Theoretical overview - looking forward Mark Strikman, Penn State

3rd workshop on quantitative Challenges in EMC and SRC research, March 22-26, 2021 Emergence of light cone (LC) dominance at high energies

Deuteron - LC - nonrelativistic correspondence

Construction of LC density matrix, spectral and decay function

Quest for 3N SRCs

Quest for nonnucleonic degrees of freedom - have no time

2

To resolve short-range structure of nuclei on the level of nucleon/hadronic constituents one needs processes which transfer to the nucleon constituents both energy and momentum larger than the scale of the NN short range correlations $q_0 \geq 1 GeV, \vec{q} \geq 1 GeV$

much larger for SRC with momenta up to 1 GeV/c

⇒ Need to treat the processes in the relativistic domain. The price to pay is a need to treat the nucleus wave function using light-cone quantization - One cannot use (at least in a simple way) nonrelativistic description of nuclei.

 \Rightarrow High energy process develops along the light cone.



Similar to the perturbative QCD the amplitudes of the processes are expressed through the wave functions on the light cone. In the nucleus rest frame $\alpha_N = (E_N - p_{N,z})/(m_A/A)$

In the reference frame of collider (LHC,RHIC) $\alpha_N = AE_N/E_A$

Note: in general no benefit for using LC for low energy processes.

Light-cone Quantum mechanics of two nucleon system

Due to the presence of a small parameter (inelasticity of NN interactions) it makes sense to consider two nucleon approximation for the LC wave function of the deuteron.

Key point is presence of the unique matching between nonrelativistic and LC wave functions in this approximation. Proof is rather involved.

First step: include interactions which do not have two nucleon intermediate states into kernel V (like in nonrel. QM) to build a Lippman-Schwinger type (Weinberg type) equation [LF & MS 81]



Second step: Impose condition that master equation should lead to the Lorentz invariance of the on-energy-shell amplitude of NN scattering

Introduce three-vector $\vec{k} = (k_3, k_t)$ with

$$\alpha = 1 - \frac{k_3}{\sqrt{k^2 + m^2}}$$

Invariant mass of two nucleon system is

$$M_{NN}^2 = 4\frac{m^2 + k_t^2}{\alpha(2 - \alpha)} = 4m^2 + 4k^2$$

$$T(k,k_{\rm f}) = V(k,k_{\rm f}) + \int V(k,k') \frac{{\rm d}^3 k'}{4\sqrt{k'^2 + m^2}} \frac{1}{k'^2 - k_{\rm f}^2} \frac{1}{(2\pi)^3} T(k',k_{\rm f})$$

Very similar structure for the equation for the scattering amplitude in NR QM and for LC. If a NR potential leads to a good description of phase shifts the same is true for its LC analog. Hence simple approximate relation for LC and NR two nucleon wave function

$$\alpha = (\sqrt{p^2 + m^2} - p_3)/(m_D/2)$$

where p is rest frame momentum of nucleon spectator in reaction h+ D-> p +X

Highly nonlinear relation between momentum k and momentum p: backward p=3m/4 $<\!\!<\!\!> k \rightarrow \infty$

backward p=0.5 GeV—> k=0.8 GeV



would be desirable to have data from Jlab (real photon, moderate x ~.1-.2)

First application of the logic of decay function - spectator mechanism of production of fast backward nucleons - observed in high energy proton, pion, photon - nucleus interactions with a number of simple regularities. We suggested - spectator mechanism - breaking of 2N, 3N SRCs. We extracted (Phys.Lett 1977) two nucleon correlation function from analysis of

 $\gamma(p) \stackrel{12}{\sim} C \rightarrow backward p+X processes [no backward nucleons are produced in the scattering off free protons!!!]$



Dynamical quantity (ones which can be directly observe)

Nonrelativistic

Light cone

momentum distribution n(k) not observable directly

LC density matrix $\rho_A(\alpha, k_t)$

calculated for A=3 Spectral function

modeled in 2N moving in mean field model

Decay function

 $D_A(k_2, k_1, E_r) = |\langle \phi_{A-1}(k_2, ...) | \delta(H_{A-1} - E_r) a(k_1) | \psi_A \rangle|^2$ FS81 -88

Ab-initio NR calculation of double momentum distribution + ansatz 2N moving in mean field are used for modeling spectral and decay functions Numerical calculations in NR quantum mechanics confirm dominance of two nucleon correlations in the spectral functions of nuclei at k> 300 MeV/c - could be fitted by a motion of a pair in a mean field (Ciofi, Simula, Frankfurt, MS - 89-91). However numerical calculations for nuclear matter ignored three nucleon correlations - 3p3h excitations. Relativistic effects maybe important rather early as the recoil modeling does involve k^2/m_N^2 effects.



Points are numerical calculation of the spectral functions of ³He and nuclear matter - curves two nucleon approximation from CSFS 91

question/concern : removing one nucleon of SRC does not destroy interactions of second nucleon of SRC with mean field - *should suppress emission from pairs with high momenta of the pair.*

Additional Ansatz - LC implementation of motion of the pair in the mean field

symmetry in LC NN fraction around a_{NN}=2

Phenomenologically such model works well for A(e, e'p), A(e,e'pp),... - several talks

further open questions:

- ***** with what is accuracy WF of pn pair $\propto \Psi^2_D(\kappa)$; FSIs Boeglin talk
- Need observables sensitive to LC dynamics study scattering off polarized deuteron (S/D ratio) or studying variation of scalar/tensor ratio for different angles and same momentum









small component in coordinate space generates dominant contribution in momentum space

Three nucleon SRCs = three nearby nucleons with large relative momenta

Since NN interaction is sufficiently singular for large momenta

 $\rho_A^N(\alpha, p_t)$ can be expanded over contributions of j-nucleon correlations $\rho_j(\alpha, p_t)$ $\rho_A^N(\alpha > 1.3, p_t) = \sum_{j=2}^{A} a_j(A)\rho_j(\alpha, p_t)$ FS 79

iterations of NN interactions (Plus 3N from 3N forces possible)



 $\rho_j(\alpha, p_t)(j-\alpha)^{n(j-1)+j-2}$, where $\rho_j(\alpha, 0) \propto (2-\alpha)^n$

α up to 2 (3) are allowed for 2N (3N) SRC (plus small mean field corrections)

NR case large k = 2N SRC, qualitative difference relativistic and nonrelativistic dynamics

Evidence from NR calculations? 3N SRC can be seen in the structure of decay of ³He (Sarsgian et al).



Figure 8: Dependence of the decay function on the residual nuclei energy and relative angle of struck proton and recoil nucleon. Figure (a) neutron is recoiling against proton, (b) proton is recoiling against proton. Initial momentum of the struck nucleon as well as recoil nucleom momenta is restricted to $p_{in}, p_r \geq 400 \text{ MeV/c}$.



Recoil energy dependence of the ratio of decay function calculated for the case of struck and recoil nucleons - $p_s \& p_r$ for struck proton and recoil proton and neutron for $p_s \& p_r >$ 400MeV/c & 180° > $\theta(p_s p_r) > 170°$

Some of experimental evidence in historic order

Plenty of data were described using few nucleon SRC approximation with 3N, 4N correlations dominating in certain kinematic ranges. Strength of 2N correlations is similar to the one found in (e,e'),(p,2p)





for backward emission at Ep=9 GeV

Observations of (p,2pn) &(e,e') at x>1 confirm the origin of SRC as the dominant source of the fast backward nucleons



recent analysis of (e,e') x> 2 (Day, Sargsian, LF , MS)

need larger Q²



Onset of 3N dominance at $\alpha \sim 1.6$



Further studies are necessary of LC_0^1 scaling of the ratios, etc. Recoil structure more complicated than in 2N case 10 A



Correlations in $p A \rightarrow p$ (backward) + p (backward) +X measurements of Bayukov et al 86



FIG. 1. Diagram of apparatus. (a)—Side view, (b)—view along the beam direction. Only the Z counters are shown.

 $p_i \approx 0.5 \, GeV, \alpha \approx 1.4, p_t \approx .25 \, GeV$

$$R_2 = \frac{1}{\sigma_{pA}^{in}} \frac{d\sigma(p+A \to pp+X)/d^3 p_1 d^3 p_2}{d\sigma(p+A \to p+X)/d^3 p_1 d\sigma(p+A \to p+X)/d^3 p_2}$$



the pattern of Ψ dependence of R_2 can be reproduced



Study 3N correlations in A(e,e' p +2 backward nucleons) &A(p,p' p +2 backward nucleons). Reminder: for the neutron star dynamics mostly isotriplet nn, nnn,.. SRC are relevant.



Start with ³He, followed by ⁴He, C. Expectations:

(a) $\alpha_{1 Back.Nucl} + \alpha_{2 Back.Nucl} + \alpha_{3 Forw.Nucl} \approx 3$ (b) ppn ~ nnp >> nnn, ppp (c) e+A \rightarrow e+ 2N +X stronger angular dependence and larger R₂(ψ =-180°) than in pA.

Can inverse kinematics help? Lumi

Summary - discoveries through precision and through new processes

- Precision theoretical and experimental studies of the lightest nuclei, including relativistic dynamics
- Tests of realistic modeling of FSI
- Tests of factorization (comparing electron, photon, nucleon- nucleus SRC sensitive processes.
- Tests of dynamic assumptions of LC many nucleon approximation.
- Separation of S and D waves in SRCs
- Calculating and looking for 3N SRC
- Looking for Δ-isobars and other nonnucleonic degrees of freedom in nuclei



a question - is w(k) /u(k) universal for k> 300 MeV/c?