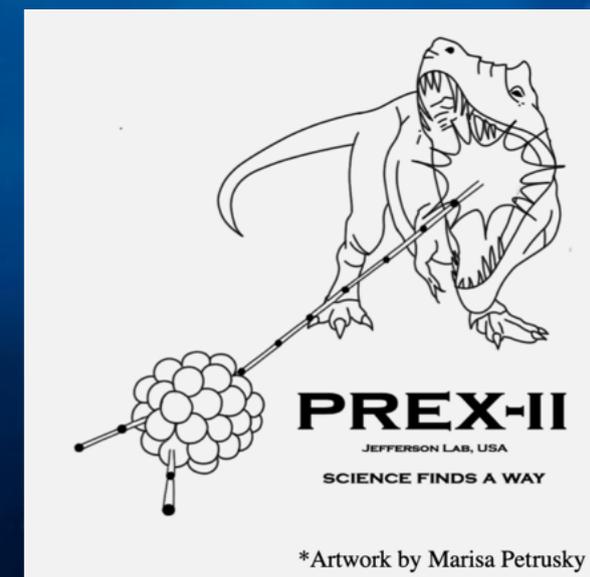
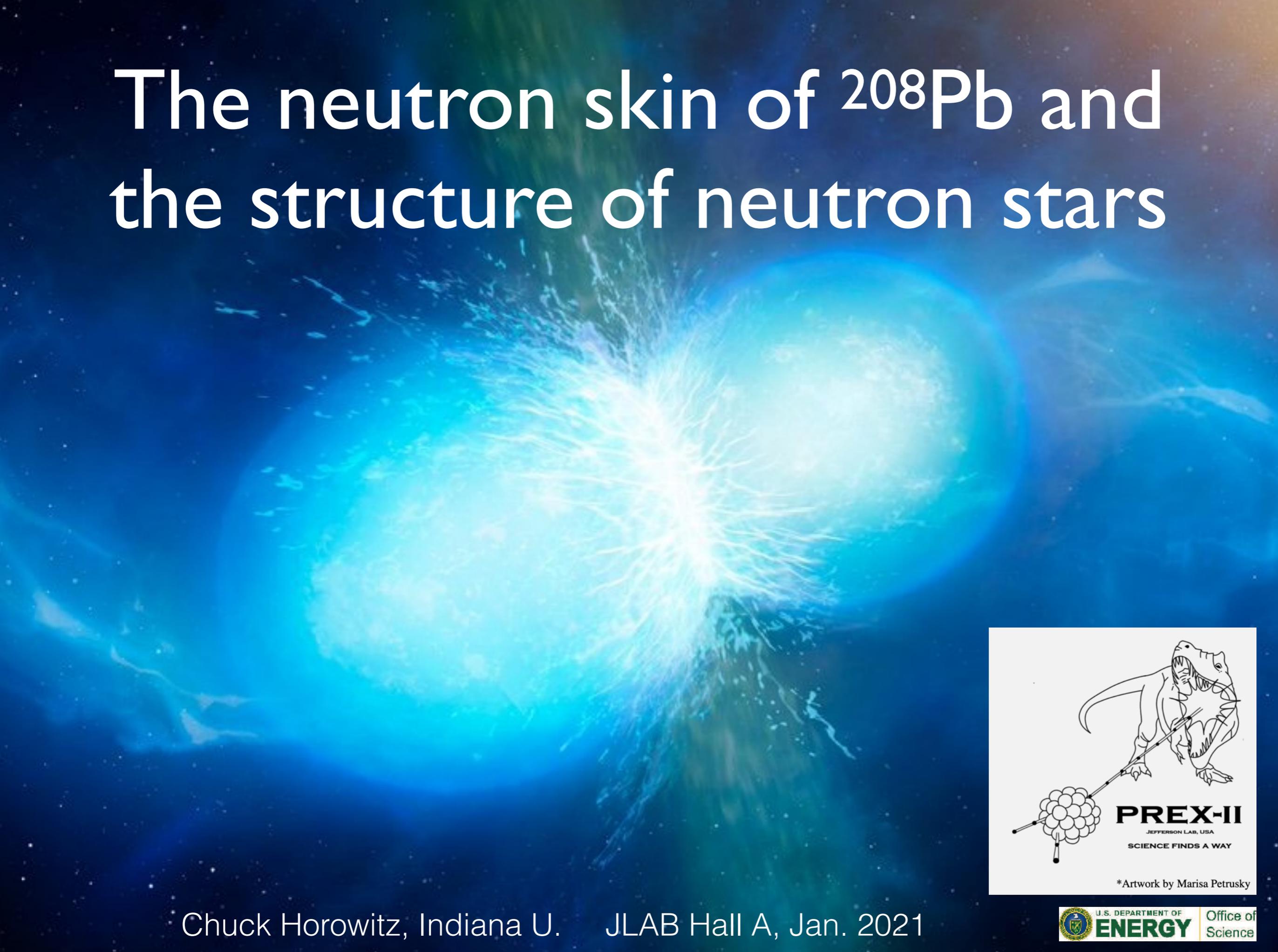


# The neutron skin of $^{208}\text{Pb}$ and the structure of neutron stars

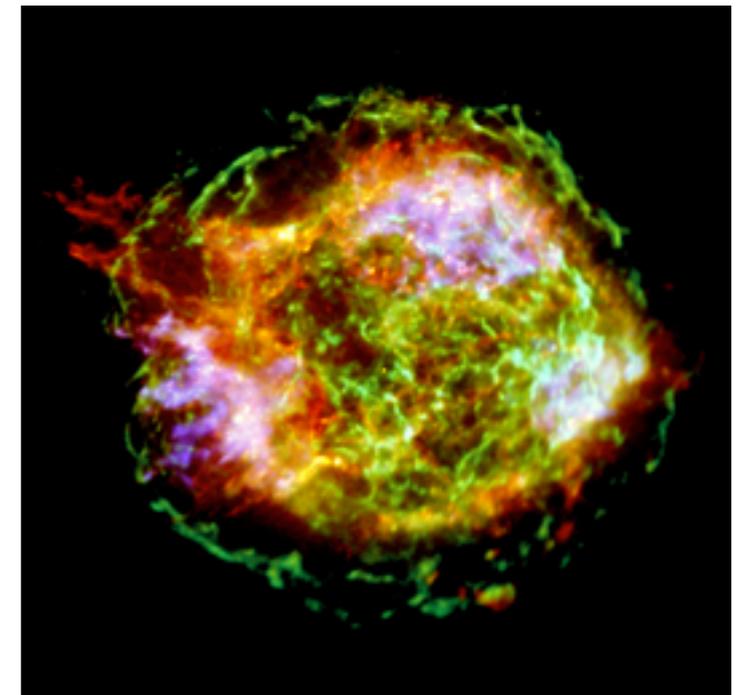


# The neutron skin of $^{208}\text{Pb}$ and the structure of neutron stars

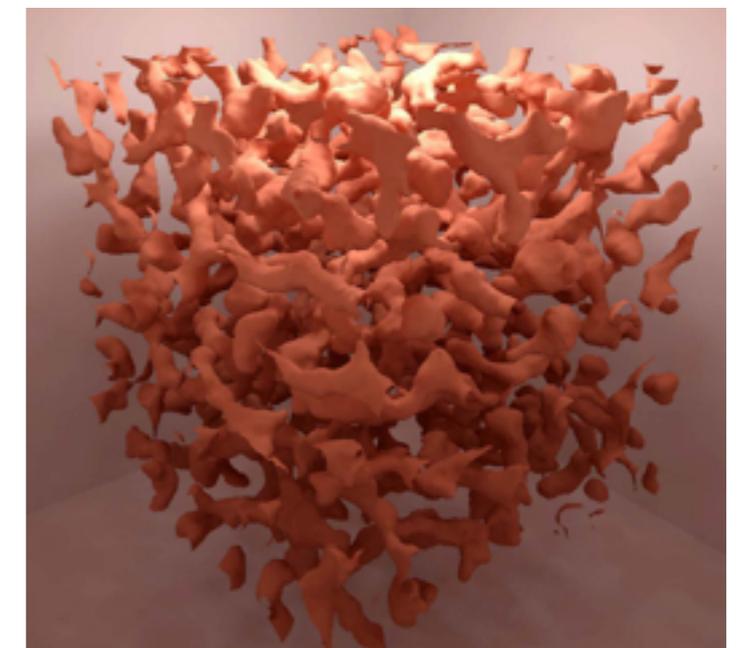
- PREX II experimental results for the neutron skin of  $^{208}\text{Pb}$  and the pressure of neutron rich matter.
- Comparison to Laser Interferometer Gravitational-Wave Observatory (LIGO) and X-ray telescope NICER results for the radius and polarizability of neutron stars.
- *Measuring nuclear density  $\rho_0$* : the interior weak and baryon densities of  $^{208}\text{Pb}$ .

# Neutron Rich Matter

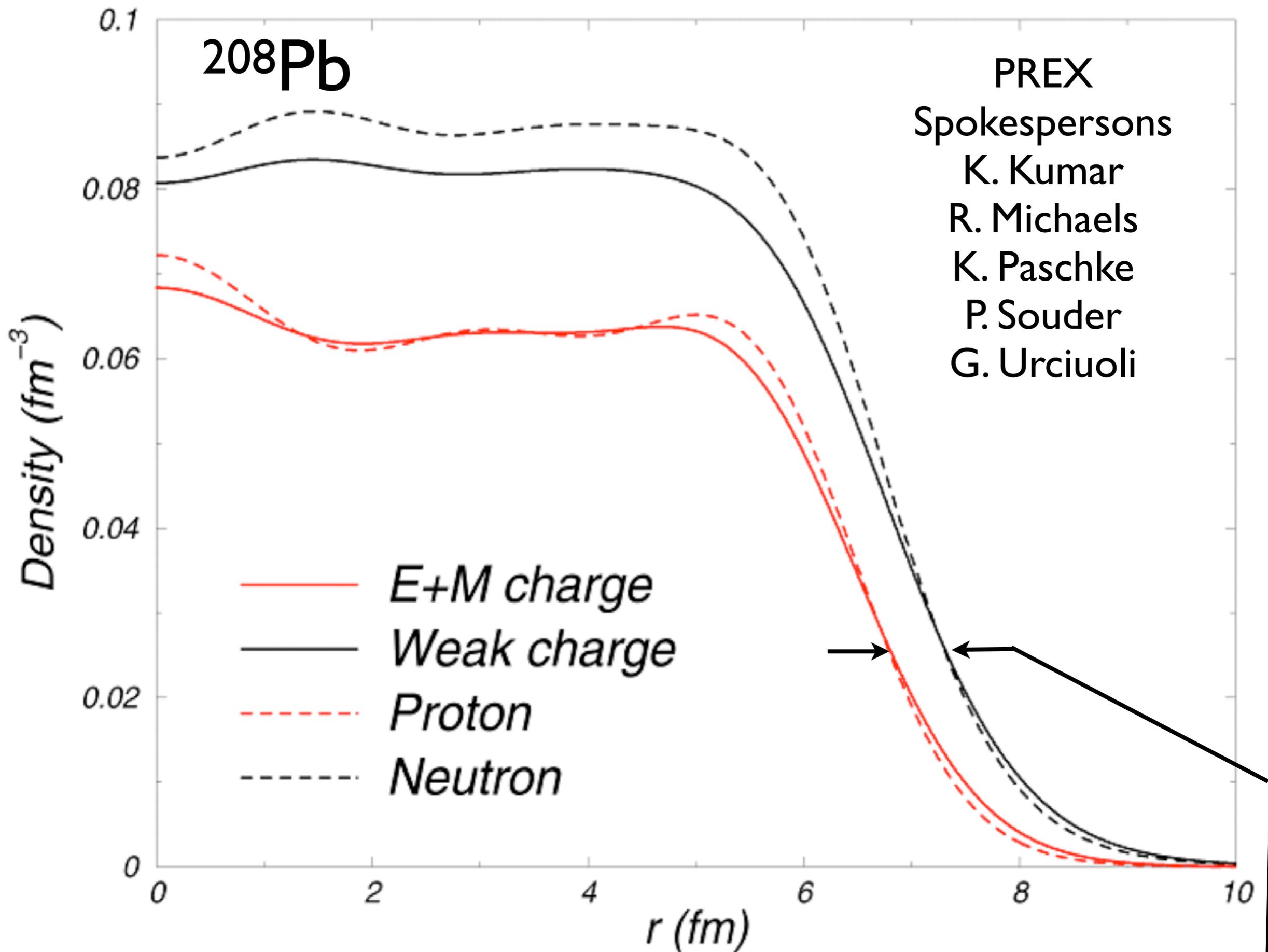
- Compress almost anything to  $10^{11}+$  g/cm<sup>3</sup> and electrons react with protons to make neutron rich matter. This material is at the heart of many fundamental questions in nuclear physics and astrophysics.
  - What are the high density phases of QCD?
  - Where did chemical elements come from?
  - What is the structure of many compact and energetic objects in the heavens, and what determines their electromagnetic, neutrino, and gravitational-wave radiations?
- Interested in neutron rich matter over a tremendous range of density and temperature where it can be a *gas, liquid, solid, plasma, liquid crystal (nuclear pasta), superconductor ( $T_c=10^{10}$  K!), superfluid, color superconductor...*
- For example a heavy nucleus such as  $^{208}\text{Pb}$  expected to have neutron rich skin.



Supernova remnant  
Cassiopea A in X-rays



MD simulation of Nuclear  
Pasta with 100,000 nucleons



- PREX measures how much neutrons stick out past protons (neutron skin).

# PREX uses Parity V. to Isolate Neutrons

- In Standard Model  $Z^0$  boson couples to the weak charge.

- Proton weak charge is small:

$$Q_W^p = 1 - 4\sin^2\Theta_W \approx 0.05$$

- Neutron weak charge is big:

$$Q_W^n = -1$$

- Weak interactions, at low  $Q^2$ , probe neutrons.

- Parity violating asymmetry  $A_{PV}$  is cross section difference for positive and negative helicity electrons

$$A_{PV} = \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R} \approx \frac{G_F Q^2 |Q_W| F_W(Q^2)}{4\pi\alpha\sqrt{2}Z F_{ch}(Q^2)}$$

- $A_{PV}$  from interference of photon and  $Z^0$  exchange.

- Determines weak form factor

$$F_W(Q^2) = \int d^3r \frac{\sin(Qr)}{Qr} \rho_W(r)$$

- Model independently map out distribution of weak charge in a nucleus.

- **Electroweak reaction free from most strong interaction uncertainties.**

# PREX in Hall A at Jefferson Lab



- **PREX**: ran in 2010. 1.05 GeV electrons elastically scattering at  $\sim 5$  deg. from  $^{208}\text{Pb}$

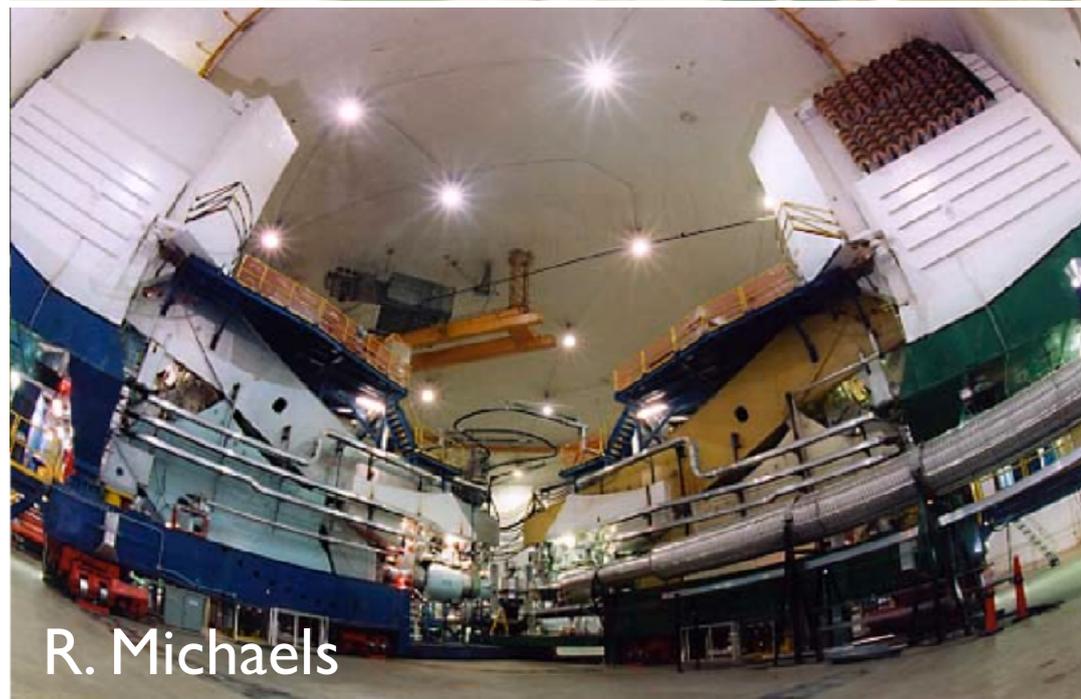
$$A_{\text{PV}} = 0.657 \pm 0.060(\text{stat}) \pm 0.014(\text{sym}) \text{ ppm}$$

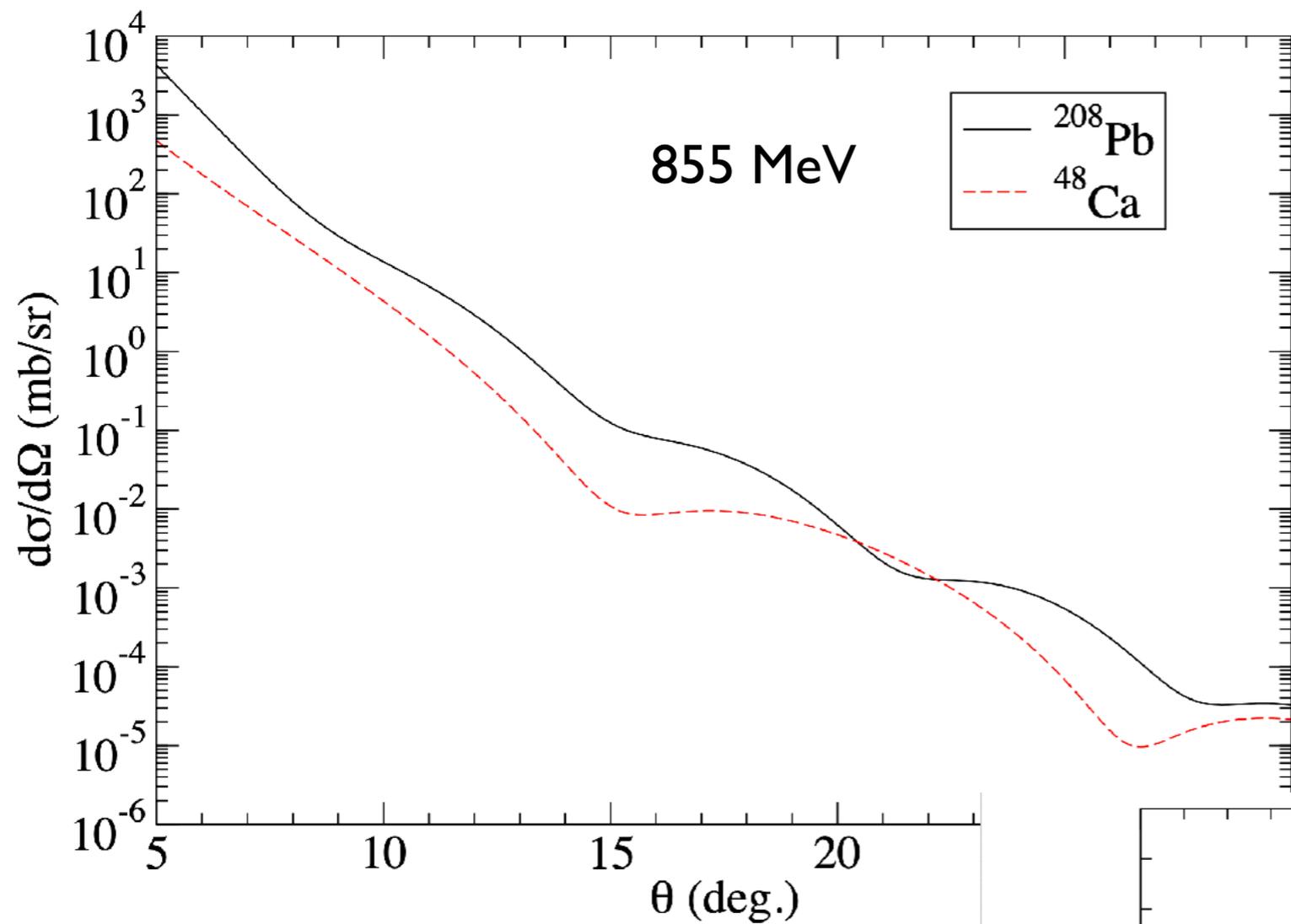
- From  $A_{\text{PV}}$  I inferred neutron skin:

$$R_n - R_p = 0.33^{+0.16}_{-0.18} \text{ fm.}$$

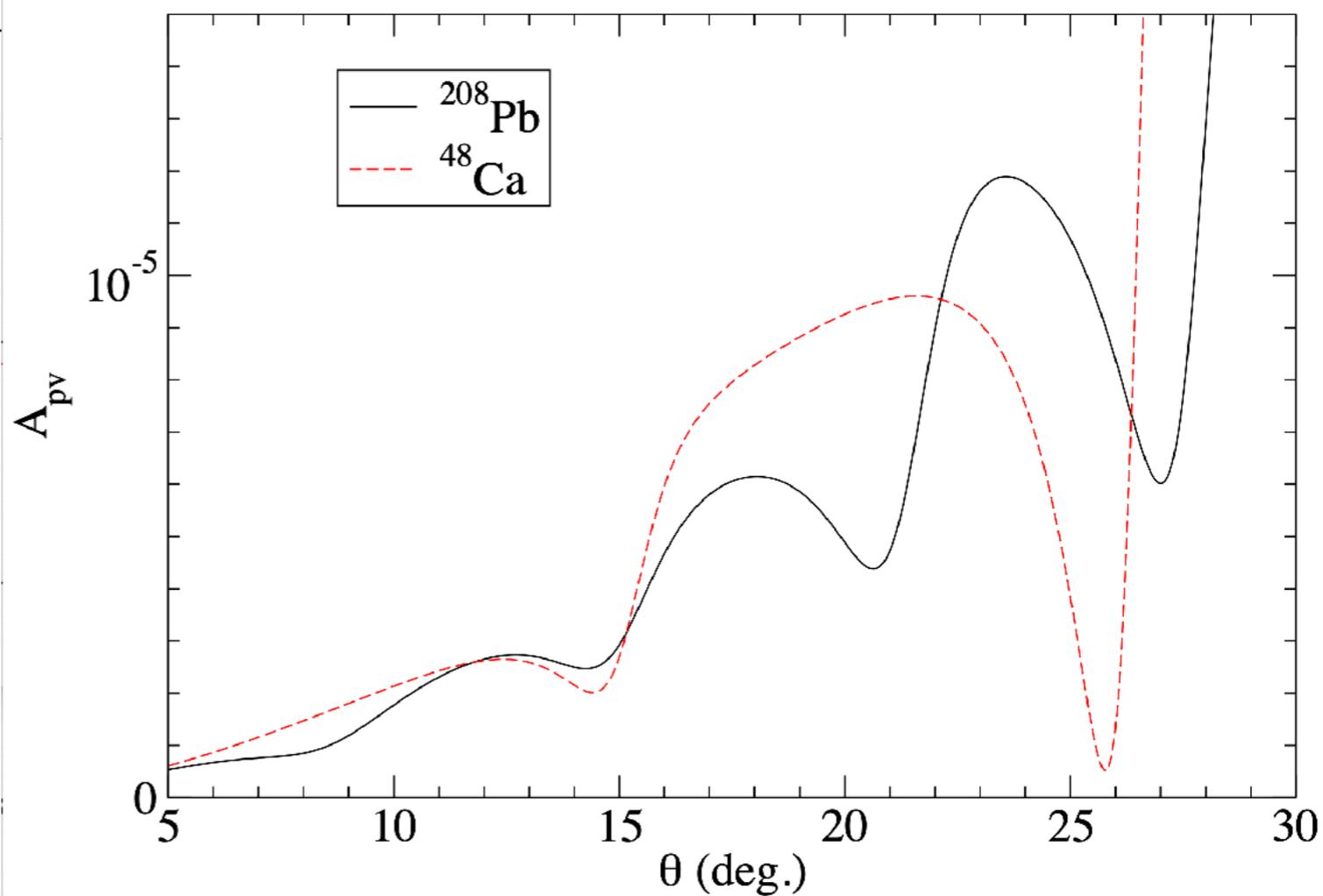
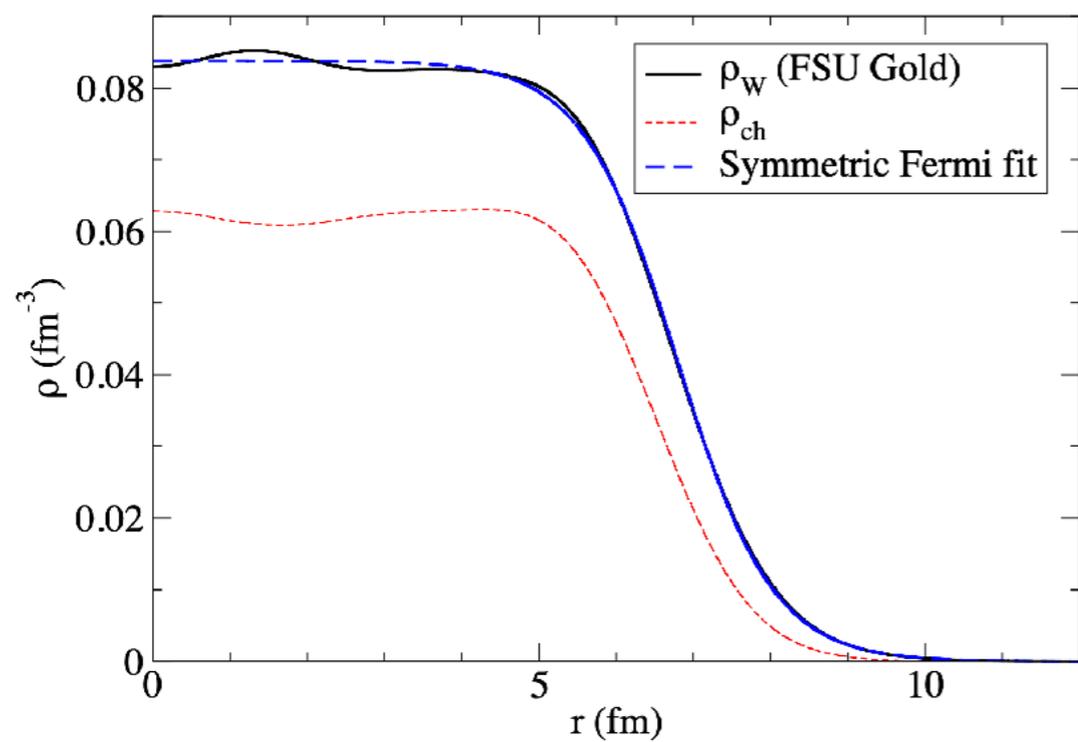
- **PREX-II**:  $^{208}\text{Pb}$  with more statistics. Ran during summer of 2019. 950 MeV electrons scattering at about 5 degrees.

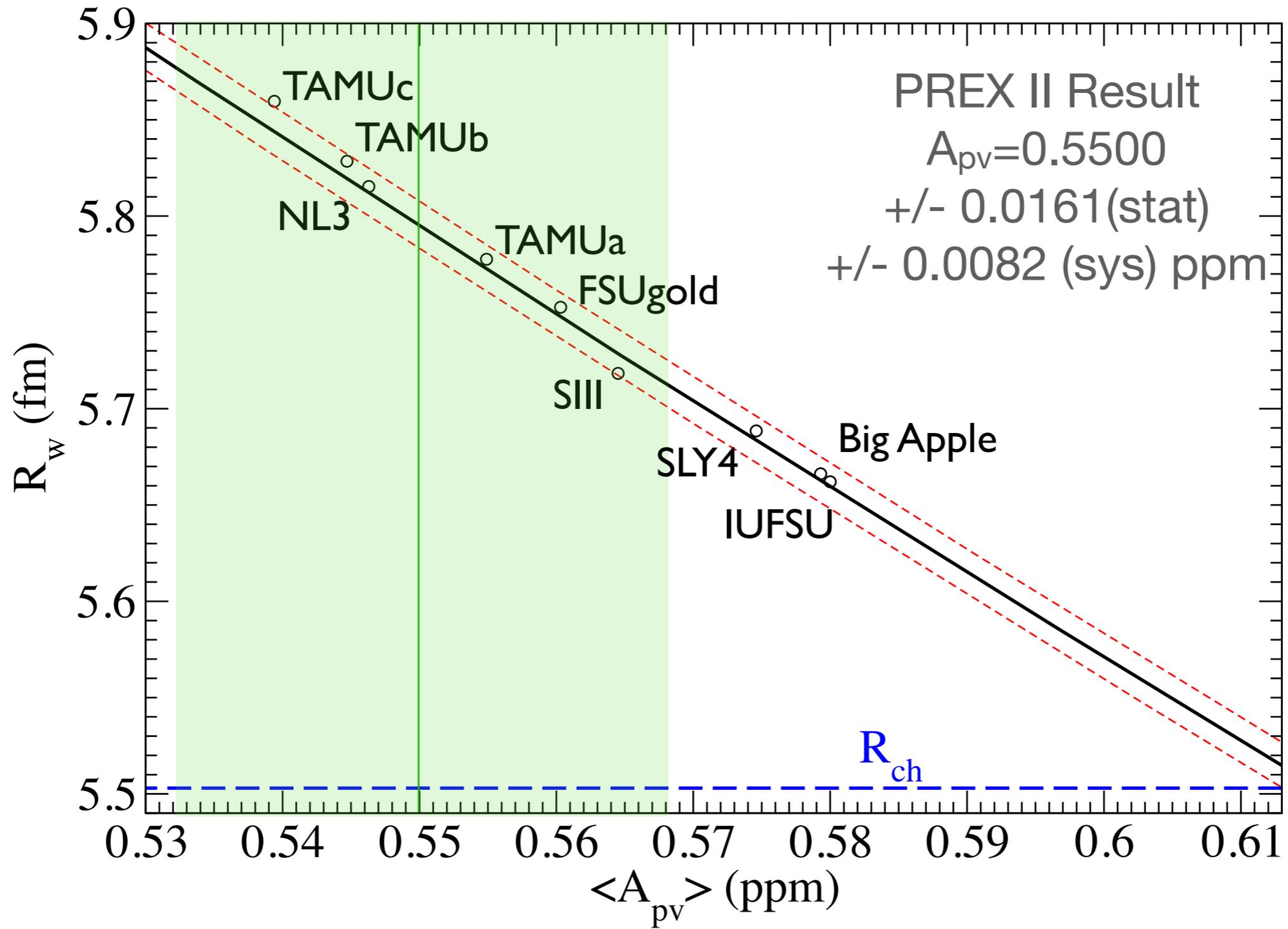
- **CREX**: Measure  $R_n$  of  $^{48}\text{Ca}$  to  $\pm 0.02$  fm. Microscopic calculations feasible for light n rich  $^{48}\text{Ca}$  to relate  $R_n$  to *three neutron forces*. Ran during Fall 2019-Spring 2020 and Summer 2020. Analysis underway, results in a year.





Numerical solution of Dirac equation for electron moving in coulomb plus small weak potentials give cross section and  $A_{pv}$ .





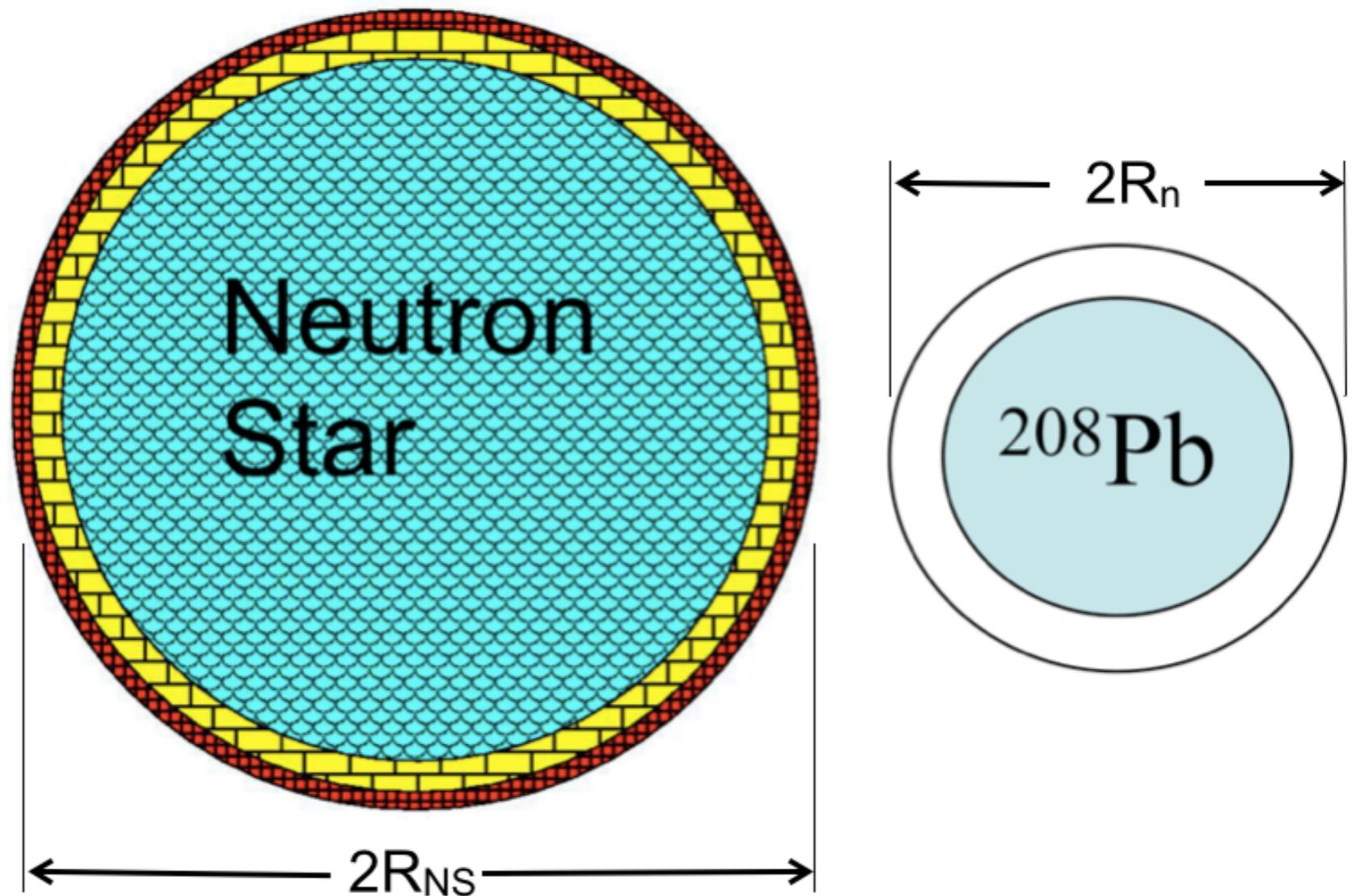
# Preliminary results

- $\langle A_{pv} \rangle = 550.0 \pm 18.1$  ppb
  - Weak charge of  $^{208}\text{Pb}$ :  $-117.9 \pm 0.3$  and numerical coulomb distortions yield  $\rightarrow$
- $R_w = 5.795 \pm 0.082$  fm [  $R_{ch} = 5.503$  fm ]
- $R_w - R_{ch} = 0.292 \pm 0.082$  fm
- $R_n - R_p = 0.278 \pm 0.078$  fm **PREX II neutron skin**
  - Model error (from uncertainty in surface thickness)  $\pm 0.013$  fm
  - Gamma-Z box diagram radiative correction uncertainty  $\pm 0.006$  fm.

Weak form factor:  $F_w(Q^2=0.00616 \text{ GeV}^2) = 0.3676 \pm 0.0125$

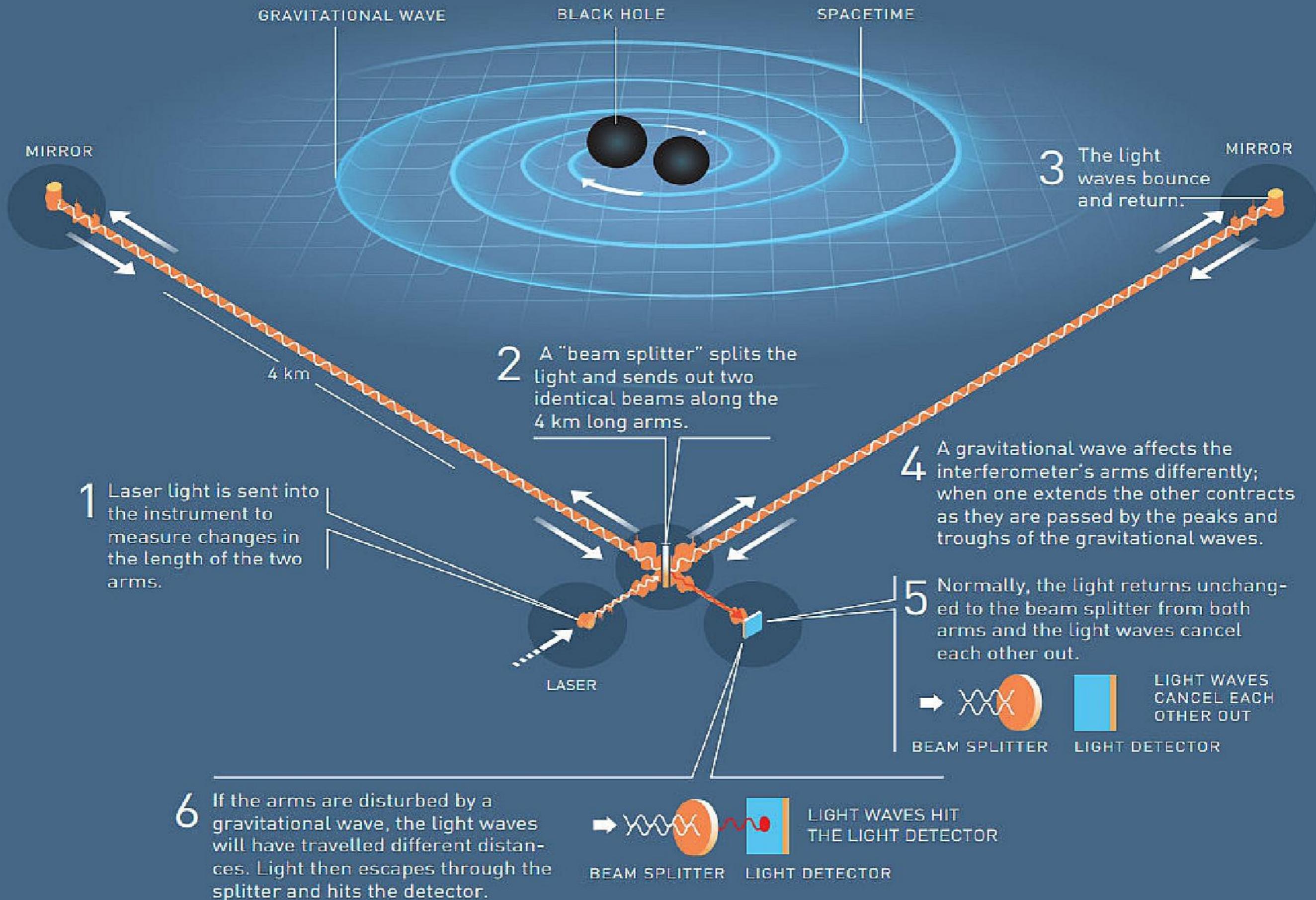
# Radii of $^{208}\text{Pb}$ and Neutron Stars

- Pressure of neutron matter pushes neutrons out against surface tension  $\implies R_n - R_p$  of  $^{208}\text{Pb}$  correlated with  $P$  of neutron matter.
- Radius of a neutron star also depends on  $P$  of neutron matter.
- Measurement of  $R_n$  ( $^{208}\text{Pb}$ ) in laboratory has important implications for the structure of neutron stars.

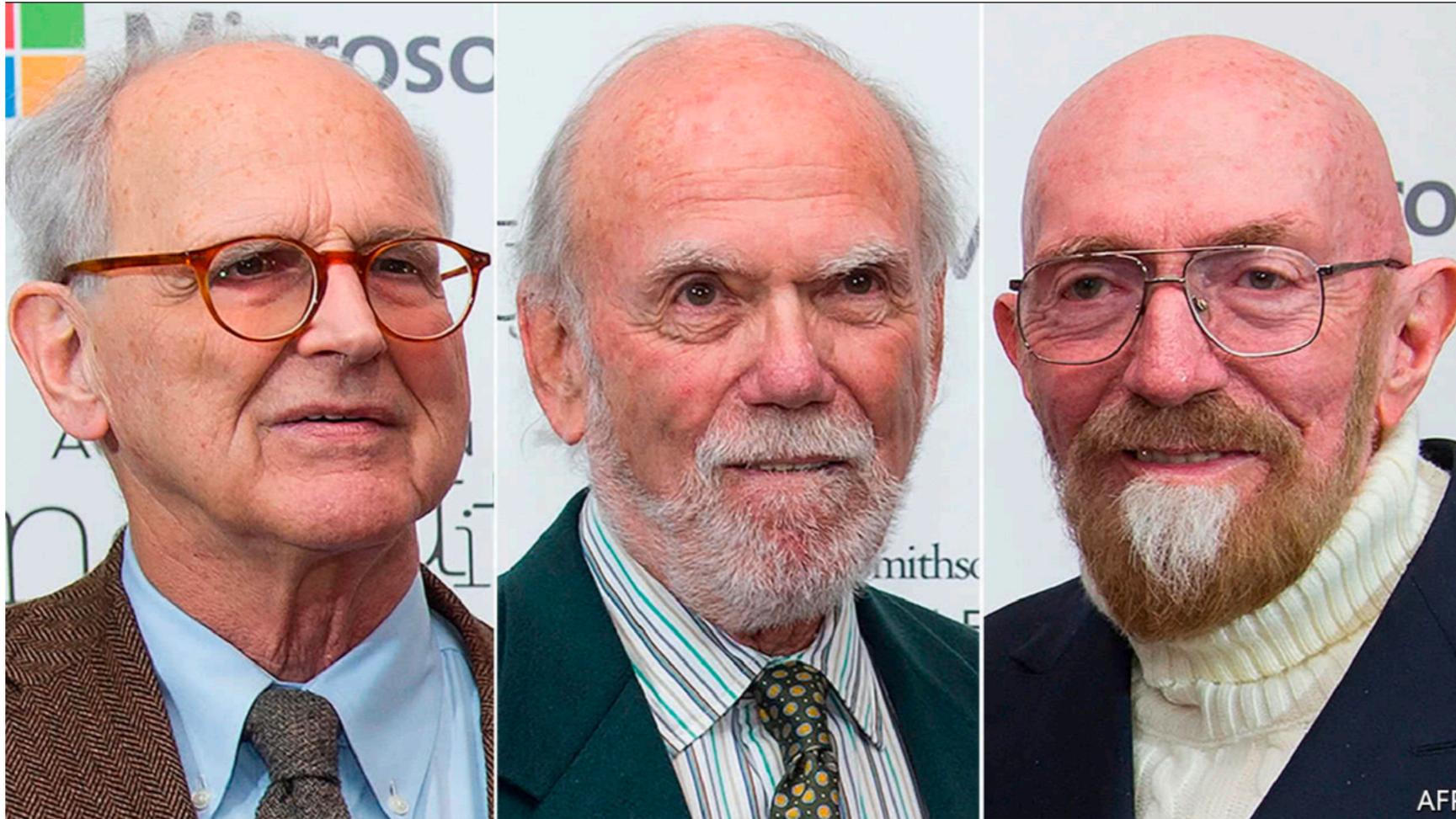


Neutron star is 18 orders of magnitude larger than Pb nucleus but has same neutrons, strong interactions, and equation of state.

# LIGO – A GIGANTIC INTERFEROMETER



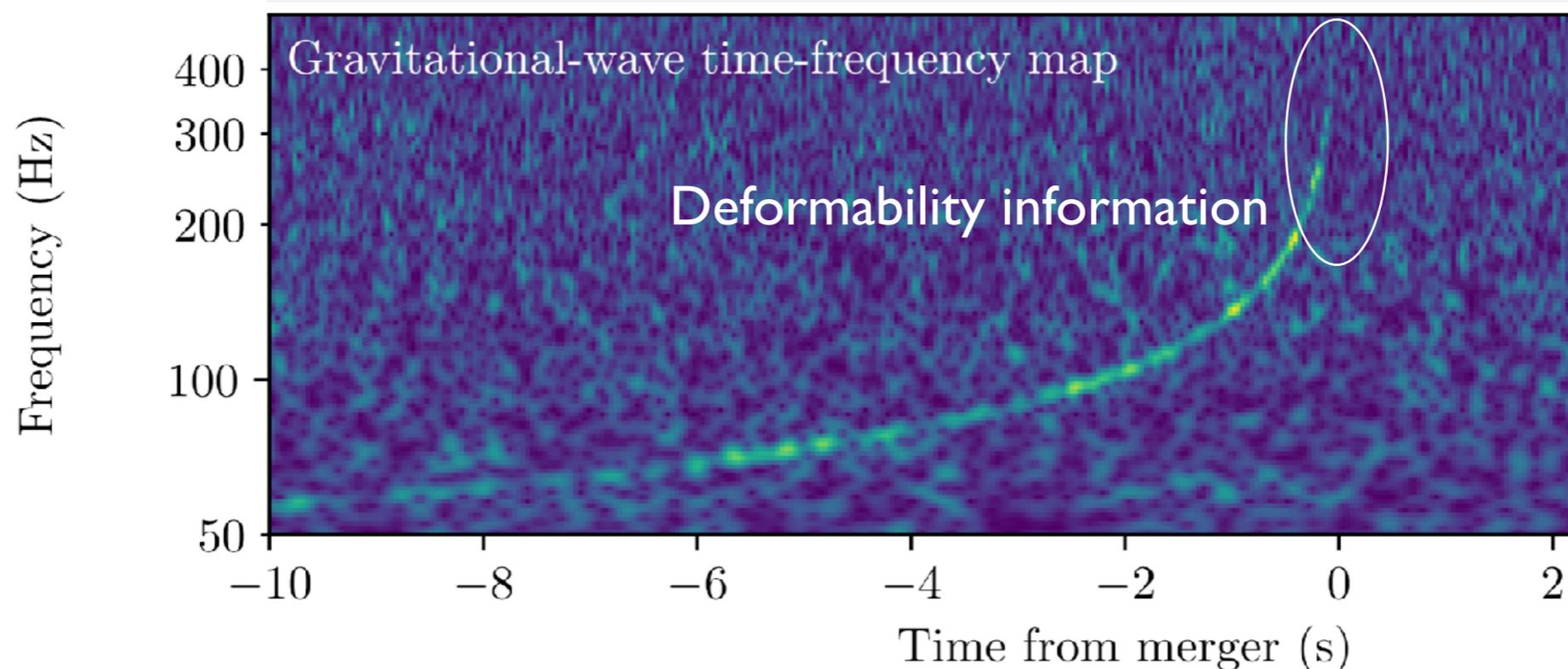
# Nobel Prize in Physics 2017



- One half to Rainer Weiss (MIT) and the other half jointly to Barry C. Barish and Kip S. Thorne (Caltech)
- “for decisive contributions to the LIGO detector and the observation of gravitational waves”

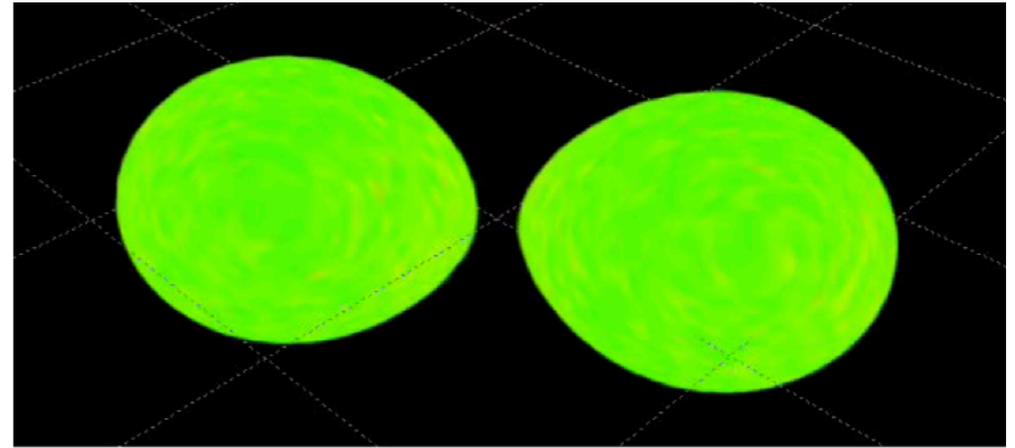
# Spectacular event GW170817

- On Aug. 17, 2017, the merger of two NS observed with GW by the LIGO and Virgo detectors.
- The Fermi and Integral spacecrafts independently detected a short gamma ray burst.
- Extensive follow up observed this event at X-ray, ultra-violet, visible, infrared, and radio wavelengths.



# Merger GW170817: deformability of NS

- Gravitational tidal field distorts shapes of neutron stars just before merger.
- Dipole polarizability of an atom  $\sim R^3$ .

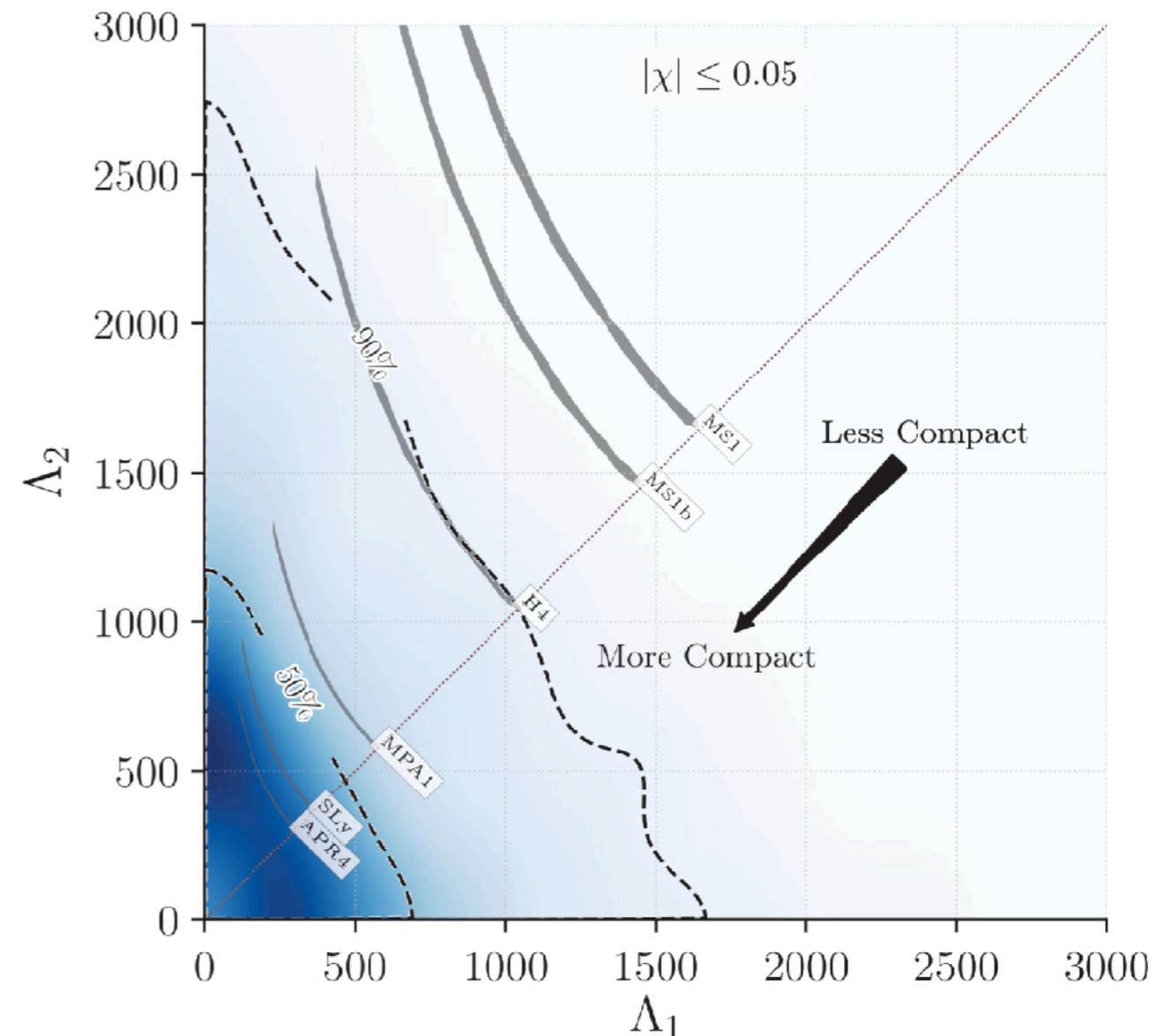


$$\kappa = \sum_f \frac{|\langle f | r Y_{10} | i \rangle|^2}{E_f - E_i} \propto R^3$$

- Tidal deformability (or quadrupole polarizability) of a neutron star scales as  $R^5$ .

$$\Lambda \propto \sum_f \frac{|\langle f | r^2 Y_{20} | i \rangle|^2}{E_f - E_i} \propto R^5$$

- GW170817 observations set upper limits on  $\Lambda_1$  and  $\Lambda_2$ .



# NICER

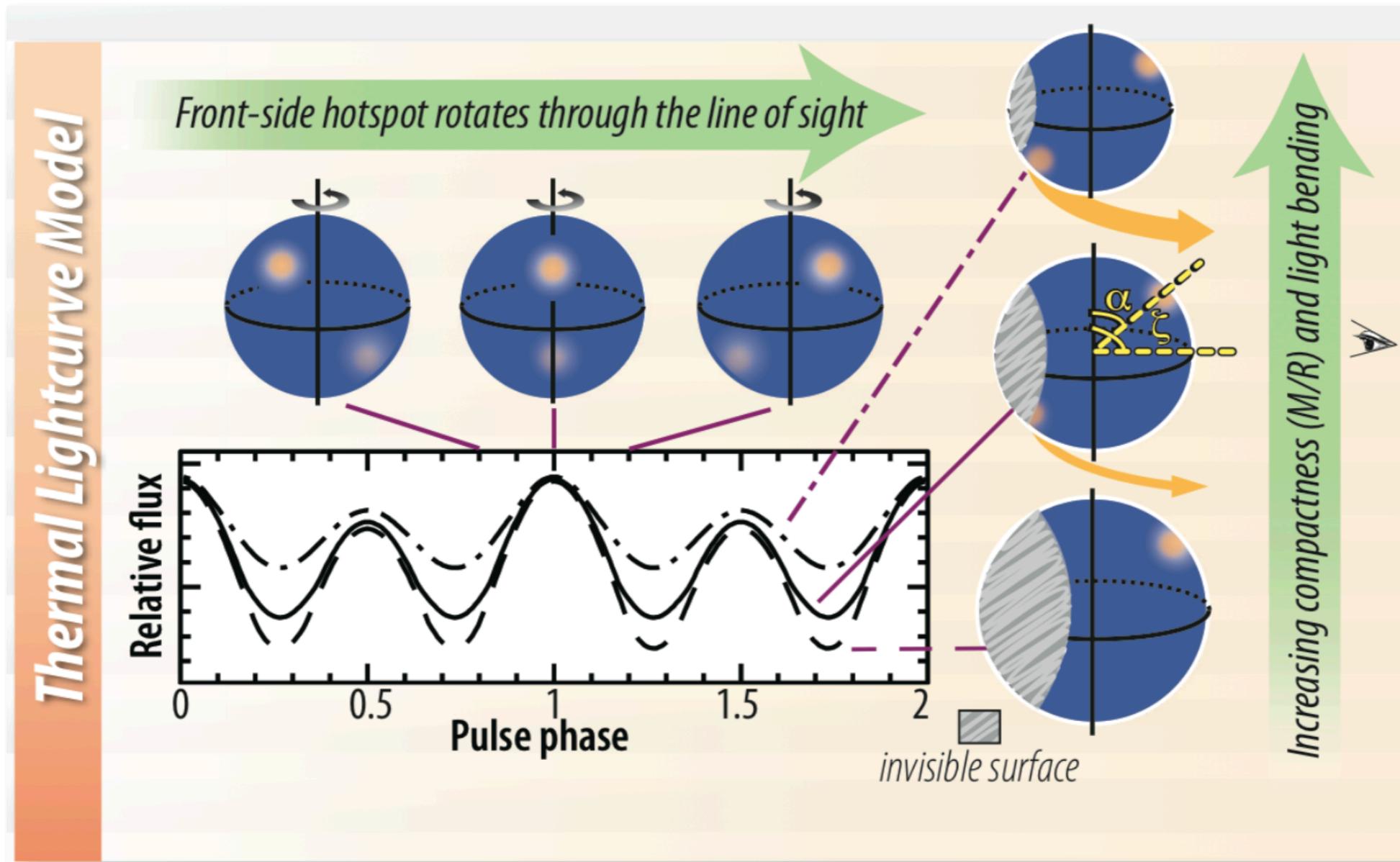
Neutron star Interior Composition Explorer

Keith Gendreau, NASA GSFC  
Principal Investigator

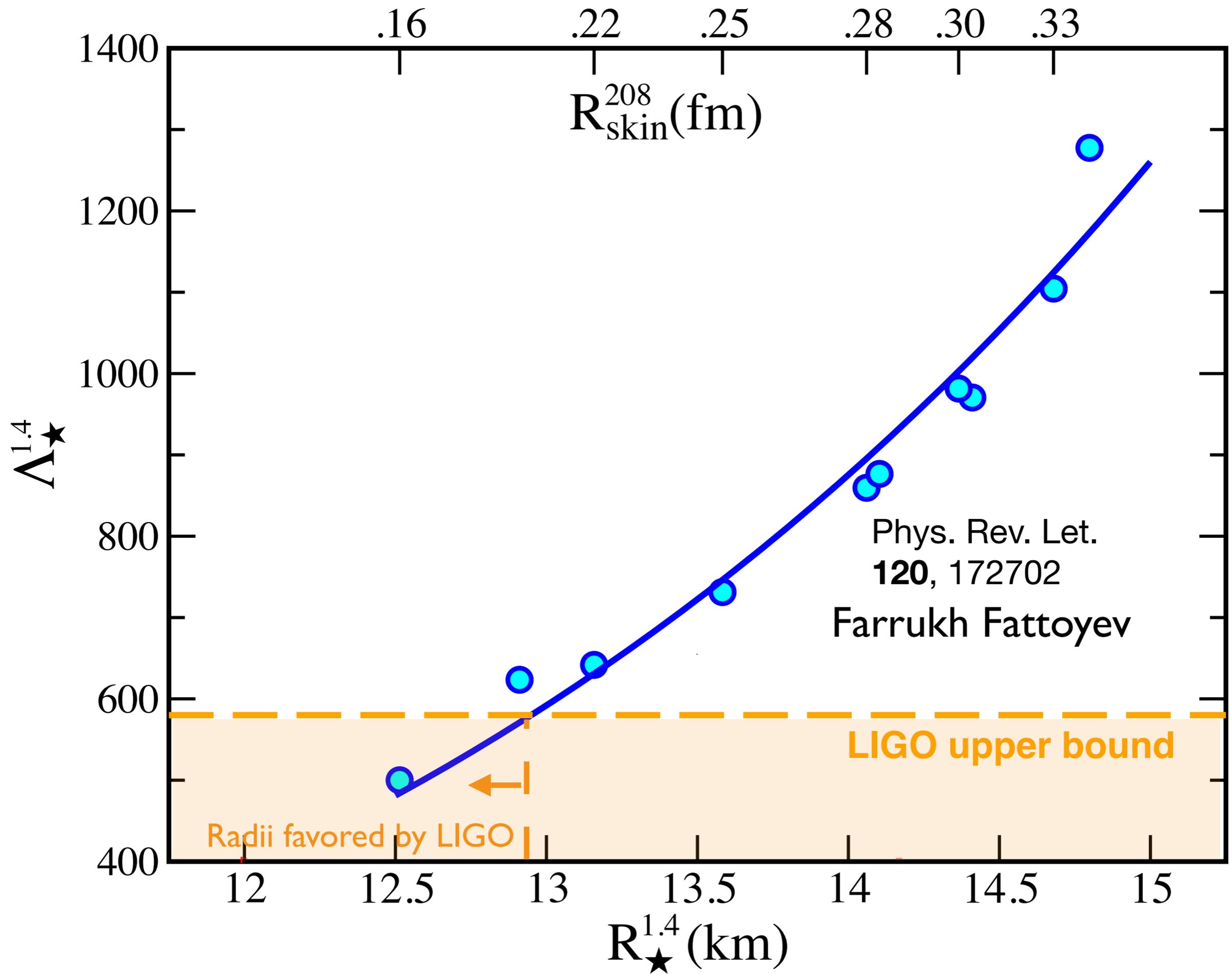


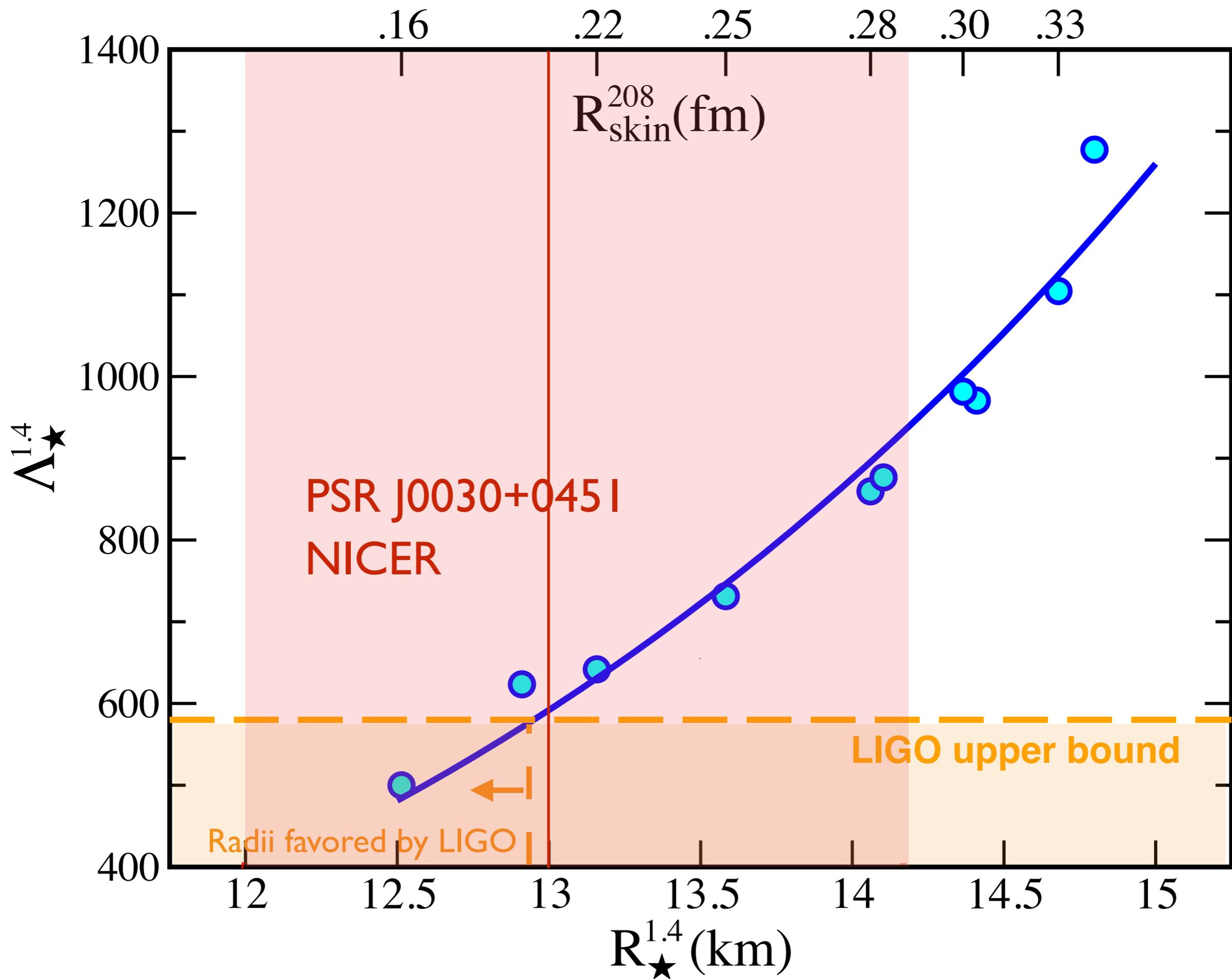
MIT KAVLI  
INSTITUTE

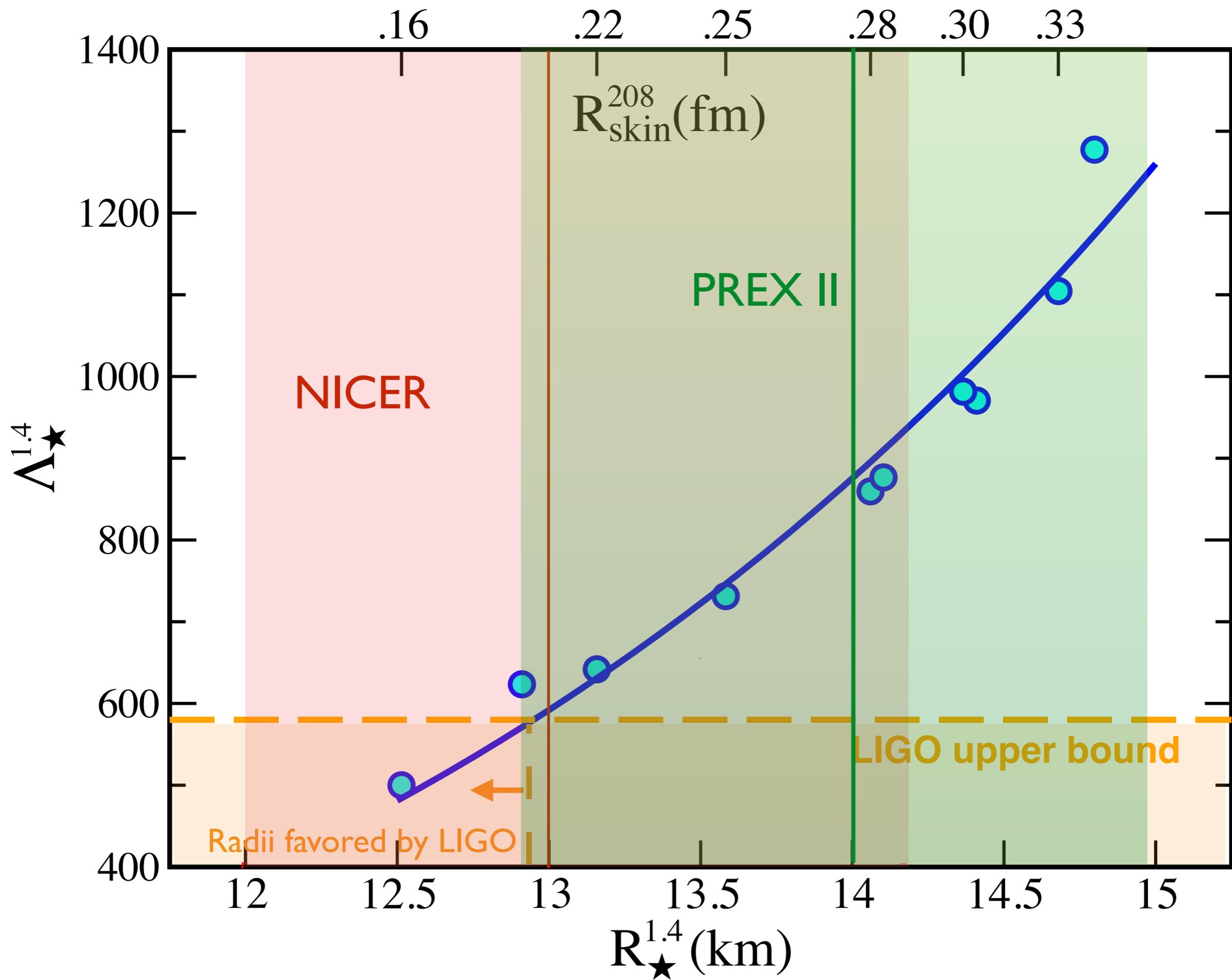
Reveal stellar structure through lightcurve modeling, long-term timing, and pulsation searches



**Lightcurve modeling** constrains the compactness ( $M/R$ ) and viewing geometry of a non-accreting millisecond pulsar through the depth of modulation and harmonic content of emission from rotating hot-spots, thanks to **gravitational light-bending**...







# Radius of $1.4M_{\text{sun}}$ Neutron Star

- LIGO: Small NS deformability suggests  $R$  is 13 km or less.
- Nicer:  $R=13 \pm \sim 1$  km
- PREX II:  $R=13$  km or more.
- Simplest solution  $\rightarrow$  Radius of NS near 13 km and neutron skin in  $^{208}\text{Pb}$  near PREX II lower limit:  $R_n - R_p \sim 0.20$  to  $0.22$  fm.

# Nuclear saturation

- In 1930s semi-empirical mass formula:  
 $BE = a_1A + a_2A^{2/3} \dots$  suggested nuclear saturation:
  - Minimum in  $E$  vs density curve for nuclear matter.
  - Baryon density  $\rho_b(r)$  approx constant in interior of heavy nucleus. *Never been directly observed.*
  - $^{40}\text{Ca}$  is largest stable  $N=Z$  nucleus, too small to have nearly constant interior density.
  - Have interior charge densities for heavier  $N > Z$  systems but not neutron or baryon densities.

*PREX II provides first clean measurement of interior baryon density of a heavy nucleus.*

# Interior weak density of $^{208}\text{Pb}$

- Two parameter (sym.) Fermi function [see PRC **102** (2020) 044321]

$$\rho_{\text{wk}}(r, c, a) = \rho_{\text{wk}}^0 \frac{\sinh(c/a)}{\cosh(r/a) + \cosh(c/a)}$$

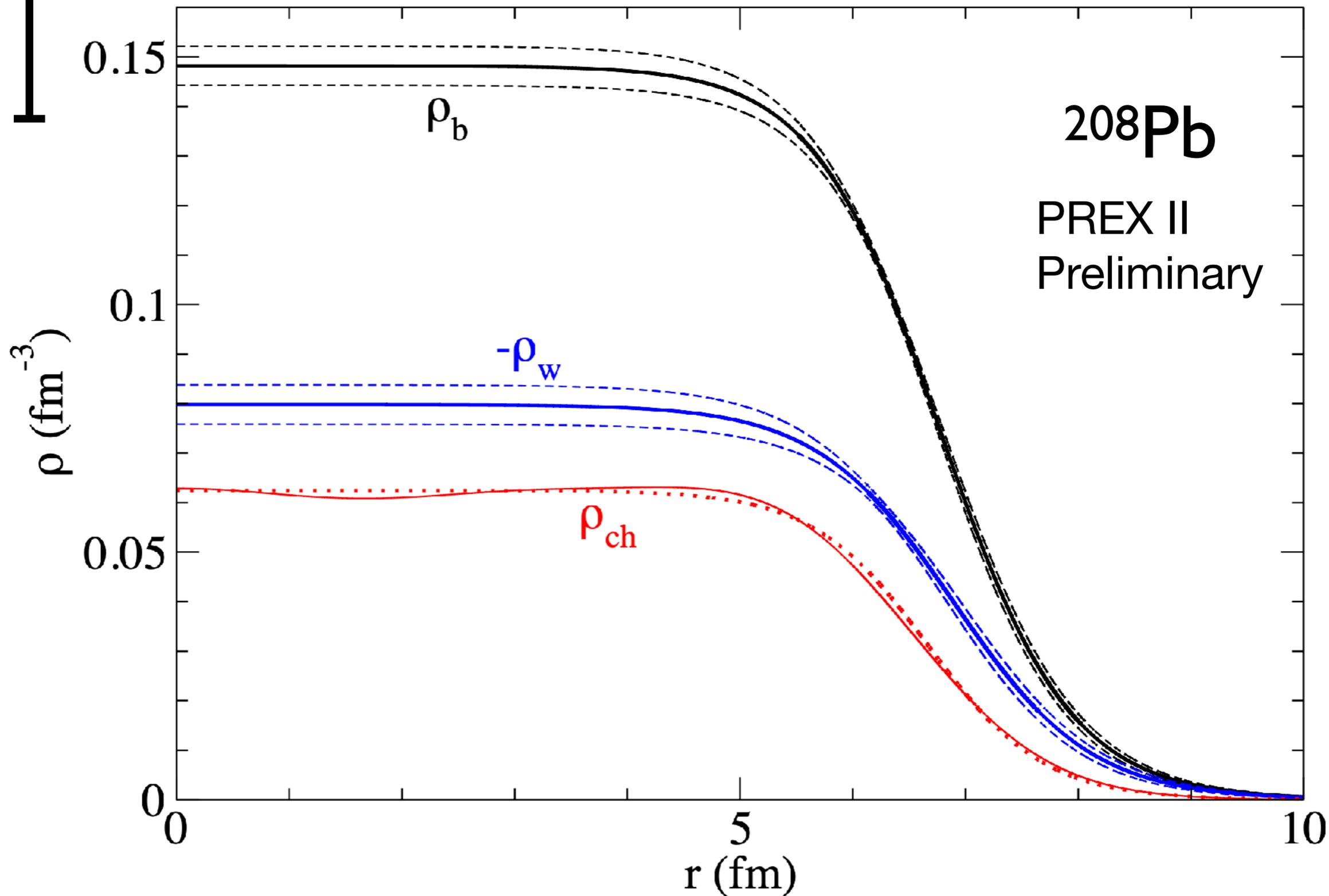
$$R_{\text{wk}}^2 = \frac{1}{Q_{\text{wk}}} \int r^2 \rho_{\text{wk}}(r) d^3r = \frac{3}{5}c^2 + \frac{7}{5}(\pi a)^2$$

$$\rho_{\text{wk}}^0 = \frac{27Q_{\text{wk}}}{4\pi(5R_{\text{wk}}^2 - 4\pi^2a^2)\sqrt{15R_{\text{wk}}^2 - 21\pi^2a^2}}$$

- Surface thickness parameter  $a$ : measure with 2<sup>nd</sup> PV exp at higher  $Q^2$ , feasible at Mainz [PRC **102**, 064308], or take from theory.
- Many nonrel. and relativistic density functionals have  $a=0.605 \pm 0.025$  fm

Drischler et al, N3LO Chiral EFT calculation of  $\rho_0$

PRC **102**, 054315

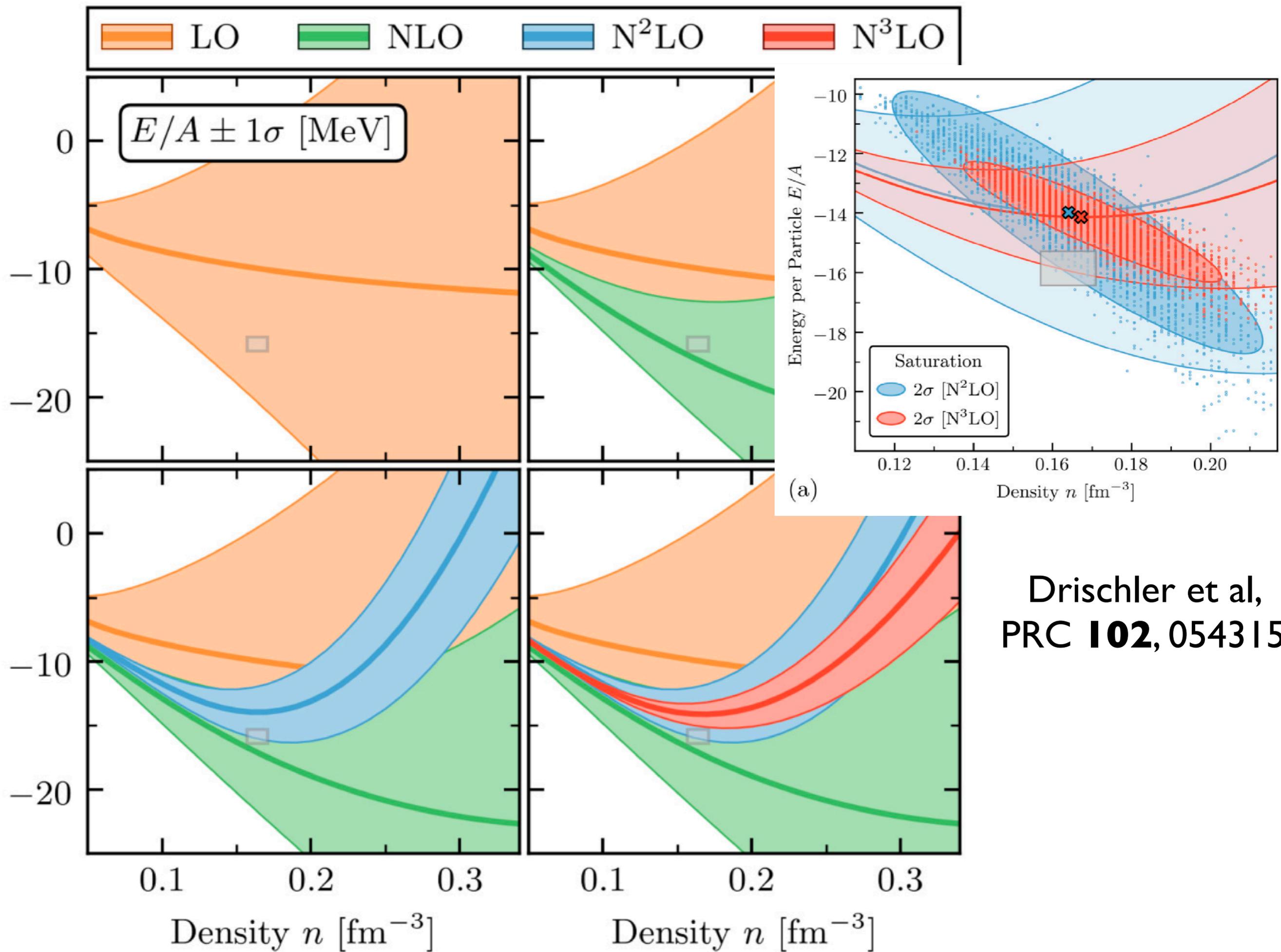


# Interior baryon density of $^{208}\text{Pb}$

- Prex II:  $R_W=5.795\pm 0.082$  fm (1.4%) implies
- $\rho_w^0 = -0.0798 \pm 0.0039(R_W) \pm 0.0013(a)$  fm $^{-3}$
- Add known charge density to get baryon density
- $\rho_b^0 = 0.1482 \pm 0.0042$  fm $^{-3}$  PREX II, preliminary
- *Fundamental nuclear structure measurement, very closely related to nuclear density  $\rho_0 \sim 0.16$  fm $^{-3}$ .*
- *Error is small only 2.8%  $\sim (1/2) 3 \Delta R_W/R_W$*

# Chiral EFT and CREX

- Chiral EFT expands 2, 3, ... nucleon interactions in powers of momentum transfer over chiral scale.
- Nucleon “hard cores” too small -> calculations with only 2 nucleon forces saturate at too high a density (up to  $2\rho_0$ )
- Three nucleon forces very important for  $\rho_0$ .
- **Three neutron forces** are hard to directly observe. They increase the pressure of neutron matter and the neutron skin thickness of both  $^{208}\text{Pb}$  and  $^{48}\text{Ca}$ .
- CREX measures neutron skin in  $^{48}\text{Ca}$ . Smaller system allows direct comparison to Chiral EFT calculations and very sensitive to 3 *neutron* forces.



Drischler et al,  
 PRC **102**, 054315

# Parity violation at Mainz

- Can measure surface thickness of  $^{208}\text{Pb}$  weak charge density at Microtron.
- At MESA (new high current low energy machine) measure:
  - Weak charge of proton (improve on  $Q_{\text{weak}}$ )
  - Weak charge of  $^{12}\text{C}$  (“Atomic PNC without the atomic structure” )
  - MREX: Neutron skin thickness of  $^{208}\text{Pb}$  (improve on PREX II by more than factor of two).

# The neutron skin of $^{208}\text{Pb}$ and the structure of neutron stars

- PREX/ CREX: K. Kumar, P. Souder, R. Michaels, K. Paschke, G. Urciuoli...
- NS deformability vs  $^{208}\text{Pb}$  skin:  
**Farrukh Fattoyev**, Jorge Piekarewicz
- Neutron skins: Concettina Sfienti
- Graduate students: **Brendan Reed**, Zidu Lin (2018), Jianchun Yin, Matt Caplan (2017)...

