

Nuclear effects in the deuteron: model-independent constraints from QCD global analysis

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https://www.jlab.org/jam

Overview

- Nucleon structure and global QCD analysis of PDFs
- Nuclear effects in DIS from the deuteron
- $\blacksquare d/u \text{ PDF ratio at large } x$
- JAM Monte Carlo analysis
- Outlook

Parton distributions in the nucleon

Inclusive high-energy particle production $AB \rightarrow CX$



Collins, Soper, Sterman (1980s)

→ <u>QCD factorization</u>: separation of hard (perturbative, calculable) from soft (nonperturbative, parametrized) physics

$$\sigma_{AB\to CX}(p_A, p_B) = \sum_{a,b} \int dx_a \, dx_b \, \underbrace{f_{a/A}(x_a, \mu)}_{\dots} \underbrace{f_{b/B}(x_b, \mu)}_{\dots} \times \sum_n \alpha_s^n(\mu) \, \hat{\sigma}_{ab\to CX}^{(n)} \left(x_a p_A, x_b p_B, Q/\mu\right)$$

→ process-independent parton distribution functions $f_{a/A}$ characterizing structure of bound state A



Parton distributions in the nucleon

Ubiquity of proton F_2 data (SLAC, EMC, NMC, BCDMS, HERA, JLab, ...) provides strong constraints on *u*-quark PDF over large *x* range



$$F_2^p \sim \frac{4}{9}xu + \frac{1}{9}xd + \cdots$$

- Absence of free-neutron data and smaller $|e_q|$ of d quarks limit precision of d-quark PDF, especially at high x
 - → <u>nuclear effects in deuteron</u> obscure free-neutron structure

- Approximate scattering from weakly-bound nuclei at $x \gg 0$ in terms of incoherent scattering from bound nucleons
 - → generalized convolution in "weak binding approximation" (WBA)

$$F_2^A(x,Q^2) = \sum_N \int \frac{d^4p}{(2\pi)^4} \mathcal{F}_0^N(\varepsilon,\mathbf{p}) \left(1 + \frac{\gamma p_z}{M}\right) \mathcal{C}_{22} \, \widetilde{F}_2^N(x/y,Q^2,p^2)$$

nuclear spectral function

bound nucleon momentum $p = (p_0; \mathbf{p}) = (M + \varepsilon; \mathbf{p}_{\perp}, p_z)$

kinematic factor
$$C_{22} = \frac{1}{\gamma^2} \left[1 + \frac{(\gamma^2 - 1)}{2y^2 M^2} (2p^2 + 3\mathbf{p}_{\perp}^2) \right] \qquad \gamma^2 = 1 + \frac{4M^2 x^2}{Q^2}$$

nuclear momentum fraction $y = \frac{M_A}{M} \frac{p \cdot q}{P \cdot q} = \frac{p_0 + \gamma p_z}{M}$

 \rightarrow factorized formula valid up to $\mathcal{O}(\mathbf{p}^2/M^2)$ corrections

WM, Schreiber, Thomas (1994) Kulagin et al. (1994)

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off-shell nucleon structure function

 \rightarrow expand to lowest order in nucleon virtuality $(p^2 - M^2)$

$$\widetilde{F}_{2}^{N}(x,Q^{2},p^{2}) = F_{2}^{N}(x,Q^{2}) \left(1 + \frac{p^{2} - M^{2}}{M^{2}} \delta f^{N}(x)\right)$$

on-shell structure function

off-shell correction

$$\delta f^N = \frac{\partial \log \widetilde{F}_2^N}{\partial \log(p^2/M^2)} \bigg|_{p^2 = M^2}$$

Write total nuclear structure function as a sum of nucleon on-shell and off-shell contributions

$$F_2^A(x,Q^2) = F_2^{A(\text{on})}(x,Q^2) + F_2^{A(\text{off})}(x,Q^2)$$

where

$$F_2^{A(\text{on})}(x,Q^2) = \sum_N \int dy \, f^{N/A}(y,\gamma) \, F_2^N(x/y,Q^2)$$
$$F_2^{A(\text{off})}(x,Q^2) = \sum_N \int dy \left[\tilde{f}^{N/A}(y,\gamma) \, F_2^N(x/y,Q^2) \right] \delta f^N(x/y)$$

Nucleon "smearing functions" (light-cone momentum distributions)

on-shell
$$f^{N/A}(y,\gamma) = \int \frac{d^4p}{(2\pi)^4} \mathcal{F}_0^N(\varepsilon,\mathbf{p}) \left(1 + \frac{\gamma p_z}{M}\right) \mathcal{C}_{22} \,\delta\left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right)$$

off-shell
$$\tilde{f}^{N/A}(y,\gamma) = \int \frac{d^4p}{(2\pi)^4} \mathcal{F}_0^N(\varepsilon,\mathbf{p}) \left(1 + \frac{\gamma p_z}{M}\right) \mathcal{C}_{22} \frac{(p^2 - M^2)}{M^2} \delta\left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{C}_{22} \left(y - 1 - \frac{\varepsilon + \gamma p_z}{M}\right) \mathcal{$$



→ off-shell smearing functions \ll on-shell smearing functions for most kinematics of interest, strongly peaked around y = 1, opposite sign

→ for deuteron,
$$f^{p/D}(y) = f^{n/D}(y) \equiv f^{N/D}(y)$$

For A = D off-shell part of structure function can be written

$$F_2^{D(\text{off})}(x,Q^2) = \int dy \,\tilde{f}^{N/D}(y) \,F_2^{p+n}(x/y,Q^2) \,\delta f^D(x/y)$$

where

$$\delta f^{D}(x) \equiv \delta f^{0}(x) = \frac{\delta f^{p/D}(x)F_{2}^{p}(x) + \delta f^{n/D}(x)F_{2}^{n}(x)}{F_{2}^{p}(x) + F_{2}^{n}(x)}$$

 \rightarrow deuteron not sensitive to possible differences $\delta f^{p/D} \stackrel{?}{\neq} \delta f^{n/D}$

To ensure conservation of valence quark (baryon) number in the deuteron, off-shell function is normalized

$$\int_0^1 dx \, q_v(x) \, \delta f^{N/D}(x) = 0$$

for q = u, d and N = p, n

CJ (CTEQ-JLab) Collaboration has performed global QCD studies focusing in particular on the high-x, low-W region, to better constrain PDFs at large x

						χ^2	
Observable	Experiment	# points	LO	NLO	NLO (OCS)	NLO (no nucl)	NLO (no nucl/D0)
IS F_2	BCDMS (p) [81]	351	426	438	436	440	427
2	BCDMS (d) [81]	254	292	292	289	301	301
	SLAC (p) [82]	564	480	434	435	441	440
	SLAC (d) [82]	582	415	376	380	507	466
	NMC (p) [83]	275	416	405	404	405	403
	NMC (d/p) [84]	189	181	172	173	174	173
	HERMES (p) [86]	37	57	42	43	44	44
	HERMES (d) [86]	37	52	37	38	36	37
	Jefferson Lab (p) [87]	136	172	166	167	177	166
	Jefferson Lab (d) [87]	136	131	123	124	126	130
F_2 tagged	Jefferson Lab (n/d) [21]	191	216	214	213	219	219
Sσ	HERA (NC $e^{-}p$) [85]	159	315	241	240	247	244
	HERA (NC e^+p 1) [85]	402	952	580	579	588	585
	HERA (NC e^+p 2) [85]	75	177	94	94	94	93
	HERA (NC e^+p 3) [85]	259	311	249	249	248	248
	HERA (NC e^+p 4) [85]	209	352	228	228	228	228
	HERA (CC e^-p) [85]	42	42	48	48	45	49
	HERA (CC e^+p) [85]	39	53	50	50	51	51
Drell-Yan	E866 (pp) [29]	121	148	139	139	145	143
	E866 (pd) [29]	129	202	145	143	158	157
harge asymmetry	CDF (e) [88]	11	11	12	12	13	14
	$DØ(\mu)$ [17]	10	18	20	19	29	28
	DO(e) [18]	13	49	29	29	14	14
	CDF (W) [89]	13	16	16	16	14	14
	DØ (W) [19]	14	35	14	15	82	_
apidity	CDF (Z) [90]	28	108	27	27	26	26
1 2	DØ(Z)[91]	28	26	16	16	16	16
et	CDF (run 2) [92]	72	29	15	15	23	25
	DØ (run 2) [93]	110	87	21	21	14	14
jet	DØ 1 [94]	16	16	7	7	7	7
-	DØ 2 [94]	16	34	16	16	17	17
	DØ 3 [94]	12	35	25	25	24	25
	DØ 4 [94]	12	79	13	13	13	13
l		4542	5935	4700	4702	4964	4817
al + norm			6058	4708	4710	4972	4826
datum			1.33	1.04	1.04	1.09	1.07



Accardi, Brady, WM, Owens, Sato (2016)

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- → off-shell effects are correlated with nuclear wave function, but best fit is for AV18 model
- \rightarrow gives small negative (positive) shift in d/N at low x (high x)

- Valence *d/u* ratio at high *x* of particular interest
 - → testing ground for nucleon models in $x \rightarrow 1$ limit
 - $d/u \rightarrow 1/2$ SU(6) symmetry
 - $d/u \rightarrow 0$ $S = 0 \ qq$ dominance (color-hyperfine interaction)
 - $d/u \rightarrow 1/5$

 $S_z = 0$ qq dominance (perturbative gluon exchange)

• $d/u \to 0.18 - 0.28$

DSE with qq correlations



 \rightarrow considerable uncertainty at high x from deuterium corrections

- Valence *d/u* ratio at high *x* of particular interest
 - → significant reduction of PDF errors with new JLab tagged neutron & FNAL W-asymmetry data







- → extrapolated ratio at x = 1 $d/u \rightarrow 0.09 \pm 0.03$ does not match any model...
- → experiments at JLab (MARATHON, BONUS, SoLID) will determine d/u up to $x \sim 0.8-0.85$

Previous analyses — AKP17

- Similar global analysis was performed by Alekhin, Kulagin, Petti (AKP17) — similar data sets & cuts (earlier analyses used data on heavy nuclei), similar nuclear theory...
 - → find qualitatively different behavior for off-shell function, and EMC ratio shape that resembles ratio for heavy nuclei!



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 - \rightarrow correspondingly larger n/p ratio at large x



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 - \rightarrow correspondingly larger n/p ratio at large x



 \rightarrow ... but curiously smaller d/u at $x \gtrsim 0.75$

Previous analyses — AKP17 vs. CJ15

- Which (if any) is correct?
 - → benchmarking efforts by Accardi & Alekhin/Kulagin...



Both use a lot of data, have a lot of phenomenological experience, but rely on single-fit technology, which can sometimes be problematic...

... is there a more robust analysis?



- JAM iterative, multi-step Monte Carlo
 - \rightarrow traditional functional form for distributions

$$f(x) = N x^{\alpha} (1-x)^{\beta} P(x)$$

but <u>sample large parameter space</u>



→ robust determination of <u>PDF uncertainties</u>





polynomial, neural net, ...

Analysis of data requires estimating expectation values E and variances V of "observables" O (functions of PDFs) which are functions of parameters

$$E[\mathcal{O}] = \int d^n a \,\mathcal{P}(\vec{a}|\text{data}) \,\mathcal{O}(\vec{a})$$
$$V[\mathcal{O}] = \int d^n a \,\mathcal{P}(\vec{a}|\text{data}) \left[\mathcal{O}(\vec{a}) - E[\mathcal{O}]\right]^2$$

"Bayesian master formulas"

■ Using Bayes' theorem, probability distribution \mathcal{P} given by $\mathcal{P}(\vec{a}|\text{data}) = \frac{1}{Z} \mathcal{L}(\text{data}|\vec{a}) \pi(\vec{a})$

in terms of the likelihood function ${\cal L}$ and priors π

Likelihood function

$$\mathcal{L}(\text{data}|\vec{a}) = \exp\left(-\frac{1}{2}\chi^2(\vec{a})\right)$$

is a Gaussian form in the data, with χ^2 function

$$\chi^{2}(\vec{a}) = \sum_{i} \left(\frac{\text{data}_{i} - \text{theory}_{i}(\vec{a})}{\delta(\text{data})} \right)^{2}$$

with priors $\pi(\vec{a})$ and evidence Z

$$Z = \int d^n a \, \mathcal{L}(\text{data}|\vec{a}) \, \pi(\vec{a})$$

 \rightarrow Z tests if *e.g.* an *n*-parameter fit is statistically different from (*n*+1)-parameter fit

		Step:	01 (Start)	02 (+HERA)	03 (W2 cut -> 4)	0	4 (+JLab)	(+test nuc.	. smearin	g)	05 (+	TMCs)	06 (+DY)	09(OS++)	10 (+Z)	11 (+W)	12 (+	+PCs)	13 (+Tevatron)	14 (+LHC)	15 (+offshell)	16 (PCs p≠n)	17 (+mix par.)	18 (+par. range)	19 (W2 cut -> 3)	22 (+Jets)	CJ15
		W2 cut :	10	10	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4	3.5	3.5	3.5	3.5	3.5	3	3	3
		TMCs:									GP	AOT	AOT	AOT	AOT	AOT	AOT	AOT	AOT	AOT	AOT	AOT	AOT	AOT	AOT	AOT	GP
v7 ⁻	[ahle	Power Corrections:															Mult. (p=n)	Add. (p=n)	Mult. (p=n)	Mult. (p=n)	Mult. (p=n)	Mult. (p≠n)	Mult. (p≠n)	Mult. (p≠n)	Mult. (p≠n)	Mult. (p≠n)	Mult. (p=n)
Λ-	abic	Offshell Corrections:																			Yes	Yes	Yes	Yes	Yes	Yes	Yes
		Deuteron Smearing:	Paris	Paris	Paris	Paris	AV18	CD-Bonn	WJC-1	WJC-2	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	Paris	AV18
		# Data Points/# Runs:	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	50	176	1
	BCDMS (p)	351	1.11	1.10	1.70	1.72	1.73	1.74	1.73	1.73	1.24	1.15	1.08	1.18	1.18	1.20	1.15	1.16	1.20	1.18	1.17	1.17	1.17	1.13	1.13	1.14	1.25
	BCDMS (d)	254	1.03	1.12	1.73	1.74	1.73	1.78	1.71	1.73	1.26	1.13	1.09	1.11	1.10	1.10	1.05	1.05	1.07	1.09	1.08	1.08	1.07	1.06	1.06	1.06	1.15
	NMC (p)	275	1.32	1.77	1.75	1.72	1.72	1.76	1.71	1.72	1.78	1.76	1.90	1.65	1.65	1.66	1.64	1.62	1.64	1.64	1.64	1.64	1.64	1.63	1.63	1.64	1.47
	NMC (d/p)	189	0.94	0.95	0.90	0.91	0.91	0.93	0.91	0.91	0.96	0.94	1.04	0.89	0.89	0.88	0.88	0.87	0.89	0.93	0.93	0.91	0.91	0.90	0.90	0.90	0.91
	SLAC (p)	564	0.85	1.20	1.67	1.66	1.66	1.70	1.67	1.66	1.19	1.07	1.28	0.97	0.97	0.95	0.78	0.80	0.76	0.78	0.78	0.79	0.79	0.80	0.80	0.78	0.77
	SLAC (d)	582	0.61	0.88	1.42	1.42	1.42	1.48	1.45	1.43	0.88	0.81	1.00	0.77	0.77	0.78	0.62	0.61	0.61	0.64	0.63	0.63	0.63	0.64	0.65	0.65	0.65
	HERA (NC e+p 1)	402		1.57	1.75	1.83	1.83	1.78	1.82	1.82	1.64	1.60	1.64	1.47	1.48	1.48	1.46	1.44	1.46	1.49	1.48	1.48	1.48	1.48	1.47	1.51	1.44
DIS	HERA (NC e+p 2)	75		1.20	1.23	1.25	1.24	1.21	1.24	1.24	1.21	1.21	1.22	1.50	1.15	1.14	1.14	1.13	1.14	1.12	1.12	1.13	1.13	1.11	1.11	1.11	1.25
	HERA (NC e+p 3)	259		1.02	1.04	1.05	1.06	1.04	1.06	1.06	1.03	1.02	1.03	1.00	1.01	1.01	1.01	1.00	1.00	1.00	1.00	1.00	1.00	1.00	1.00	1.02	0.96
	HERA (NC e+p 4)	209		1.11	1.13	1.15	1.15	1.13	1.15	1.15	1.11	1.10	1.10	1.09	1.09	1.09	1.09	1.09	1.09	1.09	1.10	1.09	1.09	1.09	1.09	1.09	1.09
	HERA (NC e-p)	159		1.77	2.13	2.14	2.14	2.14	2.13	2.13	1.82	1.73	1.68	1.62	1.62	1.63	1.54	1.54	1.57	1.61	1.60	1.61	1.60	1.59	1.59	1.55	1.52
	HERA (CC e-p)	39		1.48	1.44	1.32	1.30	1.59	1.30	1.30	1.66	1.62	1.29	1.33	1.32	1.34	1.32	1.32	1.46	1.63	1.63	1.65	1.65	1.48	1.46	1.53	1.28
	HERA (CC e-p)	42		1.17	1.07	1.10	1.09	1.08	1.09	1.10	1.13	1.19	1.25	1.06	1.05	1.01	1.04	1.04	1.23	1.16	1.16	1.16	1.15	1.16	1.16	1.08	1.14
	JLab (d)	136				0.66	0.66	0.58	0.68	0.67	0.55	0.59	0.60	0.57	0.57	0.57	0.75	0.68	0.73	0.89	0.91	0.95	0.94	0.93	0.92	0.89	0.90
	JLab (p)	136				1.12	1.12	1.13	1.12	1.11	0.93	0.94	0.95	0.94	0.94	0.94	1.08	1.02	1.09	1.24	1.25	1.21	1.21	1.22	1.79	1.79	1.22
	BONUS (n/d)	191				0.99	0.99	1.02	0.98	0.99	0.98	0.98	1.00	1.00	1.00	1.00	1.00	0.90	1.01	1.12	1.13	1.10	1.10	1.09	1.13	1.13	1.12
	E866 (nn)	121											1.49	1.16	1 16	1 17	1 14	1.15	1.20	1 24	1.24	1.25	1 24	1.25	1 23	1 24	1.15
DY	E866 (pd)	129											2.99	1.10	1.00	1 12	0.95	0.92	0.91	0.95	0.94	0.92	0.91	0.92	0.93	0.98	1.12
	CDE (7)	28											2.55	1.00	1 18	1 19	1 15	1.22	1 23	1 13	1 12	1 10	1.09	1 32	1 35	1.07	0.96
Z	D0 (7)	20													0.59	0.59	0.59	0.50	0.61	0.59	0.59	0.59	0.59	0.60	0.60	0.59	0.50
	CDE (W)	13													0.50	1.04	0.98	1.08	0.51	0.55	0.55	0.55	0.55	1 18	1.28	1 21	1.23
w		10														0.84	0.50	0.59	3.33	3 15	3.23	3 23	3 21	1.10	1.20	1.21	1.25
	D0 (w)	14														0.04	0.00	0.55	1 59	1.65	1 72	1.85	1.94	1.25	1.24	2 29	2.23
		11																	1.03	1.05	1.01	1.00	1.01	1.00	1.05	1.04	1.00
		10																	2.57	1.01	1.01	1.00	1.01	1.00	1.03	2.26	1.05
	DU (μ) ΑΤΙ ΑΕ (2012)	10																	3.57	4.68	4.52	4.62	4.59	3.02	2.99	5.20	2.00
	ATLAS (2012)	22																		0.56	0.56	0.59	0.59	0.49	0.50	0.45	
Lantan	ATLAS (2011)	22																		2.04	2.08	2.55	2.54	1.93	1.92	1.98	
Lepton	ATLAS (2010)	22																		1.92	1.93	1.99	1.98	1.90	1.95	1.87	
	CMS (sig)	22																		0.37	0.35	0.35	0.36	0.40	0.41	0.38	
	CMS (µ) (2011)	11																		4.10	3.92	3.65	3.69	3.97	4.02	4.11	
	CMS (e) (2011)	11																		1.37	1.37	1.29	1.30	1.59	1.65	1.59	
	CMS (e) (2010)	6																		0.74	0.73	0.71	0.71	0.58	0.56	0.62	
	CMS (µ) (2010)	6																		0.10	0.10	0.09	0.09	0.07	0.07	0.08	
	D0	110																								0.92	0.19
lots	CDF	76																								1.34	0.21
Jets	STAR MB	2																								0.02	
	STAR HT	8																								1.59	
		Total:	1.002	1.265	1.530	1.505	1.505	1.523	1.509	1.506	1.241	1.183	1.323	1.123	1.116	1.121	1.059	1.054	1.084	1.109	1.108	1.108	1.107	1.093	1.106	1.109	1.070

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 \rightarrow data randomized at each step

proton DIS data

deuteron DIS data



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→ impact of LHC data mostly on sea-quark PDFs

Monte Carlo analysis tells a different story...



- \rightarrow effect very small!
- \rightarrow sits between KP and CJ15 at small $x (\sim 0.1 0.3)$
- \rightarrow more consistent with CJ15 at large x (~0.5 0.8)

Monte Carlo analysis tells a different story...



→ fitted result fairly robust

→ reveals some tension between different data sets (e.g. SLAC vs. JLab)

Fitted deuteron EMC ratio has small, < 2% deviations from unity for x < 0.6</p>





Resulting neutron structure function is <u>smaller</u> at large x







d/u PDF ratio from global fit to all data is well constrained up to $x \sim 0.8$ (mostly by *W*-asymmetry data)



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→ DIS measurements on deuteron (& other light nuclei) may be more sensitive to nuclear physics than to d/u ratio!

Outlook

- Most reliable information on nucleon PDFs requires Monte Carlo analysis and modern Bayesian analysis tools
 - \rightarrow upcoming JAM21 global QCD analysis

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- Data on A = 3 nuclei may shed light on isospin dependence of nuclear (including off-shell) effects
 - → upcoming results from MARATHON experiment