

*Beam Charge Asymmetries for  
Deeply Virtual Compton Scattering  
on the proton at CLAS12*

PR12-20-009

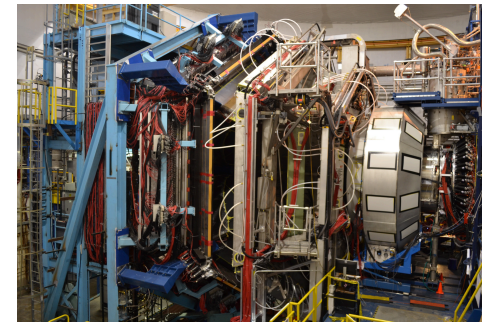
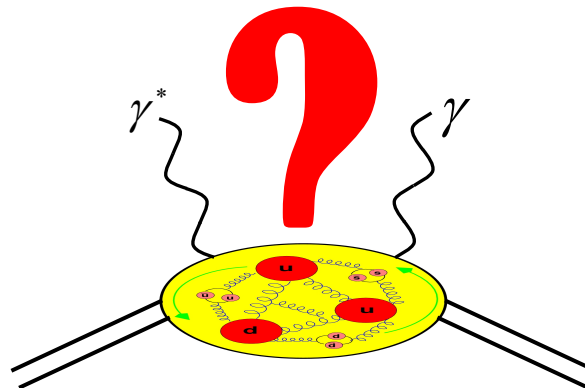
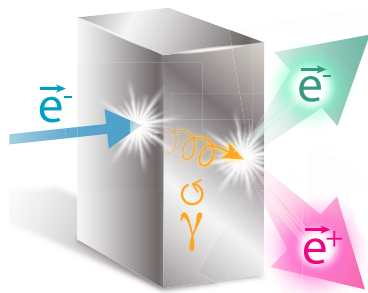
V. Burkert<sup>1</sup>, L. Elouadrhiri<sup>1</sup>, F.-X. Girod<sup>3</sup>, S. Niccolai<sup>2</sup>, E. Voutier<sup>2</sup>  
and the CLAS Collaboration

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<sup>2</sup>Université Paris Saclay, CNRS/IN2P3

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<sup>3</sup>University of Connecticut, Storrs, CT, USA



**PR12-20-009**

V. Burkert, L. Elouadrhiri, F.-X. Girod, S. Niccolai, E. Voutier et al.

« We propose to measure the unpolarized and polarized Beam Charge Asymmetries (**BCAs**) of the  $\vec{e}^{\pm}p \rightarrow e^{\pm}p\gamma$  process on unpolarized Hydrogen with **CLAS12**, using **polarized positron and electron beams** at 10.6 GeV.

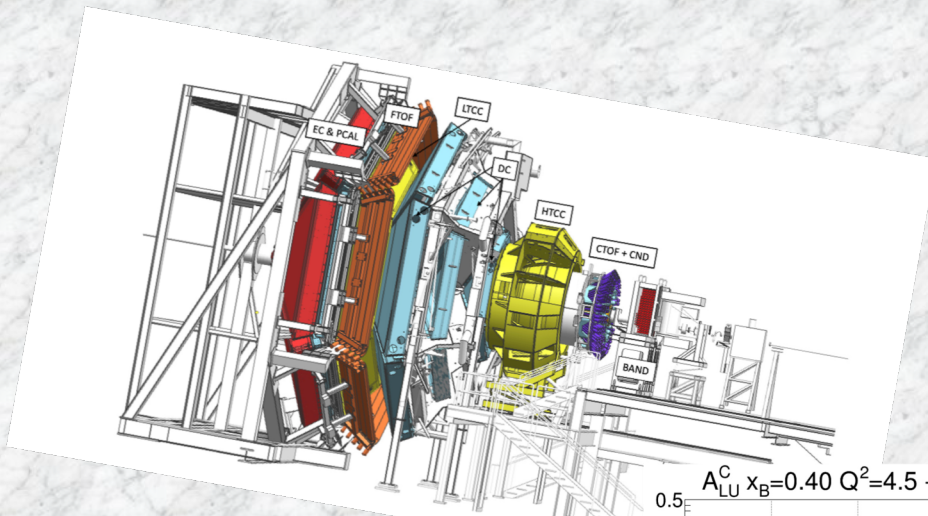
The azimuthal and  $t$ -dependences of the unpolarized and polarized BCAs will be measured over a large ( $x_B, Q^2$ ) phase space using a **2400 hours** run with a luminosity of  **$0.6 \times 10^{35} \text{ cm}^{-2}\cdot\text{s}^{-1}$** . »

☛ This proposal follows the Letter-of-Intent LOI12-18-004 discussing the perspectives of an experimental program with positron beams at JLab.

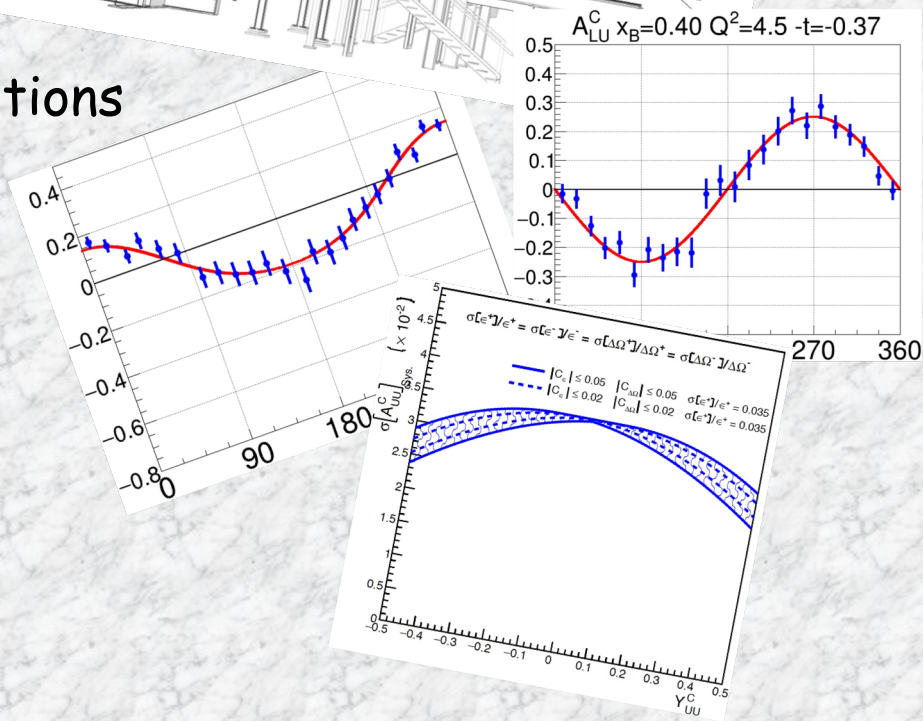
***“These measurements all have significant physics interest. The proposers should carefully evaluate feasibility and present the best case possible in a future proposal.”***

**PR12-20-009****PR12-20-012****e+@JLab White Paper**





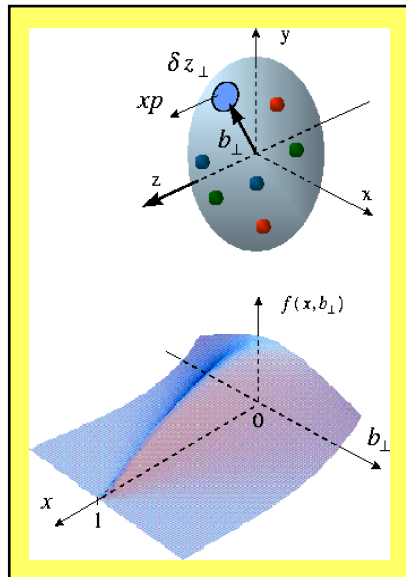
- (i) Generalized parton distributions
- (ii) Beam charge asymmetries
- (iii) Positron beams at JLab
- (iv) Experimental configuration
- (v) Experimental methodology
- (vi) Summary



## Parton Imaging

D. Müller, D. Robaschik, B. Geyer, F.M. Dittes, J. Horejsi, FP 42 (1994) 101 X. Ji, PRD 55 (1997) 7114 A. Radyushkin, PRD 56 (1997) 5524

- **GPDs parameterize** the **partonic structure** of hadrons and offer the unprecedented possibility to access the **spatial distribution** of partons.



**GPDs** encode the **correlations between partons** and contain information about the dynamics of the system like the **angular momentum** or the **distribution of the strong forces** experienced by quarks and gluons inside hadrons.

X. Ji, PRL 78 (1997) 610 M. Polyakov, PL B555 (2003) 57

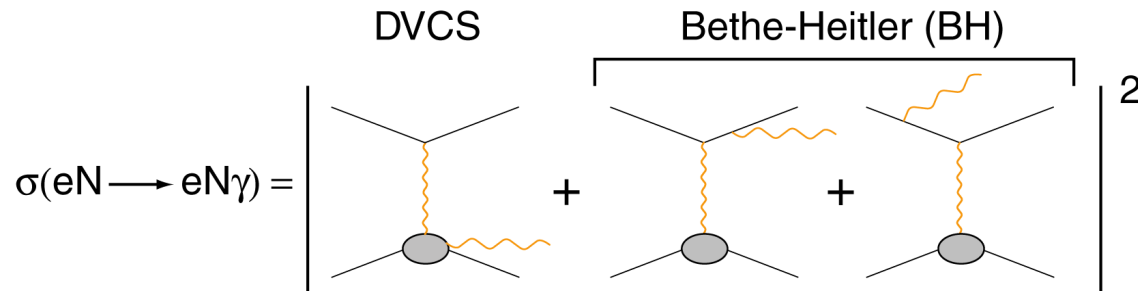
M. Burkardt, PRD 62 (2000) 071503 M. Diehl, EPJC 25 (2002) 223

**GPDs** can be interpreted as a **distribution** in the **transverse plane** of partons carrying some **fraction** of the **longitudinal momentum** of the nucleon.

A new light  
on hadron  
structure

## $\mathcal{N}(e, e'\gamma\mathcal{N})$ Differential Cross Section

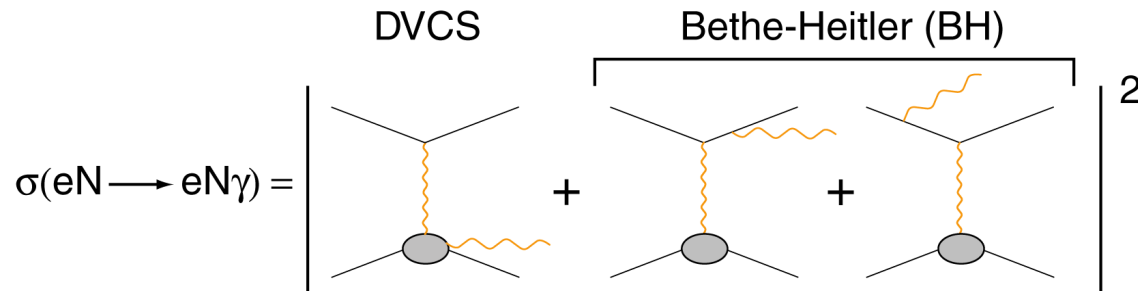
M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009



$$\sigma_{P0}^e = \sigma_{BH} + \sigma_{DVCS} + P_1 \tilde{\sigma}_{DVCS} + e_1 (\sigma_{INT} + P_1 \tilde{\sigma}_{INT})$$

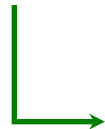
## $\mathcal{N}(e, e'\gamma\mathcal{N})$ Differential Cross Section

M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009



$$\sigma_{P0}^e = \sigma_{BH} + \sigma_{DVCS} + P_1 \tilde{\sigma}_{DVCS} + e_1 (\sigma_{INT} + P_1 \tilde{\sigma}_{INT})$$

Electron  
observables

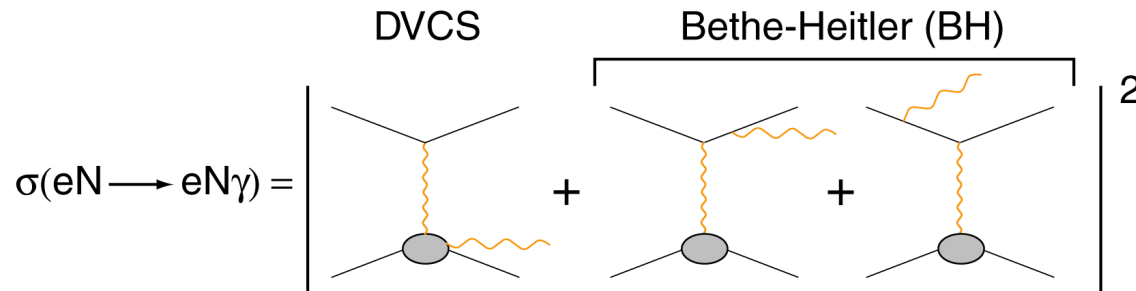


$$\sigma_{00}^- = \sigma_{BH} + \sigma_{DVCS} - \sigma_{INT}$$

$$\sigma_{+0}^- - \sigma_{-0}^- = 2\tilde{\sigma}_{DVCS} - 2\tilde{\sigma}_{INT}$$

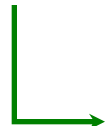
# $\mathcal{N}(e, e'\gamma\mathcal{N})$ Differential Cross Section

M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009



$$\sigma_{P0}^e = \sigma_{BH} + \sigma_{DVCS} + P_1 \tilde{\sigma}_{DVCS} + e_1 (\sigma_{INT} + P_1 \tilde{\sigma}_{INT})$$

Electron observables



$$\begin{aligned} \sigma_{00}^- &= \sigma_{BH} + \sigma_{DVCS} - \sigma_{INT} \\ \sigma_{+0}^- - \sigma_{-0}^- &= 2\tilde{\sigma}_{DVCS} - 2\tilde{\sigma}_{INT} \end{aligned}$$

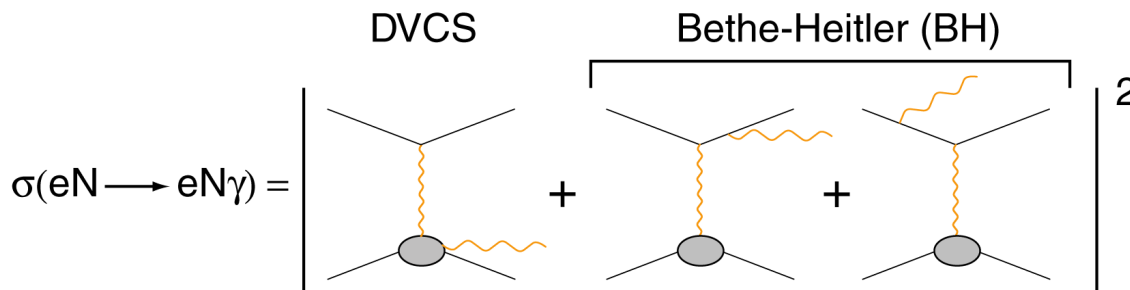
Electron & positron observables



$$\begin{aligned} \sigma_{00}^+ - \sigma_{00}^- &= 2\sigma_{INT} \\ [\sigma_{+0}^+ - \sigma_{-0}^+] - [\sigma_{+0}^- - \sigma_{-0}^-] &= [\sigma_{+0}^+ - \sigma_{+0}^-] - [\sigma_{-0}^+ - \sigma_{-0}^-] = 4\tilde{\sigma}_{INT} \end{aligned}$$

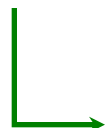
## $\mathcal{N}(e, e'\gamma\mathcal{N})$ Differential Cross Section

M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009



$$\sigma_{P0}^e = \sigma_{BH} + \sigma_{DVCS} + P_1 \tilde{\sigma}_{DVCS} + e_1 (\sigma_{INT} + P_1 \tilde{\sigma}_{INT})$$

Electron observables



$$\begin{aligned} \sigma_{00}^- &= \sigma_{BH} + \sigma_{DVCS} - \sigma_{INT} \\ \sigma_{+0}^- - \sigma_{-0}^- &= 2\tilde{\sigma}_{DVCS} - 2\tilde{\sigma}_{INT} \end{aligned}$$

Electron & positron observables



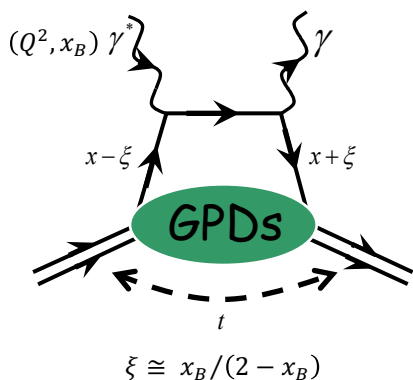
$$\begin{aligned} \sigma_{00}^+ - \sigma_{00}^- &= 2\sigma_{INT} \\ [\sigma_{+0}^+ - \sigma_{-0}^+] - [\sigma_{+0}^- - \sigma_{-0}^-] &= [\sigma_{+0}^+ - \sigma_{+0}^-] - [\sigma_{-0}^+ - \sigma_{-0}^-] = 4\tilde{\sigma}_{INT} \end{aligned}$$

**Polarized electrons and positrons** allow to **separate** the **unknown amplitudes** of the cross section for electro-production of photons.



## Compton Form Factors

- GPDs enter the  $ep\gamma$  cross section via Compton Form Factors (CFFs) representing an integral over the intermediate quark longitudinal momentum.



$$\sigma \propto \int_{-1}^{+1} dx \frac{\text{GPD}(x, \xi, t)}{x \pm \xi \mp i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{\text{GPD}(x, \xi, t)}{x \pm \xi} \pm i\pi \text{GPD}(x = \pm \xi, \xi, t)$$

↗  $\sigma_{INT, DVCS}$       ↗  $\tilde{\sigma}_{INT, DVCS}$

- At twist-2 and leading  $\alpha_{QCD}$ -order, the  $ep\gamma$  reaction accesses the four chiral even and parton helicity conserving GPDs  $\{\mathcal{H}, \tilde{\mathcal{H}}, \mathcal{E}, \tilde{\mathcal{E}}\}$  of the proton via the CFFs  $\{\mathcal{H}, \tilde{\mathcal{H}}, \mathcal{E}, \tilde{\mathcal{E}}\}$ .

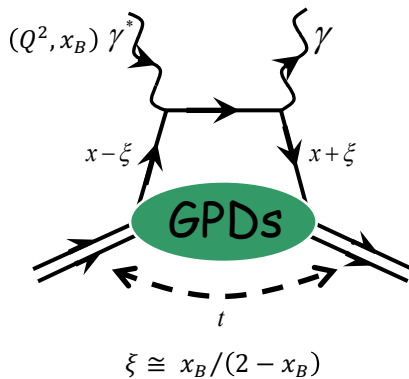
$$\mathcal{C}^{DVCS} = 4(1 - x_B)[\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*] - x_B^2[\mathcal{H}\mathcal{E}^* + \mathcal{E}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{E}}^* + \tilde{\mathcal{E}}\tilde{\mathcal{H}}^*] - \left(x_B^2 + (2 - x_B)^2 \frac{t}{4M^2}\right)\mathcal{E}\mathcal{E}^* - x_B^2 \frac{t}{4M^2}\tilde{\mathcal{E}}\tilde{\mathcal{E}}^*$$

$$\mathcal{C}^{INT} = F_1\mathcal{H} - \xi[F_1 + F_2]\tilde{\mathcal{H}} - \frac{t}{4M^2}\mathcal{E}$$



## Compton Form Factors

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$$\mathcal{C}^{\text{DVCS}} = 4(1 - x_B)[\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*] - x_B^2[\mathcal{H}\mathcal{E}^* + \mathcal{E}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{E}}^* + \tilde{\mathcal{E}}\tilde{\mathcal{H}}^*] - \left(x_B^2 + (2 - x_B)^2 \frac{t}{4M^2}\right)\mathcal{E}\mathcal{E}^* - x_B^2 \frac{t}{4M^2}\tilde{\mathcal{E}}\tilde{\mathcal{E}}^*$$

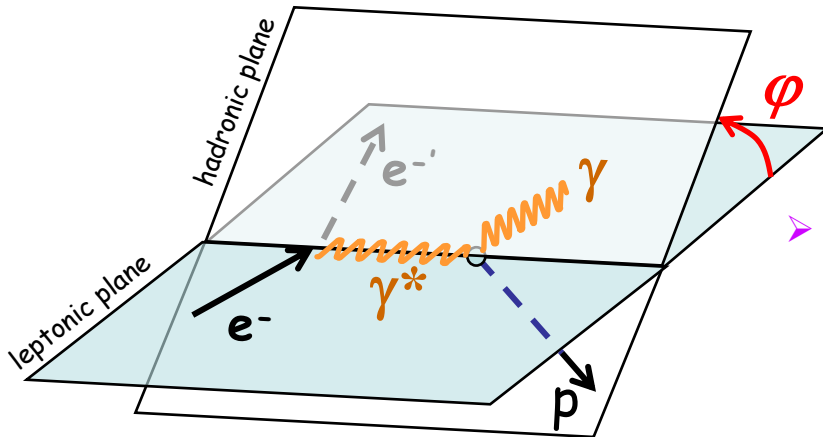
$$\mathcal{C}^{\text{INT}} = F_1\mathcal{H} - \xi[F_1 + F_2]\tilde{\mathcal{H}} - \frac{t}{4M^2}\mathcal{E}$$

Importance of the **separation** of the **DVCS** and **INT** reaction amplitudes for the **determination** of **CFFs**.

## Experimental Method

A.V. Belitsky, D. Müller, A. Kirchner, NPB 629 (2002)

$$\sigma_{P0}^e = \sigma_{BH} + \sigma_{DVCS} + P_1 \tilde{\sigma}_{DVCS} + e_1 (\sigma_{INT} + P_1 \tilde{\sigma}_{INT})$$



$$\sigma_X \equiv \frac{d^5 \sigma_X}{dQ^2 dx_B dt d\phi_e d\phi}$$

- The BH differential cross section is exactly calculable from the proton form factors ( $F_1, F_2$ ) known at small  $t$ .

$$\sigma_{BH} = \frac{1}{P_1(\varphi)P_2(\varphi)} \sum_{n=0}^2 c_n^{BH} \cos(\varphi)$$

- At twist-2 and leading  $\alpha_{QCD}$ -order, the cross section components exhibit **specific azimuthal dependences**.

$$\sigma_{DVCS} = c_0^{DVCS} \Re[C^{DVCS}]$$

$$\tilde{\sigma}_{DVCS} = 0$$

$$\sigma_{INT} = \frac{c_0^{INT} + c_1^{INT} \cos(\varphi)}{P_1(\varphi)P_2(\varphi)} \Re[C^{INT}]$$

$$\tilde{\sigma}_{INT} = \frac{s_1^{INT} \sin(\varphi)}{P_1(\varphi)P_2(\varphi)} \Im[C^{INT}]$$

## Gravitational Form Factors

V. Burkert, L. Elouadrhiri, F.-X. Girod, Nature 557 (2018) 396    M.V. Polyakov, P. Schweitzer, Int. J. Mod. Phys. A33 (2018) 1830025  
 K. Kumerički, Nature 570 (2019) E1

$$\int_{-1}^1 x H(x, \xi, t) dx = M_2(t) + \frac{4}{5} \xi^2 d_1(t)$$

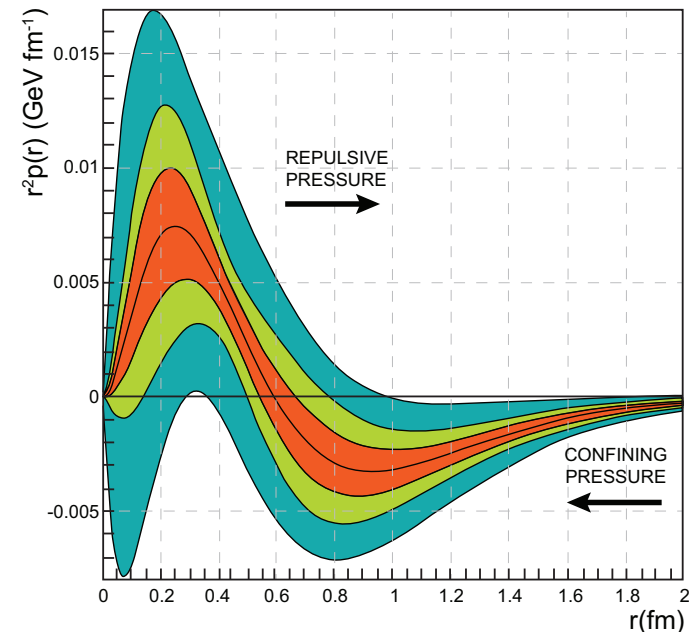
The 2<sup>nd</sup> **Mellin moment** of GPDs allows to access the pressure distribution inside hadrons through the **skewness dependency** of GPDs.

CFF 
$$\mathcal{H}(\xi, t) = \int_{-1}^1 \left[ \frac{1}{\xi - x - i\epsilon} - \frac{1}{\xi + x - i\epsilon} \right] H(x, \xi, t) dx$$

$$\Re[\mathcal{H}(\xi, t)] \stackrel{\text{LO}}{=} D(t) + \mathcal{P} \left\{ \int_{-1}^1 \left[ \frac{1}{\xi - x} - \frac{1}{\xi + x} \right] \Im[\mathcal{H}(x, t)] dx \right\}$$

$$D(t) = \frac{1}{2} \int_{-1}^1 \frac{D(z, t)}{1 - z} dz$$

$$D(z, t) = (1 - z^2) \left[ d_1(t) C_1^{3/2}(z) + \dots \right]$$



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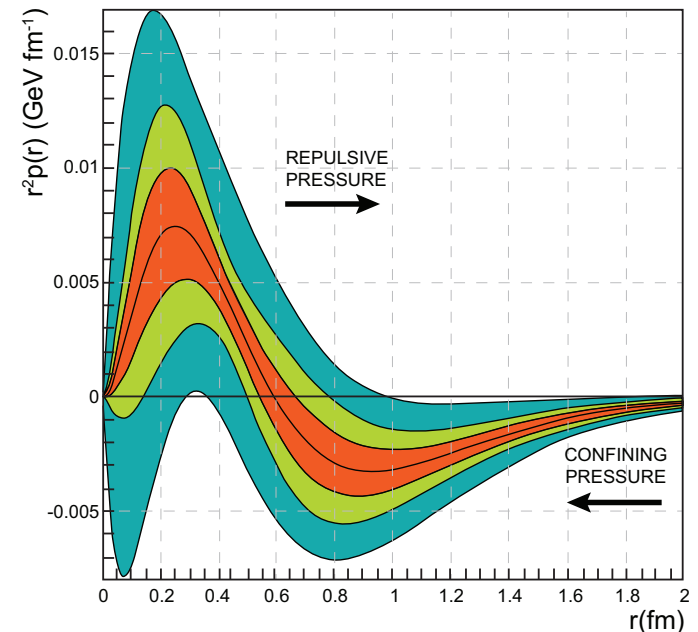
The 2<sup>nd</sup> **Mellin moment** of GPDs allows to access the pressure distribution inside hadrons through the **skewness dependency** of GPDs.

$$\text{CFF} \quad \mathcal{H}(\xi, t) = \int_{-1}^1 \left[ \frac{1}{\xi - x - i\epsilon} - \frac{1}{\xi + x - i\epsilon} \right] H(x, \xi, t) dx$$

$$\Re[\mathcal{H}(\xi, t)] \stackrel{\text{LO}}{=} D(t) + \mathcal{P} \left\{ \int_{-1}^1 \left[ \frac{1}{\xi - x} - \frac{1}{\xi + x} \right] \Im[\mathcal{H}(x, t)] dx \right\}$$

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# Gravitational Form Factors

V. Burkert, L. Elouadrhiri, F.-X. Girod, Nature 557 (2018) 396    M.V. Polyakov, P. Schweitzer, Int. J. Mod. Phys. A33 (2018) 1830025  
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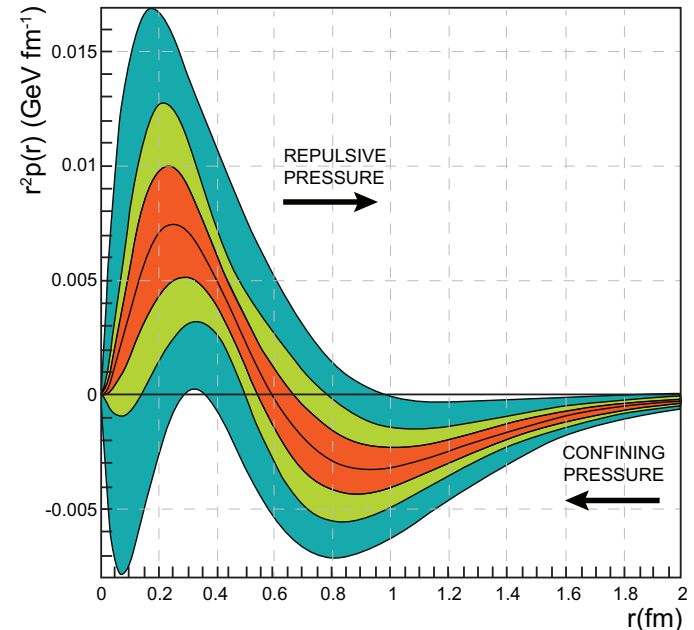
CFF 
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$$D(z, t) = (1 - z^2) \left[ d_1(t) C_1^{3/2}(z) + \dots \right]$$

**Real part of Compton form factors**  
( $\sigma_{INT}$ )



## Proposed Measurements

Using polarized electron and positron beams, we are proposing to measure

- The unpolarized beam charge asymmetry  $A_{UU}^C$ , which is sensitive to the **CFF real part**
- The polarized beam charge asymmetry  $A_{LU}^C$ , which is sensitive to the **CFF imaginary part**
- The neutral beam spin asymmetry  $A_{LU}^0$ , which is sensitive to **higher twist effects**

New GPD  
Observables  
@ JLab

$$A_{UU}^C = \frac{(Y_+^+ + Y_-^+) - (Y_+^- + Y_-^-)}{Y_+^+ + Y_-^+ + Y_+^- + Y_-^-}$$

$$= \frac{\sigma_{INT}}{\sigma_{BH} + \sigma_{DVCS}}$$

$$A_{LU}^C = \frac{(Y_+^+ - Y_-^+) - (Y_+^- - Y_-^-)}{Y_+^+ + Y_-^+ + Y_+^- + Y_-^-}$$

$$= \frac{\tilde{\sigma}_{INT}}{\sigma_{BH} + \sigma_{DVCS}}$$

$$A_{LU}^0 = \frac{(Y_+^+ + Y_-^-) - (Y_+^- + Y_-^+)}{Y_+^+ + Y_-^+ + Y_+^- + Y_-^-}$$

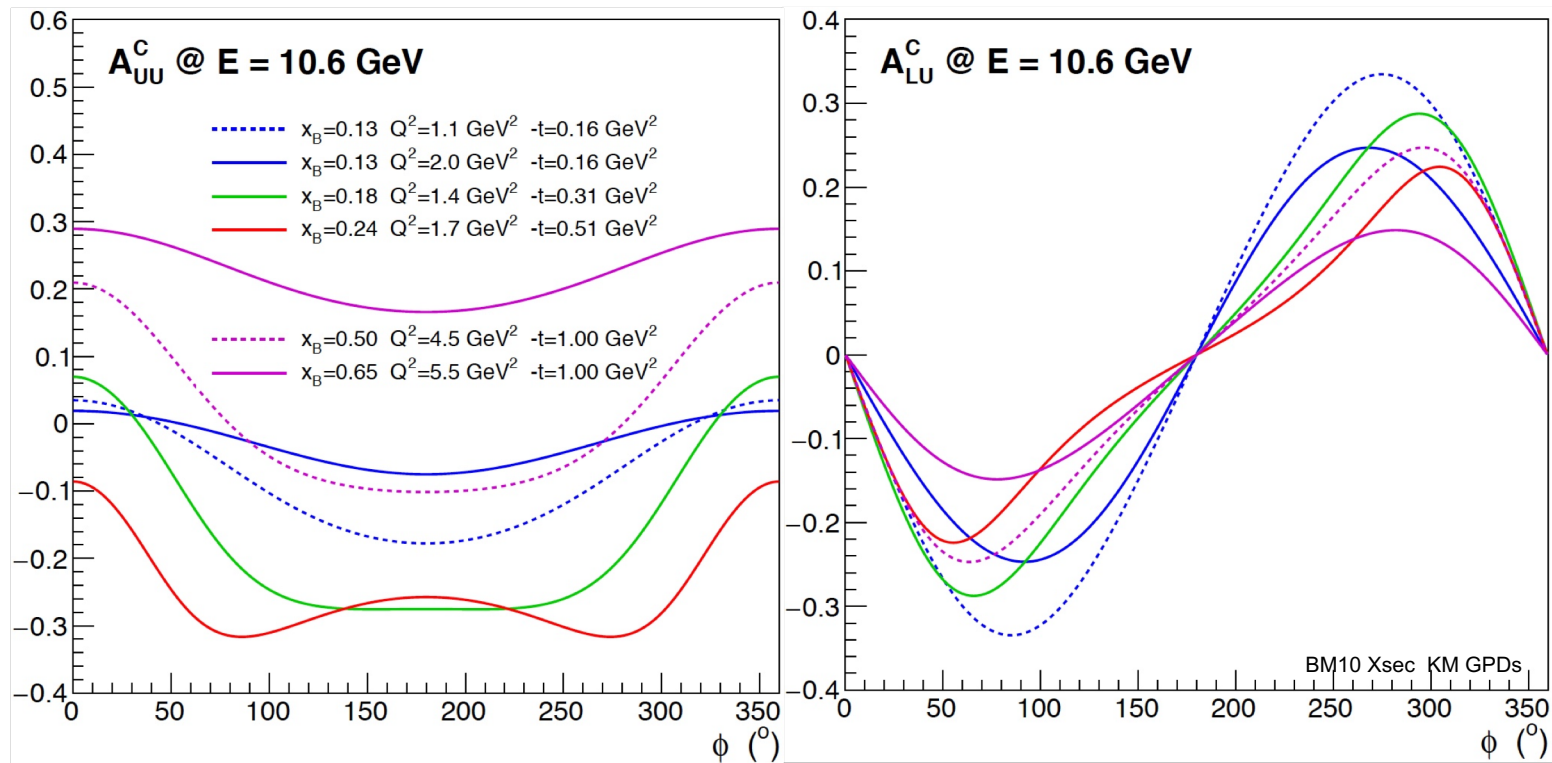
$$= \frac{\tilde{\sigma}_{DVCS}}{\sigma_{BH} + \sigma_{DVCS}}$$

$$\text{👉 } A_{LU}^C \neq A_{LU}^\pm = \frac{\pm(\tilde{\sigma}_{INT} \pm \tilde{\sigma}_{DVCS})}{\sigma_{BH} + \sigma_{DVCS} \pm \sigma_{INT}}$$

$Y_\pm^{e^\pm} = \frac{N_\pm^\pm}{Q_\pm^\pm}$  is the yield normalized the accumulated beam charge.

## Experimental Signal

A.V. Belitsky, D. Müller, PRD 82 (2010) 074010    K. Kumerički, D. Müller, NPB 841 (2010)



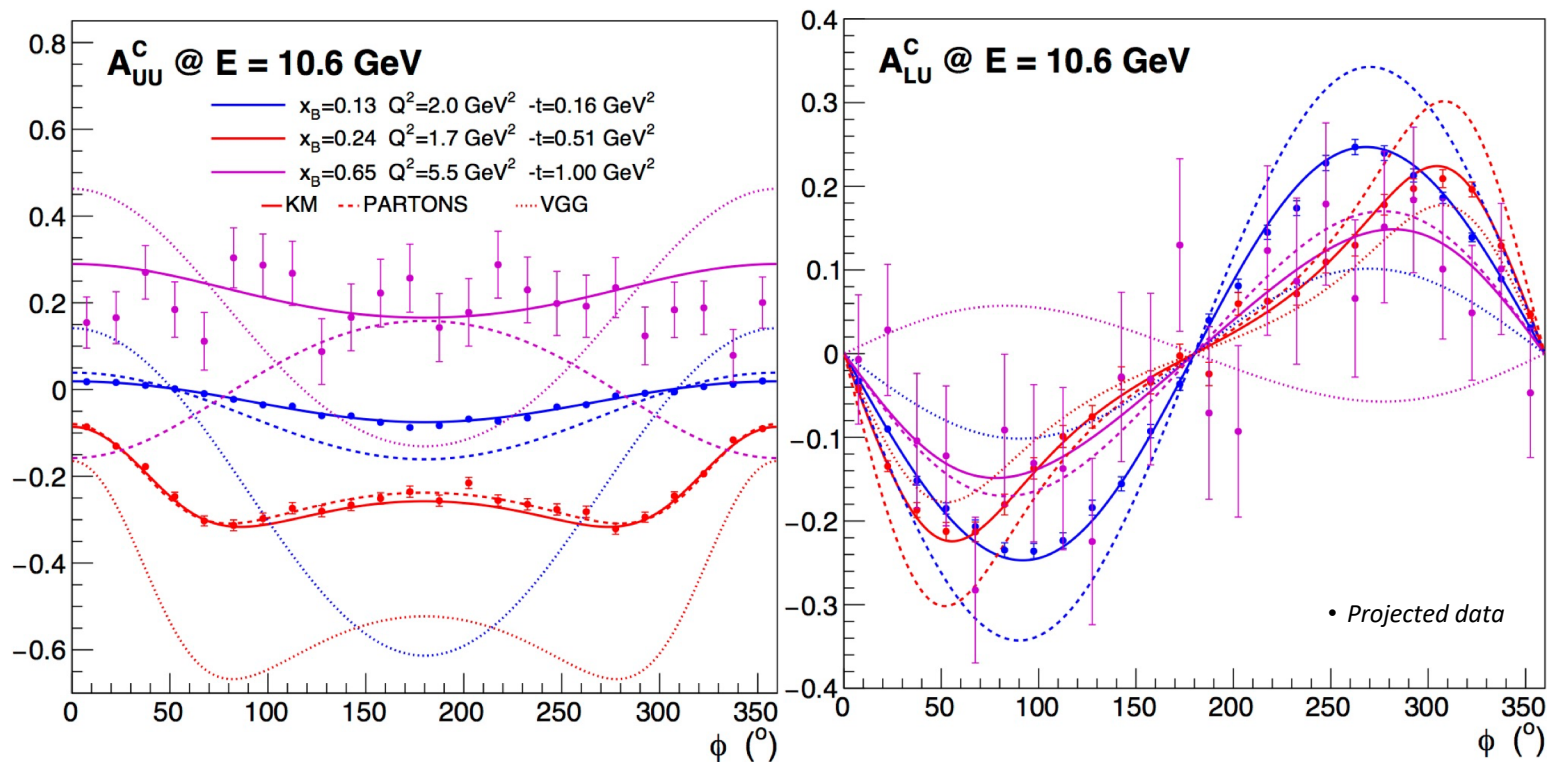
- The **magnitude** of unpolarized and polarized BCAs is **sizeable** and **kinematics dependent**.



## Partonic Content Sensitivity

A.V. Belitsky, D. Müller, PRD 82 (2010) 074010

K. Kumerički, D. Müller, NPB 841 (2010) 1    B. Berthou et al. EPJC 78 (2018) 478    M. Vanderhaeghen, P.A.M. Guichon, M. Guidal, PRD 60 (1999) 094017



- The **shape** and **magnitude** of BCAs are **sensitive** to the specific **GPD model**.

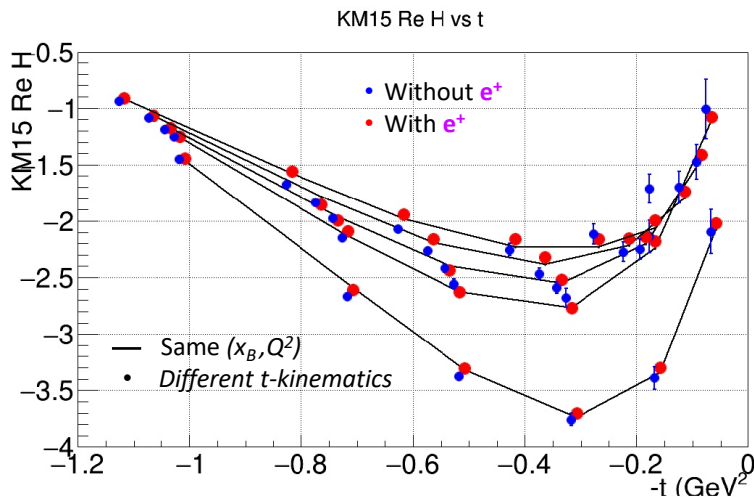
# Impact of Positron Measurements

K. Kumerički, D. Müller, NPB 841 (2010) 1

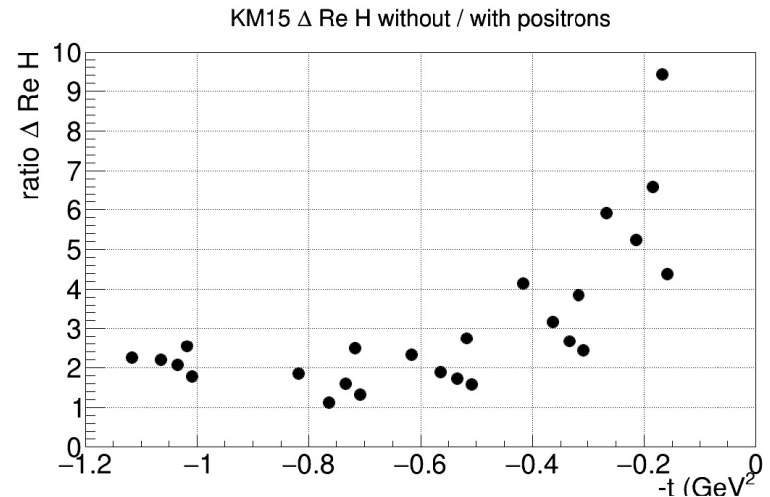
- The importance of positron beams for the determination of CFFs can be quantified in a **model-dependent** way depending on : the **cross section model**, the **GPDs model**, and the hypotheses of the **fitting approach**.

Observable	$\sigma_{UU}$	$A_{LU}$	$A_{UL}$	$A_{LL}$	$A_{UU}^C$	$A_{LU}^C$
Sytematics (%)	5	3	3	$3\oplus 3$	3	3

Fitting of  $\{\mathcal{H}, \tilde{\mathcal{H}}\}$  CFFs assuming model values for  $\{\varepsilon, \tilde{\varepsilon}\}$  CFFs.



Purity of  $e^+$  observables



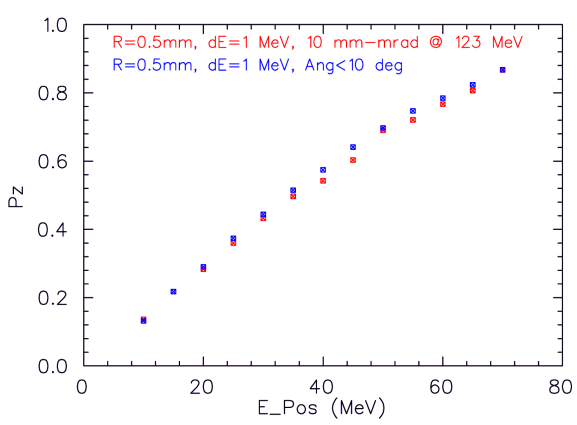
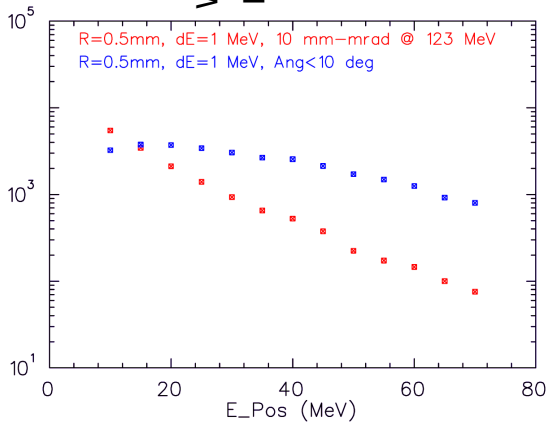
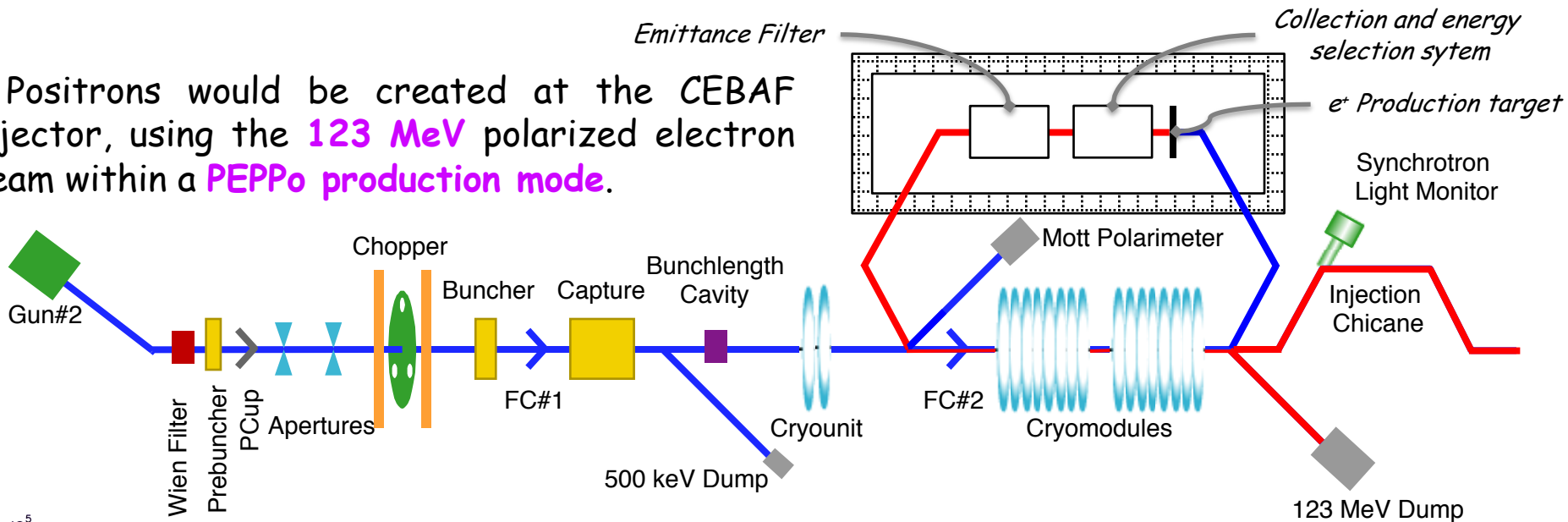
Reduction of CFF correlations

- Improvement of the **statistical** and **systematical** error of  $\Re[\mathcal{H}]$ .

# Positron Beam Concept

(PEPPo Collaboration) D. Abbott et al. , PRL 116 (2016) 214801 L. Cardman et al. (JPos17) AIP CP 1970 (2018) 050001

➤ Positrons would be created at the CEBAF injector, using the **123 MeV** polarized electron beam within a **PEPPo production mode**.

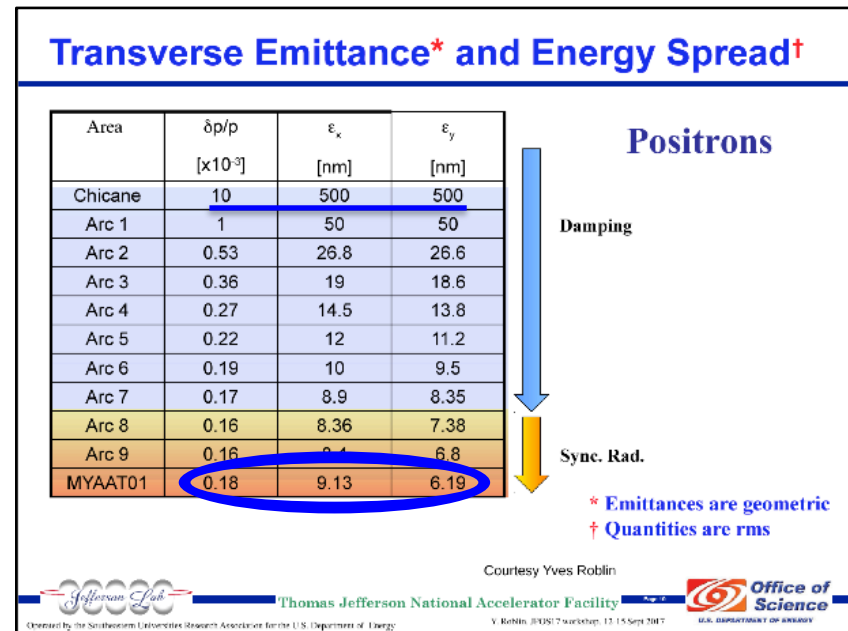
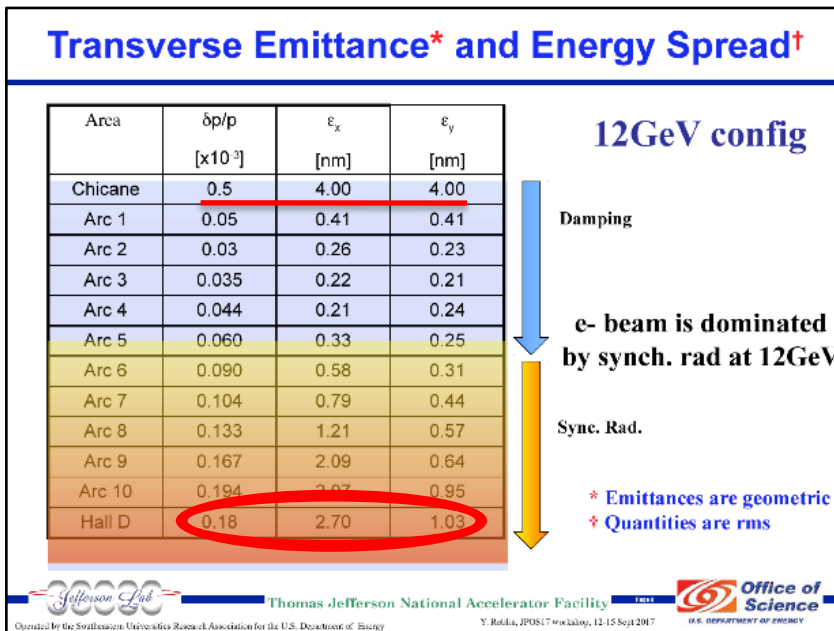


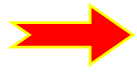
- ✓ Selecting **60 MeV** positrons maximizes the FoM (**60%**, **100 nA**).
- ✓ Selecting **6 MeV** positrons maximizes the flux (**0%**, **5 μA**).

# Beam Properties

Y. Roblin at the International Workshop on Physics with Positrons at Jefferson Lab, Newport News, September 12-15, 2017

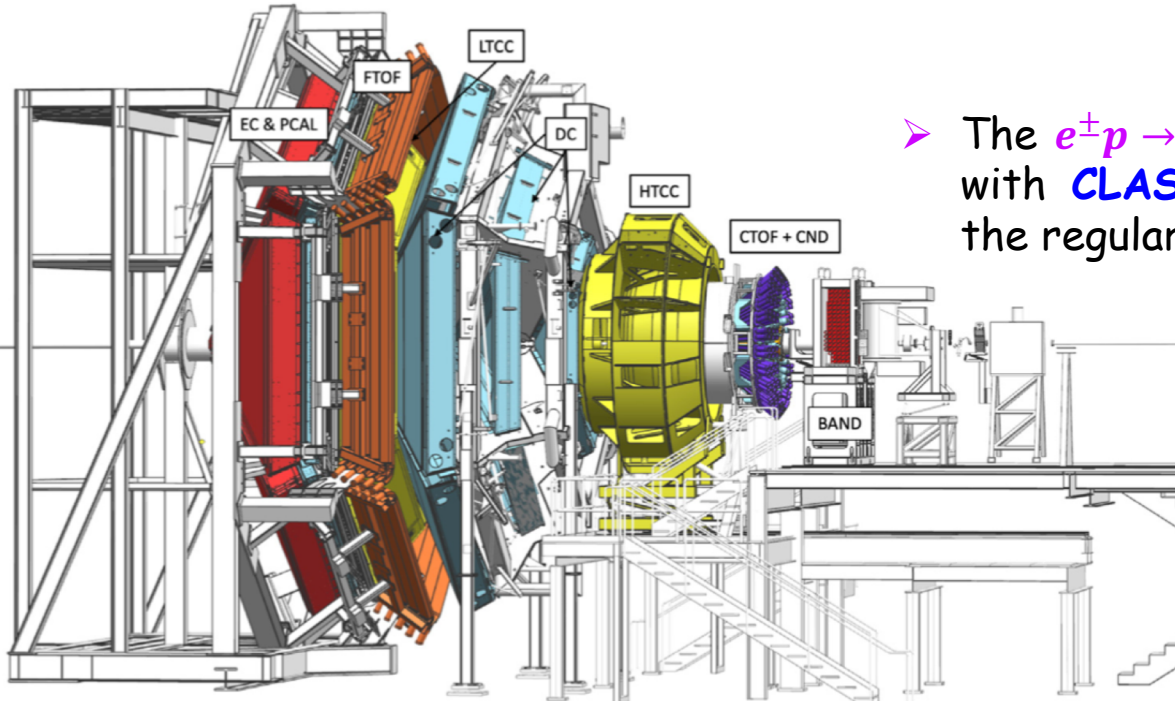
- At the source, the secondary  $e^+$  beam is much larger momentum dispersion and emittance than the primary  $e^-$  beam, but has the **same  $\delta p/p$**  with a spot size **2-3 times larger** at target.



 A **target cell** with **50% larger** (15 mm diameter) entrance and exit **windows** is anticipated to **avoid** any **interaction** with the target **structure frames**.

## CLAS12

V. Burkert et al. NIMA 959 (2020) 163419



➤ The  $e^\pm p \rightarrow e^\pm p \gamma$  reaction will be measured with **CLAS12** in **OUT-bending** mode, using the regular detector arrangement.

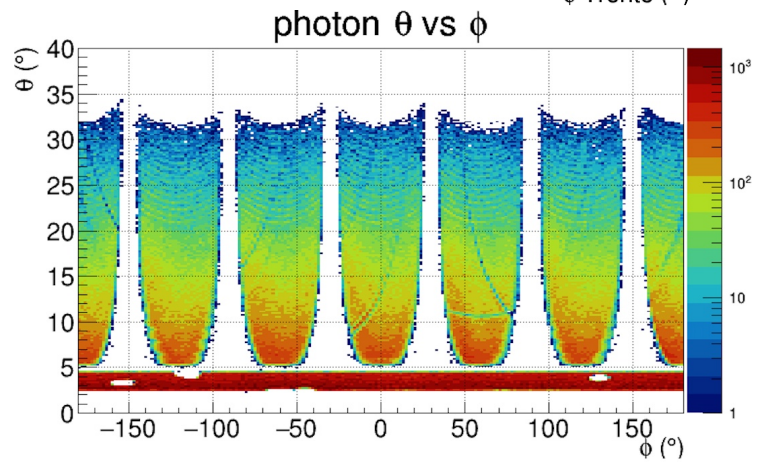
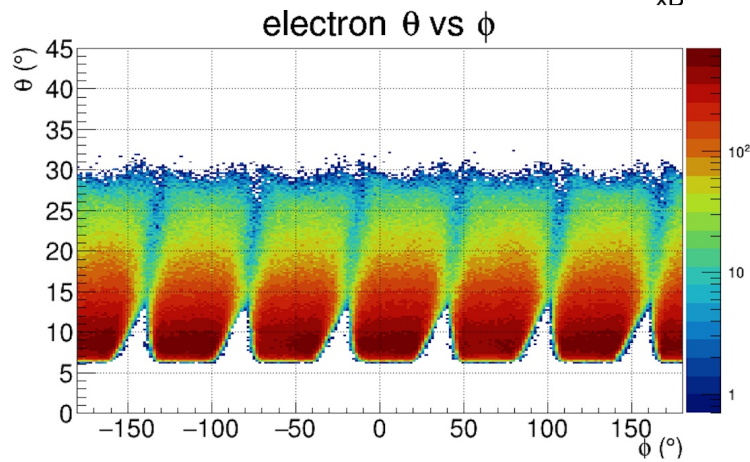
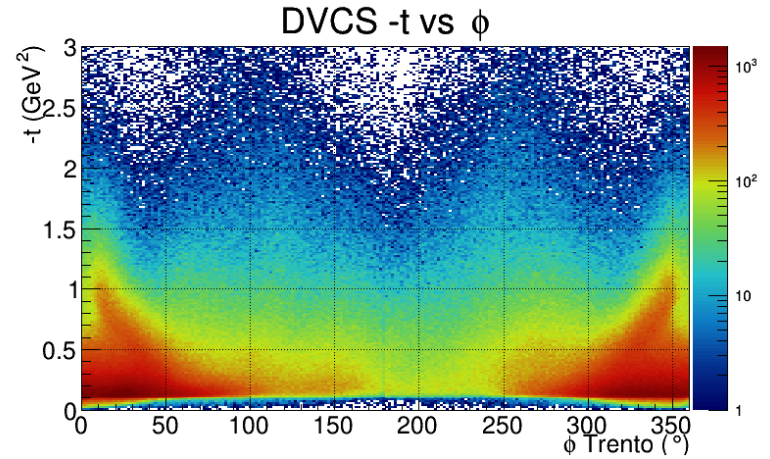
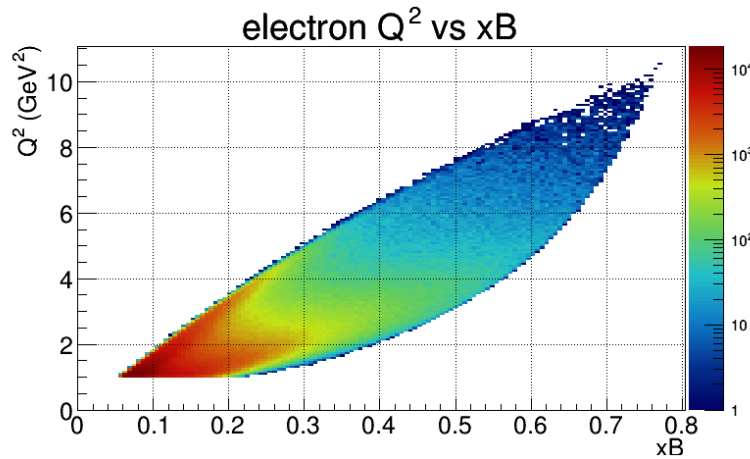
- $\mathcal{L} = 0.6 \times 10^{35} \text{ cm}^{-2} \cdot \text{s}^{-1}$
- New **5 cm LH<sub>2</sub>** target cell

- There is **no difference** between  $e^-$  and  $e^+$  beam transport in Hall B beam line, nor in beam related detector background.
- Beam diagnostics are expected to operate similarly with  $e^-$  and  $e^+$  beam.



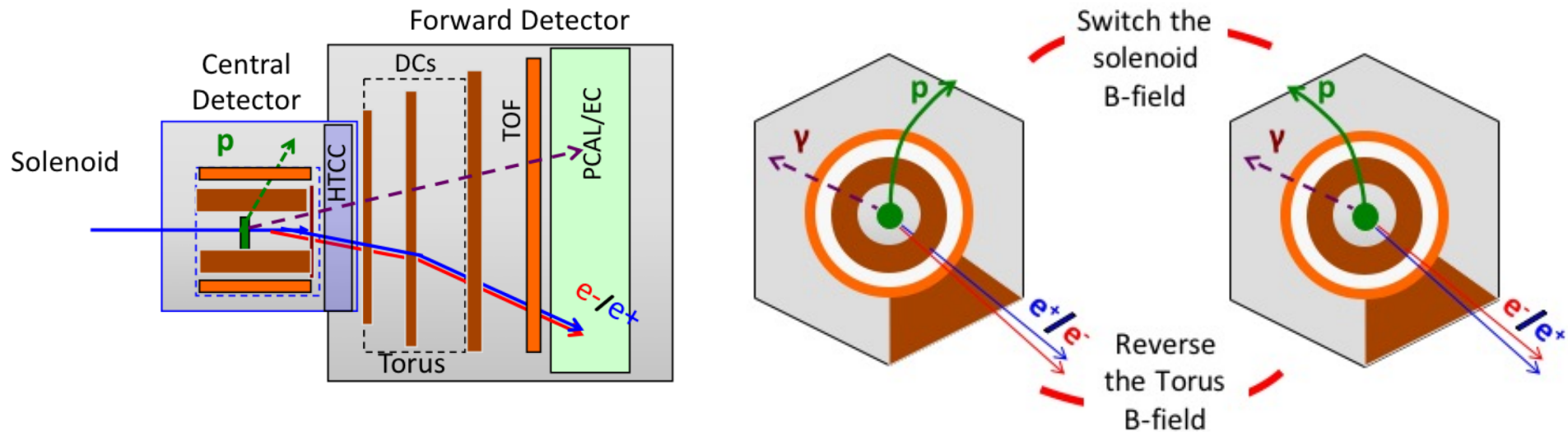
# Kinematic Coverage

From a subset of RGA data in OUT-bending mode



## False Asymmetries

- Potential false asymmetries may occur due to  $e^-$  and  $e^+$  from same vertex and kinematics passing through different part of the detector shifted in  $\phi$  in a sector.



- **Switch** the **solenoid field** to reveals false asymmetries in the Forward Detector, which may creates false asymmetries in proton tracking.
- **Measure**, simultaneously to DVCS, **elastic scattering** cross sections for  $e^-$  and  $e^+$  at low- $Q^2$  where  $2\gamma$ -effects are small.



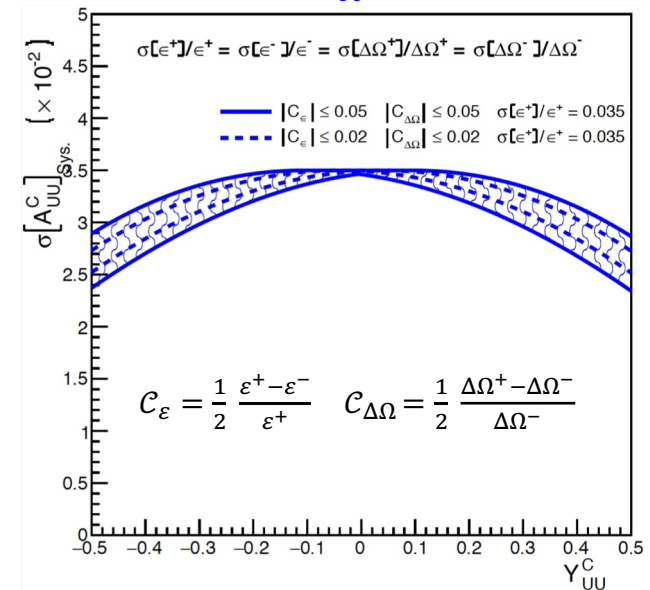
## Systematical Effects

- The measurement of **BCAs** is comparable to the measurement of **relative cross sections** where some of systematical effects cancel out while others ask for careful control and monitoring.

Minizing systematics in  $e^+/e^-$  comparison requires:

- Same **beam qualities** (energy, transverse profile, emittance...)
- Same **detector** (target, efficiency, solid angle...)
- Same **statistics** (accumulated charge, beam polarization)

Systematics projection on  $A_{UU}^C$  for 5% uncertainty on  $\varepsilon \cdot \Delta\Omega$ .



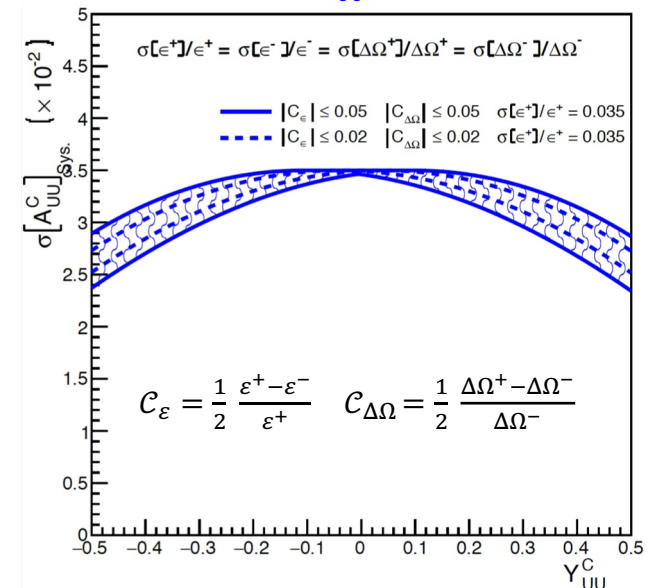
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Systematics projection on  $A_{UU}^C$  for 5% uncertainty on  $\varepsilon \cdot \Delta\Omega$ .



- Take **DVCS data** with the **secondary electron beam** (simultaneously produced at the positron production target) prior data taking with the **secondary positron beam**.

# Beam Time Request

We are asking for a total of **100 days** of beam, operating **CLAS12** with **45 nA e<sup>-</sup>** and **e<sup>+</sup>** beams polarized at **60%**, distributed in :

Purpose	Label	Beam parameters					Target	Sol. Pol.	Tor. Pol.	Time (h)
		q (e)	Nat.	E (GeV)	I (nA)	λ (%)				
ep → ep	Cal.	-	P	2.2	45	0	5 cm LH <sub>2</sub>	-	+	24
								+		24
								+		24
								-		24
ep → epγ	Phy.		S	10.6	60	60		-		480
Background	Cal.							-		48
ep → epγ	Phy.							+		480
Background	Cal.							+		48
Commissioning		+	S	2.2	45	0	5 cm LH <sub>2</sub>	+	-	72
ep → ep	Cal.							+		24
								-		24
Commissioning								10.6		60
ep → epγ	Phy.			-	480					
Background	Cal.			-	48					
ep → epγ	Phy.			+	480					
Background	Cal.			+	48					
<b>Total</b>										<b>2400</b>

- **80 days** for **physics** data taking;
- **20 days** for **commissioning** and **calibration**;

using lepton beams of different nature

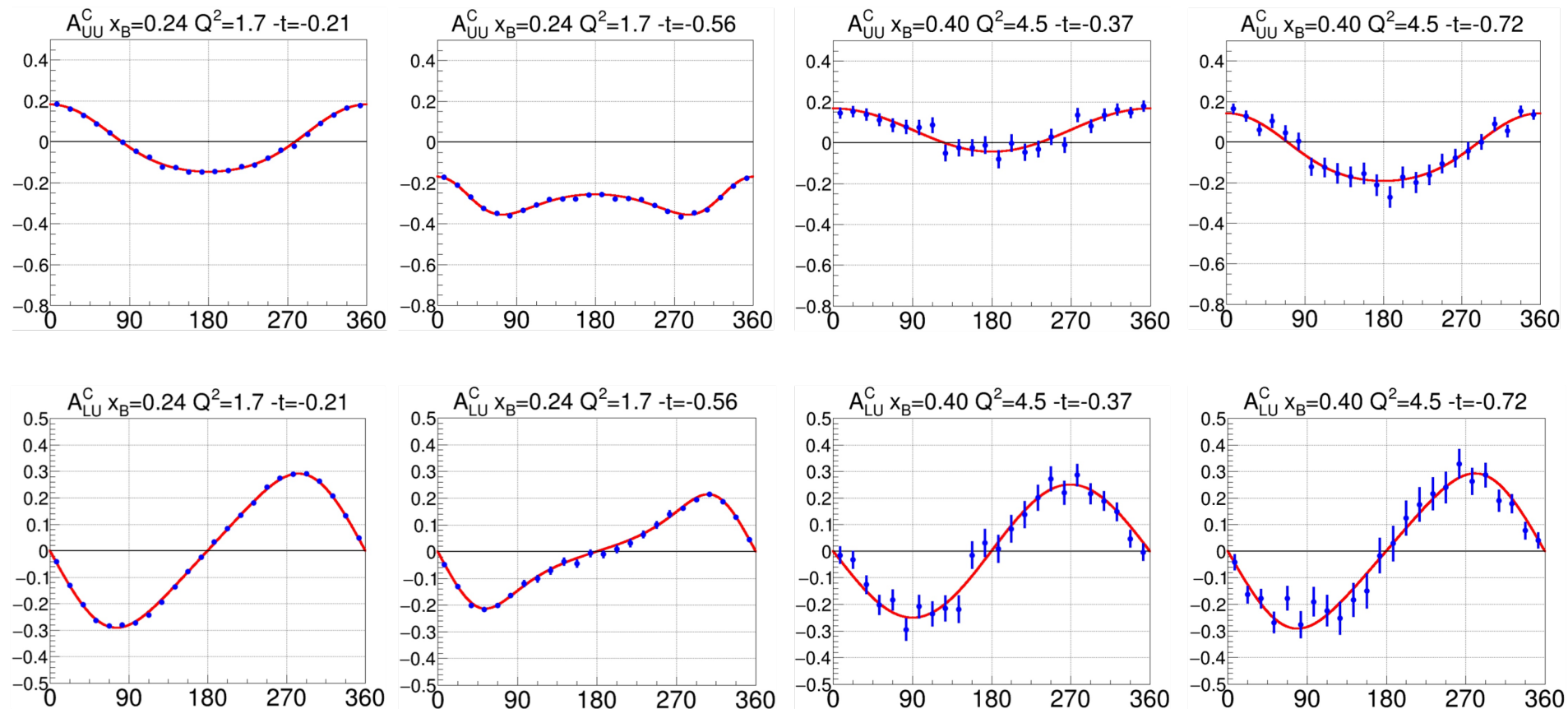
- 2 days with the CEBAF e<sup>-</sup> beam;
- 46 days with the secondary e<sup>-</sup> beam;
- 52 days with the secondary e<sup>+</sup> beam;

and different energies

- 9 days at **2.2 GeV**;
- 91 days at **10.6 GeV**.

## Experimental Projections

➤ A sample of expected experimental data... 15 bins in  $(x_B, Q^2) \times 6$  bins in  $t = 90$  azimuthal dependences per observable ( $A_{UU}^C, A_{LU}^C$ ).



$p$ -DVCS BCAs @ CLAS12

- We propose to measure **Beam Charge Asymmetries for the DVCS reaction off protons** at 10.6 GeV with **CLAS12** and using the **secondary polarized electron and positron beams** produced at a **PEPPo positron source**, over a **100 days** data taking period.
- The **separation** of the **DVCS** and **INT** reaction amplitudes will provide **unambiguous experimental signals** that **uniquely determine CFFs**, particularly  $\Re[\mathcal{H}]$  for the proposed experiment.

*The direct access to the real part of the INT amplitude via positron beams will constitute a major step forward for DVCS studies.*

This project has received funding from the European Union's Horizon 2020 research and innovation program under agreement No 824093.

*Additional Slides ...*

LOI12-18-004

J. Grames, E. Voutier et al. Jefferson Lab LOI12-18-004 (2018), arXiv:1906.09419

Letter of Intent submitted to JLab PAC46 (July 2018)

Highlighting 7 mini-LOI's Supported by 127 Members from 39 Institutions

Table with 2 columns: Physics with Positron Beams at Jefferson Lab 12 GeV and list of member institutions and contact information.

Table with 5 columns: Experiment Name, I- (nA), I+ (nA), Beam Polarization, Time (d). Rows include Two-photon exchange, Generalized Parton Distributions, and Test of the Standard Model.

These measurements all have significant physics interest. The proposers should carefully evaluate feasibility and present the best case possible in a future proposal.



e+ @ JLab White Paper

A. Accardi et al. JLab-PHY-20-3232, arXiv:2007.15081 (2020)

An experimental program with positron beams at Jefferson Lab

discussing 15 possible experiments covering:

- Generalized Parton Distributions physics,
Two-Photon Exchange physics,
Tests of the Standard Model,
and other specific measurements.

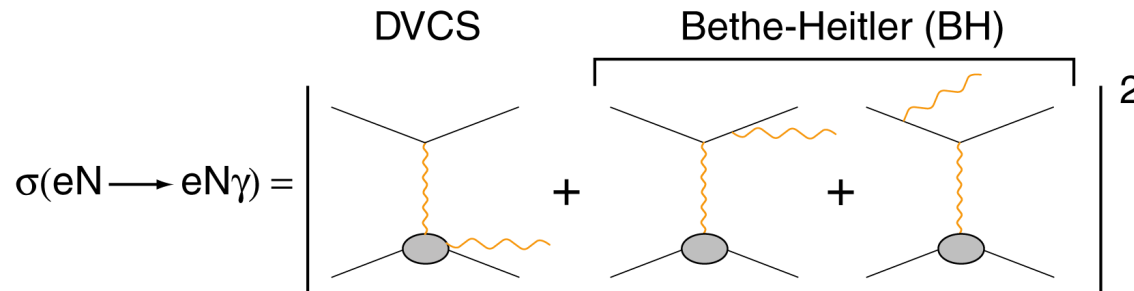
Supported by 197 Members from 59 Institutions

Left column of the white paper containing the title 'e+ @ JLab White Paper' and a list of 15 experimental topics.

Right column of the white paper containing the title 'An Experimental Program with Positron Beams at Jefferson Lab' and a list of 197 member institutions.

## $\mathcal{N}(e, e'\gamma\mathcal{N})$ Differential Cross Section

M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009

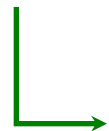


$$\sigma_{\text{PS}}^e = \sigma_{\text{P0}}^e + S \left[ P_1 \Delta\sigma_{\text{BH}} + (\Delta\tilde{\sigma}_{\text{DVCS}} + P_1 \Delta\sigma_{\text{DVCS}}) + e_1 (\Delta\tilde{\sigma}_{\text{INT}} + P_1 \Delta\sigma_{\text{INT}}) \right]$$

Electron  
observables

Additional observables

Electron & positron  
observables



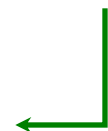
$$\begin{aligned} \sigma_{00}^- &= \sigma_{\text{BH}} + \sigma_{\text{DVCS}} - \sigma_{\text{INT}} \\ \sigma_{+0}^- - \sigma_{-0}^- &= 2\tilde{\sigma}_{\text{DVCS}} - 2\tilde{\sigma}_{\text{INT}} \end{aligned}$$

$$[\sigma_{++}^\pm - \sigma_{+-}^\pm] - [\sigma_{-+}^\pm - \sigma_{--}^\pm] = 4\Delta\sigma_{\text{BH}} + 4\Delta\sigma_{\text{DVCS}} \pm 4\Delta\sigma_{\text{INT}}$$

$$\sigma_{0+}^\pm - \sigma_{0-}^\pm = 2\Delta\tilde{\sigma}_{\text{DVCS}} \pm 2\Delta\tilde{\sigma}_{\text{INT}}$$

$$\sigma_{00}^+ - \sigma_{00}^- = 2\sigma_{\text{INT}}$$

$$[\sigma_{+0}^+ - \sigma_{-0}^+] - [\sigma_{+0}^- - \sigma_{-0}^-] = [\sigma_{+0}^+ - \sigma_{+0}^-] - [\sigma_{-0}^+ - \sigma_{-0}^-] = 4\tilde{\sigma}_{\text{INT}}$$



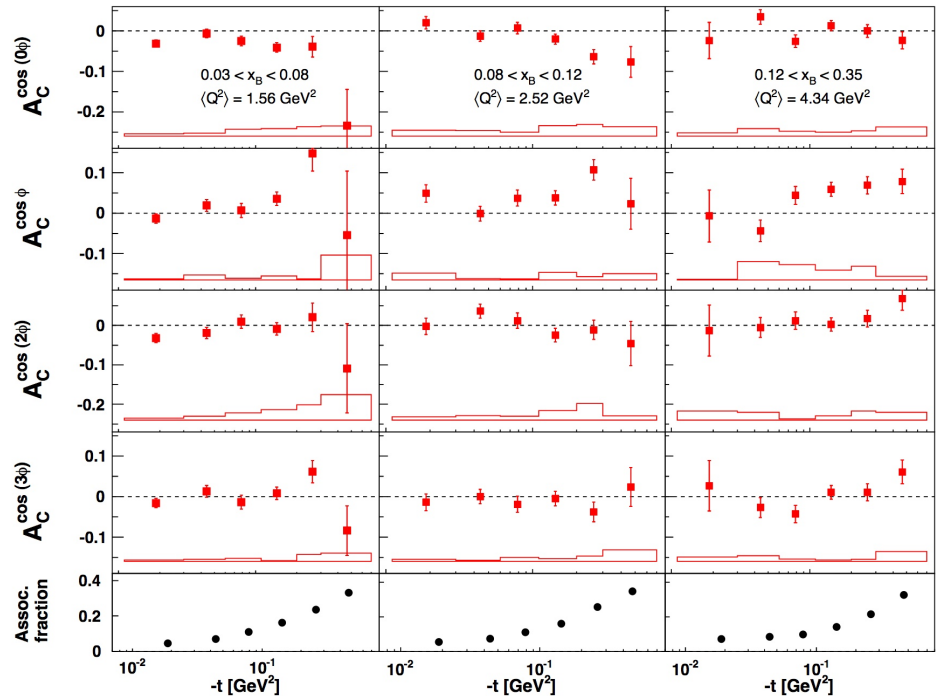
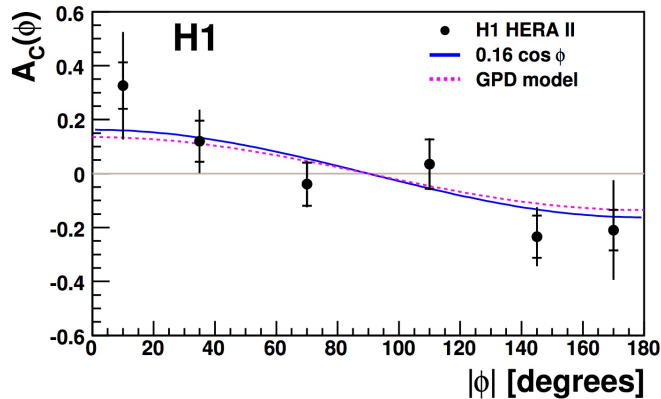
**Polarized electrons and positrons** allow to **separate** the **unknown amplitudes** of the cross section for electro-production of photons.

## Existing Data

(H1 Collaboration) F.D. Aaron et al. PLB 681 (2009) 391

(HERMES Collaboration) A. Airapetian et al. JHEP 06 (2008) 066 - 11 (2009) 083 - 07 (2012) 032

- Pioneering measurements of DVCS with electron and positron beams at **HERA** demonstrated the existence of a **BCA-signal**.

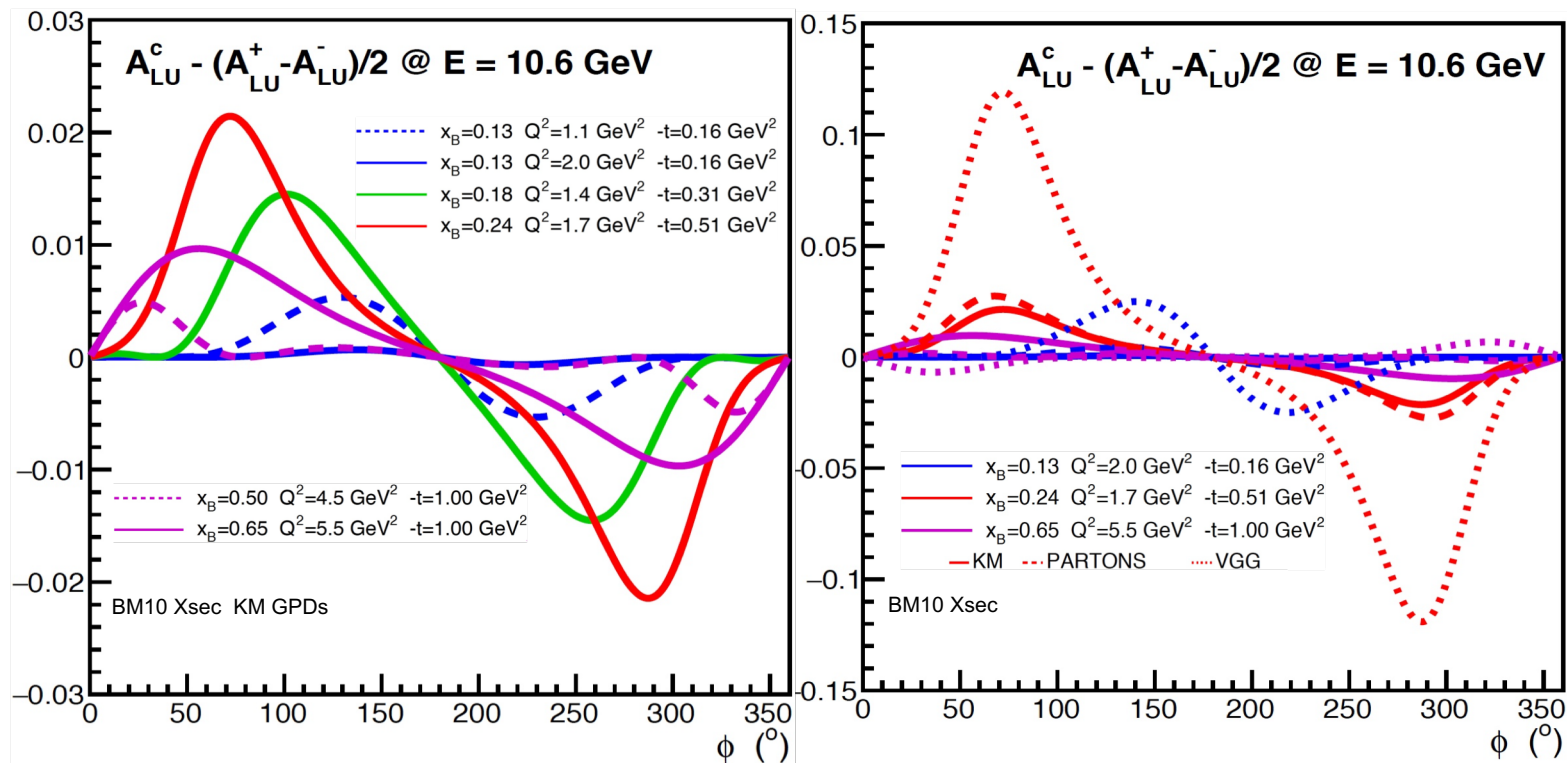


- The **COMPASS** experiment operating high energy  $\mu^\pm$  beams should release in a near future **BCA data** in the **sea-quark region**.

## Bethe-Heitler Dominance

A.V. Belitsky, D. Müller, PRD 82 (2010) 074010

K. Kumerički, D. Müller, NPB 841 (2010) 1    B. Berthou et al. EPJC 78 (2018) 478    M. Vanderhaeghen, P.A.M. Guichon, M. Guidal, PRD 60 (1999) 094017

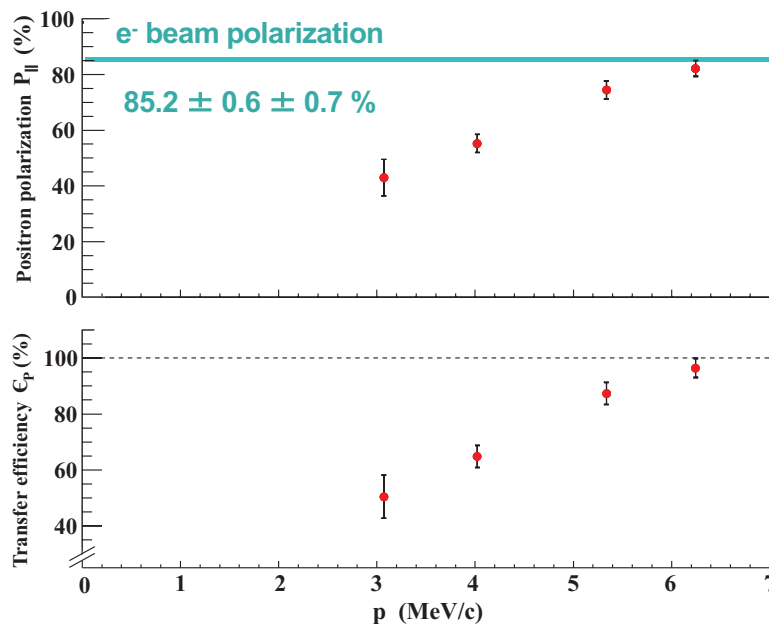
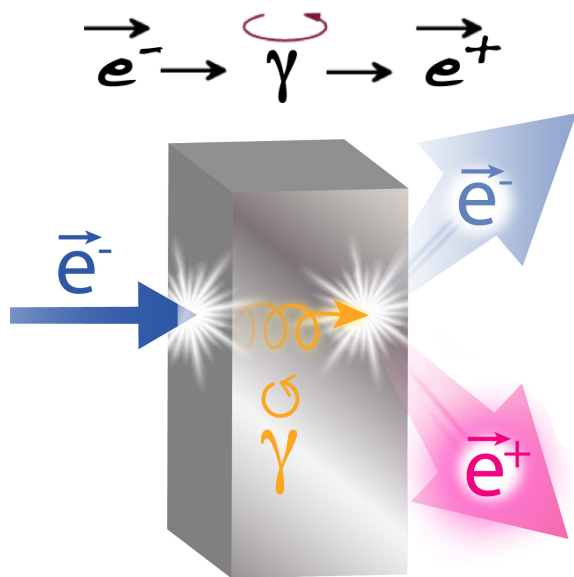


- BH-dominance  $\equiv A_{LU}^c - (A_{LU}^+ - A_{LU}^-)/2 = 0$  is a **kinematics-** and **GPD model-dependent hypothesis** eventually testable at CLAS12.

## Polarized Electrons for Polarized Positrons

(PEPPo Collaboration) D. Abbott et al. , PRL 116 (2016) 214801

- PEPPo demonstrated **efficient polarization transfer** from **8.2 MeV/c electrons to positrons**, expanding polarized positron capabilities from **GeV to MeV accelerators**.

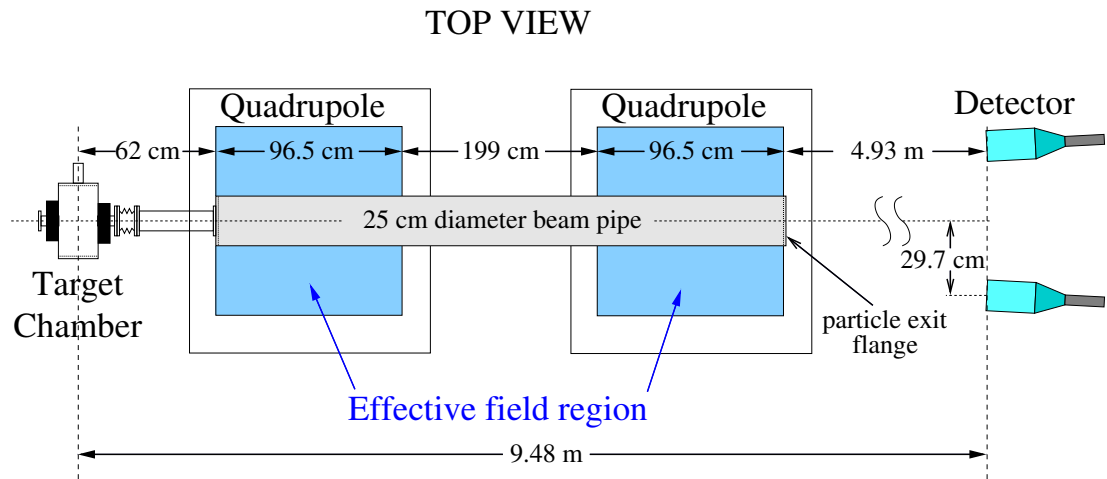


The PEPPo technique can achieve up to **100% transfer** of the electron polarization.

## Møller to Bhabha Polarimeter



**Bhabha asymmetries** are **identical** to **Møller's**, and **cross sections** are **similar magnitude** at 90°c.m.



- ❖ The transition of the **Møller** polarimeter into a **Bhabha** polarimeter will be achieved by adapting the detector configuration to allow for single ( $e^+$ ) and/or coincidence ( $e^+e^-$ ) detection at 90°c.m.

*On-going design optimization*

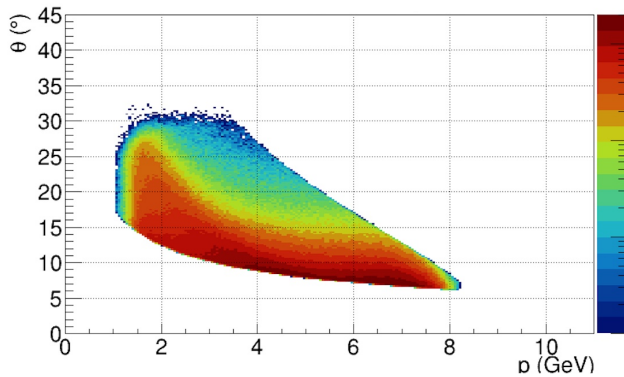


## Detector Acceptance

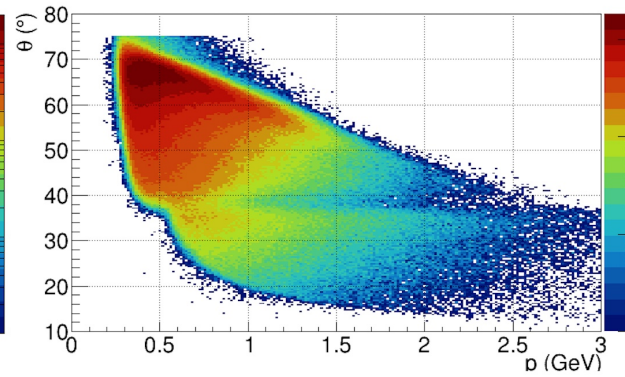
From a subset of out-bending RGA data

- The CLAS12 torus will operate in **OUT-Bending mode**.

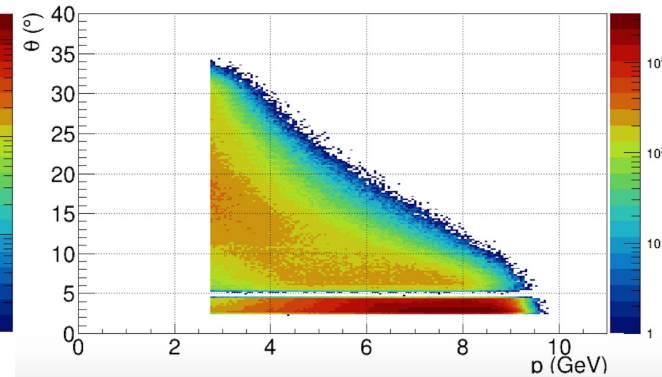
Electron (Positron)



Proton



Photon

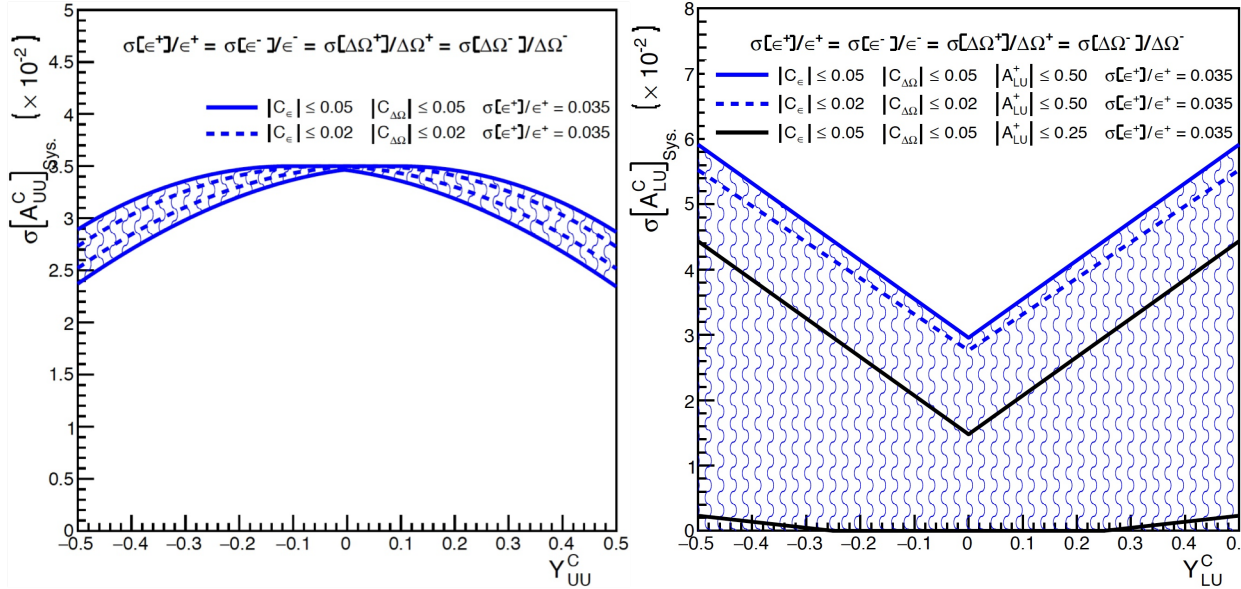


- Scattered electrons (positrons) are detected in the **Forward Detector (FD)**.
- Recoil protons are detected essentially in the **Central Detector (CD)**, but also in the **FD** at large  $-t$ .
- Produced photons are measured in the **Electromagnetic Calorimeter (Ecal)** of the FD and in the **Forward Tagger Calorimeter (FTCal)**.



## Evaluation of Systematics

- Solid angle and detector efficiency differences between  $e^-$  and  $e^+$  result in corrections ( $\eta_c$ ) to raw normalized yield asymmetries ( $y_{UU}^C, y_{LU}^C$ ), and related systematic uncertainties.



$$\eta_c = C_{\Delta\Omega} - C_\epsilon - 2C_{\Delta\Omega}C_\epsilon$$

$$C_\epsilon = \frac{1}{2} \frac{\epsilon^+ - \epsilon^-}{\epsilon^+}$$

$$C_{\Delta\Omega} = \frac{1}{2} \frac{\Delta\Omega^+ - \Delta\Omega^-}{\Delta\Omega^-}$$

$$A_{UU}^C = \frac{(1 + \eta_c)y_{UU}^C - \eta_c}{1 + \eta_c - \eta_c y_{UU}^C}$$

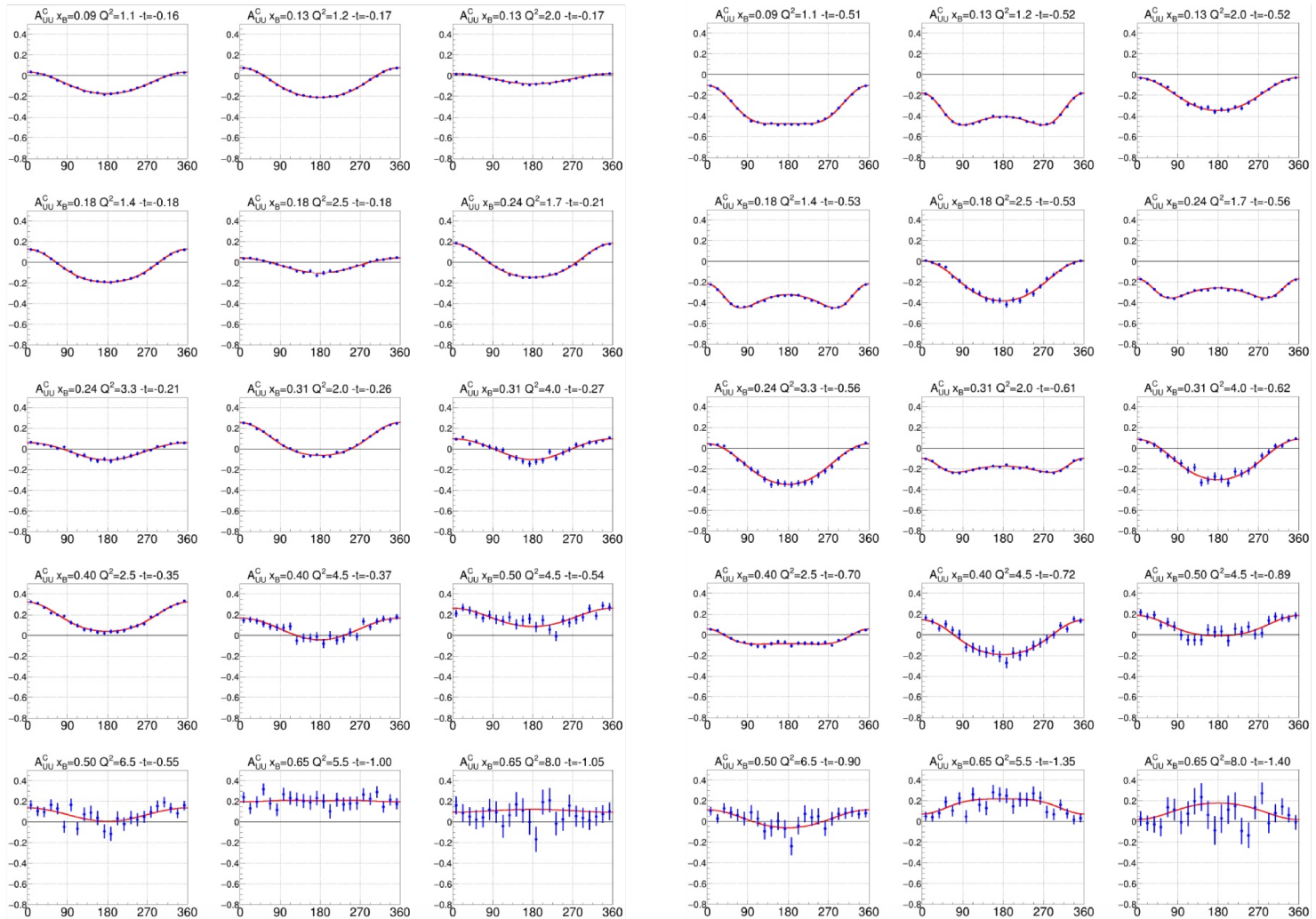
$$A_{LU}^C = [1 + \eta_c(1 + A_{UU}^C)]y_{LU}^C - \eta_c(1 + A_{UU}^C)A_{LU}^+$$

Assuming **5% systematic** uncertainty on  $\epsilon \cdot \Delta\Omega$  and **raw asymmetries** within **[-50%;50%]**:

- **Unpolarized BCA** systematics are about **0.03**;
- **Null polarized BCA** can be measured with a systematic uncertainty about **0.03**.

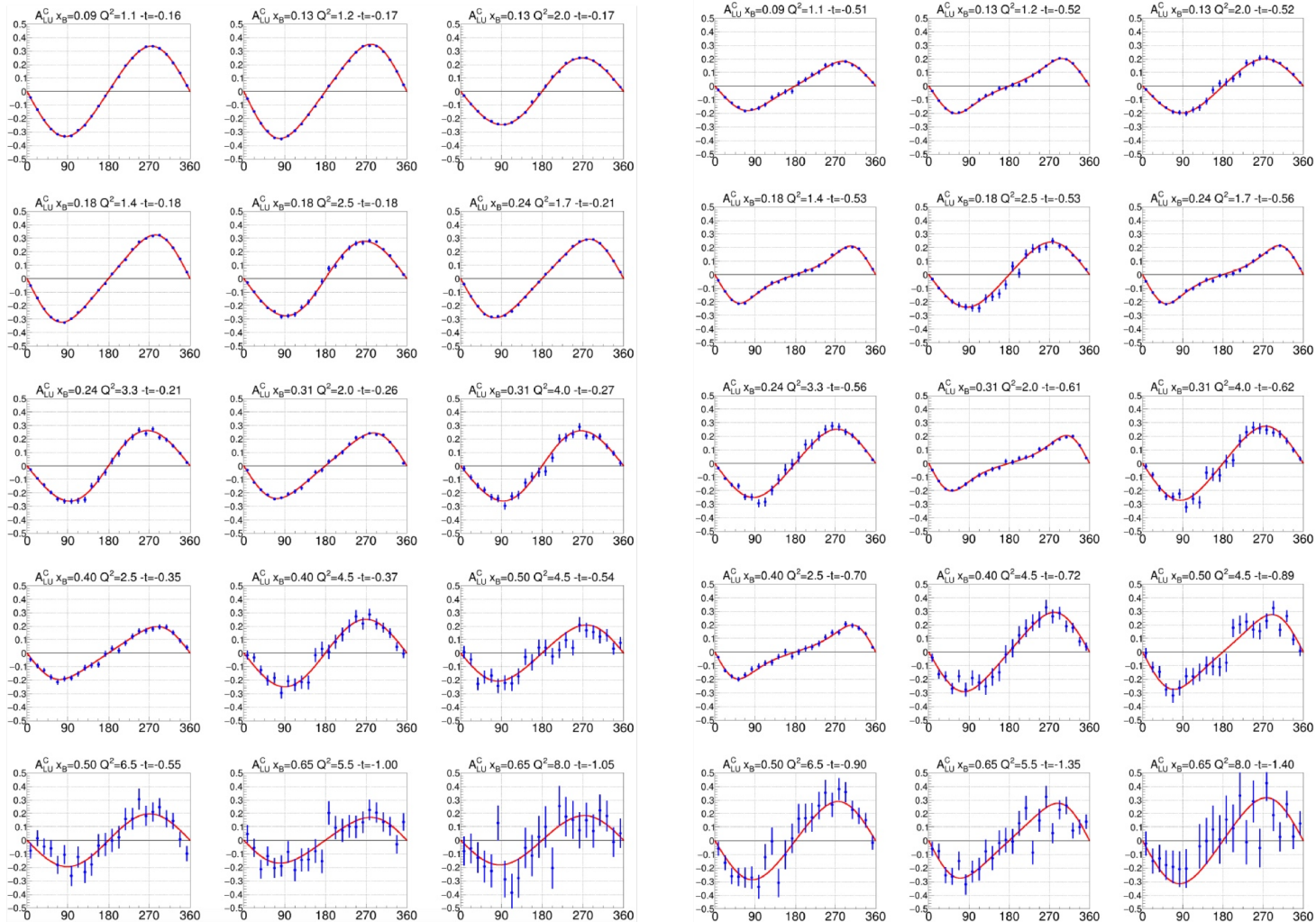
*p-DVCS unpolarized BCA @ CLAS12*

More projected data...



*p-DVCS polarized BCA @ CLAS12*

More projected data...



*p-DVCS positron BSA @ CLAS12*

More projected data...

