



**University of New Hampshire**  
**Nuclear & Particle Physics Group**

# **The $g_2p$ Experiment: A Measurement of the Proton's Spin Structure Functions 2020 Status Update**

**David Ruth**

**Hall A Collaboration Meeting**

**January 30, 2020**

**Some Slides & Figures by Ryan Zielinski**

# ESSENTIAL QUANTITIES IN $ep$ SCATTERING

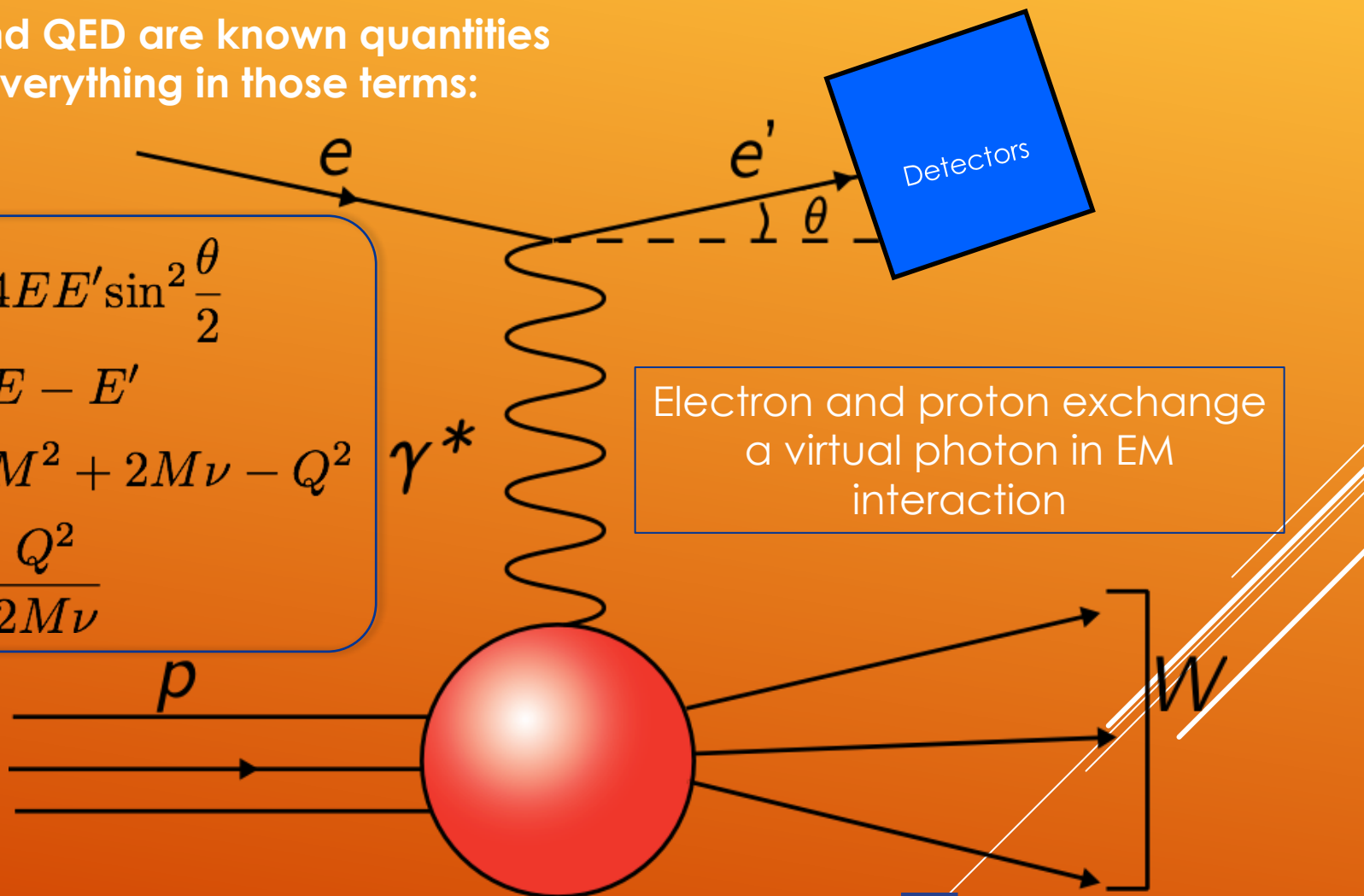
Electron and QED are known quantities  
so define everything in those terms:

$$Q^2 = 4EE' \sin^2 \frac{\theta}{2}$$

$$\nu = E - E'$$

$$W^2 = M^2 + 2M\nu - Q^2$$

$$x = \frac{Q^2}{2M\nu}$$



# INCLUSIVE $ep$ SCATTERING CROSS SECTIONS DESCRIBE NORMALIZED INTERACTION RATE

Elastic scattering: target remains in the ground state after interaction

$$E'_{\text{elas}} = \frac{E}{1 + \frac{2E}{M} \sin^2 \frac{\theta}{2}}$$

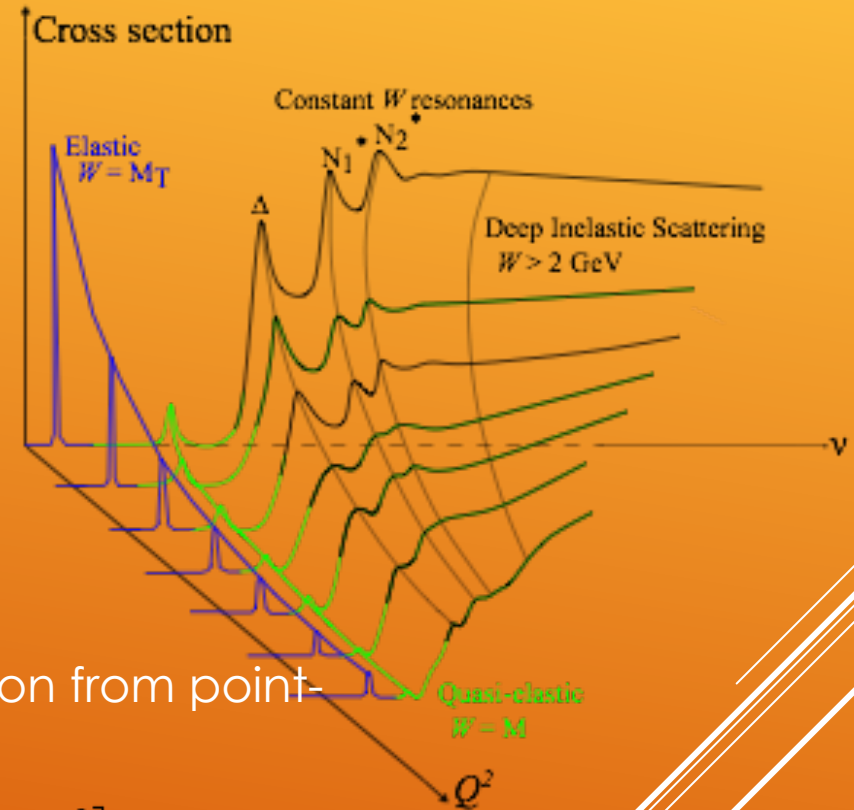
Mott cross section describes scattering from point-particle:

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} = \frac{\alpha^2}{4 E^2 \sin^4 \frac{\theta}{2}} \cos^2 \frac{\theta}{2}$$

Rosenbluth cross section describes deviation from point-particle:

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} \left[ \frac{G_E^2(Q^2) + \tau G_M^2(Q^2)}{1 + \tau} + 2\tau G_M^2(Q^2) \tan^2 \frac{\theta}{2} \right]$$

$G_E$  and  $G_M$  related to charge and current distributions



# INCLUSIVE $ep$ SCATTERING CROSS SECTIONS DESCRIBE NORMALIZED INTERACTION RATE

Inelastic scattering: Target is in excited state after interaction

Structure Functions:

Inclusive *unpolarized* cross sections

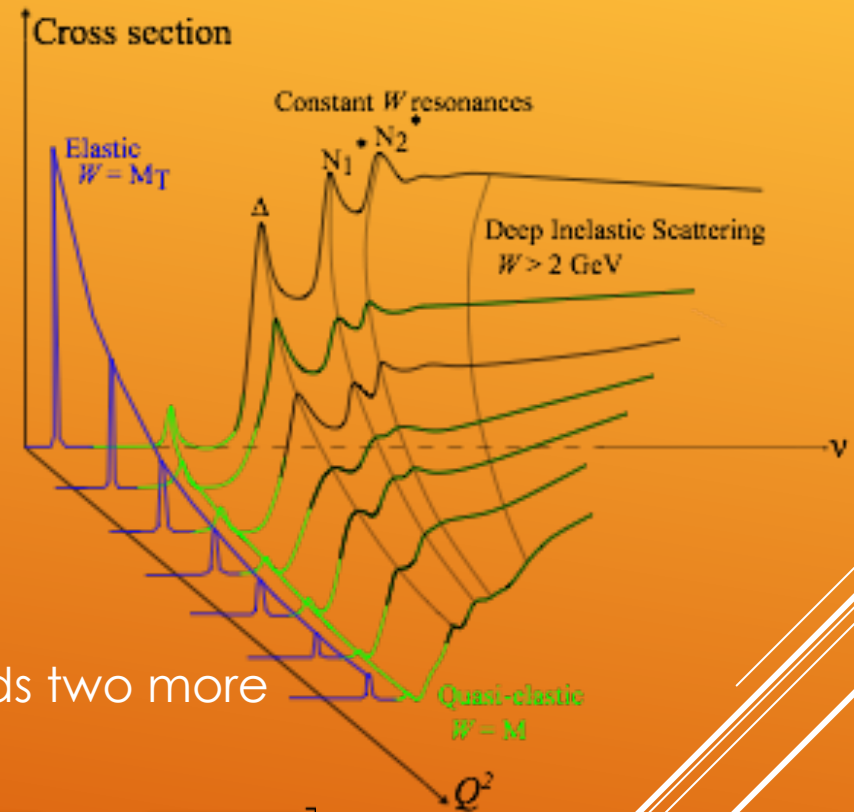
$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{\text{Mott}} \left[ \frac{1}{\nu} F_2(x, Q^2) + \frac{2}{M} F_1(x, Q^2) \tan^2 \frac{\theta}{2} \right]$$

$F_1$  and  $F_2$  related to quark/gluon distribution

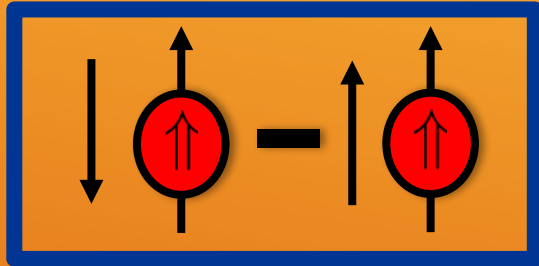
Adding a *polarized* beam and target adds two more structure functions

$$\frac{d^2\sigma^\pm}{d\Omega dE'} = \sigma_{\text{Mott}} \left[ \alpha F_1(x, Q^2) + \beta F_2(x, Q^2) \pm \gamma g_1(x, Q^2) \pm \delta g_2(x, Q^2) \right]$$

$g_1$  and  $g_2$  related to spin distribution



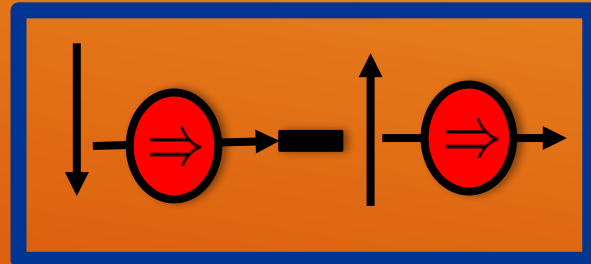
# EXTRACTING SPIN STRUCTURE BY LOOKING AT CROSS SECTION DIFFERENCES



Parallel

Inclusive *polarized* cross sections

$$\frac{d^2\sigma^{\uparrow\uparrow}}{dE'd\Omega} - \frac{d^2\sigma^{\downarrow\uparrow}}{dE'd\Omega} = \frac{4\alpha^2}{M\nu Q^2} \frac{E'}{E} \left[ g_1(x, Q^2) \{E + E' \cos\theta\} - \frac{Q^2}{\nu} g_2(\nu, Q^2) \right]$$



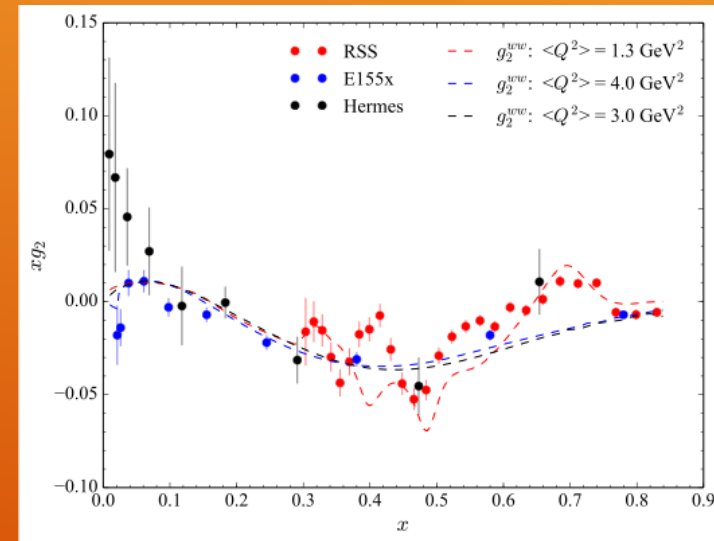
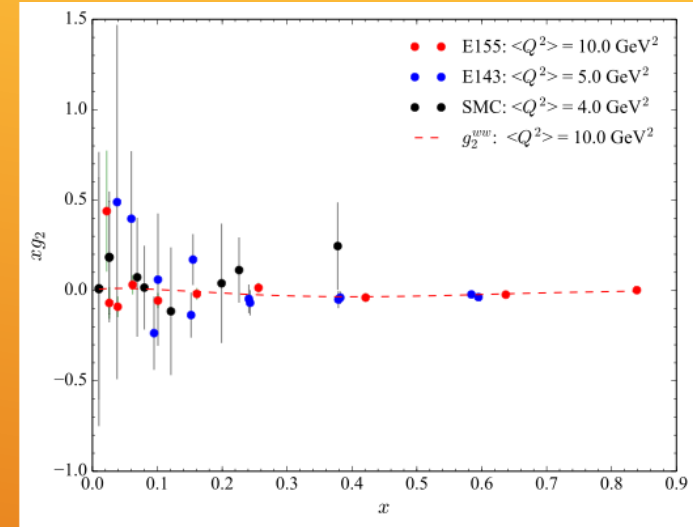
Perpendicular

$$\frac{d^2\sigma^{\uparrow\Rightarrow}}{dE'd\Omega} - \frac{d^2\sigma^{\downarrow\Rightarrow}}{dE'd\Omega} = \frac{4\alpha^2}{M\nu Q^2} \frac{E'^2}{E} \sin\theta \left[ \nu g_1(x, Q^2) + 2Eg_2(\nu, Q^2) \right]$$

Two equations, two unknowns...

# MOTIVATION: Measure a fundamental spin observable ( $g_2$ ) in the region $0.02 < Q^2 < 0.20 \text{ GeV}^2$ for the first time

- Measurements at Jefferson Lab:
  - RSS – medium  $Q^2$  ( 1-2  $\text{GeV}^2$ ) (published)
  - SANE – high  $Q^2$  (2-6  $\text{GeV}^2$ ) (analysis)
  - **$g_2\text{p}$  – low  $Q^2$  (0.02-0.20  $\text{GeV}^2$ ) (analysis)**
- **Low  $Q^2$  is difficult:**
  - Electrons strongly influenced by target field
  - Strong kinematic dependence on observables
- **Low  $Q^2$  is useful:**
  - Test predictions of Chiral Perturbation Theory ( $\chi\text{PT}$ )
  - Test sum rules and measure moments of  $g_2$
  - Study finite size effects of the proton



- **$g_2\text{p}$  experiment ran spring 2012 in Hall A**

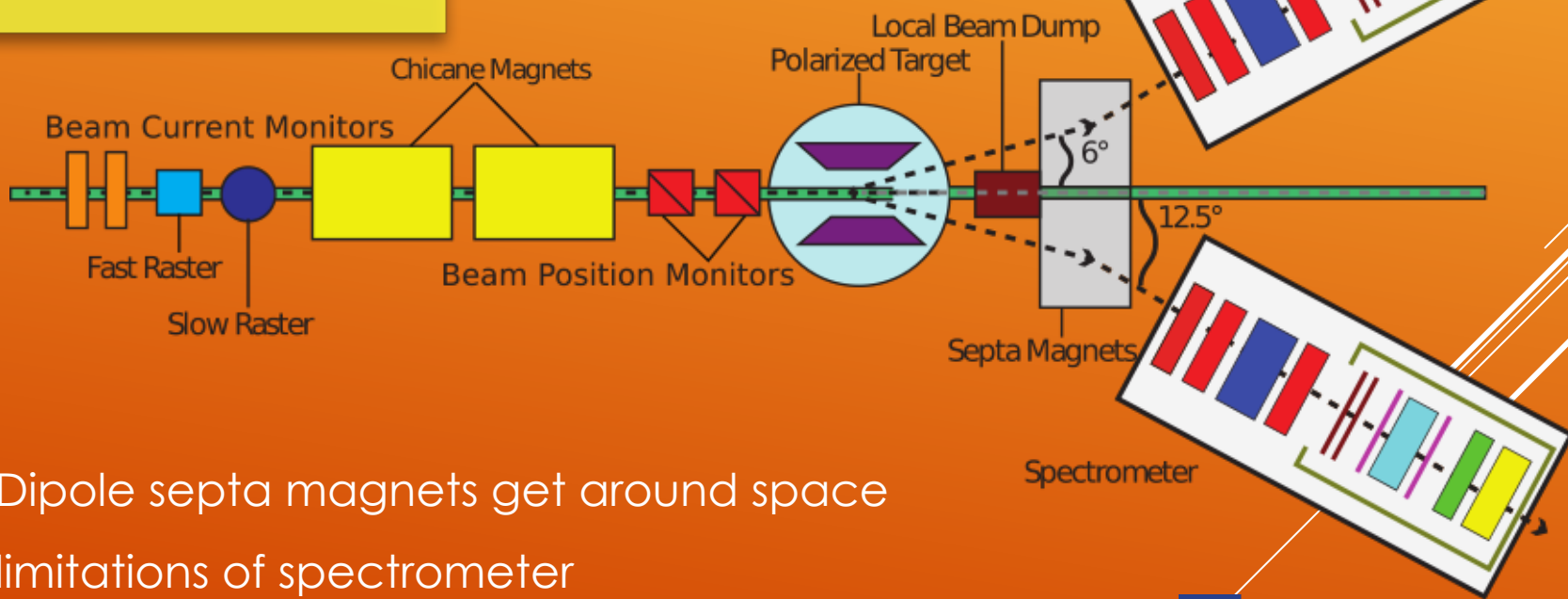


# HALL A EXPERIMENTAL SETUP:

Measuring  $g_2^p$

- Electron Beam
- Polarized Proton Target
- Spectrometer/Detectors
- Small Scattering Angle

- Transverse polarized  $NH_3$  target (2.5/5.0T)
- Dipole chicane magnets help compensate for target field bending of beam



- Dipole septa magnets get around space limitations of spectrometer

# POLARIZED PROTONS CREATED WITH DYNAMIC NUCLEAR POLARIZATION (DNP)

## Creating initial polarization:

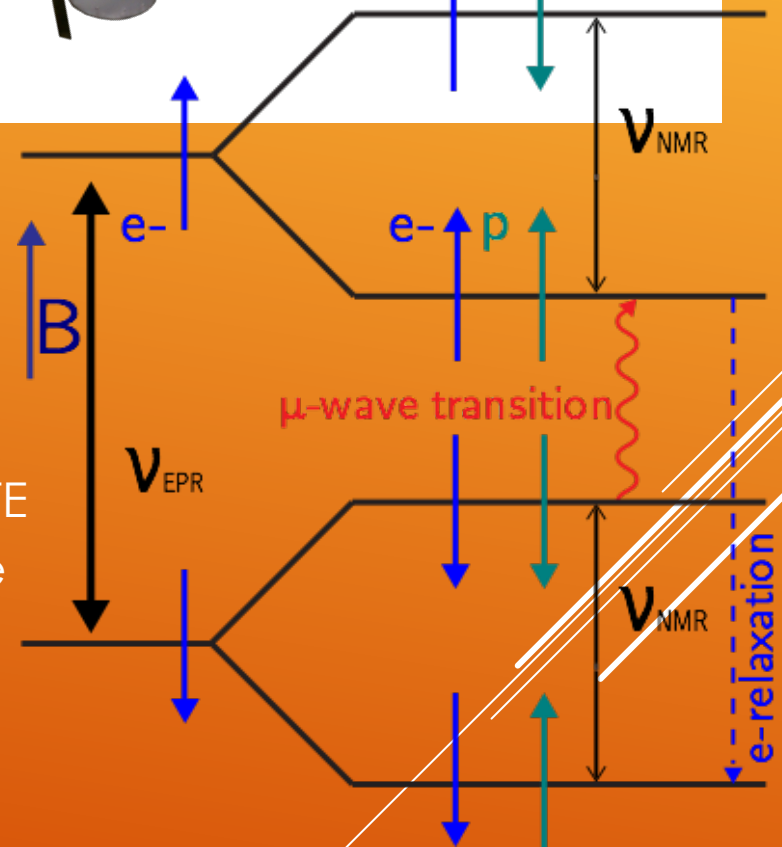
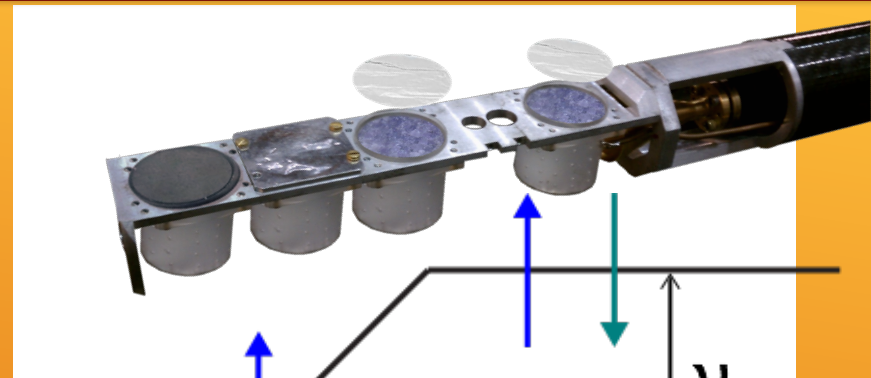
- Align spins in large B and low T
  - 5.0 T / 2.5 T @ 1 K

$$P_{TE} = \frac{N^\uparrow - N^\downarrow}{N^\uparrow + N^\downarrow} = \frac{e^{\frac{\mu B}{kT}} - e^{-\frac{\mu B}{kT}}}{e^{\frac{\mu B}{kT}} + e^{-\frac{\mu B}{kT}}}$$

- Large  $\mu_e$  ( $\sim 660\mu_p$ ) creates large electron polarization ( $\sim 99\%$  at 5T/1K)

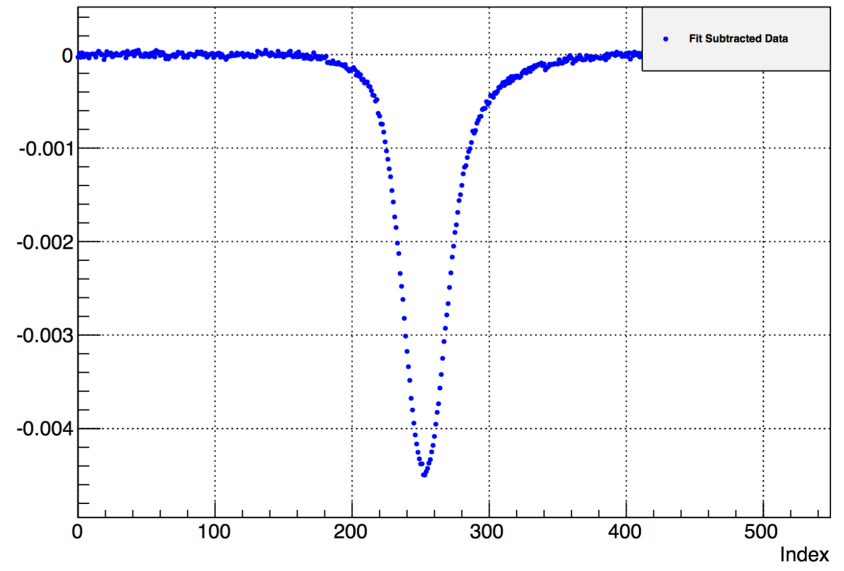
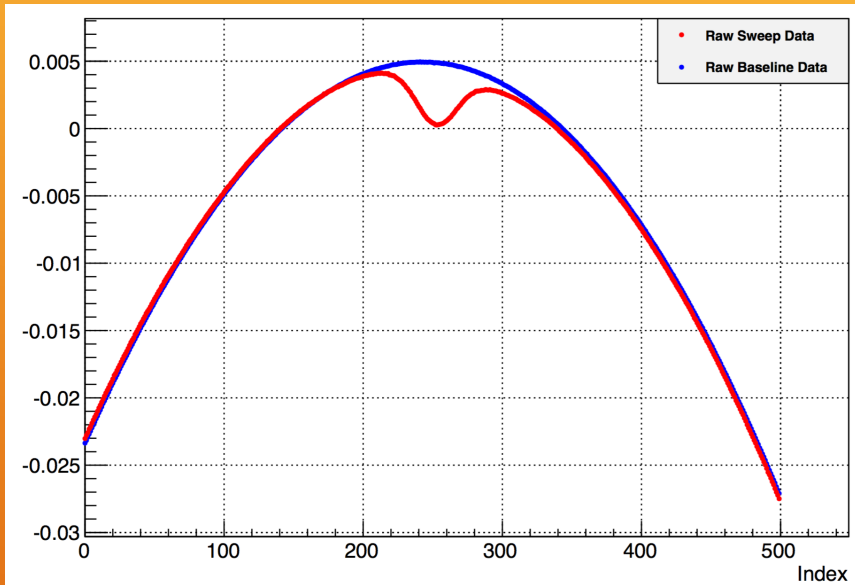
## Enhancing initial polarization:

- Proton pol. much smaller ( $\sim 0.5\%$  5T) at TE
- ep spin coupling and microwaves drive pol.
- Electrons relax much quicker than protons so polarization is sustained

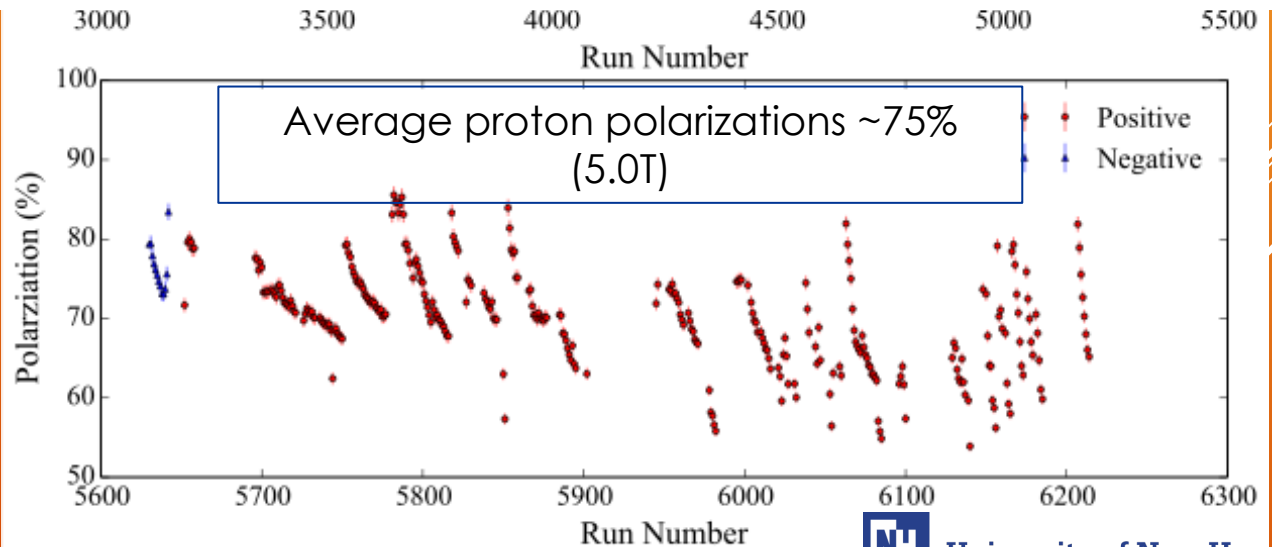




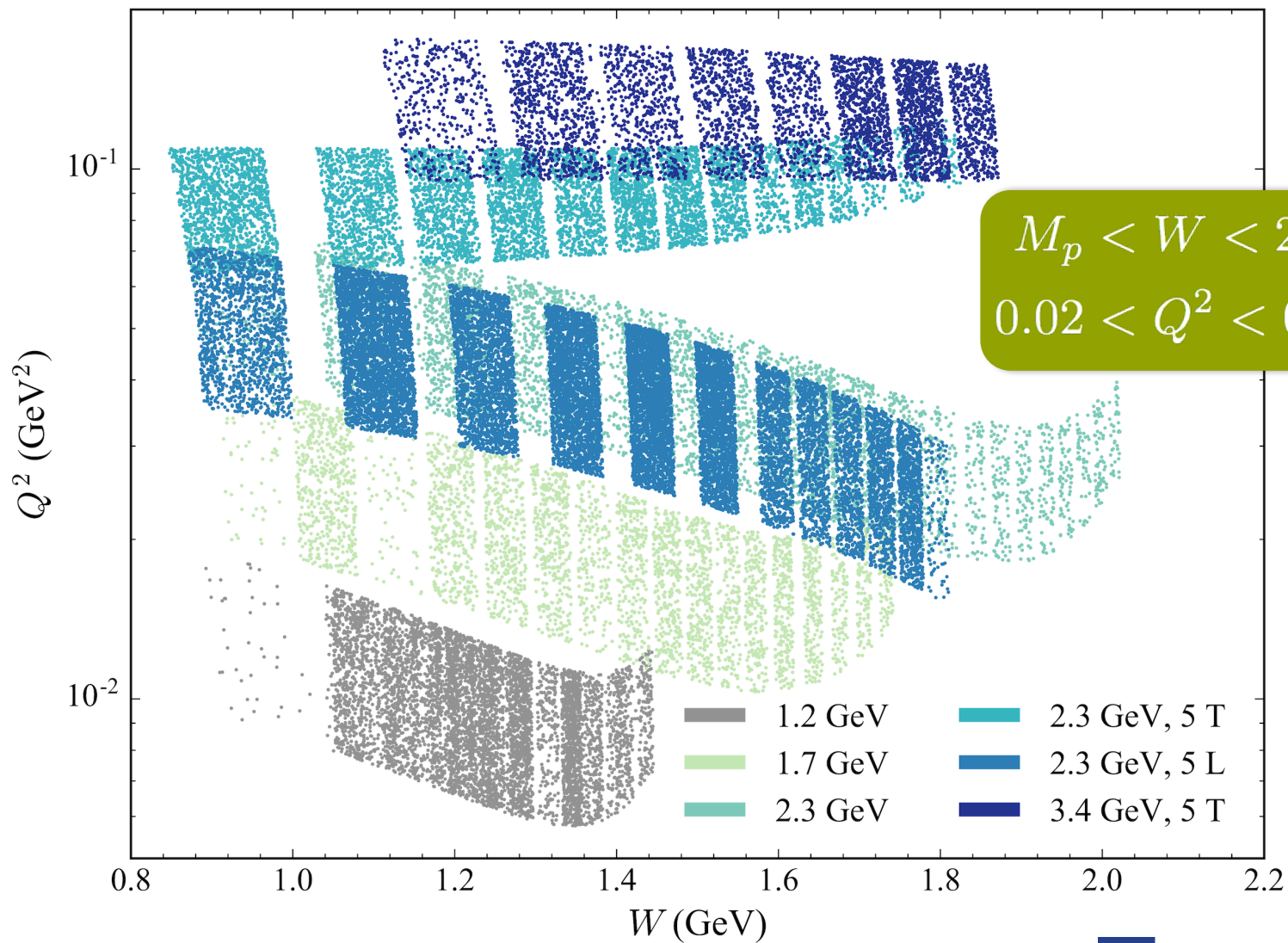
# PROTON POLARIZATION MEASURED WITH Q-METER



- LRC circuit where proton spin's couple with and change inductance



# $G_2P$ KINEMATIC COVERAGE



$M_p < W < 2 \text{ GeV}$   
 $0.02 < Q^2 < 0.20 \text{ GeV}^2$

# MEASURING $g_{1,2}$ FROM DATA

## What can we measure?

1. Helicity dependent asymmetries
2. Unpolarized cross sections
3. Polarized cross sections

$$A_{\perp} = \frac{\frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow - \uparrow\Rightarrow)}{\frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow + \uparrow\Rightarrow)}$$

2.

$$\sigma_0 = \frac{1}{2} \frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow + \uparrow\Rightarrow)$$

3.

$$\Delta\sigma_{\perp} = \frac{d^2\sigma}{d\Omega dE'}(\downarrow\rightarrow - \uparrow\rightarrow) = 2 \cdot A_{\perp} \sigma_0$$

Similar equation for parallel polarized cross section

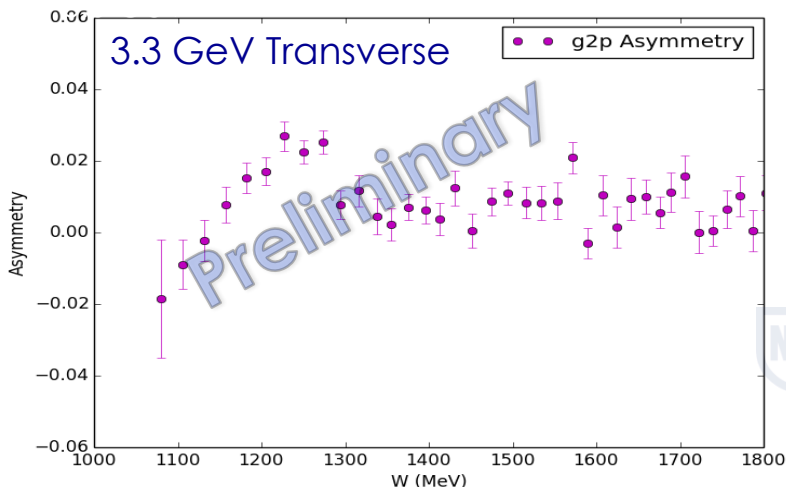
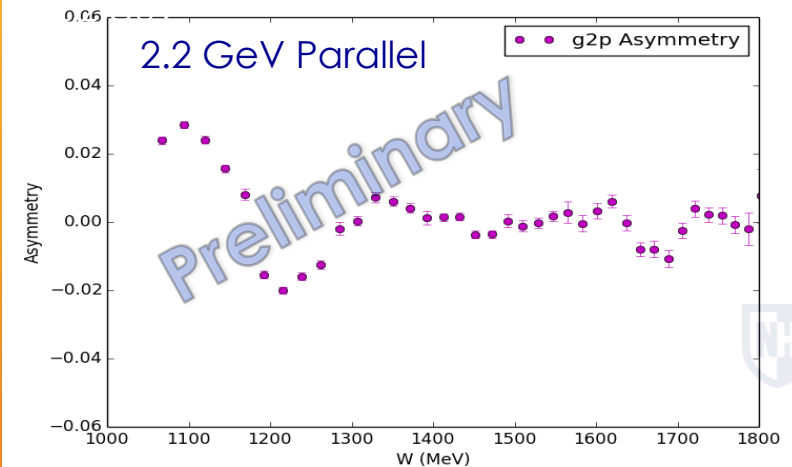
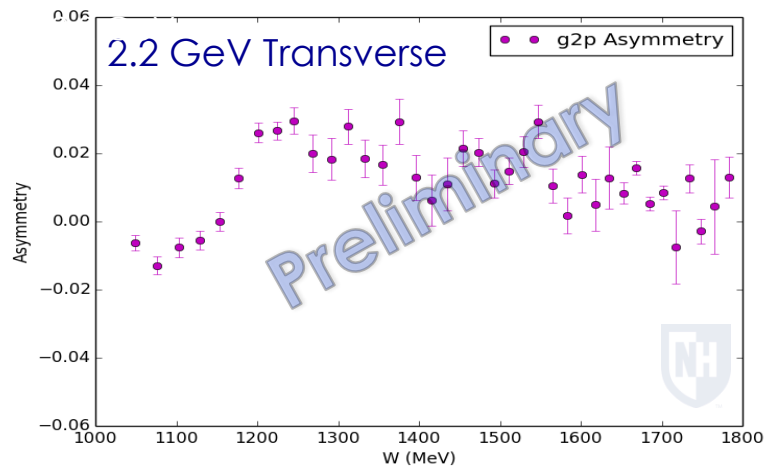
## Why do it this way?

- Asymmetries are easy to measure
- Lots of data on unpolarized cross sections so models are a possibility

Need to be mindful of contributions from scattering from anything other than protons



# 5T PROTON ASYMMETRIES



Raw Counts:

$$Y_{\pm} = \frac{N_{\pm}}{LT_{\pm}Q_{\pm}}$$

Measured Asymmetries:

$$A = \frac{Y_+ - Y_-}{Y_+ + Y_-},$$

$$A^{\text{exp}} = \frac{1}{f \cdot P_t \cdot P_b} A^{\text{raw}}$$

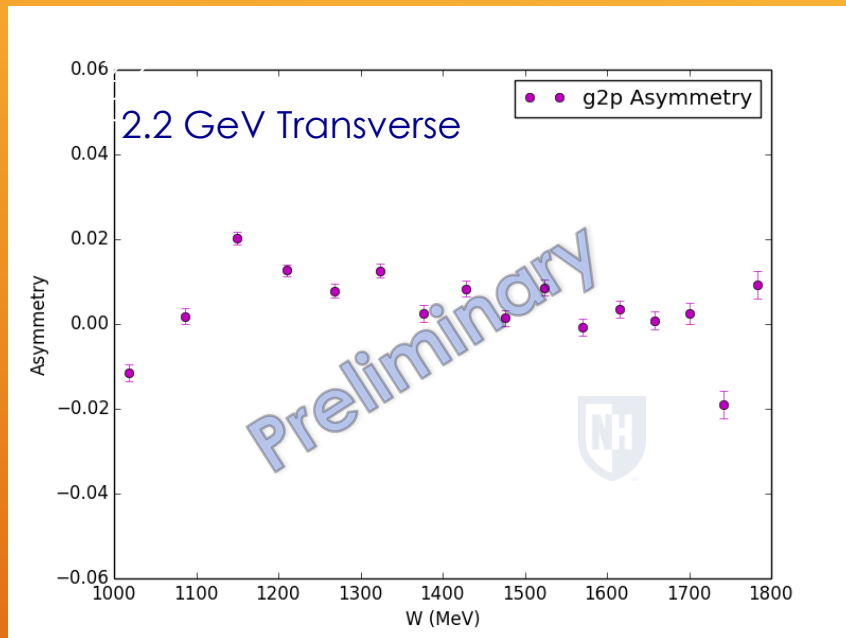
Combine both HRS for best statistics!

dilution factor

beam/target pol



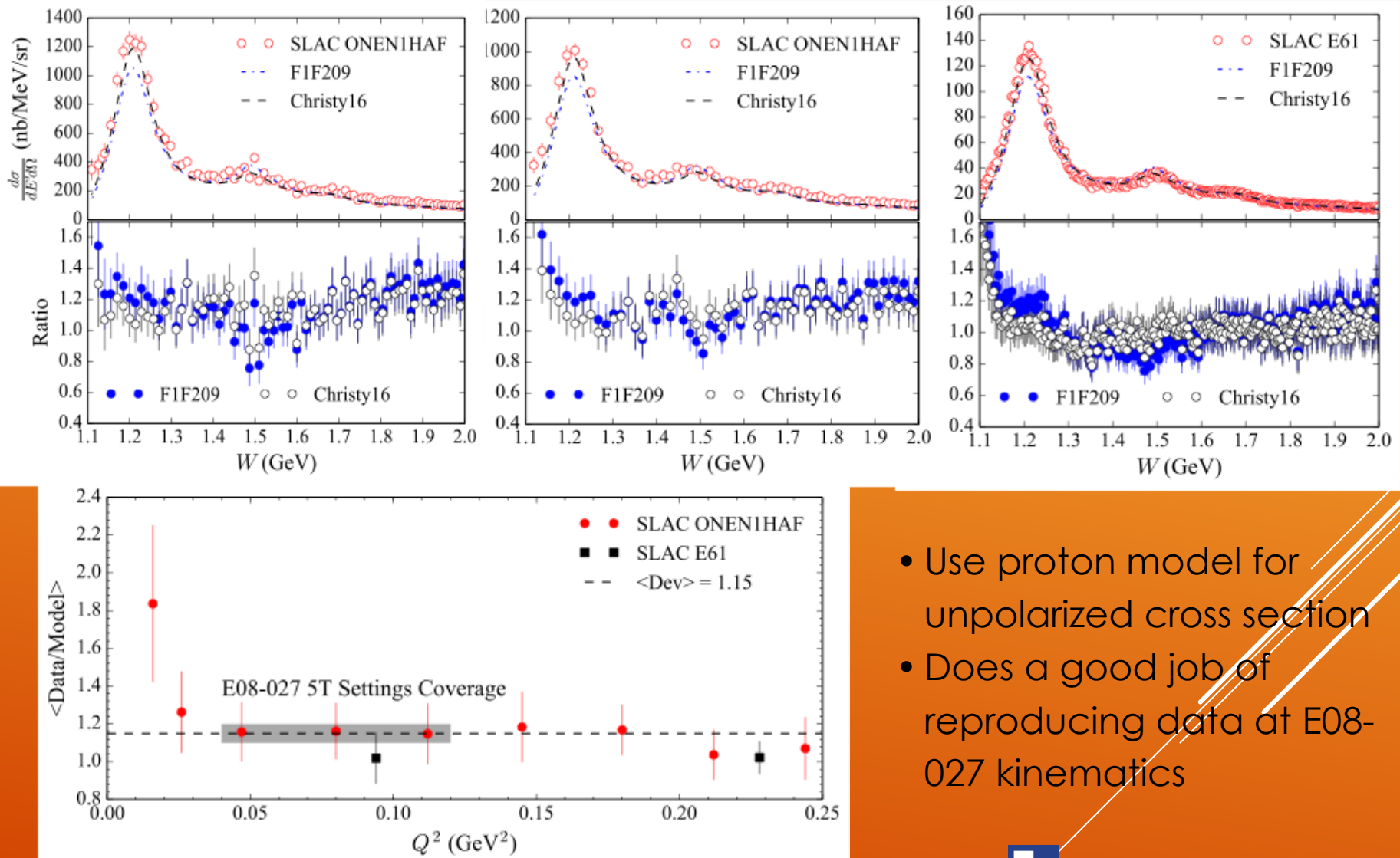
# 2.5T PROTON ASYMMETRIES



2.5T Data exists at 1.7 GeV and 1.1 GeV energy settings, but has large systematics that complicate analysis and will not be focused in initial publications



# MODEL CROSS SECTION

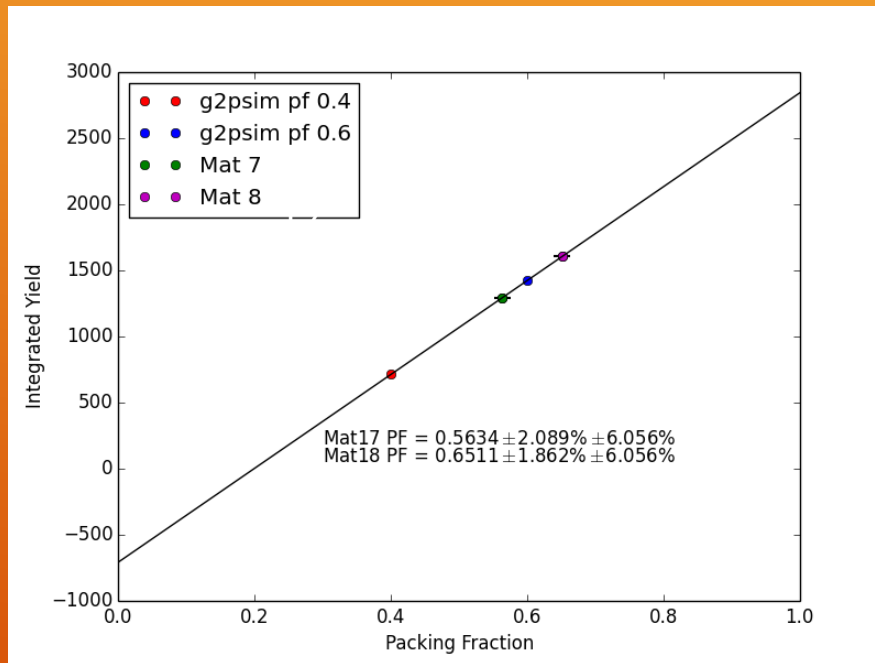
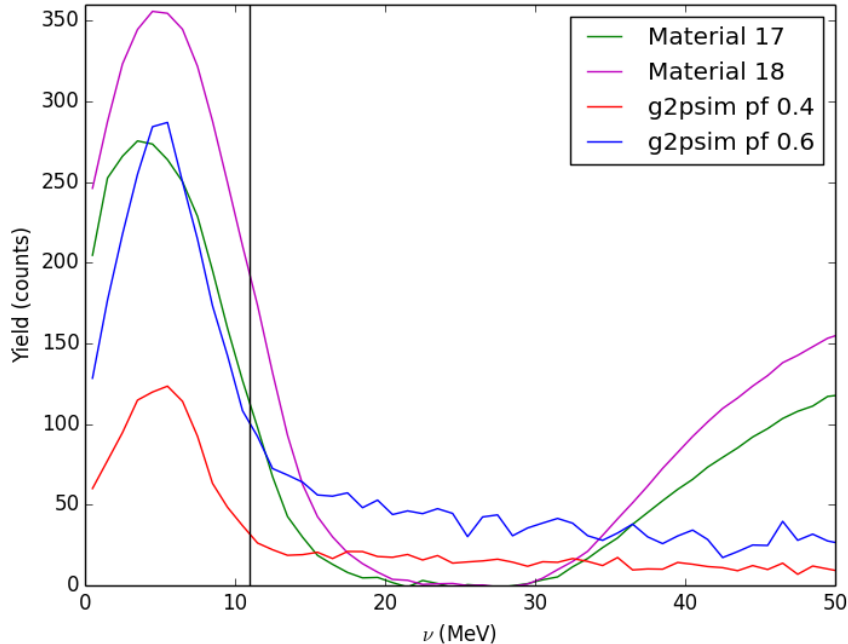


- Use proton model for unpolarized cross section
- Does a good job of reproducing data at E08-027 kinematics



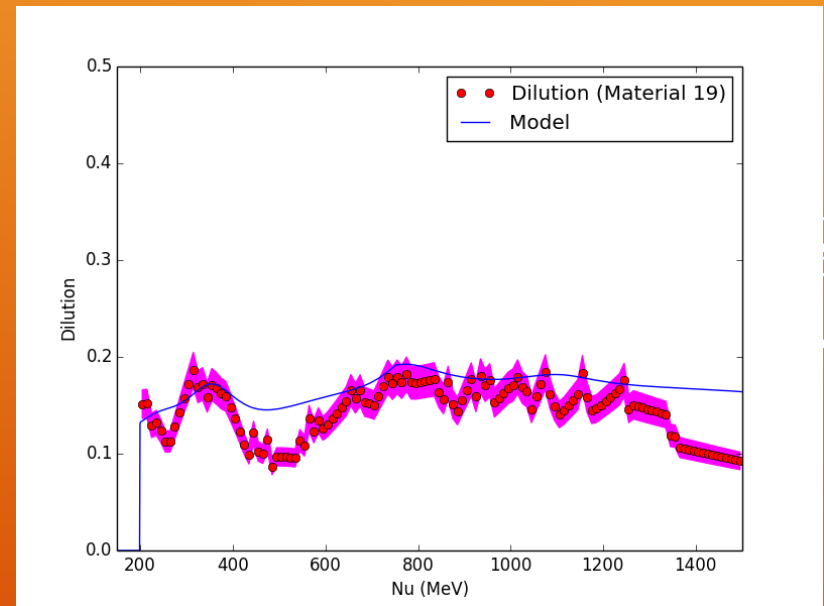
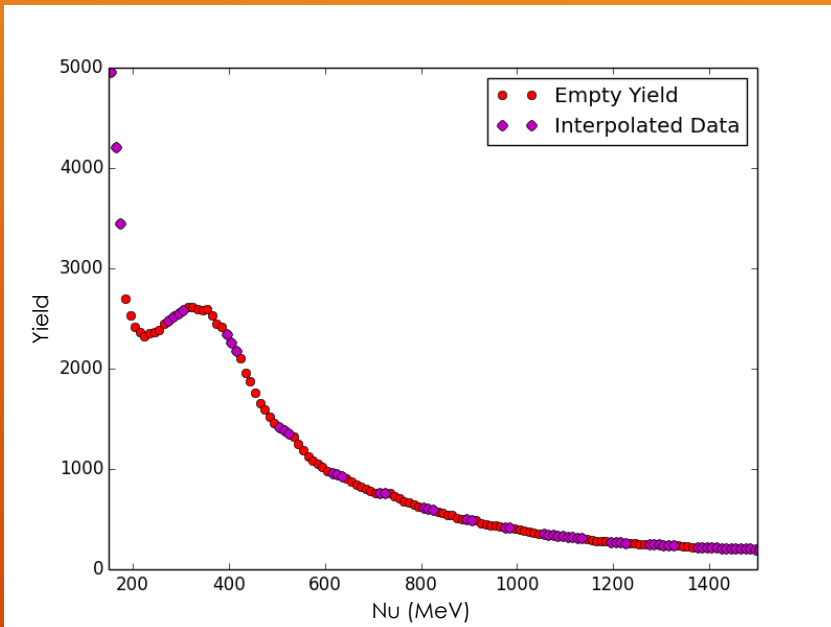
# NEW PACKING FRACTION ANALYSIS

- Packing fraction describes how much material is in the target cell, important for calculating dilution factor
- Previous packing fraction analysis yielded unrealistic results
- Analysis re-done with Oscar Rondon's method from RSS
- Because the Yield and Packing Fraction have a linear relationship, we can determine Packing Fraction by comparing integrated yield to simulated yields at known packing fractions



# NEW DILUTION ANALYSIS

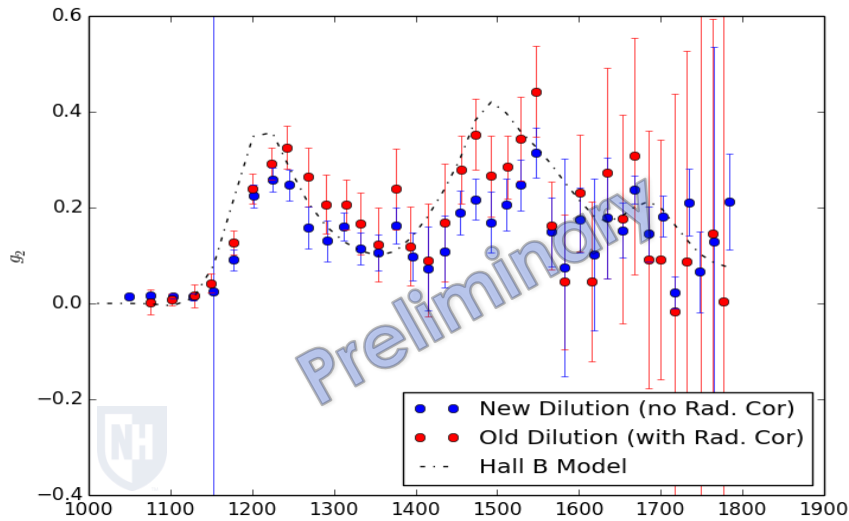
- Dilution approximates how much of data comes from other materials
- $f = \frac{\sigma_{Proton}}{\sigma_{Prod}} = 1 - \frac{Y_N + Y_{He} + Y_{Al}}{Y_{Prod}}$
- Acceptance effects on edge of momentum settings and BPM calibration issues complicated this analysis
- Sections missing above yields or including problematic data were replaced with the Bosted-Christy model, scaled to the data around it



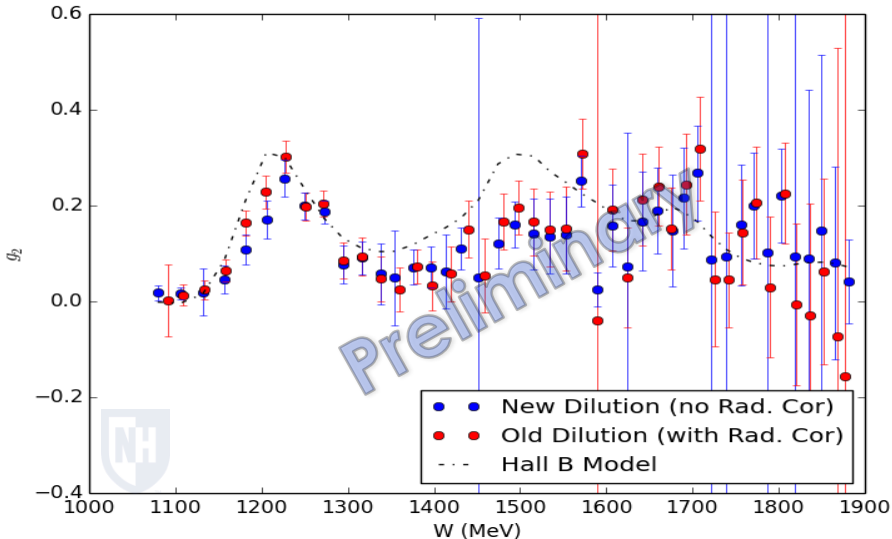


# EXTRACTING THE SPIN STRUCTURE FUNCTIONS : $G_2$

2.2 GeV Transverse



3.3 GeV Transverse



Model driven procedure for unmeasured part

$$g_2(x, Q^2) = \frac{K_1 y}{2} \left[ \Delta\sigma_{\perp} \left( K_2 + \tan\frac{\theta}{2} \right) \right] - \frac{g_1(x, Q^2)y}{2}$$

$$K_1 = \frac{MQ^2}{4\alpha} \frac{y}{(1-y)(2-y)}$$

$$K_2 = \frac{1 + (1-y)\cos\theta}{(1-y)\sin\theta}$$

Adjusting to a constant  $Q^2$

$$\delta_{\text{evolve}} = g_{1,2}^{\text{mod}}(x_{\text{data}}, Q_{\text{data}}^2) - g_{1,2}^{\text{mod}}(x_{\text{const}}, Q_{\text{const}}^2),$$

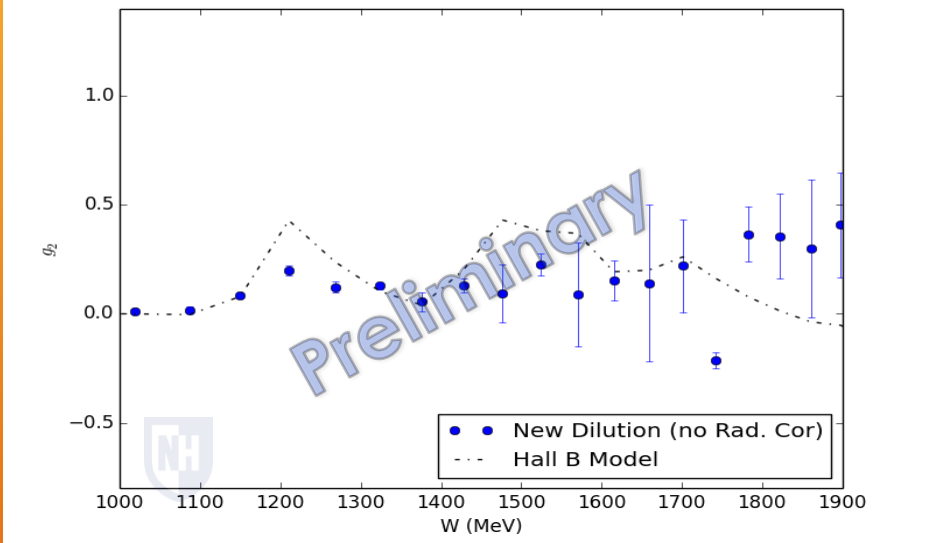
$$x_{\text{const}} = Q_{\text{const}}^2 / (W^2 - M^2 + Q_{\text{const}}^2);$$

Small effect at the transverse settings



# 2.5T G2, RADIATIVE CORRECTIONS INCOMPLETE

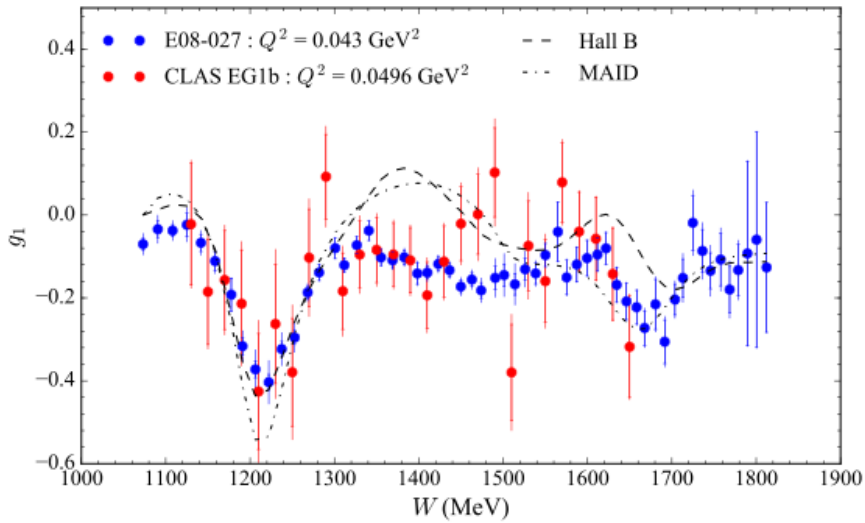
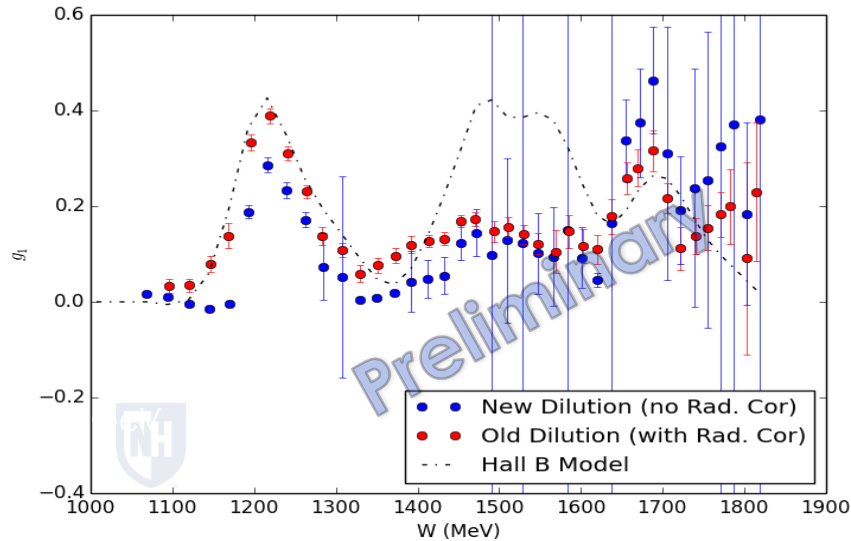
2.5T 2.2 GeV Transverse



- Structure functions with new dilution and packing fraction still need inelastic radiative corrections
- Our radiative corrections expert, Ryan Zielinski, is helping with this
- We expect R.C. to be done in the coming weeks
- Lack of R.C. is the likely reason for discrepancy from model

# EXTRACTING THE SPIN STRUCTURE FUNCTIONS: $G_1$

2.2 GeV Parallel



Model driven procedure for unmeasured part

$$g_1(x, Q^2) = K_1 \left[ \Delta\sigma_{\parallel} \left( 1 + \frac{1}{K_2} \tan \frac{\theta}{2} \right) \right] + \frac{2g_2(x, Q^2)}{K_2 y} \tan \frac{\theta}{2}$$

$$K_1 = \frac{MQ^2}{4\alpha} \frac{y}{(1-y)(2-y)}$$

$$K_2 = \frac{1 + (1-y)\cos\theta}{(1-y)\sin\theta}$$

- E08-027 data is consistent with previously published data from CLAS
- But with much better statistics!!

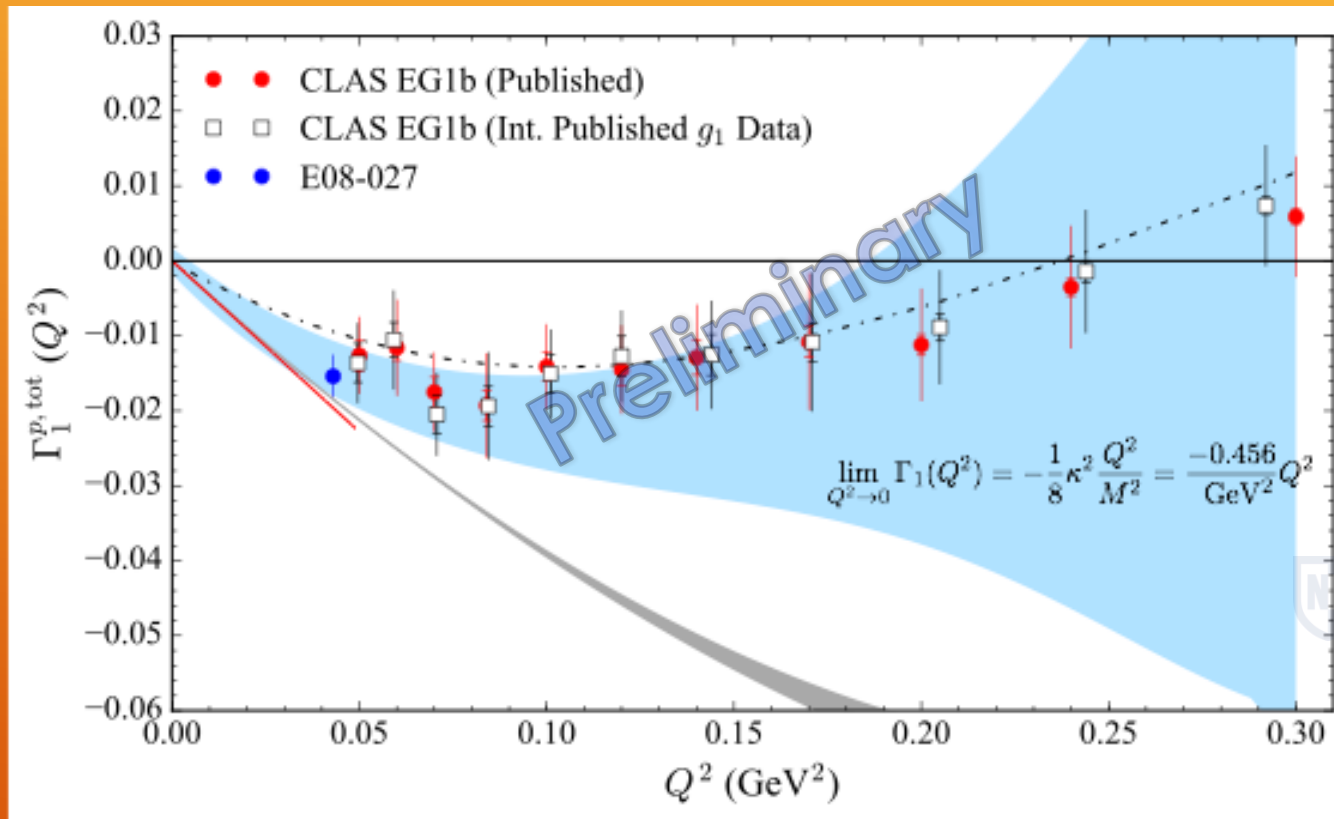
<http://clas.sinp.msu.ru/cgi-bin/jlab/db.cgi>



# FIRST MOMENT OF $g_1(x, Q^2)$

$$\Gamma_1(Q^2) = \int_0^{x_{th}} g_1(x, Q^2) dx$$

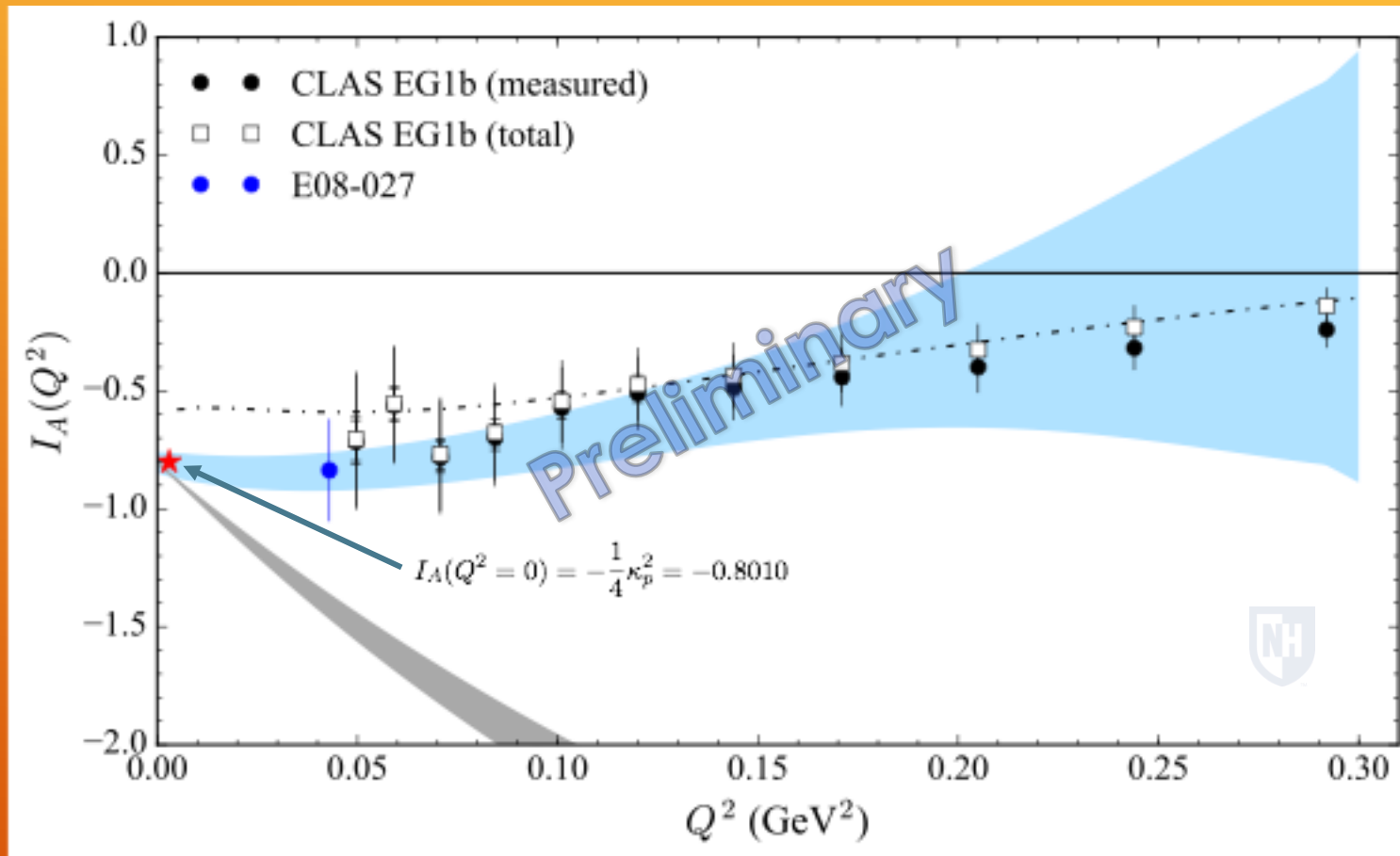
Moments provide a useful quantity that can be related back to theory predictions!



Setting	$\Gamma_1^{meas.}$	$\Gamma_1^{low x}$	$\Gamma_1^{tot.}$	$\delta_{stat}^{tot.}$	$\delta_{sys}^{tot.}$
2254 5T Long.	-0.01541	0.0003	-0.01541	0.0006	0.0028

# EXTENDED GDH SUM

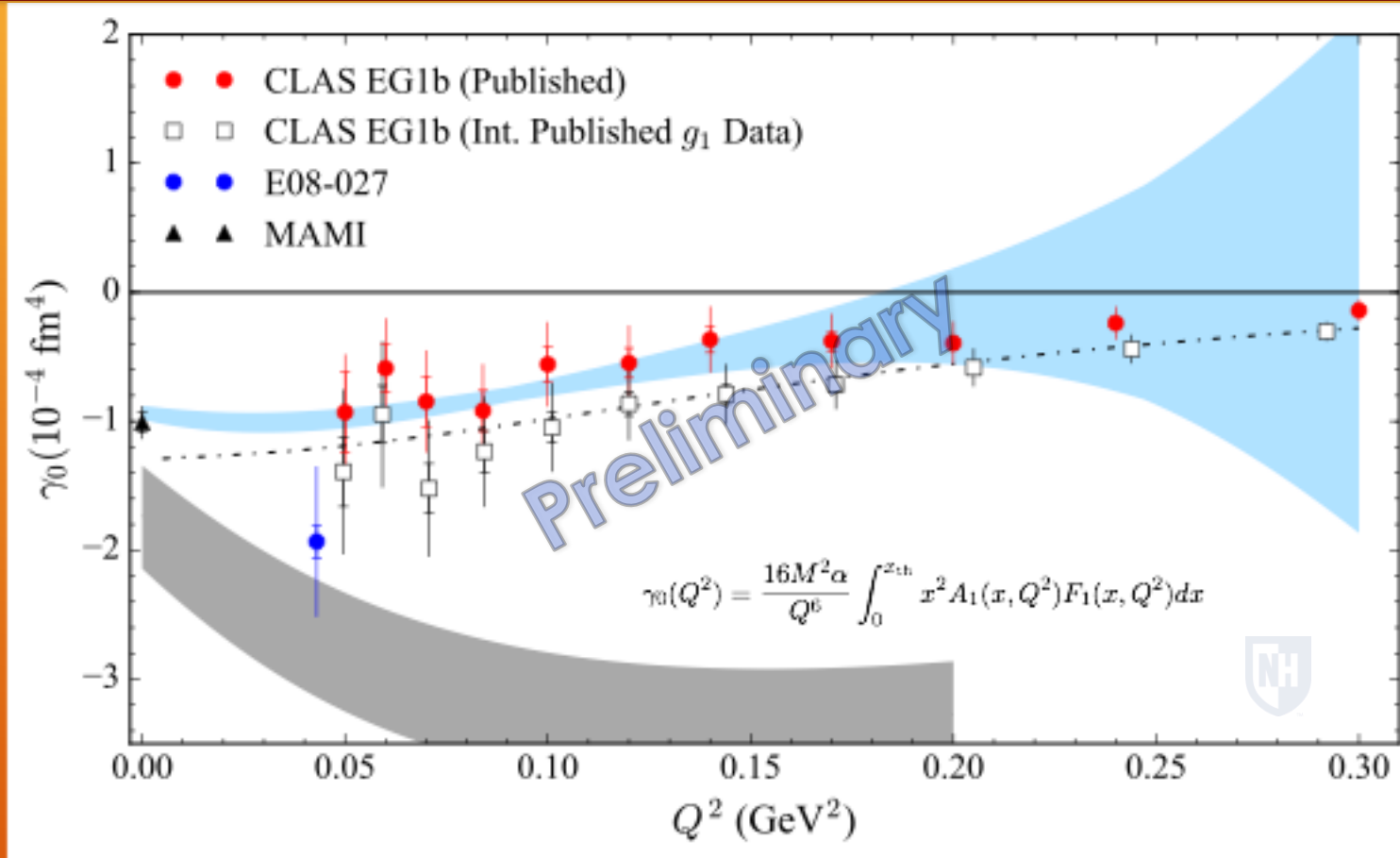
$$I_A(Q^2) = \frac{2M^2}{Q^2} \int_0^{x_{th}} \left( g_1(x, Q^2) - \frac{4M^2}{Q^2} x^2 g_2(x, Q^2) \right) dx$$



Setting	$\Gamma_A^{meas.}$	$\Gamma_A^{low x}$	$\Gamma_A^{tot.}$	$\delta_{stat}^{tot.}$	$\delta_{sys}^{tot.}$
2254 5T Long.	-0.83669	0.0119	-0.83669	0.0240	0.2155

# FORWARD SPIN POLARIZABILITY

$$\gamma_0(Q^2) = \frac{16\alpha M^2}{Q^6} \int_0^{x_{th}} x^2 \left( g_1(x, Q^2) - \frac{4M^2}{Q^2} x^2 g_2(x, Q^2) \right) dx$$



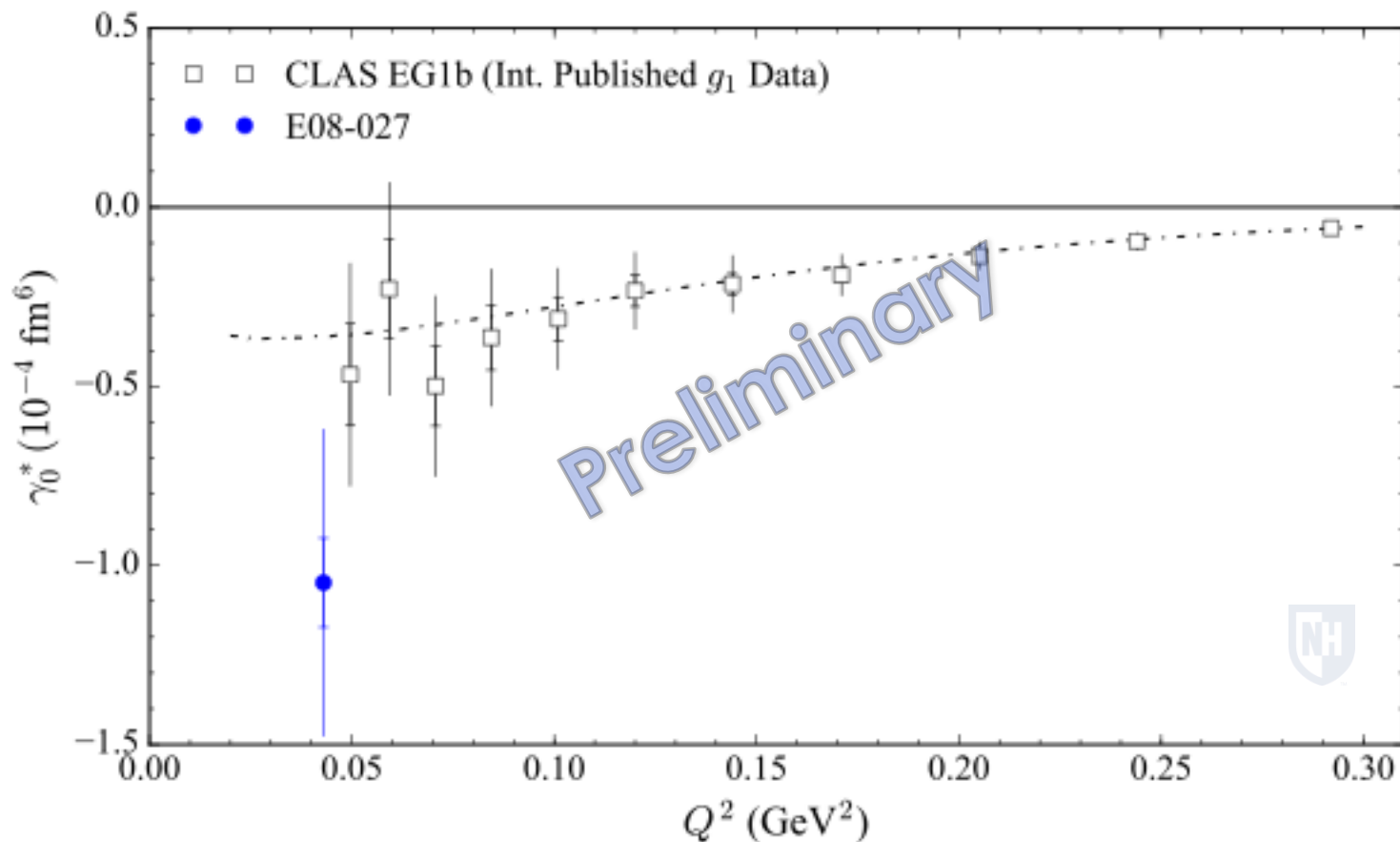
Setting	$\gamma_0^{\text{meas.}}$	$\gamma_0^{\text{low } x}$	$\gamma_0^{\text{tot.}}$	$\delta_{\text{stat}}^{\text{tot.}}$	$\delta_{\text{sys}}^{\text{tot.}}$
2254 5T Long.	-1.9352	-0.0007	-1.9352	0.1202	0.5663

$$\gamma_0(Q^2 = 0) = [-1.01 \pm 0.08 \text{ (stat)} \pm 0.10 \text{ (sys)}] \cdot 10^{-4} \text{ fm}^4$$



# HIGHER ORDER POLARIZABILITY

$$\gamma_0^*(Q^2) = \frac{64\alpha M^4}{Q^{10}} \int_0^{x_{th}} x^4 \left( g_1(x, Q^2) - \frac{4M^2}{Q^2} x^2 g_2(x, Q^2) \right) dx$$



Setting	$\gamma_0^*$ mens.	$\gamma_0^*$ low $x$	$\gamma_0^*$ tot.	$\delta_{\text{stat}}^{\text{tot.}}$	$\delta_{\text{sys}}^{\text{tot.}}$
2254 5T Long.	-1.0501	$-2 \times 10^{-5}$	-1.0501	0.1203	0.4072



# CONCLUSIONS

- Experimental measurements of proton structure are key to understanding the proton!
- The  $g_2p$  experiment was a precision measurement of proton  $g_2$  in low  $Q^2$  region **for the first time!**
- Longitudinal data agrees with previous measurements.
- Dilution and packing fraction issues have consumed the analysis for some time but are now resolved
- Still need R.C. for revised Structure Functions – once this is complete, final results will be calculated for four energy settings
- We hope to wrap up the analysis of these settings in the coming months.





# ACKNOWLEDGEMENTS

## g2p Analysis Team

### Spokespeople:

Alexandre Camsonne  
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Min Huang  
Jie Liu  
Pengjia Zhu  
Ryan Zielinski

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