

Timelike Compton Scattering with CLAS12 at Jefferson Lab

CLAS Collaboration meeting - 2nd June 2021

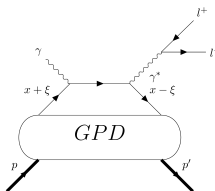
Pierre Chatagnon

INFN Genova (CMS group)

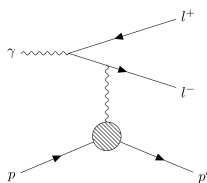
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Timelike Compton Scattering

$$\text{TCS: } \gamma p \rightarrow e^+ e^- p' \quad \text{DVCS: } ep \rightarrow e' p' \gamma$$



TCS



Bethe-Heitler

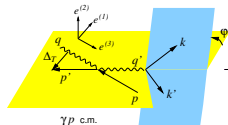
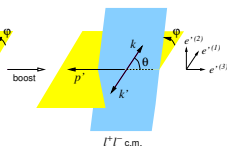
 γp c.m. $l^+ \Gamma$ c.m.

Figure in Berger et al., EPJ C, 2002

- BH cross section only depends on electromagnetic FFs
- Unpolarized interference cross section

$$\frac{d^4 \sigma_{INT}}{dQ'^2 dt d\Omega} \propto \frac{L_0}{L} \left[\cos(\phi) \frac{1+\cos^2(\theta)}{\sin(\theta)} \text{Re} \tilde{M}^{--} + \dots \right]$$

- Polarized interference cross section

$$\frac{d^4 \sigma_{INT}}{dQ'^2 dt d\Omega} = \frac{d^4 \sigma_{INT}}{dQ'^2 dt d\Omega} \Big|_{\text{unpol.}} - \nu \cdot A \frac{L_0}{L} \left[\sin(\phi) \frac{1+\cos^2(\theta)}{\sin(\theta)} \text{Im} \tilde{M}^{--} + \dots \right]$$

$$\rightarrow \tilde{M}^{--} = \frac{2\sqrt{t_0 - t}}{M} \frac{1 - \xi}{1 + \xi} \left[F_1 \mathcal{H} - \xi(F_1 + F_2) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E} \right]$$

Both $\text{Im} \mathcal{H}$ and $\text{Re} \mathcal{H}$ can be accessed in TCS

Motivations to measure TCS

Test of universality of GPDs

- TCS is parametrized by GPDs
- **Comparison between DVCS and TCS** results allows to test the **universality** of GPDs
- TCS does not involve Distribution Amplitudes unlike Deeply Virtual Meson Production
→ direct comparison between DVCS and TCS

Real part of CFFs and nucleon D-term

- As for DVCS, TCS unpolarized cross section is sensitive to $\text{Re}\mathcal{H}$, which is still not well constrained by existing data.
- The CFFs dispersion relation at leading order and leading twist :

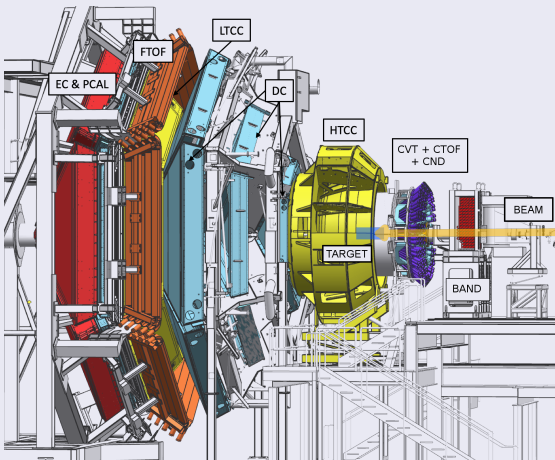
$$\text{Re}\mathcal{H}(\xi, t) = \mathcal{P} \int_{-1}^1 dx \left(\frac{1}{\xi - x} - \frac{1}{\xi + x} \right) \text{Im}\mathcal{H}(\xi, t) + D(t)$$

- $D(t)$ can be related to the **mechanical properties** of the nucleon.

Review in Polyakov, Schweitzer, *International Journal of Modern Physics A*, 2018

Experimental setup

CLAS12



● Forward Detector (6 sectors)

- Torus magnet
- Drift Chambers
- Forward Time-of-Flight
- Calorimeters (EC and PCAL)
- Cherenkov counters

● Central Detector

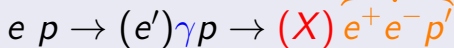
- Solenoid magnet
- Central Vertex Tracker (Silicon and micromegas)
- Central Time-of-Flight
- Central Neutron Detector

Figure in Burkert et al., *NIM A*, 2020

Data set used in this work

- Fall 2018 run period
- LH_2 target / 10.6 GeV beam / RG-A
- Inbending torus magnetic field
- Accumulated charge: ~ 150 mC (~ 200 fb $^{-1}$)

Analysis strategy



Final state selection from PID

Exclusivity cuts

$$P_X = p_{beam} + p_{target} - p_{e^+} - p_{e^-} - p_{p'}$$

$$|M_X^2| < 0.4 \text{ GeV}^2$$

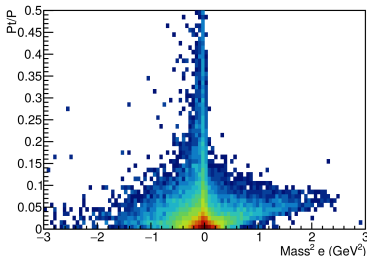
Quasi-real photoproduction

$$\frac{P_{tX}}{P_X} < 0.05$$

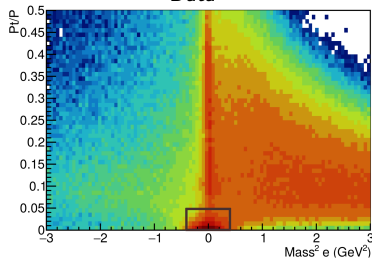
$$\rightarrow Q^2 < 0.1 \text{ GeV}^2$$

after momentum corrections and fiducial cuts [▶ Here](#)

Simulation



Data



Positron identification

Definitions

Signal: e^+ identified as e^+

Background: π^+ identified as e^+

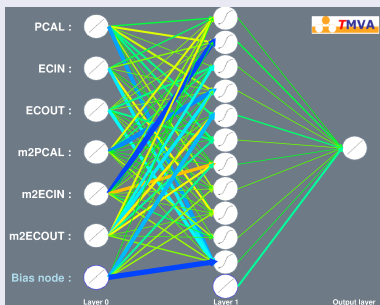
Strategy and discriminating variables

Positron: electromagnetic shower

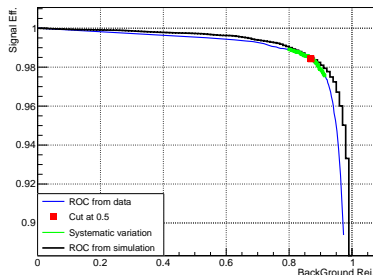
Pion: Minimum Ionizing Particle (MIP)

$$SF_{\text{EC Layer}} = \frac{E_{\text{dep}}(\text{EC Layer})}{P}$$

$$M_2 = \frac{1}{3} \sum_{U,V,W} \frac{\sum_{\text{strip}} (x-D)^2 \cdot \ln(E)}{\sum_{\text{strip}} \ln(E)} \rightarrow \mathbf{6 \text{ variables}}$$



Output: **Signal** \rightarrow 1 **Background** \rightarrow 0



B/S from 50% to 5%

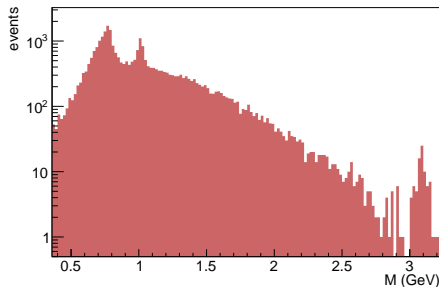
TCS with CLAS12 (P. Chatagnon)

• **Signal in data** \Rightarrow Outbending electrons

• **Background in data** $\Rightarrow ep \rightarrow e\pi^+_{PID=e^+_{6/16}}$

Data/Simulation comparison

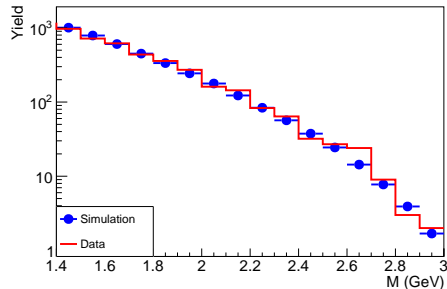
- Vector mesons peaks are visible in data:
 ω (770 MeV), ρ (782 MeV),
 Φ (1020 MeV) and J/ψ (3096 MeV)



- Data/BH comparison in the high mass region, no evident high mass vector meson production (ρ (1450 MeV, 1700 MeV))

Phase space of interest

- $1.5 \text{ GeV} < M_{e^+e^-} < 3 \text{ GeV}$
- $0.15 \text{ GeV}^2 < -t < 0.8 \text{ GeV}^2$
- $4 \text{ GeV} < E_\gamma < 10.6 \text{ GeV}$.



Data/simulation are matching at 15 % level, up to normalization factor

Acceptance

Acceptance calculation using BH-weighted events

$$Acc_B = \frac{N_B^{REC}}{N_B^{GEN}}$$

$$N_B^{REC} = \sum_{REC \in B} \text{Eff}_{corr} w$$

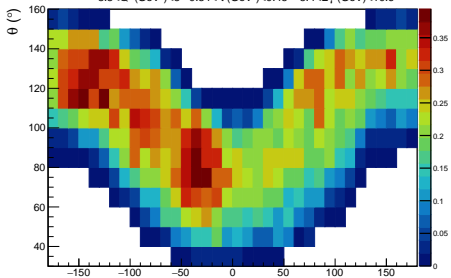
$$N_B^{GEN} = \sum_{GEN \in B} w$$

Multidimensional binning of the acceptance

4 bins in $-t$, 3 bins in E_γ and Q'^2 , $10^\circ \times 10^\circ$ bins in the ϕ/θ plane. Bins with $\frac{\Delta Acc}{Acc} > 0.5$ and $Acc < 0.05$ are discarded (ΔAcc is statistical error).

Large region with no acceptance
($\phi \sim 0^\circ/\theta \sim 180^\circ$ and $\phi \sim 180^\circ/\theta \sim 0^\circ$)

$3.5 < Q'^2 (\text{GeV}^2) < 5$ $0.34 < -t (\text{GeV}^2) < 0.48$ $8.4 < E_\gamma (\text{GeV}) < 10.6$



Efficiency corrections

- **Data-driven correction** for the proton detection efficiency derived using $ep \rightarrow e' \pi^+ \pi^- (p')$ reaction
- Efficiency correction from **background merging** using random trigger events

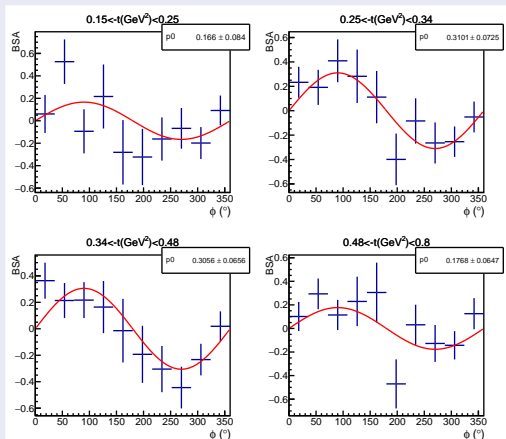
Observable 1: Photon polarization asymmetry (BSA)

Access to the imaginary part of CFFs

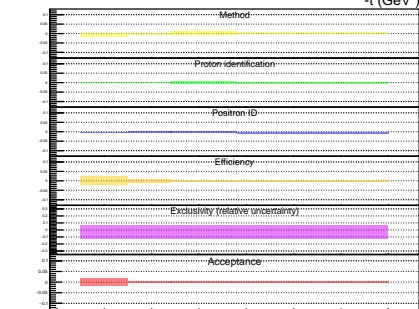
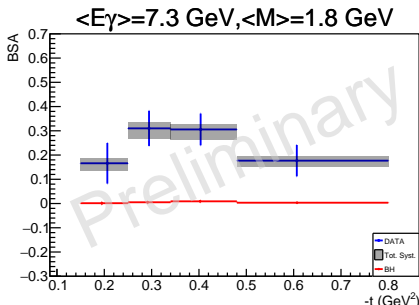
$$BSA = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} = \frac{-\frac{\alpha^3}{4\pi s^2} \frac{1}{-t} \frac{m_p}{Q'} \frac{1}{\tau\sqrt{1-\tau}} \frac{L_0}{L} \sin\phi \frac{(1+\cos^2\theta)}{\sin(\theta)}}{d\sigma_{BH}} \text{Im}\tilde{M}^{--}$$

Experimental measurement

- $BSA(-t, E_\gamma, M; \phi) = \frac{1}{Pol_{eff}} \frac{N^+ - N^-}{N^+ + N^-}$ where $N^\pm = \sum \frac{1}{Acc} Pol_{transf}$.
- Pol_{transf} is the **transferred polarization** from the electron to the photon
- Pol_{eff} is the **polarization of the CEBAF electron beam** (85%)
- The ϕ -distribution is fitted with a sine function



Systematics



Method

- Calculated from generated BH events, and full-chain simulated events.

Proton

- Apply χ^2 cut for the proton identification

Positron Identification

- Vary the positron ID cut (0.5 ± 0.3 ; max. significance region)

Efficiency

- Calculate observable with/without data-driven proton efficiency

Exclusivity cuts

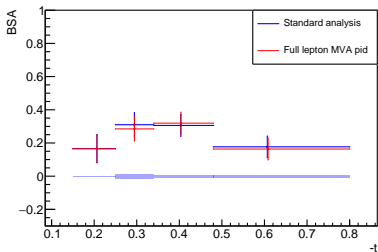
- Vary the values of the exclusivity cuts:
 $|Pt/P| < 0.05 \pm 0.01, |M_X^2| < 0.4 \pm 0.1 \text{ GeV}^2$
Fully integrated relative uncertainty

Acceptance

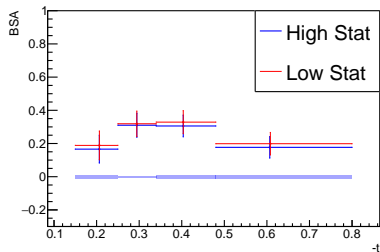
- Calculate observable with acceptance produced using BH-weighted events or unity weights
- Neighboring bins uncertainties are averaged
- Then added in quadrature

Additional systematics checks

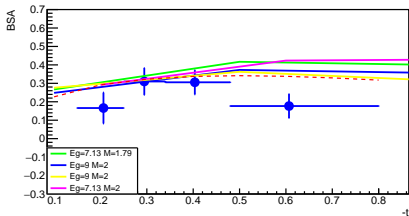
- Stability of the MVA at low momenta for all leptons



- MC statistic: acceptance calculated with 16M/36M generated events



- Study the stability of the model prediction within the integrated kinematic range

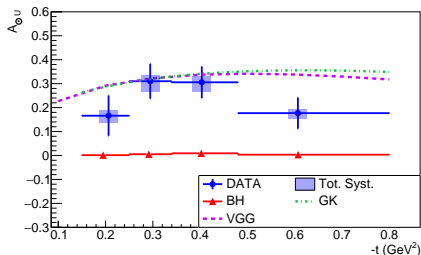


BSA selected results

- First time measurement
- A **sizeable asymmetry** is measured (above the expected vanishing BSA of BH)
→ **signature of TCS**
- Theoretical predictions were provided by M.Vanderhaeghen (using the VGG model) and P.Sznajder (using the GK model)
- Size of the asymmetry is **well reproduced** by VGG and GK models
→ **model dependent hints for universality of GPDs**

$$\langle M \rangle = 1.8 \text{ GeV}; \langle E_\gamma \rangle = 7.29 \text{ GeV};$$

$$\langle \theta \rangle = 92^\circ$$



Observable 2: Forward-Backward asymmetry

- Concept explored for J/Ψ production (Gryniuk, Vanderhaeghen, *Phys. Rev. D*, 2016).
- Exploratory studies for TCS performed in my thesis.
- Very first predictions for TCS have been published very recently (Heller, Keil, Vanderhaeghen, *Phys. Rev. D*, 2021).
- Use the different parity of the TCS and BH amplitudes under the inversion of the leptons directions
 $k \leftrightarrow k' \iff (\theta, \phi) \leftrightarrow (180^\circ - \theta, 180^\circ + \phi)$

BH cross section

$$\frac{d\sigma_{BH}}{dQ^2 dt d\Omega} \propto \frac{1+\cos^2\theta}{\sin^2\theta} \xrightarrow{FB} \frac{d\sigma_{BH}}{dQ^2 dt d\Omega}$$

Int. cross section

$$\frac{d^4\sigma_{INT}}{dQ'^2 dt d\Omega} \propto \frac{L_0}{L} \cos(\phi) \frac{1+\cos^2(\theta)}{\sin(\theta)} \xrightarrow{FB} -\frac{d\sigma_{INT}}{dQ^2 dt d\Omega}$$

A_{FB} formula

$$A_{FB}(\theta_0, \phi_0) = \frac{d\sigma(\theta_0, \phi_0) - d\sigma(180^\circ - \theta_0, 180^\circ + \phi_0)}{d\sigma(\theta_0, \phi_0) + d\sigma(180^\circ - \theta_0, 180^\circ + \phi_0)} = \frac{-\frac{\alpha_{em}^3}{4\pi s^2} \frac{1}{-t} \frac{m_P}{Q'} \frac{1}{\tau\sqrt{1-\tau}} \frac{L_0}{L} \cos\phi_0 \frac{(1+\cos^2\theta_0)}{\sin(\theta_0)}}{d\sigma_{BH}} \text{Re}\tilde{M}^{--}$$

- Access to the real part of the CFFs with no integration over angles
- Removes large dependencies on angular acceptance \rightarrow direct comparison with models
- But smaller phase space \rightarrow lower statistics

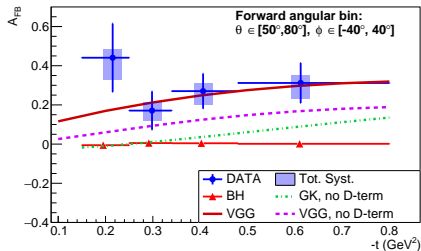
A_{FB} selected results

- A_{FB} measured in two mass regions:
 $M \in [1.5 \text{ GeV}, 3 \text{ GeV}]$ and
 $M \in [2 \text{ GeV}, 3 \text{ GeV}]$
- The measured A_{FB} is non-zero:
evidence for signal beyond pure BH contribution
- Three model predictions
 - 1 VGG without D-term
 - 2 VGG with D-term

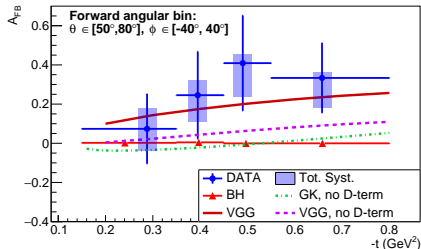
D-term in Pasquini et al., *Physics Letters B*, 2014

 - 3 GK without D-term
- Measured asymmetry is better reproduced by the VGG model **including the D-term** in both mass bins

$\langle M \rangle = 1.8 \text{ GeV}; \langle E_\gamma \rangle = 7.24 \text{ GeV}$



$\langle M \rangle = 2.25 \text{ GeV}; \langle E_\gamma \rangle = 8.13 \text{ GeV}$



Documentation and code

Analysis review

Review #1461061: "Timelike Compton scattering with CLAS12"

→ ended Monday 31st of May

We are ready for ad-hoc review ! Request sent yesterday

Analysis note

Final version of the analysis note available [▶ here](#)

Analysis scripts

(Personal) Github repository containing all the scripts used for the analysis is being set up [▶ here](#)

