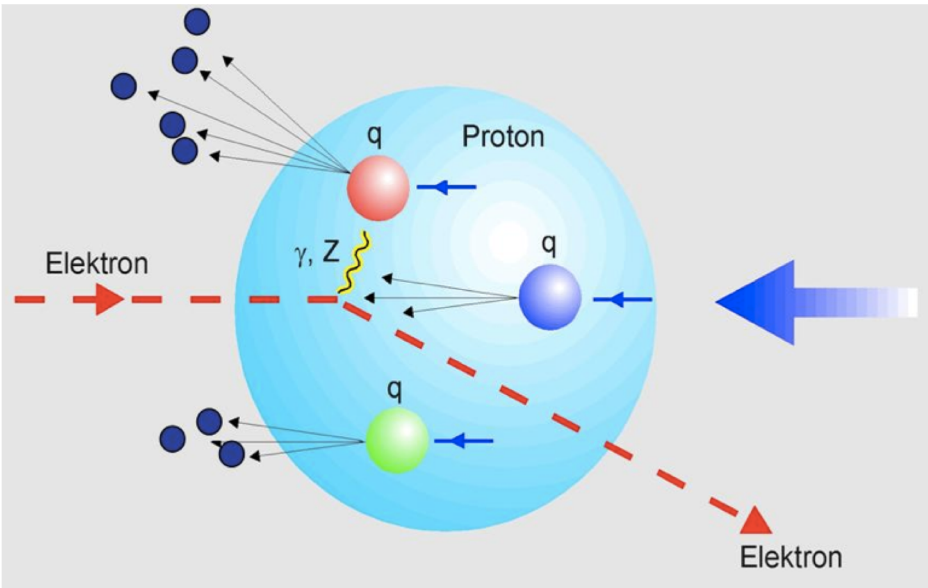
$$\mu_n = \frac{1}{3}[4\mu_d - \mu_u] \qquad \frac{\mu_u}{\mu_d} \approx \frac{2/3}{-1/3} = -2$$

$$\Rightarrow \begin{cases} \left(\dfrac{\mu_n}{\mu_p}\right)_{\mathrm{QM}} = -\dfrac{2}{3} \\ \left(\dfrac{\mu_n}{\mu_p}\right)_{\mathrm{Exp}} = -0.68497945(58) \end{cases}$$

How to "see" substructure of a nucleon?

❑ Modern Rutherford experiment – Deep Inelastic Scattering:

SLAC 1968: $e(p) + h(P) \to e'(p') + X$



✧ **Localized probe:**

$$Q^2 = -(p - p')^2 \gg 1 \text{ fm}^{-2}$$

$$\Rightarrow \quad \frac{1}{Q} \ll 1 \text{ fm}$$

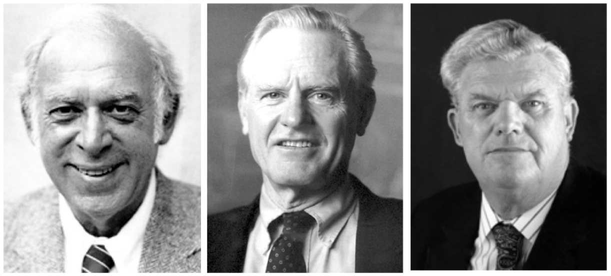
✧ **Two variables:**

$$Q^2 = 4EE' \sin^2(\theta/2)$$

$$x_B = \frac{Q^2}{2m_N \nu}$$

$$\nu = E - E'$$

→ **Discovery of spin ½ quarks, and partonic structure!**



What holds the quarks together?

Nobel Prize, 1990

⇒ **The birth of QCD (1973)**

– Quark Model + Yang-Mill gauge theory

Quantum Chromo-dynamics (QCD)

= A quantum field theory of quarks and gluons =

- **Fields:**

$\psi_i^f(x)$ — **Quark fields: spin-½ Dirac fermion (like electron)**

Color triplet: $i = 1, 2, 3 = N_c$

Flavor: $f = u, d, s, c, b, t$

$A_{\mu,a}(x)$ — **Gluon fields: spin-1 vector field (like photon)**

Color octet: $a = 1, 2, ..., 8 = N_c^2 - 1$

- **QCD Lagrangian density:**

$$\begin{aligned}\mathcal{L}_{QCD}(\psi, A) = \sum_f \overline{\psi}_i^f &\left[(i\partial_\mu \delta_{ij} - gA_{\mu,a}(t_a)_{ij})\gamma^\mu - m_f \delta_{ij}\right] \psi_j^f \\ &-\frac{1}{4}\left[\partial_\mu A_{\nu,a} - \partial_\nu A_{\mu,a} - gC_{abc}A_{\mu,b}A_{\nu,c}\right]^2 \\ &+ \text{gauge fixing} + \text{ghost terms}\end{aligned}$$

- **QED – force to hold atoms together:**

$$\mathcal{L}_{QED}(\phi, A) = \sum_f \overline{\psi}^f \left[(i\partial_\mu - eA_\mu)\gamma^\mu - m_f\right] \psi^f - \frac{1}{4}\left[\partial_\mu A_\nu - \partial_\nu A_\mu\right]^2$$

QCD is much richer in dynamics than QED

Gluons are dark, but, interact with themselves, NO free quarks and gluons

Gauge property of QCD

❑ Gauge Invariance:

$$\psi_i(x) \rightarrow \psi'_j(x) = U(x)_{ji}\,\psi_i(x)$$

$$A_\mu(x) \rightarrow A'_\mu(x) = U(x)\,A_\mu(x)\,U^{-1}(x) + \frac{i}{g}\left[\partial_\mu U(x)\right]U^{-1}(x)$$

where $A_\mu(x)_{ij} \equiv A_{\mu,a}(x)(t_a)_{ij}$

$$U(x)_{ij} = \left[e^{i\,\alpha_a(x)\,t_a}\right]_{ij} \qquad \text{Unitary} \quad [\,\text{det=1, SU(3)}]$$

❑ Color matrices:

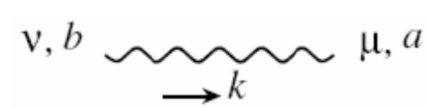
$$[\,t_a, t_b\,] = i\,C_{abc}\,t_c$$

Generators for the fundamental representation of SU3 color

❑ Gauge Fixing:

$$\mathcal{L}_{gauge} = -\frac{\lambda}{2}(\partial_\mu A_a^\mu)(\partial_\nu A_a^\nu)$$

Allow us to define the gauge field propagator:



$$G_{\mu\nu}(k)_{ab} = \frac{\delta_{ab}}{k^2}\left[-g_{\mu\nu} + \frac{k_\mu k_\nu}{k^2}\left(1 - \frac{1}{\lambda}\right)\right]$$

with $\lambda = 1$ **the Feynman gauge**

Ghost in QCD

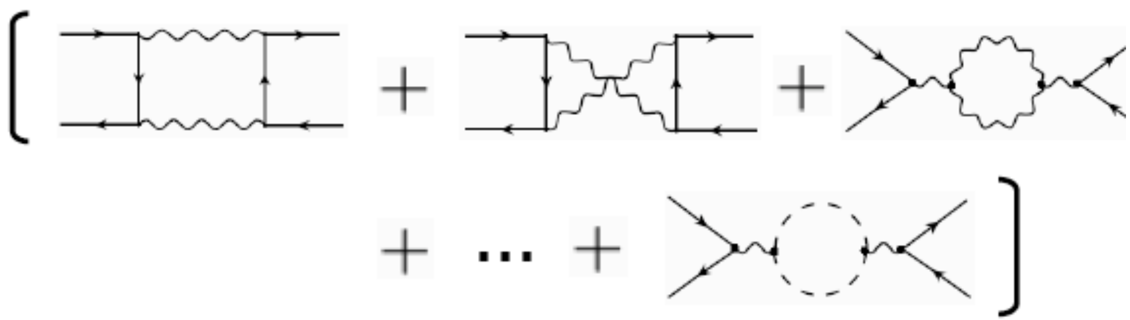
- **Ghost:**

Ghost

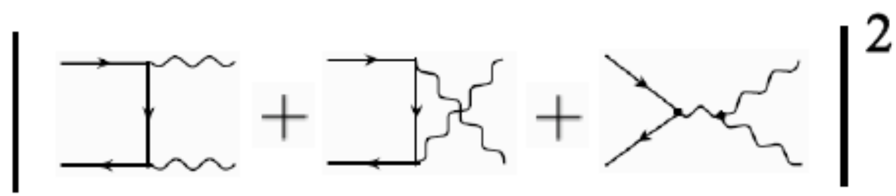
$$\mathcal{L}_{ghost} = (\partial_\mu \bar{\eta}_a(x))(\partial^\mu \eta_a(x) - g\, C_{abc}\, A_b^\mu(x)\, \eta_c(x))$$

so that the optical theorem (hence the unitarity) can be respected

2 Im



$= \Sigma$



Sum over all physical polarizations

Fail without the ghost loop

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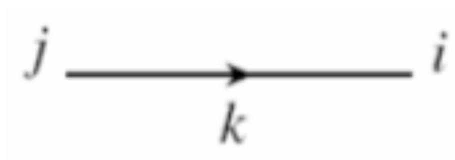
Feynman rules in QCD

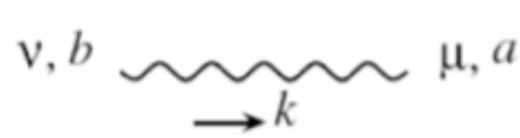
❑ Propagators:

Quark:



$$\frac{i}{\gamma \cdot k - m}\,\delta_{ij}$$

Gluon:



$$\frac{i\delta_{ab}}{k^2}\left[-g_{\mu\nu} + \frac{k_\mu k_\nu}{k^2}\left(1 - \frac{1}{\lambda}\right)\right]$$

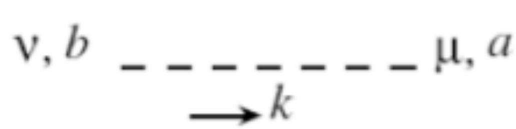
for a covariant gauge

$$\frac{i\delta_{ab}}{k^2}\left[-g_{\mu\nu} + \frac{k_\mu n_\nu + n_\mu k_\nu}{k \cdot n}\right]$$

for a light-cone gauge

$$n \cdot A(x) = 0 \quad \text{with} \quad n^2 = 0$$

Ghost::

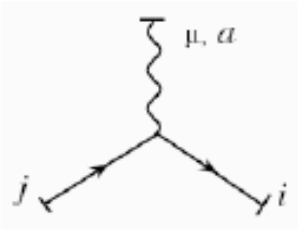


$$\frac{i\delta_{ab}}{k^2}$$

Feynman rules in QCD

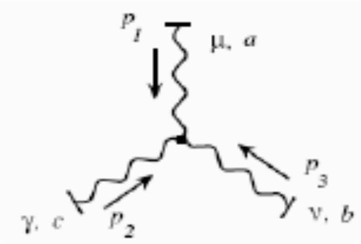
❑ **Interactions:**

$$-g\bar{\psi}\gamma^{\mu}A_{\mu,a}t_a\psi$$



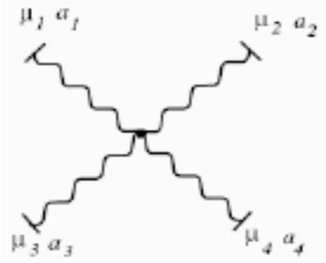
$$-ig(t_a)_{ij}\gamma_{\mu}$$

$$\frac{1}{2}gC_{abc}(\partial_{\mu}A_{\nu,a} - \partial_{\nu}A_{\mu,a})A_b^{\mu}A_c^{\nu}$$



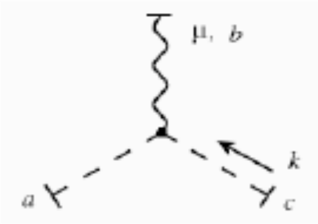
$$-gC_{abc}\left[g_{\mu\nu}(p_1 - p_2)_{\gamma} + g_{\nu\gamma}(p_2 - p_3)_{\mu} + g_{\gamma\mu}(p_3 - p_1)_{\nu}\right]$$

$$-\frac{g^2}{4}C_{abc}C_{ab'c'} * A_b^{\mu}A_c^{\nu}A_{\mu,b'}A_{\nu,c'}$$



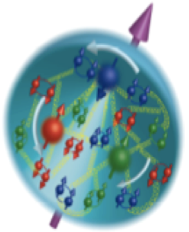
$$-ig^2\left[C_{ea_1a_2}C_{ea_3a_4} * (g_{\mu_1\mu_3}g_{\mu_2\mu_4} - g_{\mu_1\mu_4}g_{\mu_2\mu_3}) + \ldots\right]$$

$$\partial_{\mu}\bar{\eta}_a\left(gC_{abc}A_b^{\mu}\right)\eta_c$$



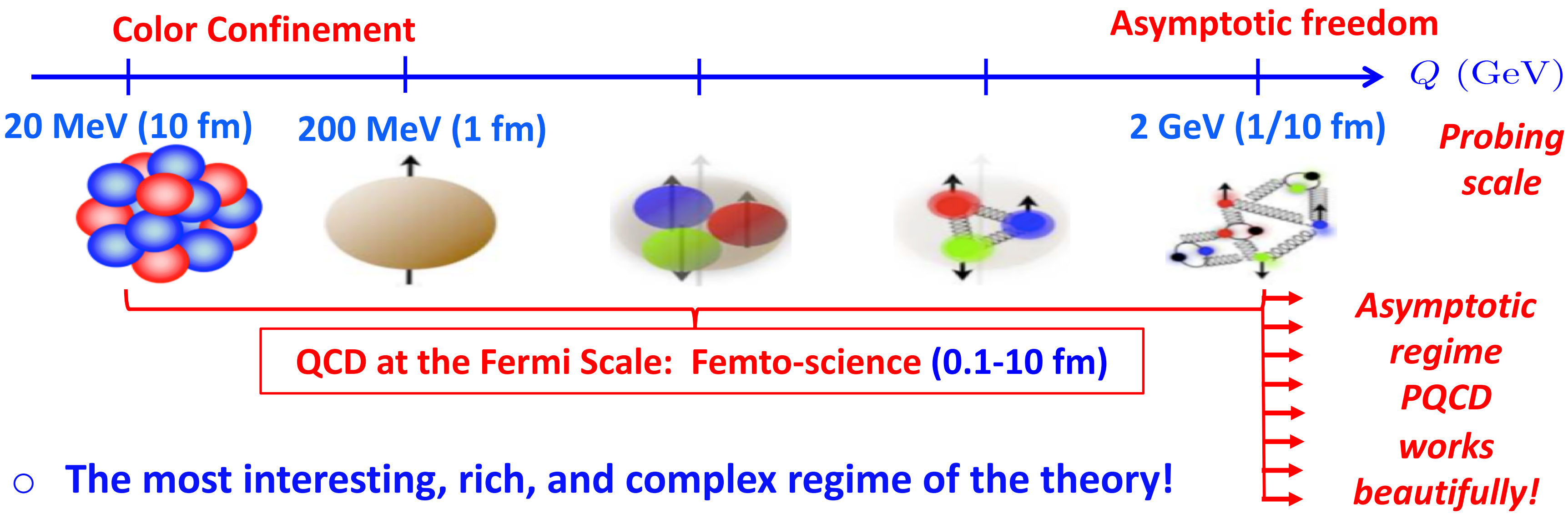
$$gC_{abc}k_{\mu}$$

QCD color is fully entangled



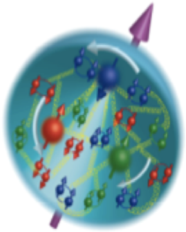
❑ QCD color confinement:

- *Do not see any quarks and gluons in isolation*
- *The structure of nucleons and nuclei – emergent properties of QCD*



- **The most interesting, rich, and complex regime of the theory!**
- **Emergent phenomena depend on the scale at which we probe them!**

QCD color is fully entangled



❑ QCD color confinement:

- *Do not see any quarks and gluons in isolation*
- *The structure of nucleons and nuclei – emergent properties of QCD*

Color Confinement — **Asymptotic freedom**

Q (GeV)

20 MeV (10 fm) — **200 MeV (1 fm)** — **2 GeV (1/10 fm)**

Probing scale

QCD at the Fermi Scale: Femto-science (0.1-10 fm)

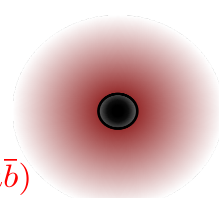
Asymptotic regime PQCD works beautifully!

- **The most interesting, rich, and complex regime of the theory!**
- **Emergent phenomena depend on the scale at which we probe them!**

❑ QCD is non-perturbative:

- *Any cross section/observable with identified hadron is not perturbatively calculable!*
- *Color is fully entangled!*

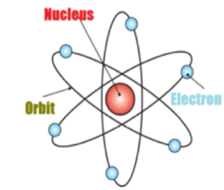
B-meson



$B^+(u\bar{b})$

Brown-Muck

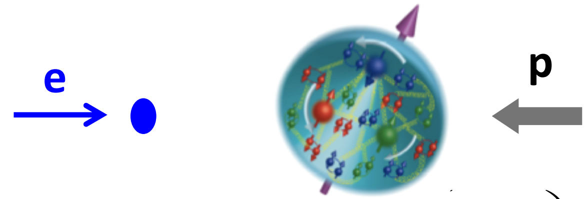
Atomic structure



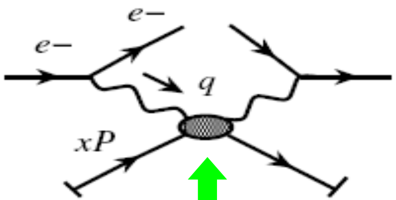
Quantum orbits

Theoretical Approaches - Approximations

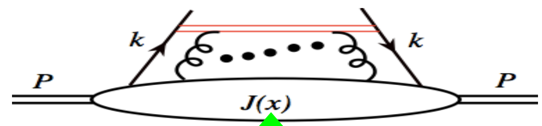
- **Perturbative QCD Factorization:**
  - *– Approximation at Feynman diagram level*



$\sigma_{\rm tot}^{\rm DIS} \sim$



$\otimes$



$+ \, O\left(\dfrac{1}{QR}\right)$

Probe Hard-part	Structure Parton-distribution	Approximation Power corrections

- **Effective field theory (EFT):**
  - *– Approximation at the Lagrangian level*

  Soft-collinear effective theory (SCET), Non-relativistic QCD (NRQCD), Heavy quark EFT, chiral EFT(s), …

- **Lattice QCD:**
  - *– Approximation mainly due to computer power*

  Hadron structure, hadron spectroscopy, nuclear structure, phase shift, …

- **Other approaches:**

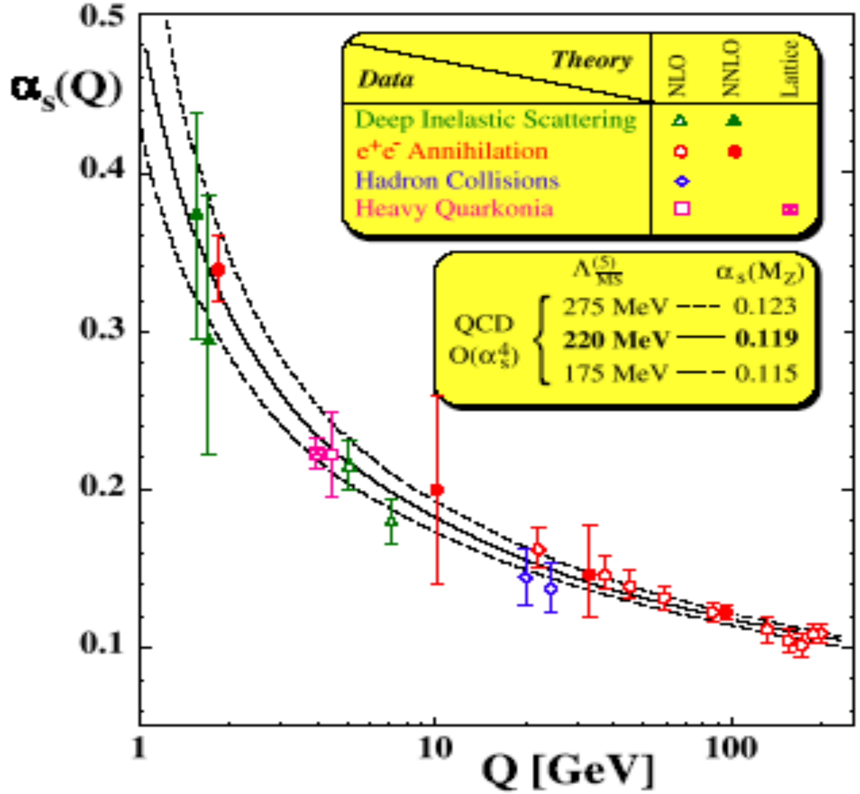
  Light-cone perturbation theory, Dyson-Schwinger Equations (DSE), Constituent quark models, AdS/CFT correspondence, …

QCD Asymptotic Freedom

- Interaction strength:

$$\alpha_s(\mu_2) = \frac{\alpha_s(\mu_1)}{1 - \frac{\beta_1}{4\pi}\alpha_s(\mu_1)\ln\left(\frac{\mu_2^2}{\mu_1^2}\right)} \equiv \frac{4\pi}{-\beta_1 \ln\left(\frac{\mu_2^2}{\Lambda_{\mathrm{QCD}}^2}\right)}$$

μ_2 and μ_1 not independent



Asymptotic Freedom ⇔ antiscreening

$$\text{QCD: } \frac{\partial \alpha_s(Q^2)}{\partial \ln Q^2} = \beta(\alpha_s) < 0$$

Compare

$$\text{QED: } \frac{\partial \alpha_{EM}(Q^2)}{\partial \ln Q^2} = \beta(\alpha_{EM}) > 0$$

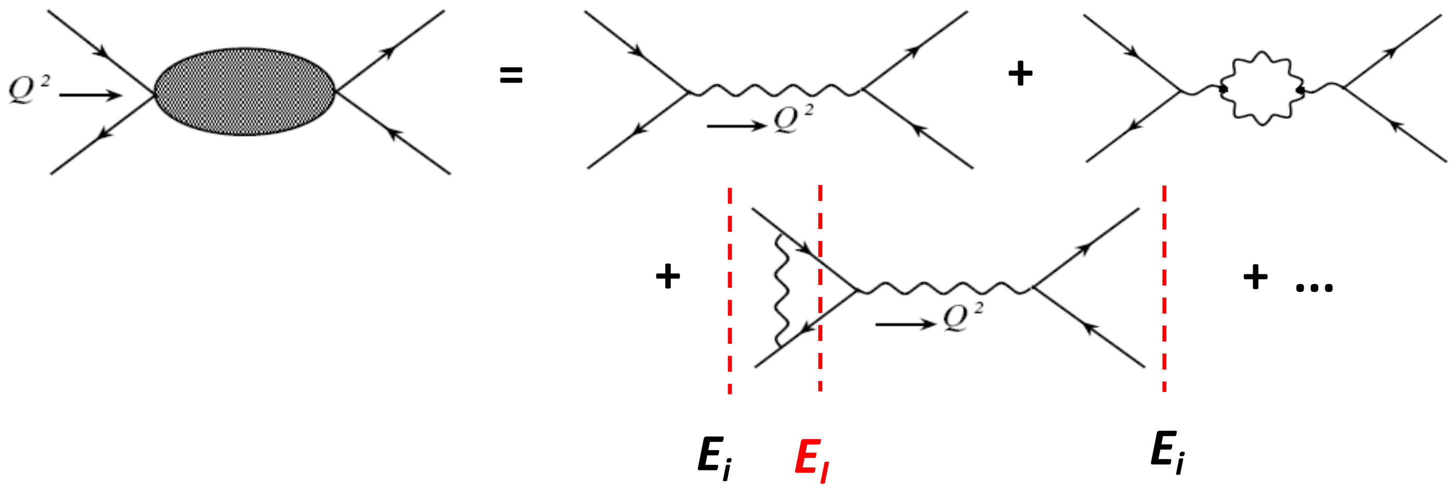
D.Gross, F.Willczek, Phys.Rev.Lett 30, (1973)
H.Politzer, Phys.Rev.Lett. 30, (1973)

→ Discovery of QCD Asymptotic Freedom

Nobel Prize, 2004

Renormalization, why need?

❑ Scattering amplitude:



$$= \int \langle PS \rangle_I \left(\frac{1}{E_i - E_I} + \ldots \right) + \ldots \Rightarrow \infty$$

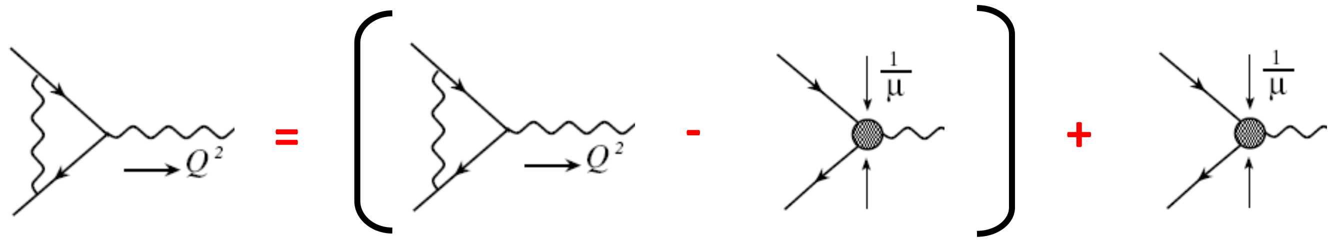
UV divergence: **result of a “sum” over states of high masses**

Uncertainty principle: **High mass states = “Local” interactions**

No experiment has an infinite resolution!

Physics of renormalization

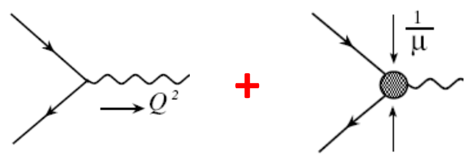
- UV divergence due to "high mass" states, not observed



"Low mass" state

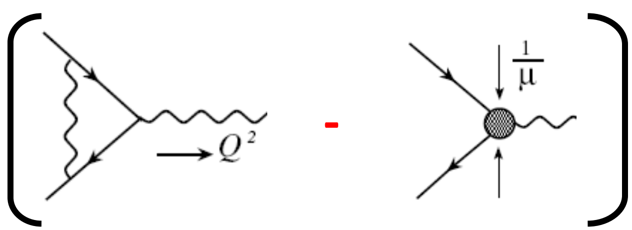
"High mass" states

- Combine the "high mass" states with LO

LO: 

= $g(\mu)$

Renormalized coupling

NLO: 

+ ...

No UV divergence!

- Renormalization = re-parameterization of the expansion parameter in perturbation theory

Renormalization Group

- **Physical quantity should not depend on renormalization scale μ ➡ renormalization group equation:**

$$\mu^2 \frac{d}{d\mu^2}\sigma_{\text{Phy}}\left(\frac{Q^2}{\mu^2}, g(\mu), \mu\right) = 0 \quad \Longrightarrow \quad \sigma_{\text{Phy}}(Q^2) = \sum_n \hat{\sigma}^{(n)}(Q^2,\mu^2)\left(\frac{\alpha_s(\mu)}{2\pi}\right)^n$$

- **Running coupling constant:**

$$\mu\frac{\partial g(\mu)}{\partial \mu} = \beta(g) \qquad \alpha_s(\mu) = \frac{g^2(\mu)}{4\pi}$$

- **QCD β function:**

$$\beta(g) = \mu\frac{\partial g(\mu)}{\partial\mu} = +g^3\frac{\beta_1}{16\pi^2} + \mathcal{O}(g^5) \qquad \beta_1 = -\frac{11}{3}N_c + \frac{4}{3}\frac{n_f}{2} < 0 \quad \text{for } n_f \le 6$$

- **QCD running coupling constant:**

$$\alpha_s(\mu_2) = \frac{\alpha_s(\mu_1)}{1 - \frac{\beta_1}{4\pi}\alpha_s(\mu_1)\ln\left(\frac{\mu_2^2}{\mu_1^2}\right)} \Rightarrow 0 \quad \text{as } \mu_2 \to \infty \quad \text{for } \beta_1 < 0$$

Effective Quark Mass

- **Running quark mass:**

$$m(\mu_2) = m(\mu_1) \exp\left[-\int_{\mu_1}^{\mu_2} \frac{d\lambda}{\lambda} [1 + \gamma_m(g(\lambda))] \right]$$

Quark mass depend on the renormalization scale!

- **QCD running quark mass:**

$$m(\mu_2) \Rightarrow 0 \qquad \text{as } \mu_2 \to \infty \qquad\qquad \text{since} \quad \gamma_m(g(\lambda)) > 0$$

- **Choice of renormalization scale:**

$\mu \sim Q$ **for small logarithms in the perturbative coefficients**

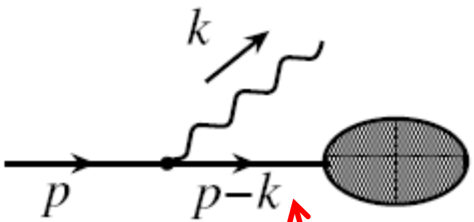
- **Light quark mass:** $m_f(\mu) \ll \Lambda_{\text{QCD}} \quad \text{for } f = u, d, \text{ even } s$

QCD perturbation theory ($Q \gg \Lambda_{QCD}$)
is effectively a massless theory

Infrared and collinear divergences

❑ **Consider a general diagram:**

$p^2 = 0, \quad k^2 = 0$ **for a massless theory**



Singularity

✧ $k^\mu \to 0 \;\Rightarrow\; (p-k)^2 \to p^2 = 0$

$\Longrightarrow$ **Infrared (IR) divergence**

✧ $k^\mu \,||\, p^\mu \;\Rightarrow\; k^\mu = \lambda p^\mu \quad \text{with} \quad 0 < \lambda < 1$

$\Rightarrow\; (p-k)^2 \to (1-\lambda)^2 p^2 = 0$

$\Longrightarrow$ **Collinear (CO) divergence**

IR and CO divergences are generic problems of a massless perturbation theory

Infrared Safety

- **Infrared safety:**

$$\sigma_{\text{Phy}}\left(\frac{Q^2}{\mu^2},\alpha_s(\mu^2),\frac{m^2(\mu^2)}{\mu^2}\right)\Rightarrow\hat{\sigma}\left(\frac{Q^2}{\mu^2},\alpha_s(\mu^2)\right)+\mathcal{O}\left[\left(\frac{m^2(\mu^2)}{\mu^2}\right)^{\kappa}\right]$$

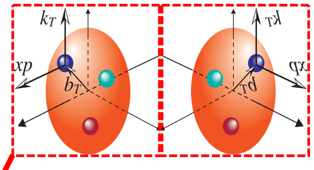
Infrared safe = κ > 0

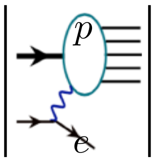
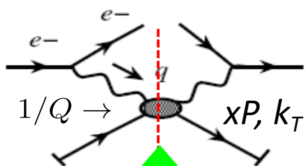
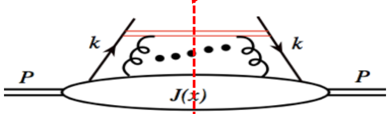
Asymptotic freedom is useful
only for
quantities that are infrared safe

- **Cross section with identified hadron(s):**
  - *Can not be calculated perturbatively!*
  - *Solution – QCD factorization:*
    - *to isolated what can be calculated perturbatively,*
    - *to represent the leading non-perturbative information by universal functions*
    - *to justify the approximation to neglect other nonperturbative information*

QCD Factorization

- ❑ **Factorization is an approximation!**

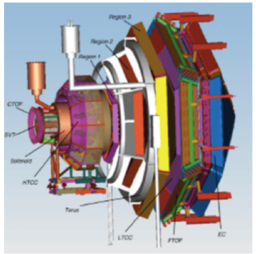


$\sigma_{\rm DIS}(x, Q^2) = \Big|$  $\Big|^2 \approx$  $\otimes$  $+ \, O\left(\frac{1}{QR}\right)$

Factorization Theory	Hard-part Probe	Parton-distribution Structure	Power corrections Approximation

- ❑ **QCD global analyses:**

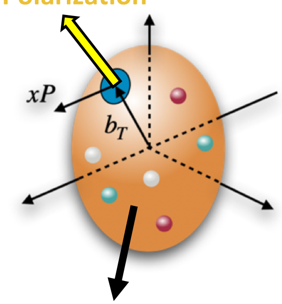
Experimental data



Jefferson Lab, KEK, CERN, BROOKHAVEN NATIONAL LABORATORY, Fermilab

Factorization Theory

$= \int \cdots \int$ Theory Calculation $\otimes$ Quantum Correlation Functions Unknown



Bayesian inference
Theory/Phenomenology
AI/ML, Computation

Extract QCFs

Jefferson Lab

Foundation of QCD perturbation theory

- ❑ **Renormalization**
  - – **QCD is renormalizable**

Nobel Prize, 1999
't Hooft, Veltman

- ❑ **Asymptotic freedom**
  - – **weaker interaction at a shorter distance**

Nobel Prize, 2004
Gross, Politzer, Welczek

- ❑ **Infrared safety and factorization**
  - – **calculable short distance dynamics**
  - – **pQCD factorization – connect the partons to physical cross sections**

J. J. Sakurai Prize, 2003
Mueller, Sterman

Look for infrared safe and factorizable observables!

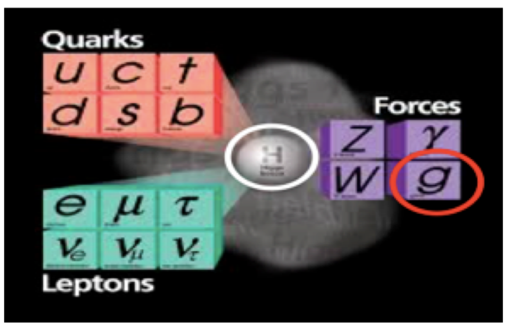
QCD is everywhere in our universe

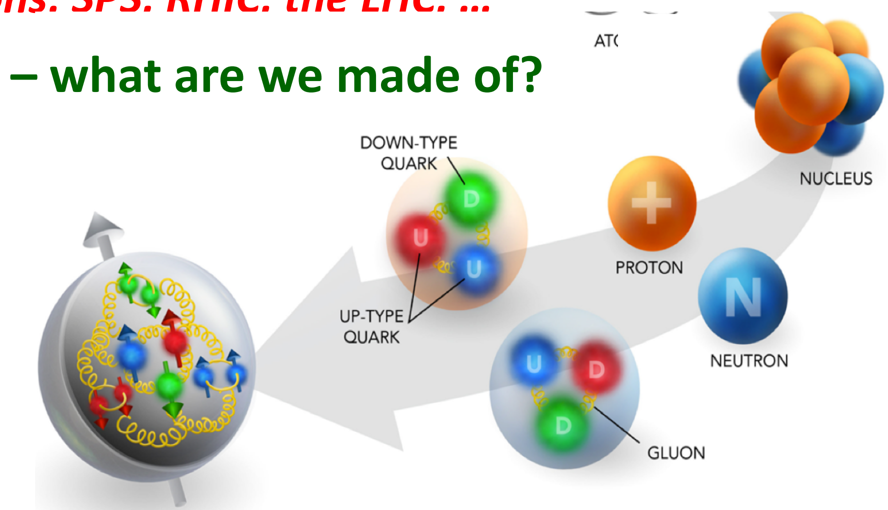
- ❑ **Understanding where did we come from?**

Global Time: →

- ○ **QCD at high temperature, high densities, phase transition, …**
- ○ ***Facilities – Relativistic heavy ion collisions: SPS. RHIC. the LHC. …***

- ❑ **Understanding the visible world at 3°K – what are we made of?**





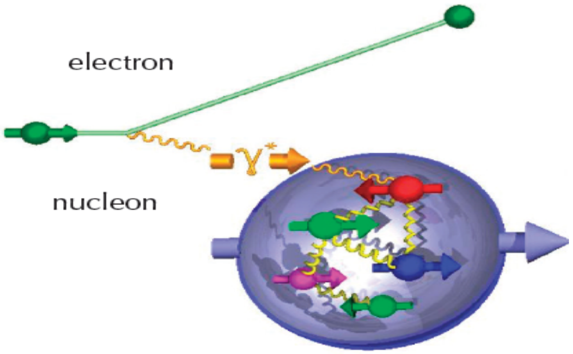
- ○ **How to understand the emergence and properties of nucleon and nuclei (elements of the periodic table) in terms of elements of the modern periodic table?**
- ○ **How does the glue bind us all?**
- ○ ***Facilities – CEBAF, EIC, …***

Nuclear Femtography

Search for answers to these questions at a Fermi scale!

Lepton-hadron scattering facility

- The new generation of “Rutherford” experiment:



- ✧ A controlled “probe” – virtual photon
- ✧ Can either break or not break the hadron

One facility covers all!
(JLab, COMPASS, EIC, …)

✧ Inclusive events: e+p/A → e’+X

Detect only the scattered lepton in the detector

(Modern Rutherford experiment!)

✧ Semi-Inclusive events: e+p/A → e’+h(p,K,p,jet)+X

Detect the scattered lepton in coincidence with identified hadrons/jets

(Initial hadron is broken – confined motion! – cleaner than h-h collisions)

✧ Exclusive events: e+p/A → e’+ p’/A’+ h(p,K,p,jet)

Detect every things including scattered proton/nucleus (or its fragments)

(Initial hadron is NOT broken – tomography!
– almost impossible for h-h collisions)