Virtuality Based Models of the EMC Effect

Gerald A. Miller University of Washington Strikman & Frankfurt several papers RMP with Or Hen, Eli Piasetzky, Larry Weinstein RMP 89 (2017) 045002 G A Miller PRC 102, 055206 ; 2008.06524 G A Miller New stuff Will focus on 0.3 <x<0.7 Remarkable experimental progress



Higinbotham, Miller, Hen, Rith CERN Courier 53N4('13)24

Ordinary nuclear physics does not explain EMC Effect

Drell-Yan on nuclei constrains models

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Virtuality is the EMC effect

- No virtuality = no EMC effect
- Average virtuality is small, average binding energy =8 MeV per nucleon,
- EMC effect is in the fluctuations
- Virtuality connects (e,e') high x to DIS

Free nucleon



Bound nucleon



A-2

Schematic two-component nucleon model

Suppression of Point Like Configurations

Blob-like config:BLC Point-like config: PLC

PLC smaller, fewer quarks high x Medium interacts with BLC energy denominator increases **PLC Suppressed**

 $|\epsilon_M| < |\epsilon|$

Dependence of the wave function of a bound nucleon on its momentum and the EMC effect C. Ciofi degli Atti, L. L. Frankfurt, L. P. Kaptari, and M. I. Strikman Phys. Rev. C 76, 055206



Questions -is there a BLC, PLC, what is x-dependence of BLC, PLC, how do neutrons and protons differ?

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Remainder of talk: towards quantitative treatment

Flavor-Independent, Two-State BLC (Blob-like configuration)-PLC Point-like configuration Model

Free nucleon :
$$H_0 = \begin{bmatrix} E_B & V \\ V & E_P \end{bmatrix}$$
, $B = BLC$, $P = PLC$, $E_P > E_B$

Eigenstates :
$$|N\rangle = \frac{1}{\sqrt{1+\epsilon^2}}(|B\rangle + \epsilon |P\rangle), \quad |X\rangle = \frac{1}{\sqrt{1+\epsilon^2}}(-\epsilon |B\rangle + |P\rangle).$$

$$\epsilon = -\frac{2V}{\Delta} \frac{1}{1 + \sqrt{1 + \frac{4V^2}{\Delta^2}}}, \ \Delta = E_P - E_B > 0$$

Deep inelastic scattering operator = \mathcal{O}_{DIS} , acts only 1 quark:

$$q(x) = \frac{1}{1+\epsilon^2} \left(\langle B|\mathcal{O}_{\rm DIS}|B\rangle + \epsilon^2 \langle P|\mathcal{O}_{\rm DIS}|P\rangle \right), \quad \langle P|\mathcal{O}_{\rm DIS}|P\rangle \equiv f(x) \langle B|\mathcal{O}_{\rm DIS}|B\rangle,$$
$$q(x) = \frac{1}{1+\epsilon^2} \langle B|\mathcal{O}_{\rm DIS}|B\rangle (1+\epsilon^2 f(x)).$$

Basic idea: ϵ reduced in nuclei

CONCLUSION AND OUTLOOK Universality of Generalized Parton Distributions in Light-Front Holographic QCD

Guy F. de Téramond, Tianbo Liu, Raza Sabbir Sufian, Hans Günter Dosch, Stanley J. Brodsky, and Alexandre Deur (HLFHS

Phys. Rev. Lett. 120, 182001 – Published 4 May 2018

ABSTRACT

(https://www.altmetric.com/details.php?domain=journals.aps.org&citation_id=32401716)

The structure of generalized parton distributions is determined from light-front holographic QCD up to a universal reparametrization function w(x) which incorporates Regge behavior at small x and inclusive counting rules at $x \to 1$. A simple ansatz for w(x) that fulfills these physics constraints with a single-parameter results in precise descriptions of both the nucleon and the pion quark distribution functions in comparison with global fits. The analytic structure of the amplitudes leads to a connection with the Veneziano model and hence to a nontrivial connection with Regge theory and the hadron spectrum. w(x) is not unique.



Based on 2-component models

Holographic QCD inputs Free nucleon $q_3 = 3$ -quark system (PLC) $P \ q_4 = BLC \ B, \ F_3(Q^2) \sim \frac{1}{Q^4} F_4(Q^2) \sim \frac{1}{Q^6}$

 $u(x) = \frac{3}{2}q_3(x) + \frac{1}{2}q_4(x), \ d(x) = q_4(x),$

excellent reproduction of measured structure functions, elastic form factors obtained!



 $q_3(x)/q_4(x) = 1/(1 - w(x)) = f(x)$

 $q_p(x) = 0.5q_4(x) + 0.5q_3(x), \quad q_n(x) = 0.75q_4(x) + 0.25q_3(x), \quad \epsilon_p^2 = 1, \ \epsilon_n^2 = 1/3, \ P_B^p = 0.5, \ P^n(B) = 0.25q_4(x) + 0.25q_4(x) + 0.25q_4(x), \quad \epsilon_p^2 = 1, \ \epsilon_n^2 = 1/3, \ P_B^p = 0.5, \ P^n(B) = 0.25q_4(x) + 0.25q_4(x) + 0.25q_4(x), \quad \epsilon_p^2 = 1, \ \epsilon_n^2 = 1/3, \ P_B^p = 0.5, \ P^n(B) = 0.25q_4(x) + 0.25q_4(x) + 0.25q_4(x), \quad \epsilon_p^2 = 1, \ \epsilon_n^2 = 1/3, \ P_B^p = 0.5, \ P^n(B) = 0.25q_4(x) + 0.25q_4($

Medium effects- flavor independent intro-Nucleon in Nucleus feels interaction U

$$H_1 = \begin{bmatrix} U & 0 \\ 0 & 0 \end{bmatrix}, H = \begin{bmatrix} E_B - |U| & V \\ V & E_P \end{bmatrix},$$

Only BLC feels interaction. Energy denominator increased, PLC probability goes down Medium-modified nucleon & excited state: $|N\rangle_M$ and $|X\rangle_M$, $\Delta \to \Delta + |U|$

$$|N\rangle_M = \frac{1}{\sqrt{1+\epsilon_M^2}} (|B\rangle + \epsilon_M |P\rangle), \ |X\rangle_M = \frac{1}{\sqrt{1+\epsilon_M^2}} (-\epsilon_M |B\rangle + |P\rangle)$$

$$q_M(x) = \frac{1}{1 + \epsilon_M^2} \langle B | \mathcal{O}_{\text{DIS}} | B \rangle (1 + \epsilon_M^2 f(x)), \ \frac{q_M(x)}{q(x)} - 1 = -2|\epsilon| \frac{|U|}{\sqrt{\Delta^2 + 4V^2}} \frac{f(x) - 1}{(1 + \epsilon^2)^2}$$

Next: relate U & virtuality = V. Use Schroedinger equation: $\mathcal{V} = \frac{-2|U|}{M} \mathcal{V} \equiv (p^2 - M^2)/M^2$

$$q_M(x) = q(x) + \frac{\mathcal{V}M}{\sqrt{\Delta^2 + 4V^2}} \frac{|\epsilon|}{(1+\epsilon^2)^2} (f(x) - 1)q(x),$$

- $\mathcal{V} < 0$ determined by kinematics
- $q_M(x) < q(x)$ if f(x) > 1 expected from idea that BLC important for high x
- \bullet x-dependence of modification independent of nuclear A
- 2 parameters Δ , V for neutron and 2 for proton

$$q_{Mp}(x) = q_p(x) + 0.686 \frac{\mathcal{V}_p}{\overline{\Delta}_n} (f(x) - 1) q_p(x),$$
$$q_{Mn}(x) = q_n(x) + 0.441 \frac{\mathcal{V}_n}{\overline{\Delta}_n} (f(x) - 1) q_n(x),$$

Neutron mass=proton mass ϵ_p, ϵ_n : 1 parameter 9/10 B. Schmookler, M. Duer, A. Schmidt, O. Hen, S. Gilad, E. Piasetzky, M. Strikman, L.B. Weinstein et al.

2004.12065



s. The EMC effect for different nuclei, as observed in (a) ratios dification of SRC pairs, as described by Eq. 2. Different colors or scale on the right The open circle show Stor of the SRC The nucleus-independent (universal) behavior of the SRC nodel, is clearly observed. The error bars on the symbols show ties, both at the 1σ or 68% confidence level and the gray bands are not isoscalar corrected.

Spares follow

Quark structure of nucleon Frankfurt-Schematic Strikman two-component **PLC** BLC nucleon model: $+ \epsilon$ (\cdot, \cdot) Blob-like config:BLC gives high x Point-like config: PLC q(x)PLC does not Cioffi degli Atti '07 Free nucleon : $H_0 = \begin{bmatrix} E_B & V \\ V & E_P \end{bmatrix}, V > 0$ interact with nucleus Α $U = \langle v(\mathbf{p}, E) \rangle / 2M$ ³H e -34.59 $|N\rangle = |B\rangle + \epsilon |P\rangle, \ \epsilon = \frac{V}{E_B - E_P} < 0$ ⁴He -69.40 ^{12}C -82.28 In nucleus (M) : $H = \begin{bmatrix} E_B - |U| & V \\ V & E_P \end{bmatrix}$ 160-79.68 ^{40}Ca -84.54 56 Fe -82.44 $|N\rangle_M = |B\rangle + \epsilon_M |P\rangle, |\epsilon_M| < |\epsilon|, \text{PLC suppressed}, \epsilon_M - \epsilon > 0 \text{ amplitude effect!}$

 $|N\rangle_M - |N\rangle \propto (\epsilon_M - \epsilon) \propto U = \frac{p^2 - m^2}{2M}$ Shroedinger eq. $q_M(x) = q(x) + (\epsilon_M - \epsilon)f(x)q(x), \ \frac{df}{dx} < 0, \ x \ge 0.3$ PLC suppression $R = \frac{q_M}{q}; \ \frac{dR}{dx} = (\epsilon_M - \epsilon)\frac{df}{dx} < 0$ Reproduces EMC effect - like every model Why this model??? Large effect if $v = p^2 - m^2$ is large, it is

large values from two nucleon correlations Simula

-92.20

12

EFT: Chen et al '16

 $^{208}\mathrm{Pb}$

Implications of model

The two state model has a ground state $|N\rangle$ and an excited state $|N^*\rangle$ $|N\rangle_M = |N\rangle + (\epsilon_M - \epsilon)|N^*\rangle$

The nucleus contains excited states of the nucleon

These configurations are the origin of high x EMC ratios

Previously missing in models of the EMC effectsame model predicts some other effect

A(e,e') at x>1 shows dominance of 2N SRC $x = \frac{Q^2}{2}$ x goes from 1 to A

x=1 is **exact** kinematic limit **for all Q**² for the scattering off a free nucleon; x=2 (x=3) is **exact** kinematic limit **for all Q**² for the scattering off a A=2(A=3) system (up to <1% correction due to nuclear binding)



How/why nucleons in nuclei cluster

one pion exchange between n and p





$1 < x < \mathbf{2}$ leading term:





DIS

Hen et al 2013



- Proper treatment of known effects: binding, Fermi motion, pionic- NO nuclear modification of Internal nucleon/pion quark structure
 - Quark based- high momentum suppression implies larger confinement volume
- a bound nucleon is larger than free one- a mean field effect Cloet Thomas

b multi-nucleon clusters - beyond the mean field

One thing I learned since '85

 Nucleon/pion model is not cool Deep Inelastic scattering from nucleinucleons only free structure function



 Hugenholz van Hove theorem nuclear stability implies (in rest frame) P⁺=P⁻=M_A

average nucleon k⁺
 k⁺=M_N-8 MeV, Not much spread

F_{2A}/A~F_{2N} no EMC effect

Momentum sum rulematrix element of energy momentum tensor

Common cause of dR/dx and a₂(A): large virtuality



Given Q^2 , x, p_{\perp} U=v/(2M)4-momentum conservation determines $2\frac{p^+}{P_D^+} \equiv \alpha$ and $v = p^2 - M^2$



|U| is large v is large
can only get this from
short range correlation
large v is responsible for
both dR/dx and a₂(A)

Sees wave function at $\alpha \approx 1.2$

The word **both** had been largely missing from models of EMC effect

many models have been ad hoc. The PLC suppression model is not.

Implications for nuclear physics

- Nucleus modifies nucleon electroweak form factors
- Nucleon excited states exist in nuclei
- Medium modifications in deuteron influence extracted neutron F₂
- spectator tagging

Logic/Summary

Data

DIS-large x (e,e') Plateau large x (e,e',NN)

Interpret:

valence quark momentum decrease in A

2 baryon clusters

QCD

nucleon wf has BLC,PLC etc PLC -high x PLC suppressed

Large virtuality

Short-ranged interactions

np dominance

Logic/Summary

EMC effect and large x plateau have same cause

Data

Interpret:

QCD

DIS-large x (e,e') Plateau large x (e,e',NN)

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Logic/Summary		EMC large have	effect and x plateau same cause
Data	DIS-large x (e	e,e') Plateau large x	(e,e',NN)
Interpret:	valence quark momentum decrease in A	2 baryon clusters	
QCD	nucleon wf has BLC,PLC etc PLC -high x PLC suppressed		
	Large virtua	lity	

Short-ranged interactions

np dominance

22





















Summary of Correlations J Ryckebusch pic



Summary of Correlations J Ryckebusch pic



Two nucleon correlations



FIG. 4: The nucleon momentum distributions $n_0(k)$ (dashed line) and n(k) (solid line) plotted versus momentum in fm⁻¹ for the deuteron, ⁴He, ¹²C and ⁵⁶Fe. Figure adapted from (Ciofi degli Atti and Simula,



Final summary

- EMC effect is related to NN correlations in two theories. Mechanism: PLC suppression enhanced by correlations
- Correlations account for high x plateau seen in several experiments
- Correlations are important in nuclear shadowing, important for EIC studies of nuclear gluons

Shadowing & Anti-shadowing

Frankfurt Strikman and Guzey Physics Reports 512 (2012) 255–393 no parton saturation

 γ^*

Green

nucleons



Kowalski Lappi Venugopalan PRL 100, 022303 use CG gluon saturation; many recent papers & discussion of detailed



But nuclear wave functions enter in all approaches

THEFT START IT'S

All approaches need two-nucleon density: $\rho^{(2)}(\mathbf{r_1}, \mathbf{r_2}) \equiv \langle A | \sum_{i \neq j} \delta(\mathbf{r_1} - \mathbf{r_i}) \delta(\mathbf{r_2} - \mathbf{r_j}) | A \rangle$ Compute thickness function $T^{(2)}(b) = \int_{-\infty}^{\infty} dz_1 \int_{-\infty}^{z_1} dz_2 \,\rho^{(2)}(b_1 = b, z_1; b_2 = b, z_2)$ Engel, Carlson, Wiringa '11 Usual approximation $\rho^{(2)}(b_1 = b, z_1; b_2 = b, z_2) \approx \rho(b, z_1)\rho(b, z_2)$ $T^{(2)}(b) = \frac{1}{2} \left(\int_{-\infty}^{\infty} dz \rho(b, z) \right)^2 = \frac{1}{2} T(b)^2$ $\rho^{(2)}(b_1 = b, z_1; b_2 = b, z_2) = \rho(b, z_1)\rho(b, z_2)(1 + C(|z_1 - z_2|_{H^2}))^{0.5}$ $T^{(2)}(b) \approx T(b)^2 \frac{l_c}{R_A}, \ l_c = 2 \int_0^\infty dz \ C(z)$ 10-20% reduction depending on 2 r (fm) $l_c/2$

Shadowing effects are overestimated by significant amounts in all approaches that neglect effects of correlations

Deep Inelastic Scattering



$$x = rac{Q^2}{2P \cdot q} = rac{k^0 + k^3}{P^0 + P^3} = rac{k^+}{P^+}$$

The 1982 EMC effect involves deep inelastic scattering from nuclei

EMC= European Muon Collaboration



Implication 1 for EIC?

Why are EMC ratios independent of Q²?

- Is the medium modification for matrix elements yielding higher-twist effects same as for leading twist?
 M. Strikman
- Can EIC add by examining Q² dependence
- Large x is on the kinematic edge, but perhaps can do during a phase in which energy is ₃₀ ramped up

The EMC Effect

Cern Courier Nov. 1982







The EMC Effect

Cern Courier Nov. 1982



How does the nucleus emerge from QCD, a theory of quarks and gluons?





CERN Courier 53N4('13)24

Progress in Quark Nuclear Physics



- Proper treatment of known effects: binding, Fermi motion, pionic- NO nuclear modification of internal nucleon/pion quark structure
- Quark based- high momentum suppression implies larger confinement volume
- a bound nucleon is larger than free one- a mean field effect
- b multi-nucleon clusters beyond the mean field

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EMC – "Everyone's Model is Cool (1985)"

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