

# $\pi^+$ cross section measurement with Eg2 experimental data

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Advisors:

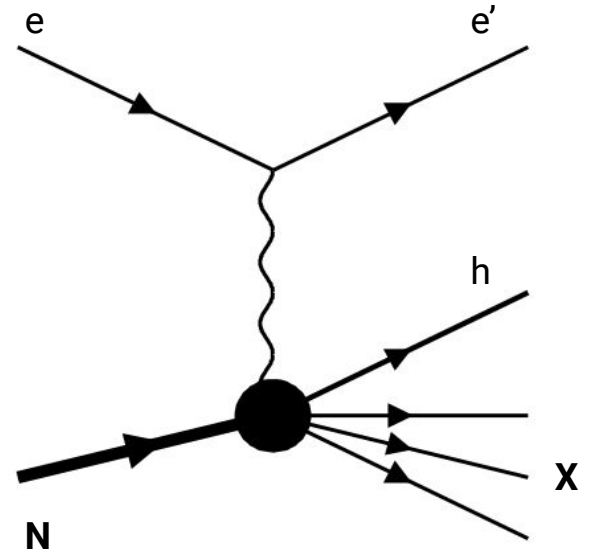
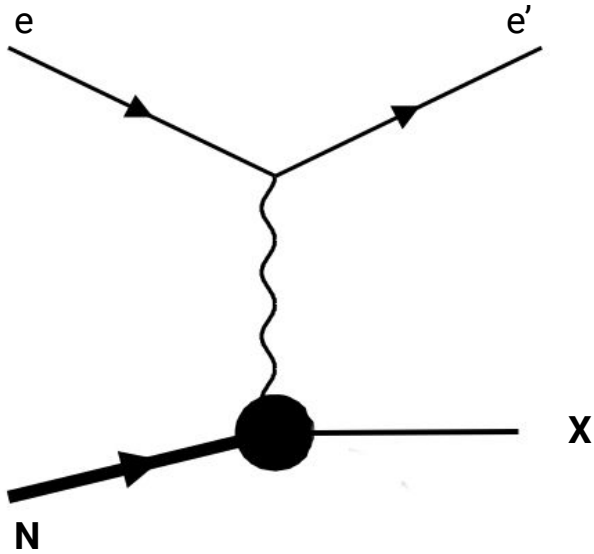
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# Outline

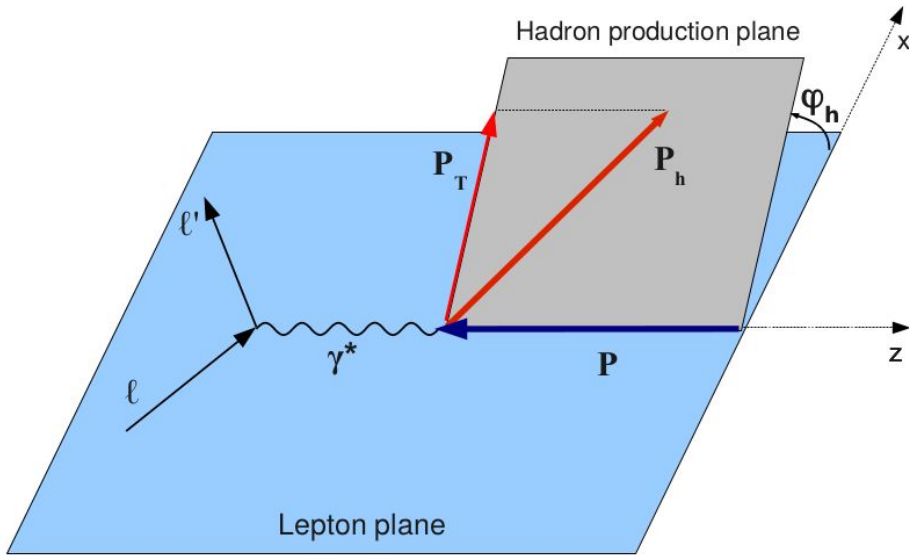
1. Brief conceptual introduction
2. Eg2 experiment
3. Motivation
4. Calculations
5. Results

# SIDIS



- $N$  -> Nucleon or nucleus
- $h$  -> Detected hadron
- $X$  -> The unknown

# SIDIS



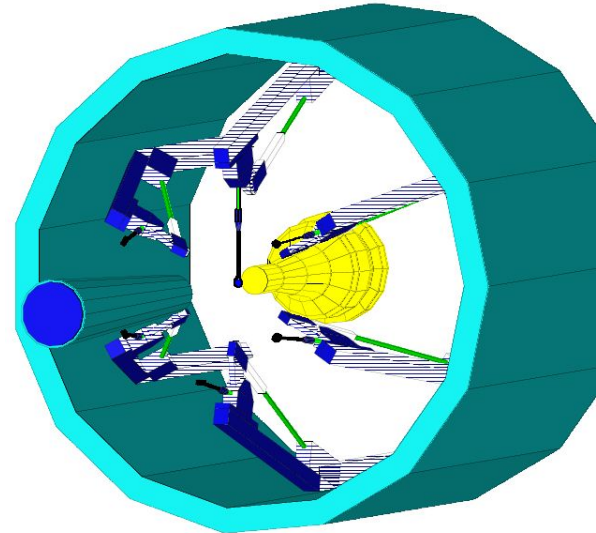
- Usually defined with five kinematical variables
  - $Q^2$  : Invariant mass virtual photon
  - $X_b$  : Bjorken scaling variable
  - $Z_h$  : Fraction of momentum that the produced hadron receives from virtual photon
  - $P_T$  : Transverse momentum of produced hadron
  - $\phi_h$  Angle between the hadronic production plane and the leptonic plane
- More details on nucleon structure dynamics

# Eg2 experiment

- Double target system
- The liquid cryo target and the solid target are simultaneously exposed to the beam
- Same luminosity for both targets (approx.)



*Eg2 double target*

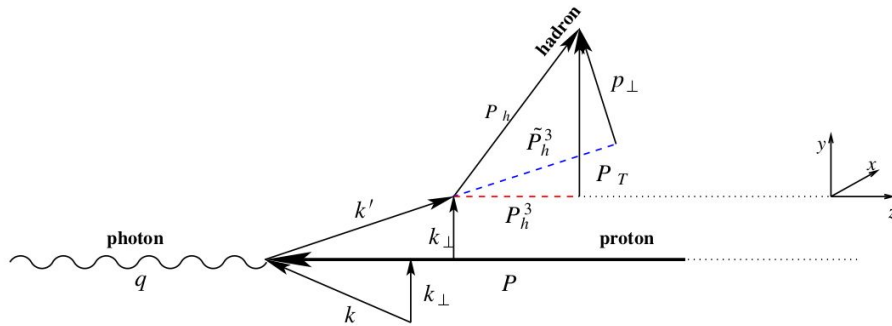


*Eg2 double target simulation*

# Motivation

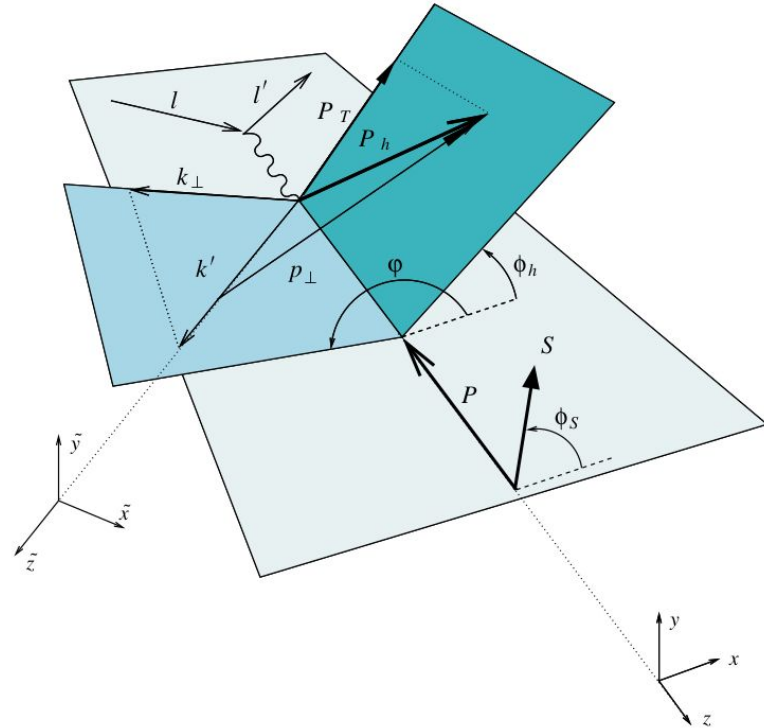
**k** : Intrinsic quark momentum -> Partonic motion inside the nucleon

Why is it important? Introduces important effects



**k**: Intrinsic quark momentum, pre scattering

**k'**: Intrinsic quark momentum, post scattering, pre hadronization



# Motivation

**k** : **Intrinsic quark momentum** -> Partonic motion inside the nucleon

Why is it important? Introduces Cahn effect

$$\frac{d\sigma}{dQ^2 dX_b dZ_h dP_t^2 d\phi_h} \propto A + B \cdot \cos(\phi_h) + C \cdot \cos(2\phi_h)$$

Azimuthal modulations are due to taking into account this intrinsic partonic motion [1]:

- **k** will be better represented by its transversal component with respect to the virtual photon direction: **kt**
- **kt** also can be found in models of structure functions
- We can extract a representative quantity for **kt** from the eq. above

**Measuring  $\langle kt^2 \rangle$  allows to understand more the azimuthal**

**behavior of the measured hadron**

# Cross section measurement

**First step** towards a  $\langle kt^2 \rangle$  measurement: Cross section measurement

- 5 dimensional analysis
  - Avoid model dependence bias in the analysis
- Binning
  - $Q^2$  : 10 bins in [1,4]
  - $X_b$  : 8 bins in adaptative range
  - $Z_h$  : 8 bins in [0.2,0.8]
  - $P_t^2$ : 8 bins in [0,1]
  - $\phi_h$  : 12 bins in [-180,180]
  - Total : 61440 five dimensional bins



# Cross section measurement

Differential cross section calculation:

$$\left( \frac{d\sigma}{dQ^2 dX_b dZ_h dP_t^2 d\phi_h} \right)_{exp} = \frac{N_{corr}(Q^2, X_b, Z_h, P_t^2, \phi_h)}{\Delta Q^2 \Delta X_b \Delta Z_h \Delta P_t^2 \Delta \phi_h} \cdot \frac{1}{\mathcal{L}_{int}}$$

Where:

$$N_{corr}(Q^2, X_b, Z_h, P_t^2, \phi_h) = \frac{N_{data}(Q^2, X_b, Z_h, P_t^2, \phi_h)}{Acc_{5dim} \cdot RC_{5dim}}$$

$$\mathcal{L}_{int} = (\# \text{ Electrons in FC}) \times \left( \frac{\# \text{ Nucleonic targets}}{\text{Area}} \right)$$

# Acceptance

The acceptance is defined as:

$$Acc_{5dim} = \left( \frac{\# \text{ Reconstructed events}}{\# \text{ Generated events}} \right)_{(Q^2, X_b, Z_h, P_t^2, \phi_h)}$$

- To obtain generated events **PYTHIA 6.319** was used
- To obtain the reconstructed events **GSIM** was used
  - Ideal detailed simulation of the detector geometry
  - Afterwards post processor is used to reproduce the inefficiencies of the detector.
- Simulations were made by Hayk Hakobyan

# Radiative Corrections

The Radiative Correction factor is defined as:

$$RC_{5dim} = \left( \frac{\sigma_{obs}}{\sigma_{Born}} \right)_{(Q^2, X_b, Z_h, P_t^2, \phi_h)}$$

- To calculate this factor a C++ version, developed by Hayk Hakobyan, of the HAPRAD code was used
- HAPRAD is based on the QED calculations of Akushevich et al.[2]
- It's specially designed for the unpolarized SIDIS case

# Results

As a start we calculate experimentally:

$$\frac{d\sigma}{dQ^2 dX_b dZ_h}$$

**Why?** Because theoretically is more simple than the five-fold differential cross section, and there's results to check if our methodology is right

**How?** Integrating experimentally in  $P_t^2$  and  $\phi_h$

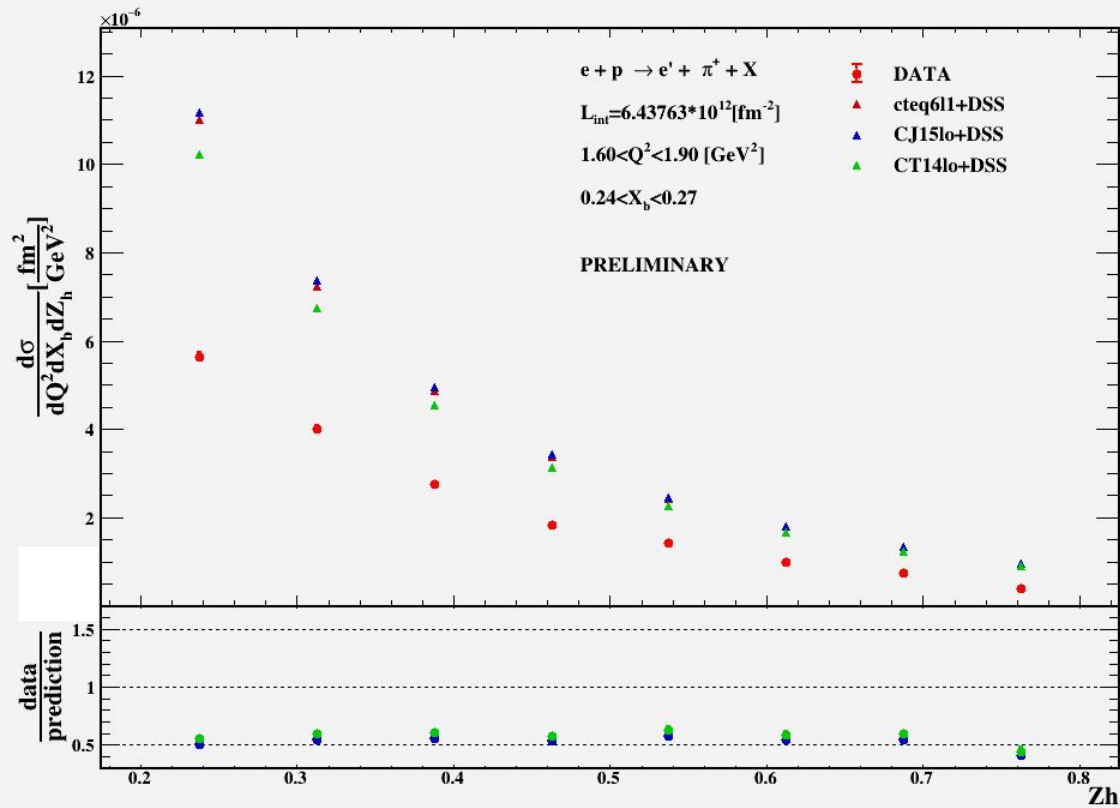
# Results

We'll compare our results to the theoretical expression derived from [3,4]:

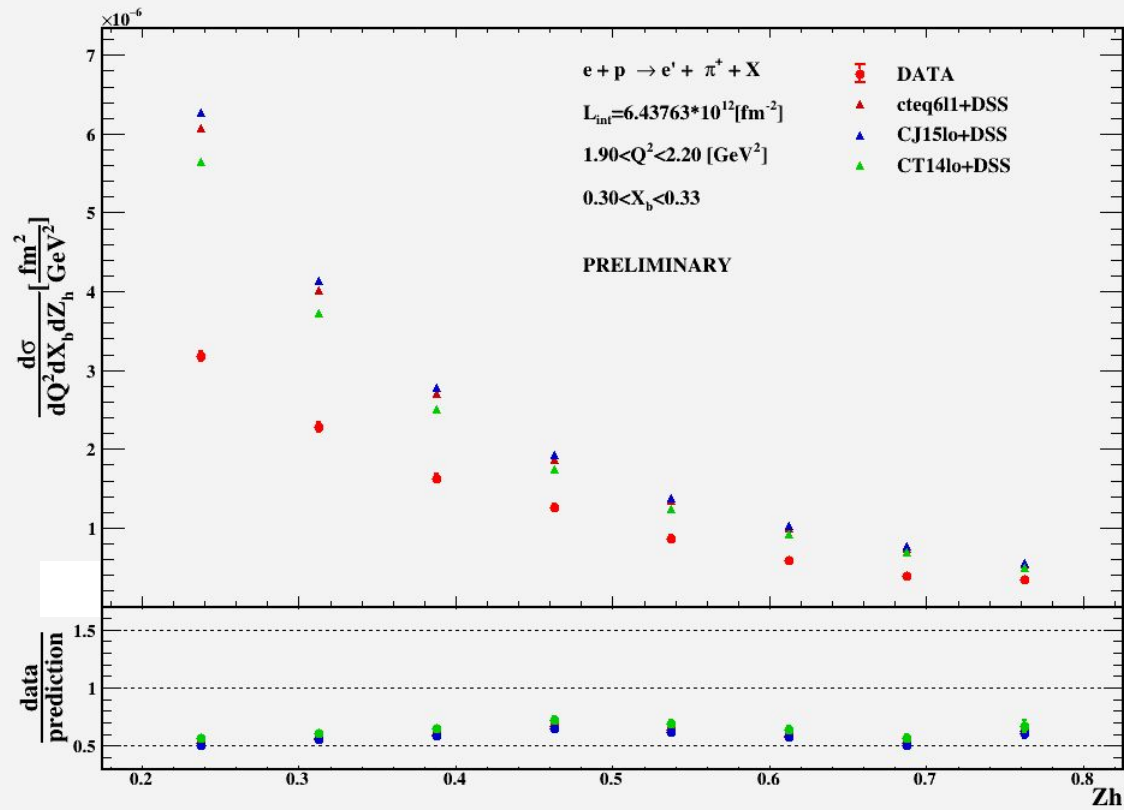
$$\frac{d\sigma}{dQ^2 dX_b dZ_h} = \frac{4\pi\alpha}{Q^4} \left( 1 - y - \frac{M \cdot X_b \cdot y}{2E} + \frac{y^2}{2\kappa \cdot R} \right) \sum_q e_q^2 f_q(X_b) \cdot D_q^{\pi^+}(Z_h)$$

- Only statistical errors will be included
- Only central values in the theory

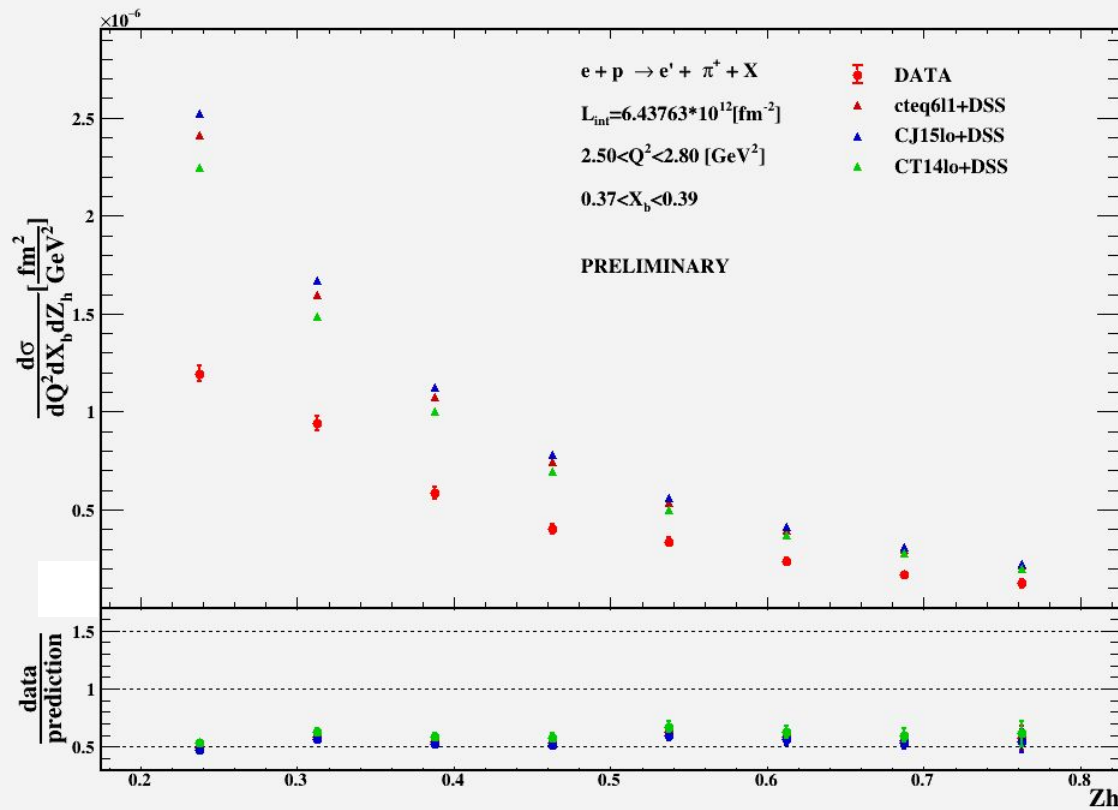
# Results: Proton target



# Results: Proton target



# Results: Proton target





# References

1. Anselmino, M. et al., *“Role of Cahn and Sivers effects in Deep Inelastic Scattering”*, 2005.
2. Akushevich, I. et al., *“Radiative Effects in the Process of Hadron Electroproduction”*, 1999.
3. Levelt, J., Mulders, P.J., *“Quark Correlation Functions in Deep Inelastic Semi-Inclusive Processes”*, 1993.
4. Osipenko, M. et al., *“Measurement of Semi-Inclusive  $\pi^+$  Electroproduction off the Proton”*, 2009.