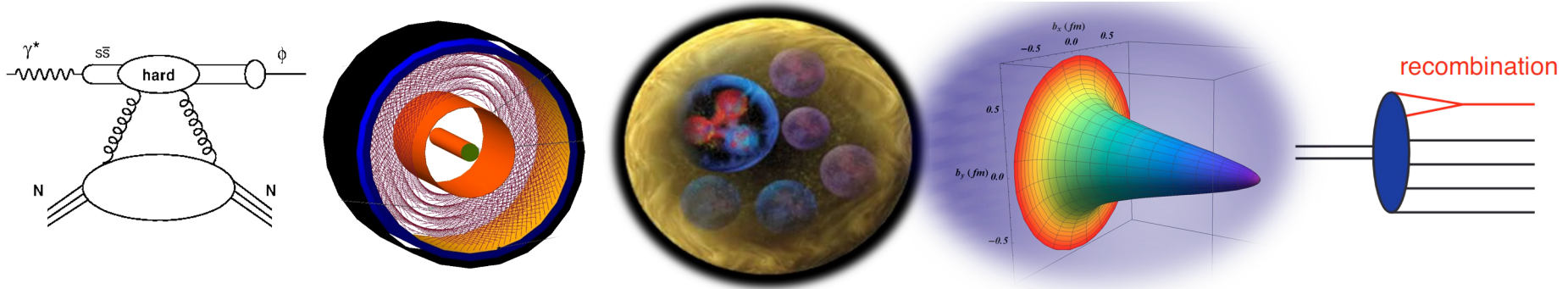


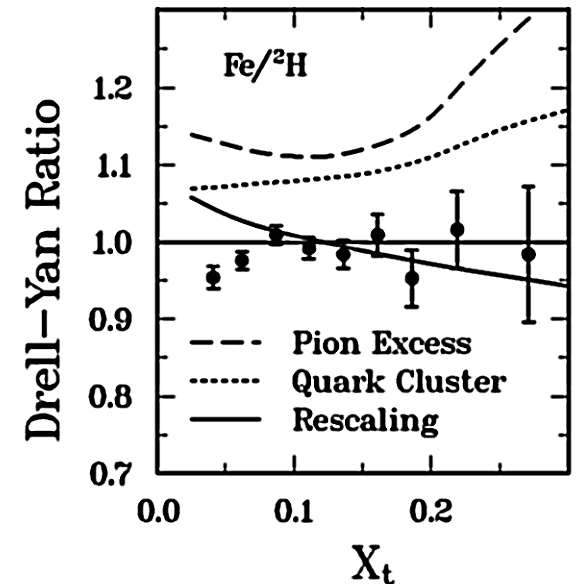
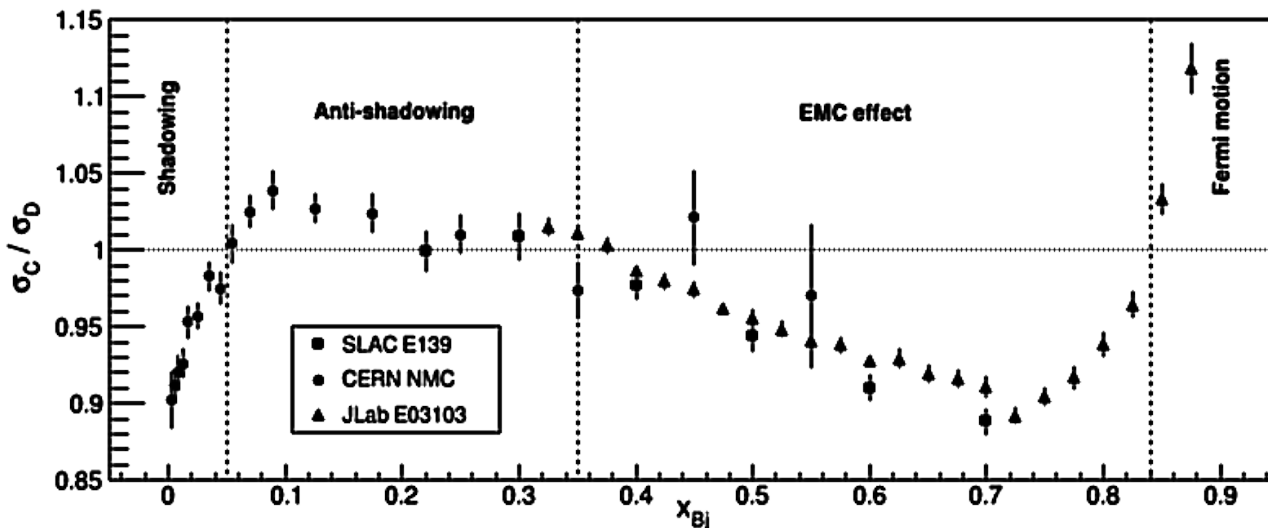
Nuclear TMDs with CLAS12



Proposition for a RG-D Run Group Proposal

R. Dupré for the authors
*W. Armstrong, L. El Fassi,
ZE. Meziani, H. Szumila-Vance et al.*

The Nuclear Effects



We discovered nuclear effects at the quark level

- Shadowing, anti-shadowing and EMC effect

The EMC effect remains a mystery to this day

- Meson content induced by NN interaction
- 6, 9, 12-quark clusters
 - Both are excluded by Drell-Yan measurements
- Nucleon size might change → bound FF
 - Difficult to prove due to FSI effects
- Q^2 - or x -rescaling with widely different physical meaning

Resolving the EMC Effect Mystery

Higher precision

- **Performed in JLab Hall-C**
- **Tough to compete with CLAS12 on this front**

New processes

- **Tagging/SRC (ALERT, BAND, Bonus)**
- **Nuclear DVCS (ALERT)**
- **Nuclear TMDs (This talk !)**

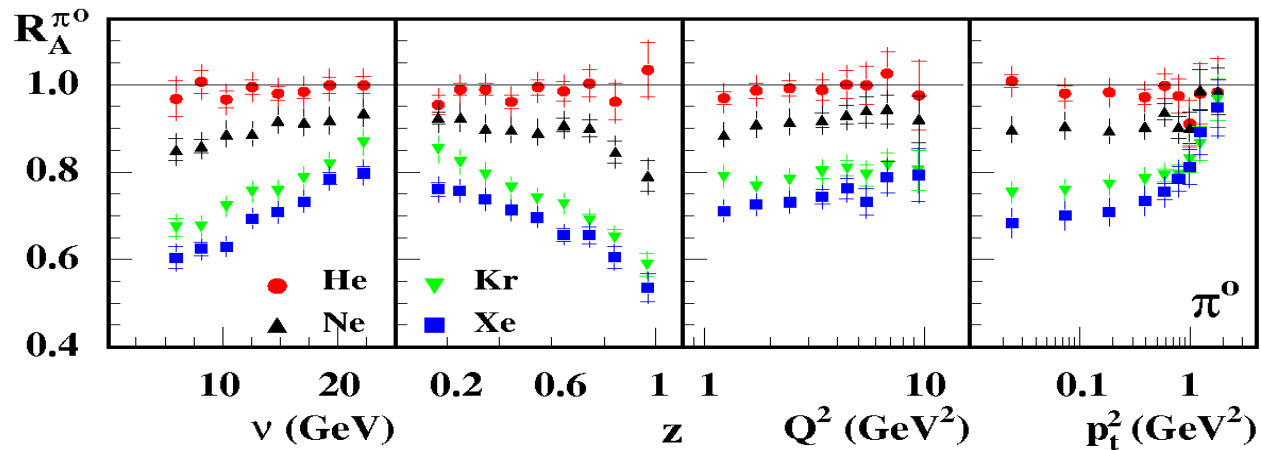
Why nuclear TMDs ?

- **The missing piece of the nucleus description**
 - *We observed surprising behavior for the GPDs*
- **Involves the fragmentation functions**
 - *These are affected by the medium*
 - *In particular with the transverse momentum broadening*
 - *The TMD framework can help treat consistently the data*
- **It will modernize the way we study nuclear SIDIS**

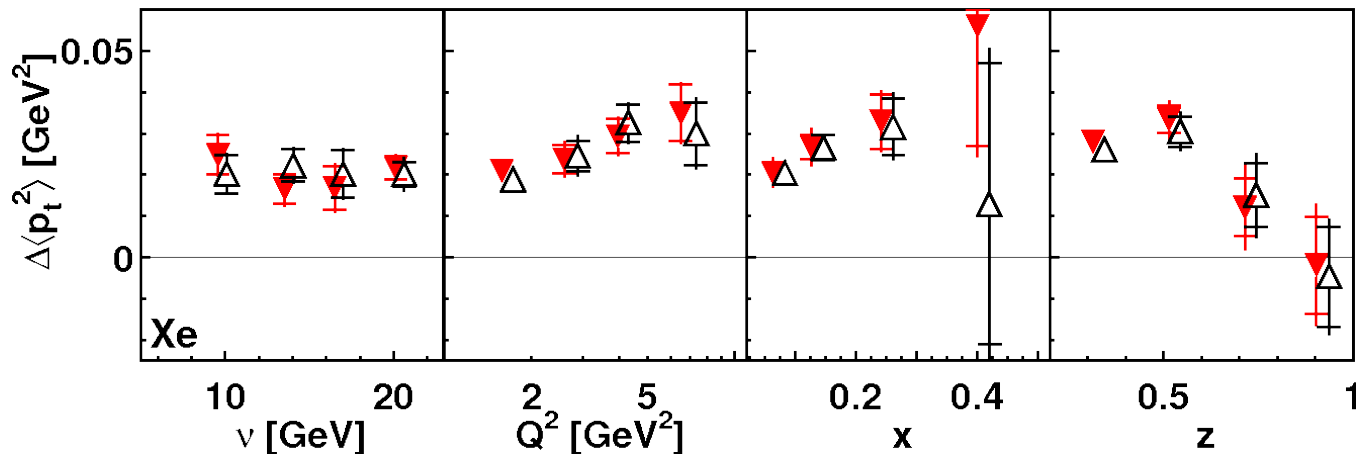
The HERMES data

Multiplicity ratio

– Hadron absorption



Transverse momentum broadening



Extracting Signal of the TMDs

TMD extraction is simple, in principle

- Each function has a different modulation
- It is a bit more complicated
 - Resolving the convolution with fragmentation functions

Experimental needs

- High acceptance
 - CLAS12 !
- Polarized beam
 - Easy in JLab
- Polarized targets
 - Probably not anytime soon for nuclear targets

$$\begin{aligned}
 \frac{d\sigma}{dx_B dy d\phi_S dz d\phi_h dP_{h\perp}^2} &= \frac{\alpha^2}{x_B y Q^2} \frac{y^2}{2(1-\varepsilon)} \\
 &\times \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 &\quad + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 &\quad + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 &\quad + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 &\quad + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 &\quad \quad + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 &\quad \quad \left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \\
 &\quad + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 &\quad \quad \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \left. \right\}.
 \end{aligned}$$

Nuclear TMD

Very little theory on the topic

- **Using the model from Liang et al. (PRD 77 125010)**

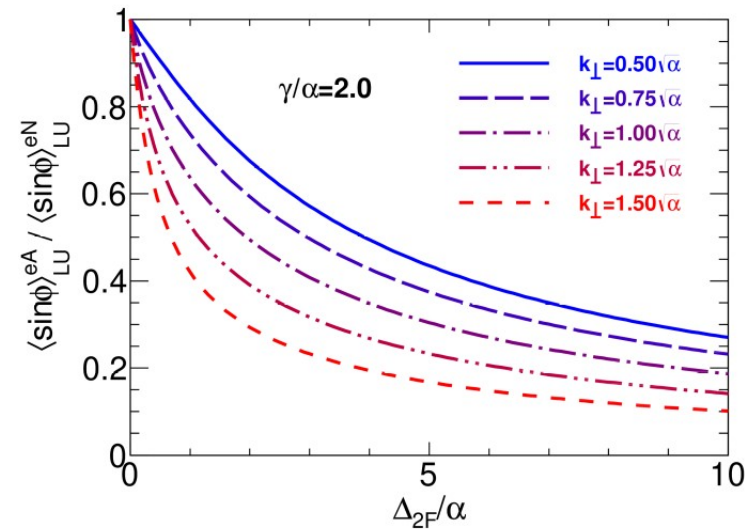
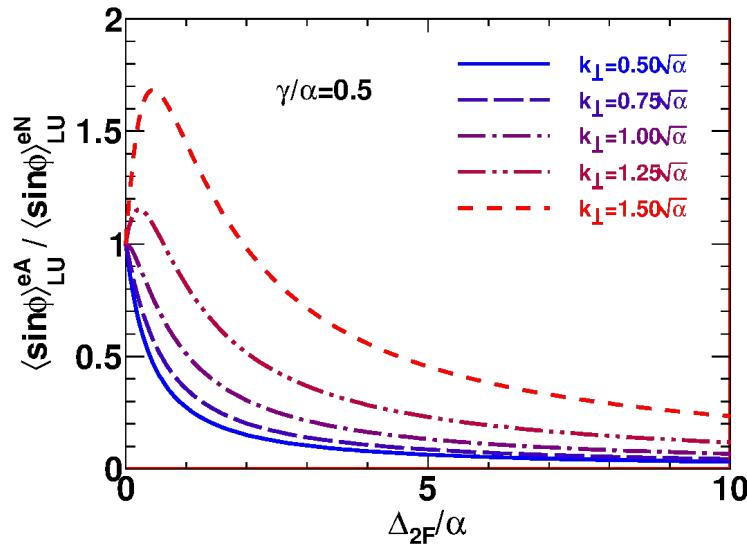
$$f_q^A(x, k_\perp) \approx \frac{A}{\pi \Delta_{2F}} \int d^2 \ell_\perp e^{-(\vec{k}_\perp - \vec{\ell}_\perp)^2 / \Delta_{2F}} f_q^N(x, \ell_\perp)$$

$$\Delta_{2F} = \int d\xi_N^- \hat{q}_F(\xi_N)$$

$$\hat{q}_F(\xi_N) = \frac{2\pi^2 \alpha_s}{N_c} \rho_N^A(\xi_N) [x f_g^N(x)]_{x \rightarrow 0}$$

- **Uses the transport coefficient of the nuclear matter**
 - *This nuclear property is poorly known, we have here the chance to measure it properly*
 - *It is directly linked to the saturation scale and should vary with A*
- **Asymmetries are generated at the partonic level**
 - *Links cleanly to the definition of the transport coefficient*

Using TMDs for Hadronization



Model give wide predictions

- α and γ are the width of the TMDs

Usual hadronization measurements use outdated methods

- We should use the TMD framework to study semi-inclusive DIS on nuclei
- The sin and cos moments give direct parton level sensitivity to the transport coefficient

Two independent transport coefficient measurements

- To be compared with the absorption and the transverse momentum broadening

Projections

We ran full simulations

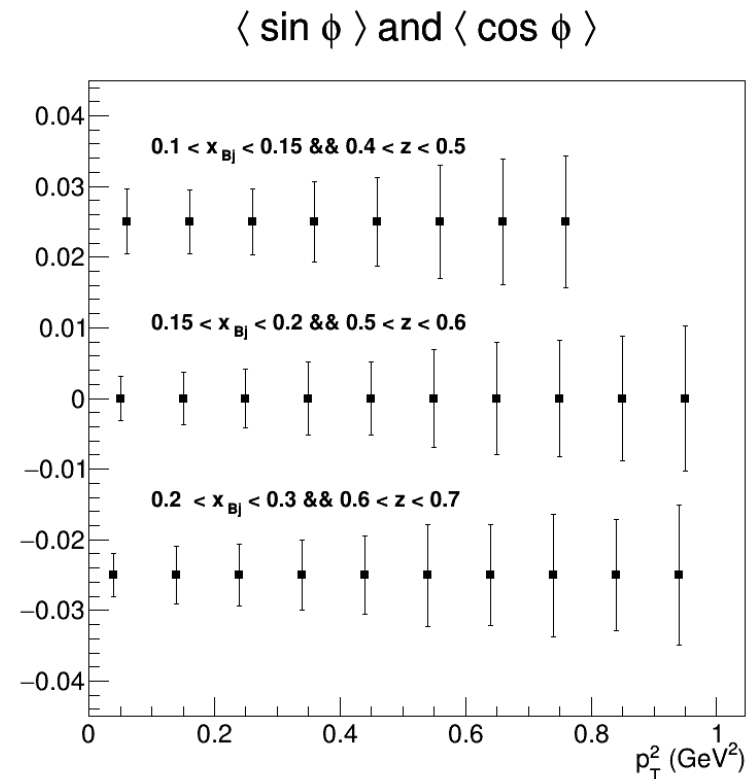
- HERMES generator `gmc_trans`
- Pass it to GEMC & Full CLAS12 reconstruction

Results for the sin and cos moment

- Identical because only statistical
- Smaller reach in P_t than for the RG-A proton target
- Logical since the beam time is much smaller

This is enough for a first measurement

- Ratios to the proton measurement are key
 - *The amount of systematic uncertainty that can be canceled will be key*
- We will use the different targets to test if a A dependence can be detected



Summary

We have studied nuclear hadronization and the EMC effect for a long time

- Both topics remain very active to this day
- Using TMD framework is a modernization of the hadronization studies as it is accounting for all components of the SIDIS cross section

The TMD framework will help progress in our understanding

- Asymmetries are generated at the partonic and level allowing a clean interpretation of the data
- Different asymmetries allow to cross check the results from the same data set with different observables
- Different nuclei allow to assess the expected variation of the transport coefficient

The Proposal

- Run group proposal that will run within the RG-D
- We only request the addition of polarization
- We know these data are prone to data mining, it is important to highlight this potential for the jeopardy process