ALPs @ GeV scale & their photoproduction

1st Workshop on New Light Physics and Photon-beam Experiments

DA, Yotam Soreq, Mike Williams, arXiv: 1811.03474, PRL 123 Cristiano Fanelli, same people, arXiv: 1903.03586, PRL 123

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Why photon-beam experiments?

Light

High

Energy

ALPs = axion like particles

- PNGB \to All interactions are suppressed by some cutoff Λ >> $\rm m_a$.
- We focused on

$$\mathscr{L}_{e\!f\!f\!.} = -\; \frac{4\pi\alpha_s}{\Lambda} c_g a G \tilde{G} \; + \; \frac{1}{4\Lambda} c_\gamma a F \tilde{F}$$

with $c_g = 1, c_\gamma = 0$ or $c_g = 0, c_\gamma = 1$

 $\& 0.1 \,\text{GeV} ≤ m_a ≤ 3 \,\text{GeV}.$

• Why are we interested in that?

Motivation A:

• ALPs are generic in many BSM scenarios.

• We have great experiments which can look precisely for that.

• If we can \rightarrow We should!

Motivation B: $c_g = 0, c_{\gamma} = 1$

• This is a blind spot of experiments





Motivation C:



Bauer et.al, 1708.00443

 Although a→gg is well known, what are the exclusive final states?

Coupling QCD-Scale ALPs to Gluons

1811.03474, PRL 123 (2019) 3, 031803

DA, Yotam Soreq, Mike Williams

ALPs to hadrons

- Interaction eigenstate \neq Mass eigenstate.
- Using U(3) flavor symmetry one can calculate the ALP-PS mixing

 $\langle a\pi \rangle$, $\langle a\eta \rangle$, $\langle a\eta' \rangle$.

• Using crossing-symmetry and a data driven method we developed a machinery to calculate the ALP decay rates to hadrons.

Some results



Results



Photoproduction of Axion-like particles

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Photoproduction of ALPs

JLAB photon-beam on fixed target experiments:

- (old ~ 2004) PrimEx 5.2 GeV γ @ Lab frame.
- (currently) GlueX 8.5-11.1 GeV γ @ Lab frame.

PrimEx

• Measurement of f_{π} via Primakoff production



• Our proposal => Bump hunter $a \rightarrow \gamma \gamma$.

Production Primakoff vs. Strong





• Can be distinguished using differential cross-section.

Case I: $c_{\gamma}=1$, $c_{g}=0$

• Theoretically the Nuclear Form-Factor (FF) is messy.



• To make long story short: The larger $m_{\rm N}$ the better - coherent scattering.

Case I: $c_{\gamma}=1$, $c_{g}=0$

• Main result - Forward direction:



with $\mathbf{P} = \pi, \eta$.

- Nuclear FF cancels in the ratio.
- Photon beam flux, yields, BG are all data driven.

One bin of PrimEx



Results



Case II: $c_{\gamma}=0$, $c_{g}=1$

• For $m_{\pi} \lesssim m_a \lesssim m_{\eta}$:

$$\frac{d\sigma_{\gamma N \to aN}^{Strong}}{dt} \simeq \left(\frac{f_{\pi}}{f_{a}}\right)^{2} \left[\left|\langle a\pi \rangle\right|^{2} \frac{d\sigma_{\gamma N \to \pi N}^{Strong}}{dt} + \left|\langle a\eta \rangle\right|^{2} \frac{d\sigma_{\gamma N \to \eta N}^{Strong}}{dt}\right]$$

 \bullet For strong production we need two measurements(roughly speaking ω & ρ exchange) for a data driven study.

• Those are strong π & η production.

Case II: $c_{\gamma}=0$, $c_{g}=1$

• For $m_{\pi} \leq m_a \leq m_{\eta}$:

$$\frac{d\sigma_{\gamma N \to aN}^{Strong}}{dt} \simeq \left(\frac{f_{\pi}}{f_{a}}\right)^{2} \left[\left|\langle a\pi \rangle\right|^{2} \frac{d\sigma_{\gamma N \to \pi N}^{Strong}}{dt} + \left|\langle a\eta \rangle\right|^{2} \frac{d\sigma_{\gamma N \to \eta N}^{Strong}}{dt}\right]$$

- We neglect interference. We gain more in the experimental error than lose in the theoretical error.
- For strong production BG &signal scale the same (incoherent). Proton target is preferred.

Results



Summary - QCD scale ALPS

- A new method to determine hadronic interaction strength with ALPs.
- Data-driven method for ALP photoproduction search at photon-beam experiments.
- World leading limits for ALP-gloun and ALP-photon couplings.
- Huge improvement with full data and dedicated searches.

Thank you!