

# New discoveries of the exotic states at LHCb

MISHA MIKHALENKO

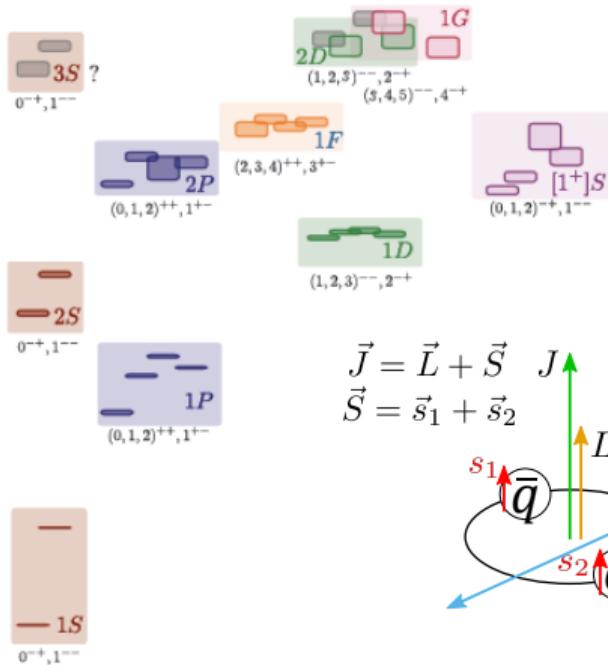
ON BEHALF OF LHCb

ORIGINS Excellence Cluster

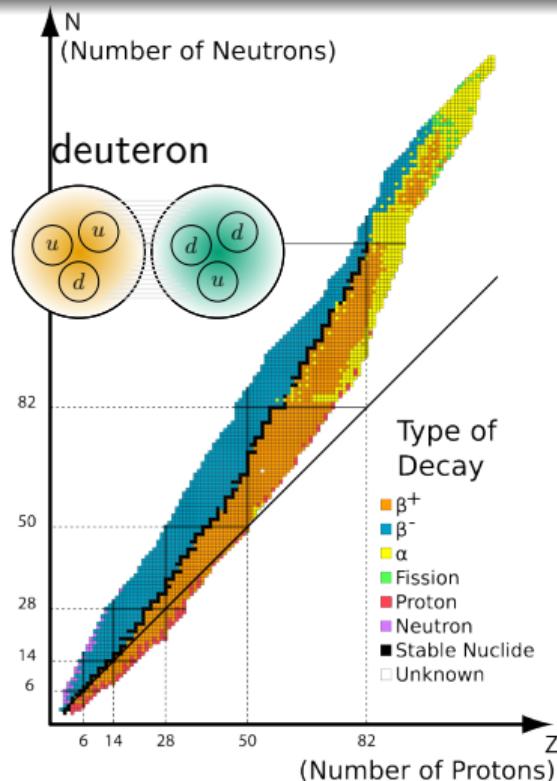
September 6<sup>nd</sup>, 2022  
QNP2022

# Two ways of building complexity

## Hadron variety

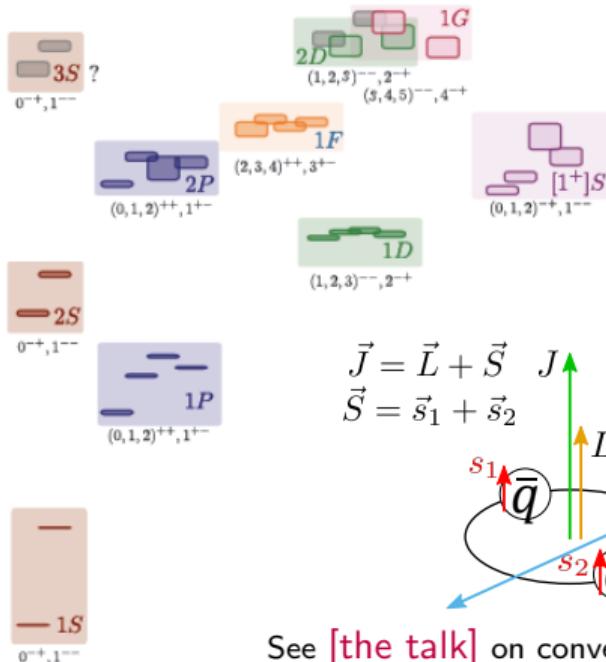


## Atomic variety



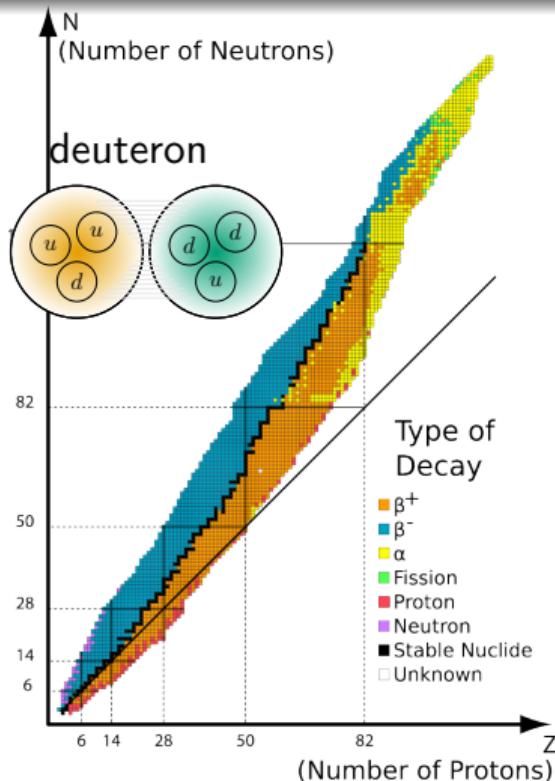
# Two ways of building complexity

## Hadron variety

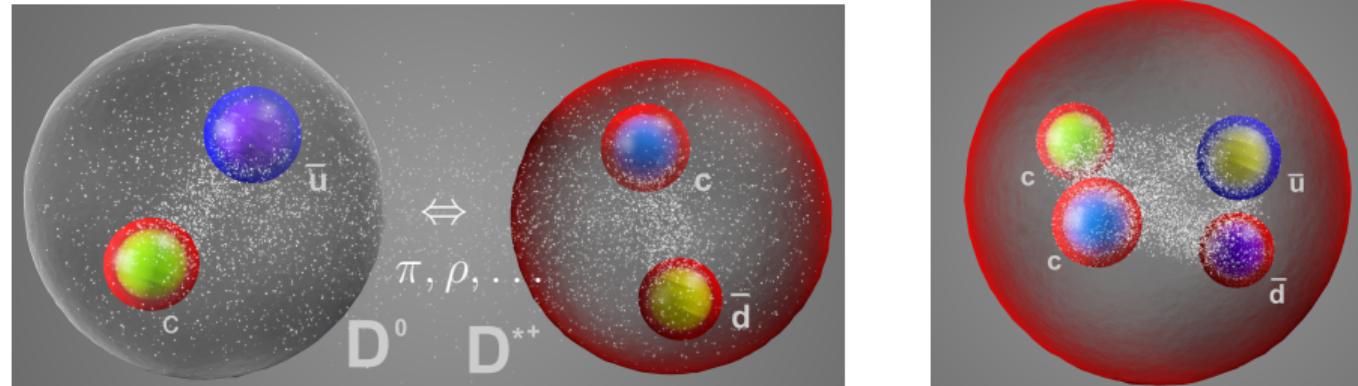


See [the talk] on conventional states at LHCb  
by Aleksei Dzyuba on Thursday

## Atomic variety



# Example: innerworking of $T_{cc}^+$



## Molecular configuration:

- two mesons are well separated,
- bound by forces similarly to el.mag. van der Waals,
- entirely coupled to  $D^{*+}D^0$ ,
- lifetime is limited by  $D^{*+}$ ,
- ? spatially-extended object.

## Compact configuration:

- genuine QCD state,
- compact ( $cc$ ) core,
- there is no limit on lifetime,  
depends on how much it couples to continuum,
- ? typical hadronic size of 1 fm.

# Summary of Pentaquarks studies

(\*) will be discussed today

$$X_b \rightarrow (J/\psi p) \dots$$

$$\Lambda_b^0 \rightarrow (J/\psi p) K^- \quad (*)$$

$$\Lambda_b^0 \rightarrow (J/\psi p) \pi^-$$

$$B_s^0 \rightarrow (J/\psi p) \bar{p}$$

Thresholds:  $\Sigma_c^{(*)+} \overline{D}^{(*)0} / \Sigma_c^{(*)++} D^{(*)-}$

$$X_b \rightarrow (J/\psi \Lambda) \dots$$

$$\Xi_b^- \rightarrow (J/\psi \Lambda) K^- \quad (*)$$

$$B^- \rightarrow (J/\psi \Lambda) \bar{p} \quad (*)$$

Thresholds:  $\Xi_c^{(*)0} \overline{D}^{(*)0} / \Xi_c^{(*)+} D^{(*)-}$

$$P_{\psi}^N:$$

$$P_{\psi S}^{\Lambda}:$$

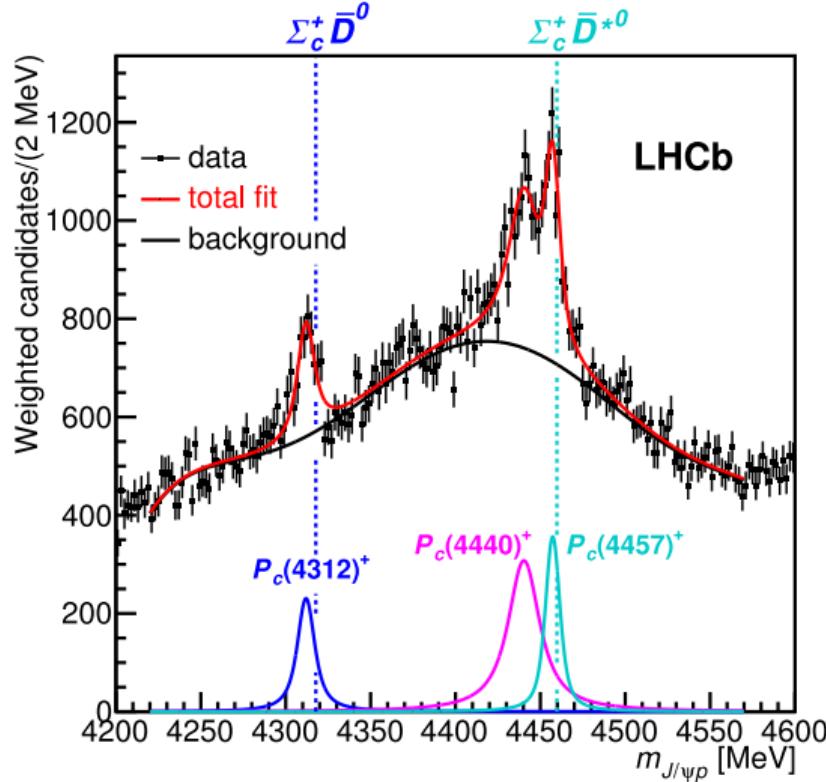
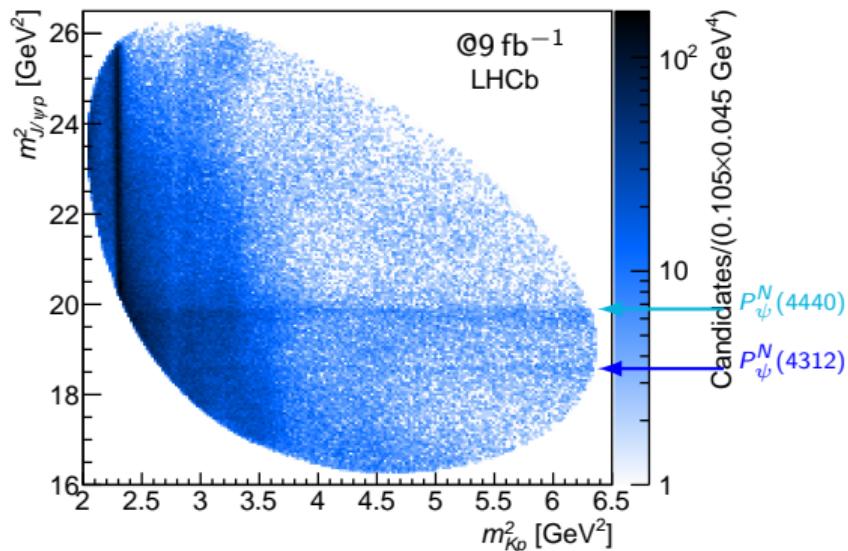
LHCb proposal for the new name convention of exotic hadrons [[arXiv:2206.15233](https://arxiv.org/abs/2206.15233)]

$$\Lambda_b^0 \rightarrow \underbrace{J/\psi p}_{P_\psi^N} K^-$$

# The first pentaquarks

[PRL 115 (2015), 072001; PRL 122 (2019) 22, 222001]

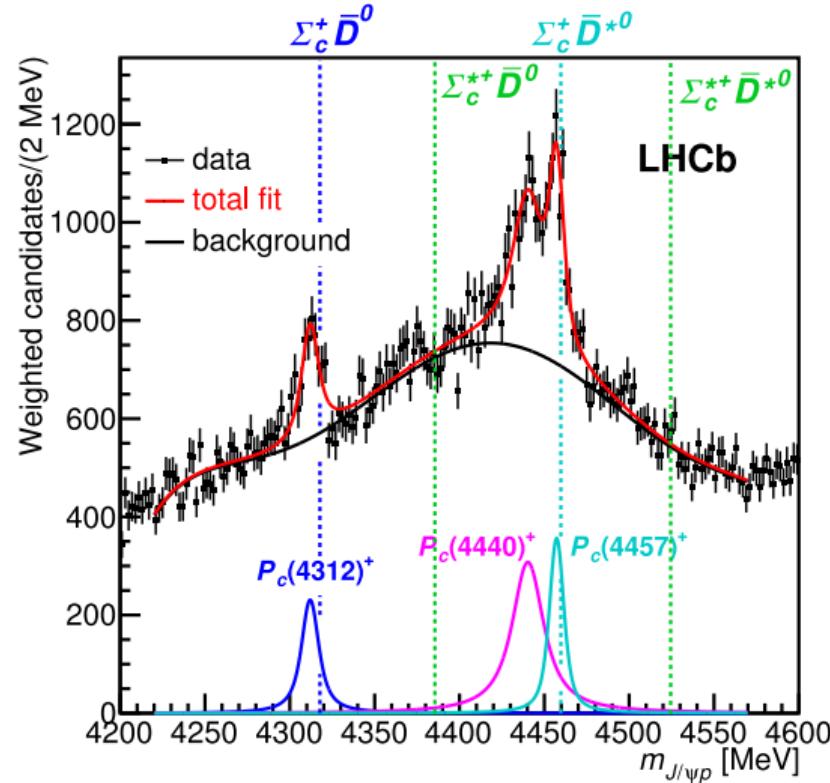
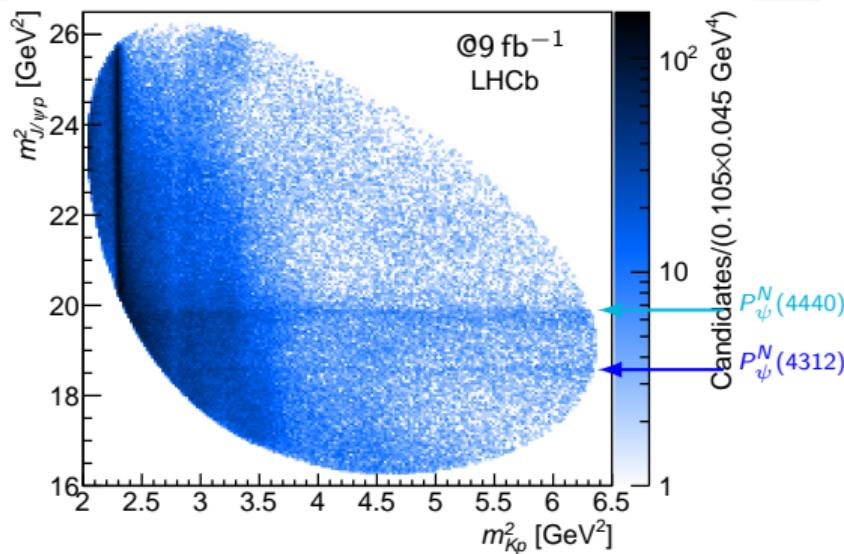
- Close to the  $\Sigma_c \bar{D}^{(*)}$  threshold,
- Multiplicity matches spin combination:  
 $1/2 \otimes 1 = 1/2 \oplus 3/2$
- Narrow(!): 10, 20, and 5 MeV for  $\Gamma_{\text{BW}}$



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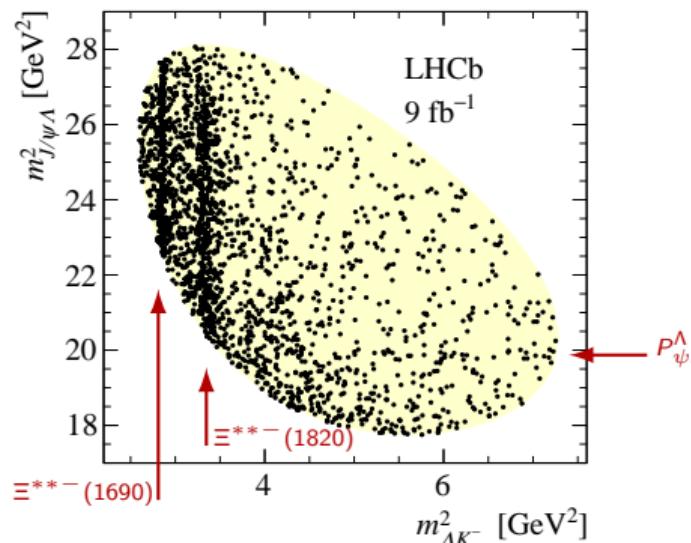
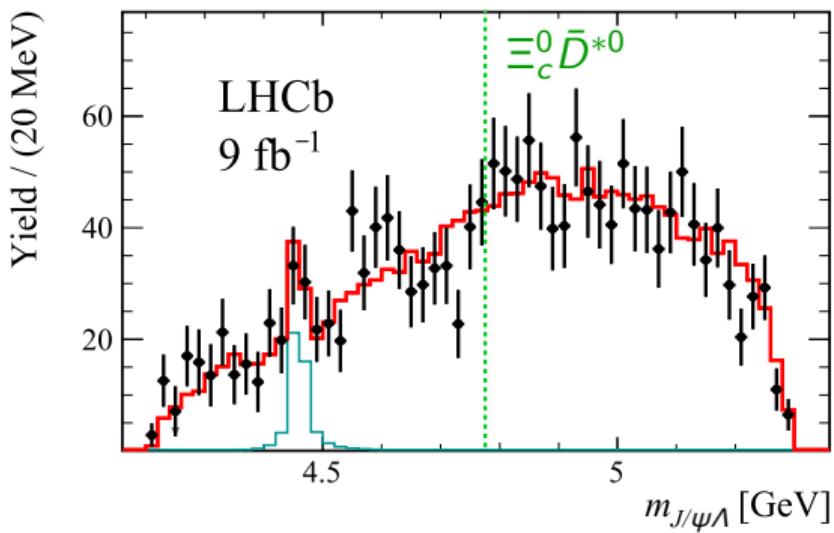


$$\Xi_b^- \rightarrow \underbrace{J/\psi \Lambda}_{P_{\psi s}^\Lambda} K^-$$

# Hint for the strange partners

$\Xi_b^- \rightarrow J/\psi(\rightarrow \mu^+\mu^-)\Lambda(\rightarrow p\pi^-)K^-$  data sample [Sci.Bull. 66 (2021) 1278-1287]

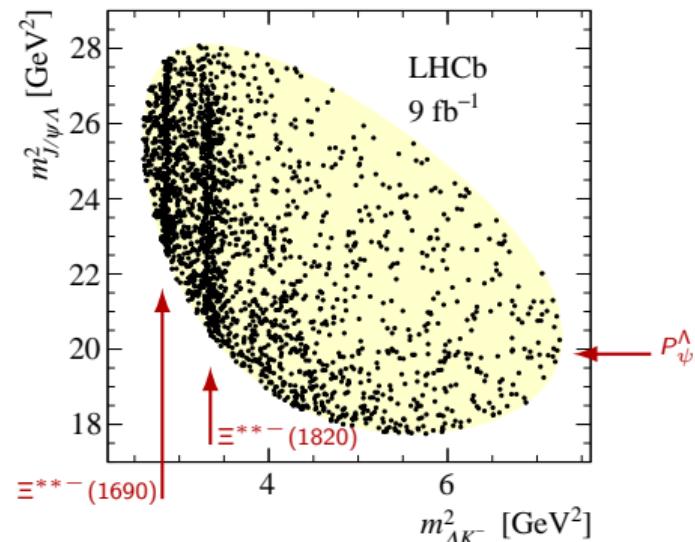
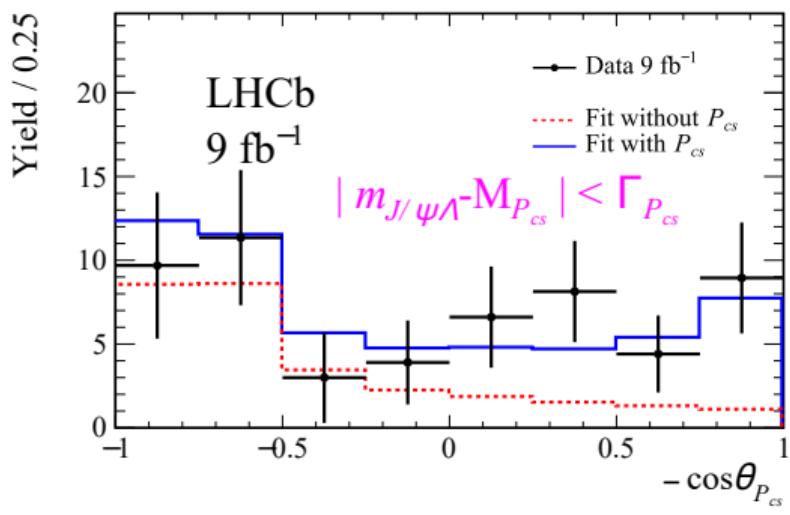
- Full data sample 1750 signals with purity 80%.
- The decay is dominated by the  $\Xi$  resonances
- $P_{\psi s}^\Lambda(4459)$ :  $m = 4458.8 \pm 2.9^{+4.7}_{-1.1}$  MeV,  $\Gamma = 17.3 \pm 6.5^{+8.0}_{-5.7}$  MeV with  $4.3\sigma$  significance



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$$B^- \rightarrow \underbrace{J/\psi \Lambda}_{P_{\psi s}^\Lambda} \bar{p}$$

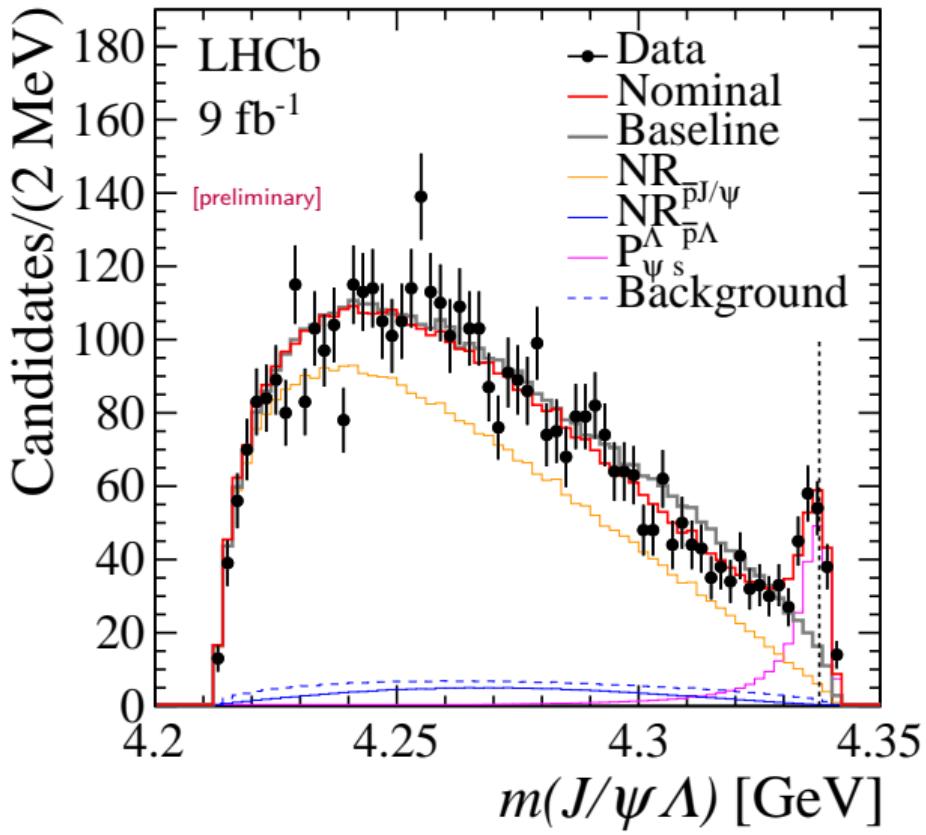
$$B^- \rightarrow J/\psi(\rightarrow \mu^+ \mu^-) \Lambda(\rightarrow p\pi)\bar{p}$$

[LHCb-PAPER-2022-031 (in preparation)]

- Amplitudes:

- ▶ NR( $J/\psi p$ ),  $84.0 \pm 2.2\%$
- ▶ NR( $\Lambda \bar{p}$ ),  $11.3 \pm 1.3\%$
- ▶ New  $P_{\psi s}^\Lambda$ ,  $12.5 \pm 0.7\%$ ,
- ▶ with parameters
  - ★  $m(P_{\psi s}^\Lambda) = 4338.2 \pm 0.7$  MeV
  - ★  $\Gamma(P_{\psi s}^\Lambda) = 7.0 \pm 1.2$  MeV

- $J^P = 1/2^-$  is preferred
- BW mass is close to  $\Xi_c \bar{D}$  thresholds:
  - ▶ 0.8 MeV above  $\Xi_c^+ D^-$
  - ▶ 2.9 MeV above  $\Xi_c^0 \bar{D}^0$



# Tetraquarks candidates

(\*) will be discussed today

$$J/\psi\pi^+$$

$$J/\psi K^+$$

$$J/\psi\phi$$

$$J/\psi J/\psi$$

$$T_\psi^b \quad (Z_c)$$

3900, 4430, ...

$$T_{\psi s}^\theta \quad (Z_{cs})$$

4000, 4220

$$X \quad (T_{\psi\phi})$$

4140, 4274, 4500, ...

$$T_{\psi\psi} \quad (T_{c\bar{c}\bar{c}\bar{c}})$$

6900, ... (!)

$$D^0 D^0 \pi^+$$

$$D^+ K^- \quad (*)$$

$$D_s^\pm \pi^+ \quad (*)$$

$$T_{cc}$$

3874

$$T_{cs} \quad (X)$$

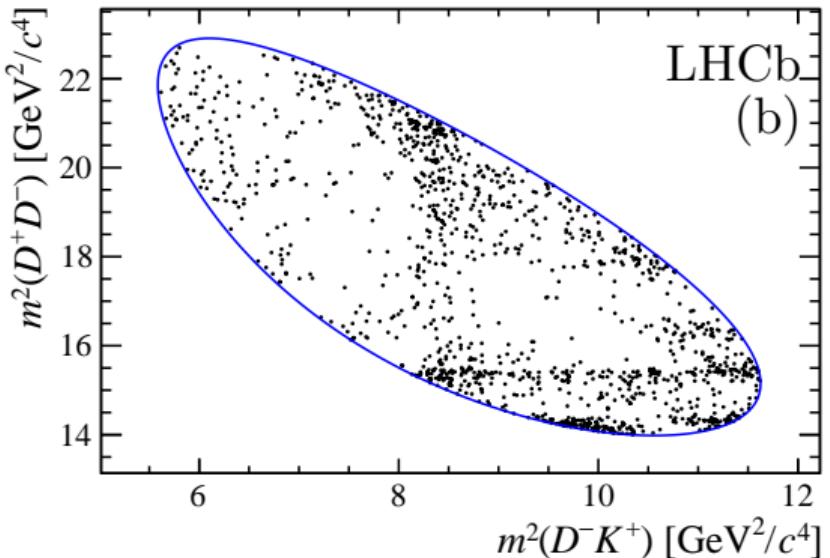
2900

$$T_{c\bar{s}} \quad (X)$$

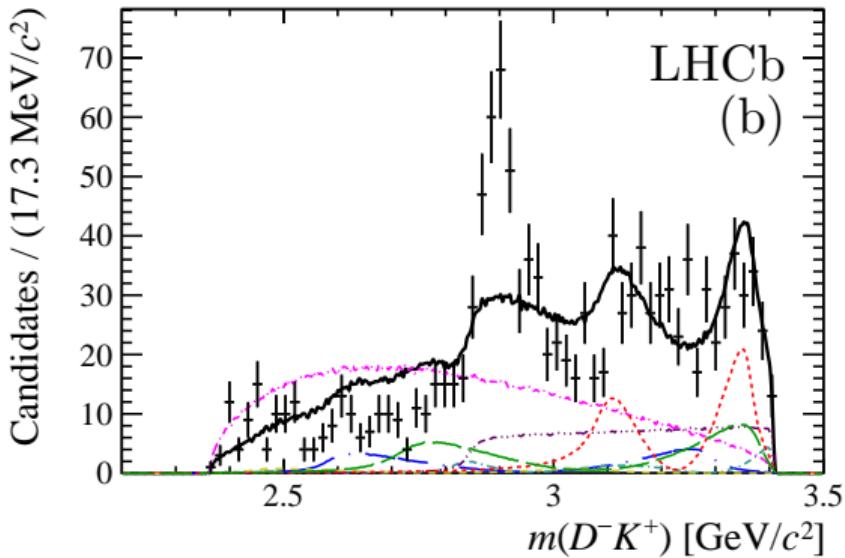
2900

$$B^+ \rightarrow D^+ \underbrace{D^- K^+}_{T_{CS}}$$

# Dalitz plot for $B^+ \rightarrow D^+ D^- K^+$

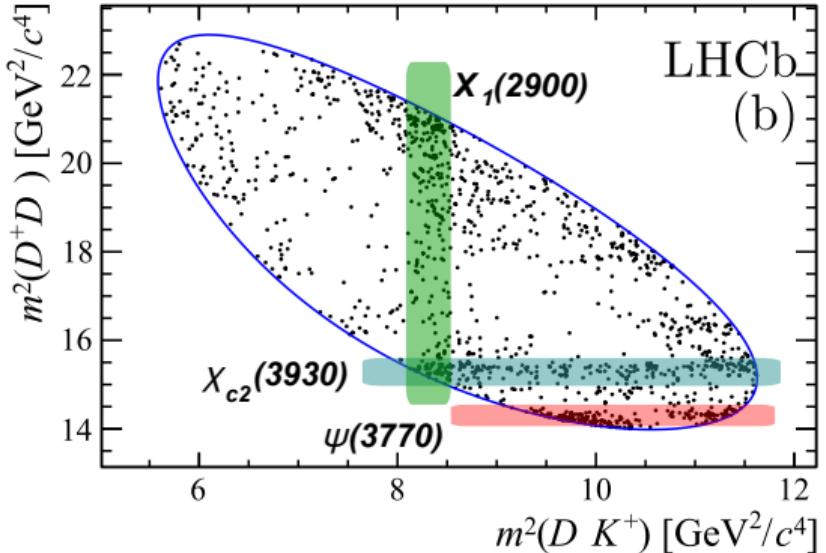


[LHCb, PRD102 (2020) 112003]

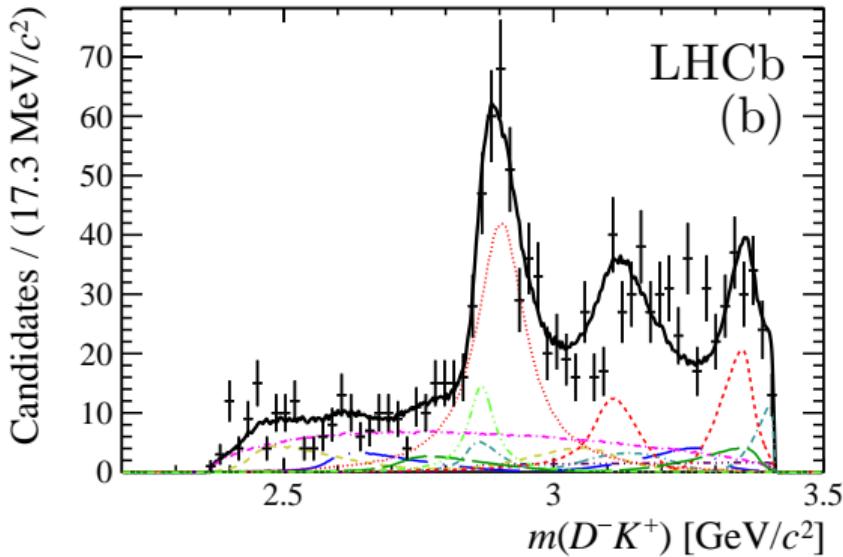


- Horizontal bands are resonances in  $D^+D^-$
- Hint for a vertical band around 8.5 GeV $^2$  in  $m^2(D^-K^+)$
- Exotic candidate  $T_{cs}(2900)$ :  $[\bar{c}\bar{s}ud]$
- Both quantum numbers  $J^P = 0^+$  and  $1^-$  are required in the fit

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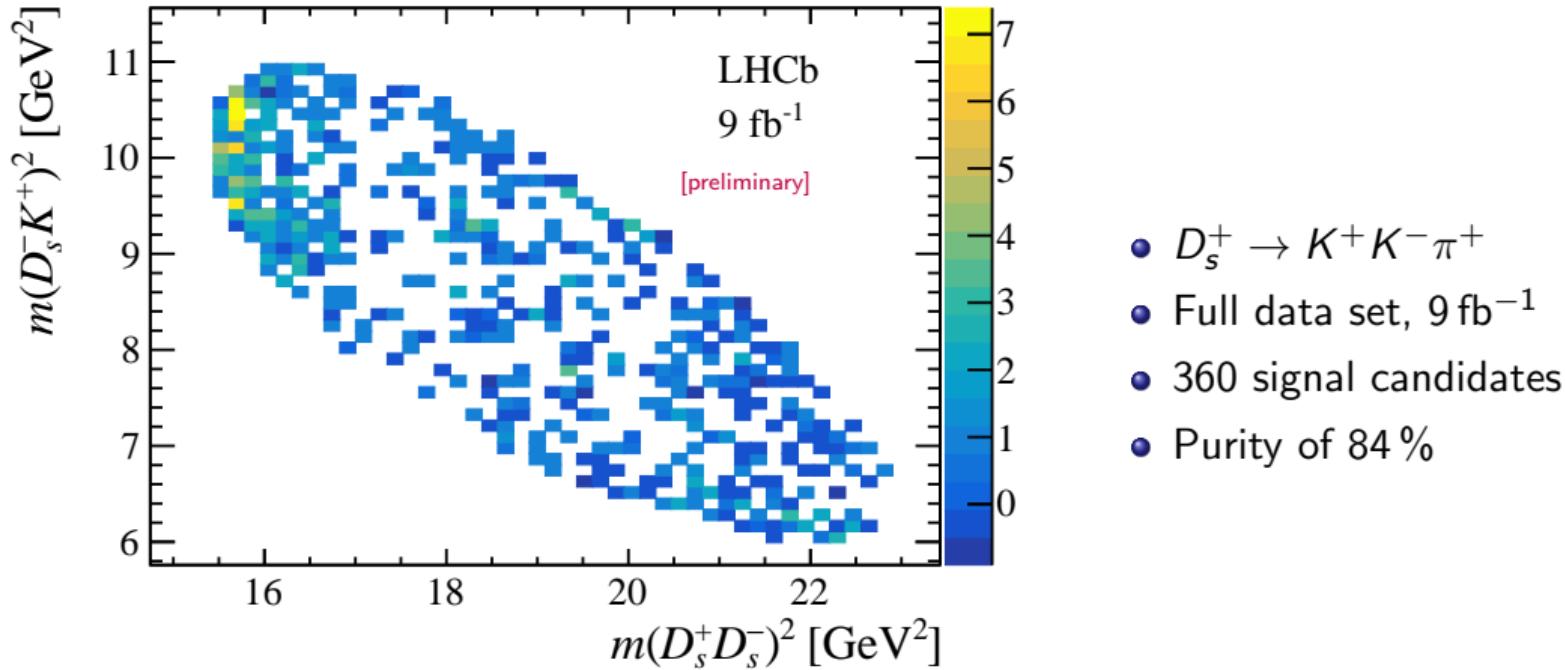


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$$B^+ \rightarrow \underbrace{D_s^+ D_s^-}_{\chi_{c0}/T_{\psi\phi}} K^+$$

# Threshold enhancement at $D_s^+ D_s^-$ in $B^+ \rightarrow D_s^+ D_s^- K^+$ decays

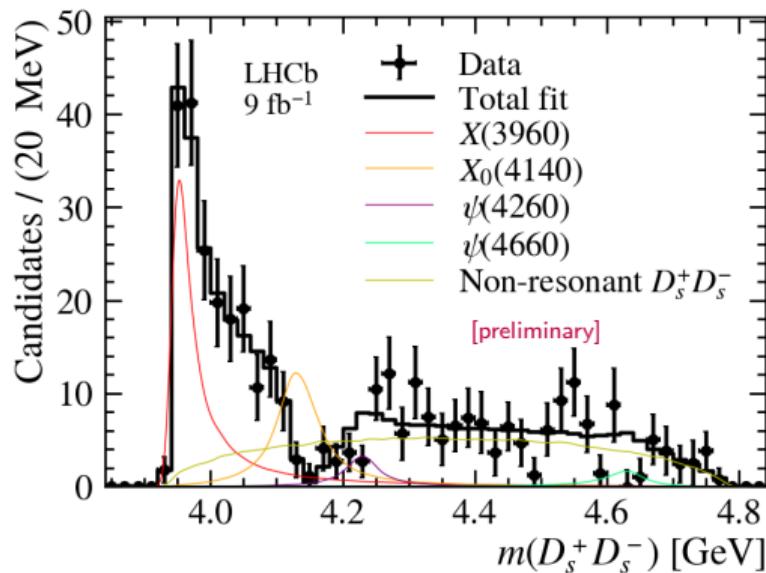
[LHCb-PAPER-2022-018, 019 (in preparation)]



# $B^+ \rightarrow D_s^+ D_s^- K^+$ amplitude analysis

Main features of the data:

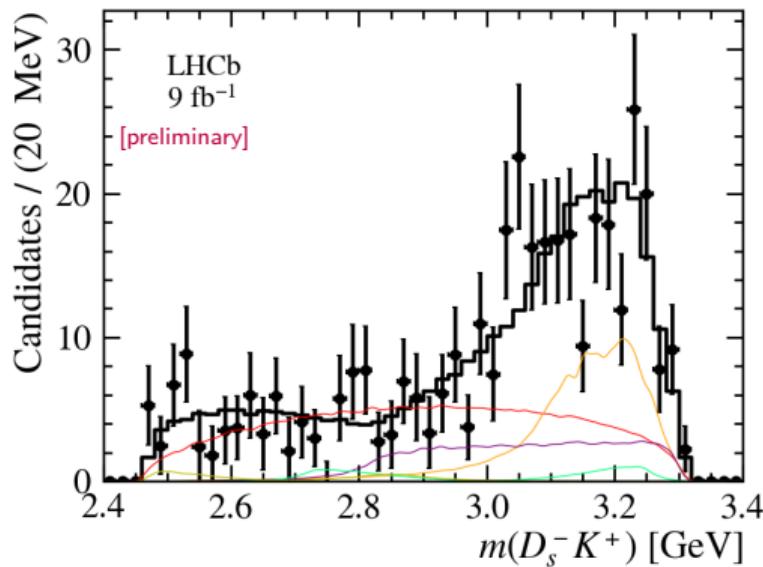
- Enhancement at the  $D_s^+ D_s^-$  threshold
- Followed by a dip at 4.15 GeV.



[LHCb-PAPER-2022-018, 019 (in preparation)]

Baseline model:  $D_s^+ D_s^-$  resonances

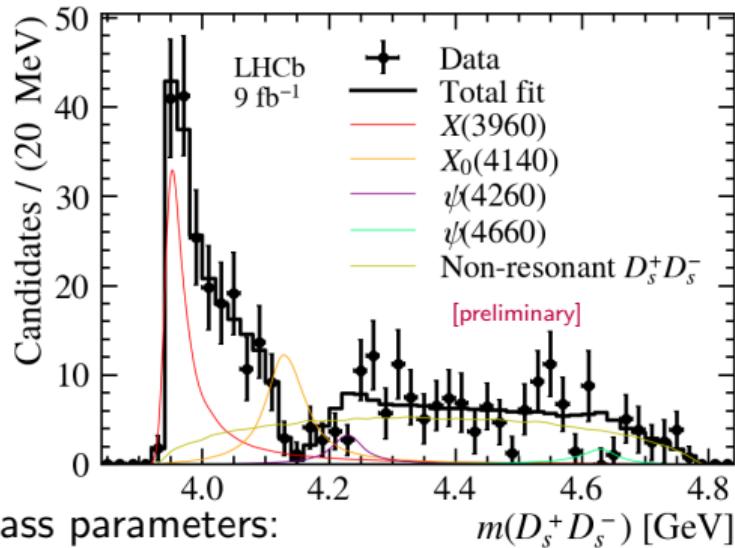
- $1^{--}$ :  $\psi(4260) \sim 4\%$ ,  $\psi(4660) \sim 2\%$
- $0^{++}$ :  $X(3960) \sim 24\%$ ,  $X(4140) \sim 18\%$ , NR  $\sim 50\%$



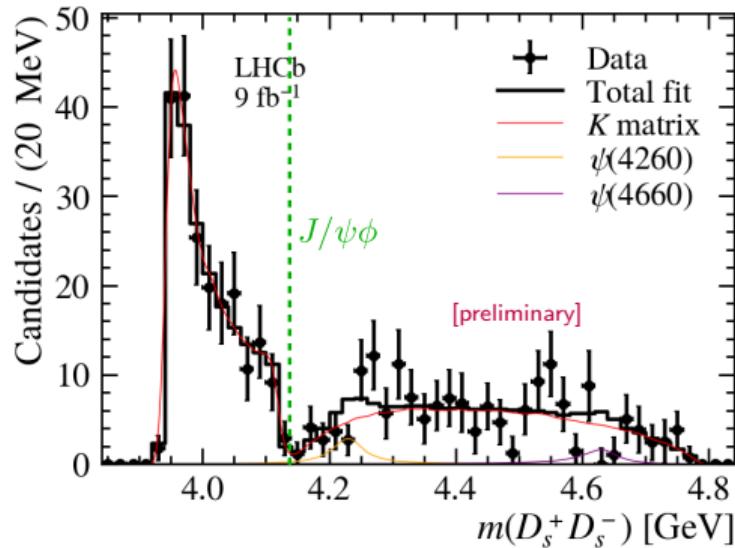
# Alternative model using $K$ -matrix

[LHCb-PAPER-2022-018, 019 (in preparation)]

Three interfering components in  $0^{++}$  are replaced by the  $K$ -matrix.



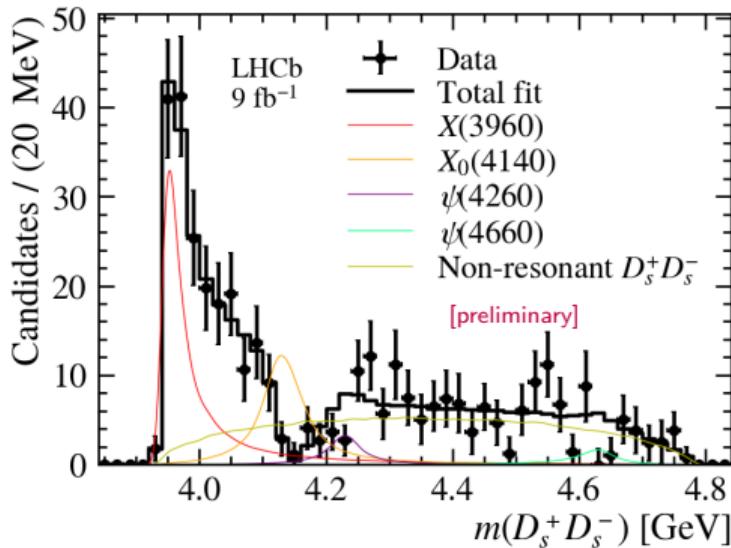
- 3960:  $m = 3956 \pm 5 \pm 11 \text{ MeV}$ ,  
 $\Gamma = 43 \pm 13 \pm 8 \text{ MeV}$
- 4140:  $m = 4133 \pm 6 \pm 11 \text{ MeV}$ ,  
 $\Gamma = 67 \pm 17 \pm 7 \text{ MeV}$



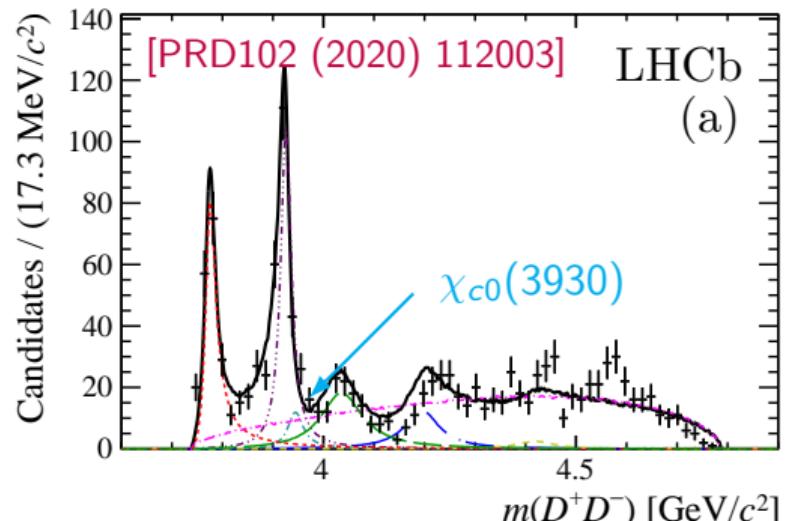
- Coupled channels  $D_s^+ D_s^- / J/\psi\phi$
  - One  $K$ -matrix pole + bgd term
- Gives equally good fit

# Is $X(3960)$ the same as $\chi_{c0}(3930)$ from $D^+D^-$ ?

$B^+ \rightarrow (D_s^+ D_s^-) K^+$  by LHCb:



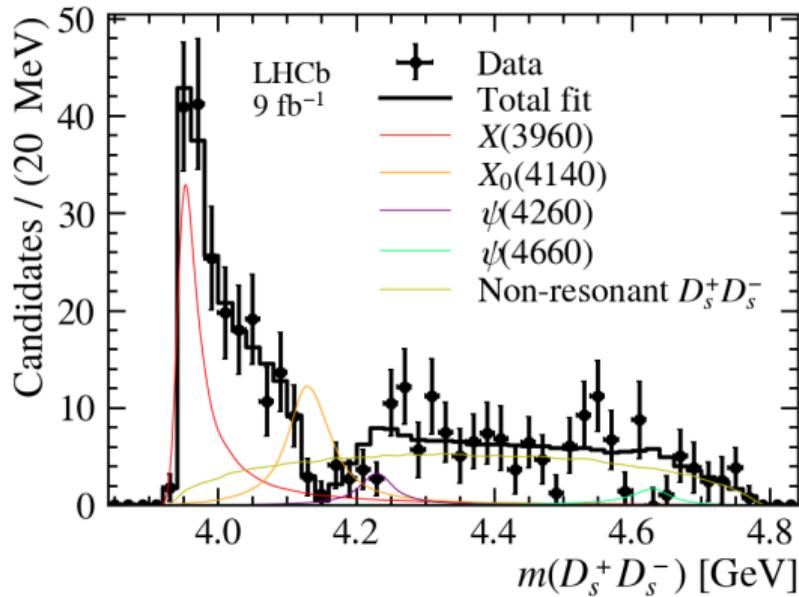
$B^+ \rightarrow (D^+ D^-) K^+$  by LHCb:



- Assuming to be the same,  $\mathcal{B}(\chi_{c0} \rightarrow D^+ D^-)/\mathcal{B}(\chi_{c0} \rightarrow D_s^+ D_s^- P) \sim 0.3$   
large molecular component, or large tetraquark component,  $T_{\psi\phi}$
- [JHEP 06 (2021) 035] finds a state coupled to  $D_s^+ D_s^-$  on the lattice

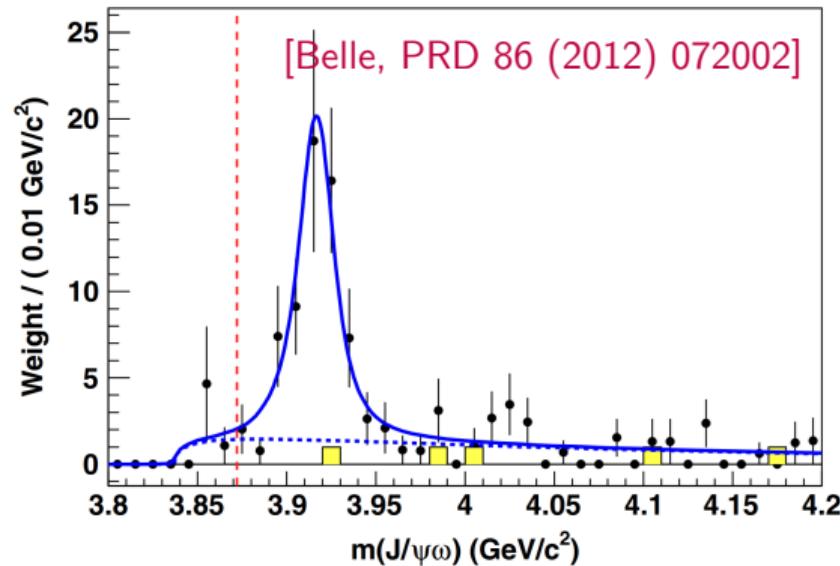
# Is $X(3960)$ the same as $\chi_{c0}(3915)$ ?

$B^+ \rightarrow (D_s^+ D_s^-) K^+$  by LHCb:



[LHCb-PAPER-2022-018, 019 (in preparation)]

$\gamma\gamma \rightarrow J/\psi\omega$  by Belle:



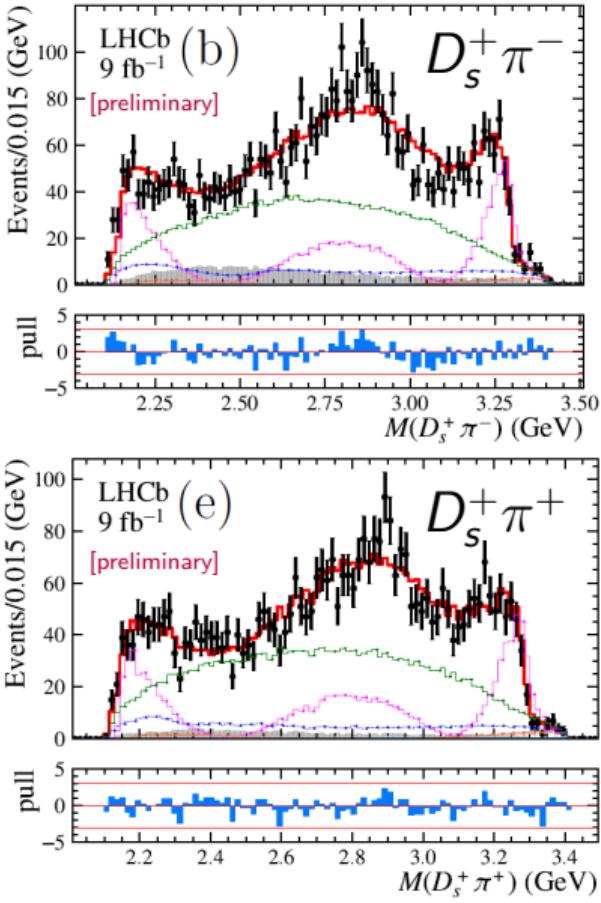
- Belle sees a clean state in  $J/\psi\omega$  with  $J^P = 0^+$
- The  $D_s^+ D_s^-$  signal might be a tail of the  $\chi_{c0}(3915)$  state

$$B^+ \rightarrow D^- \underbrace{D_s^+ \pi^+}_{T_{c\bar{s}}^a}$$

$$B^0 \rightarrow \bar{D}^0 \underbrace{D_s^+ \pi^-}_{T_{c\bar{s}}^a}$$

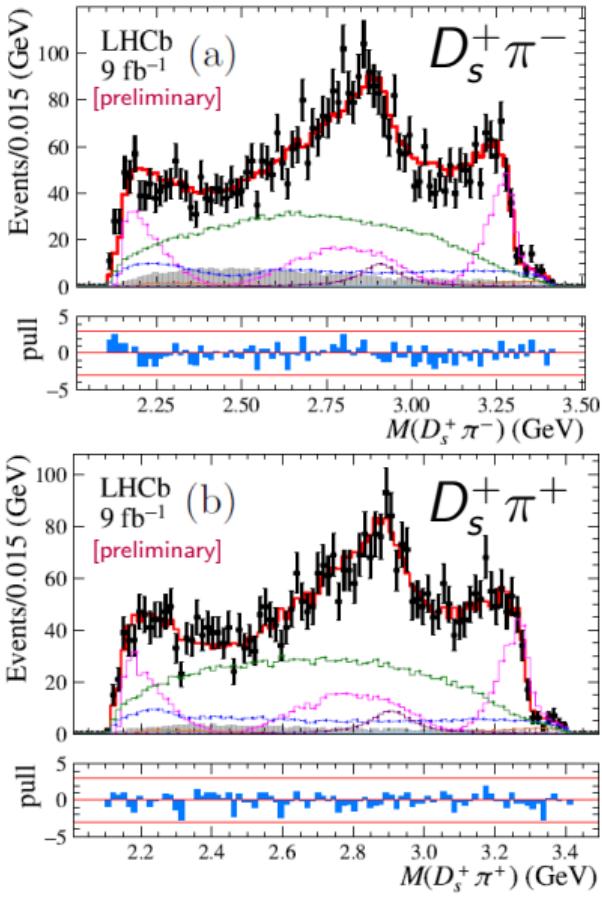
# $T_{c\bar{s}}^a(2900)$ in the $D_s^\pm \pi^\pm$ system [LHCb-PAPER-2022-026 (in preparation)]

- 4420  $B^0$  decays and 3940  $B^-$  decays, including charge-conjugated reactions
- Simultaneous fit using the isospin symmetry
- Main components in  $B^0/B^+$  model:
  - ▶  $D^* \sim 17/14\%$
  - ▶  $D_2^* \sim 22/23\%$
  - ▶  $D\pi$  S-wave  $\sim 45/48\%$ .
- $T_{c\bar{s}}^a \sim 2\%$  needed ( $> 5\sigma$ ),  $J^P = 0^+$  is favored ( $7.5\sigma$ )
- Mass and width are close to these of  $T_{c\bar{s}}^a(2900)$ 
  - ▶  $T_{c\bar{s}}^{a0}$ :  $m = 2892 \pm 14 \pm 15$  MeV,  
 $\Gamma = 119 \pm 26 \pm 12$  MeV;
  - ▶  $T_{c\bar{s}}^{a++}$ :  $m = 2921 \pm 17 \pm 19$  MeV,  
 $\Gamma = 137 \pm 32 \pm 14$  MeV



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## Summary and conclusion

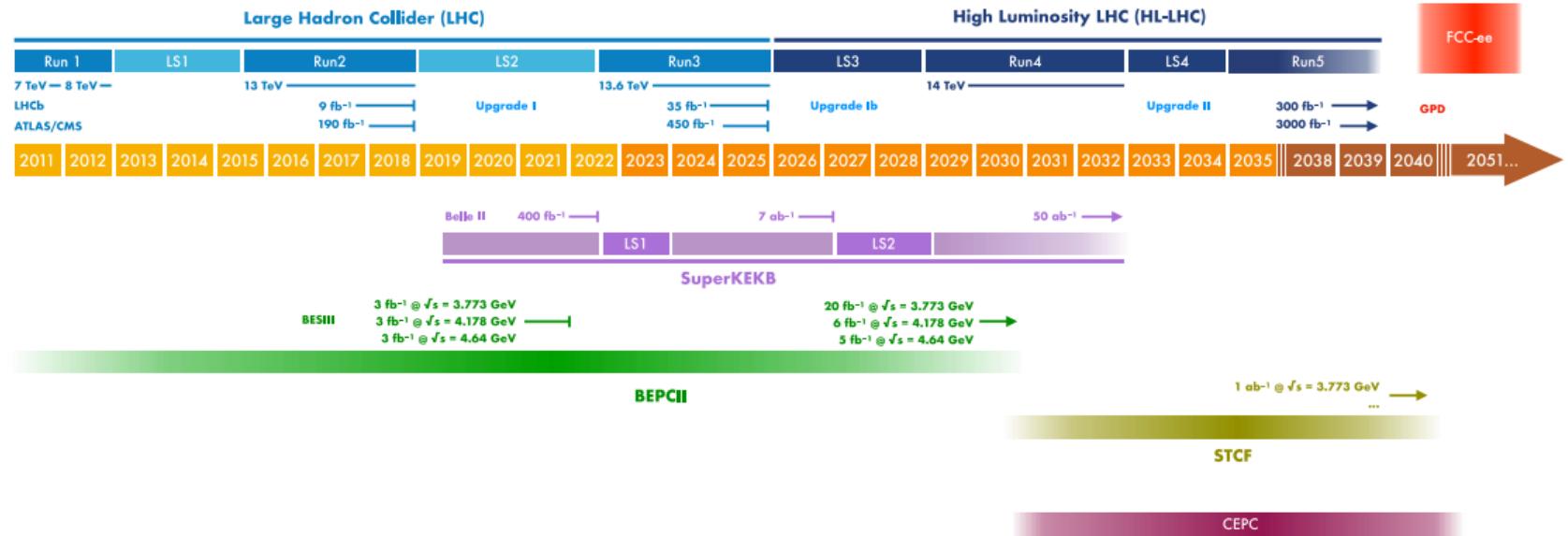
- With excellent performance of LHCb, we are exploring the uncharted waters
- This summer, more puzzle pieces for the quest of understanding nature complexity

Hints for the bigger picture:

- new  $P_{\psi s}^\Lambda$  in  $B$  decays are close to  $\Xi_c D$  threshold  $\longleftrightarrow P_\psi^N$  are connected to  $\Sigma_c \bar{D}^{(*)}$   
Note, not the SU(3) flavour symmetry(!). Check  $\Lambda_c \bar{D}^{(*)}$  and  $\Xi'_c \bar{D}^{(*)}$  thresholds
- new  $T_{c\bar{s}}^a$  appear to be similar to  $T_{cs}$   
HQSS is in action? Where is the good/bad diquark mass difference?
- new  $\chi_{c0}/T_{\psi\phi}$  close to the  $D_s^+ D_s^-$  threshold.  
A hadronic molecule or the scalar charmonium?

# Updated timeline for LHC

[W. Altmannshofer, F. Archilli, arXiv:2206.11331]

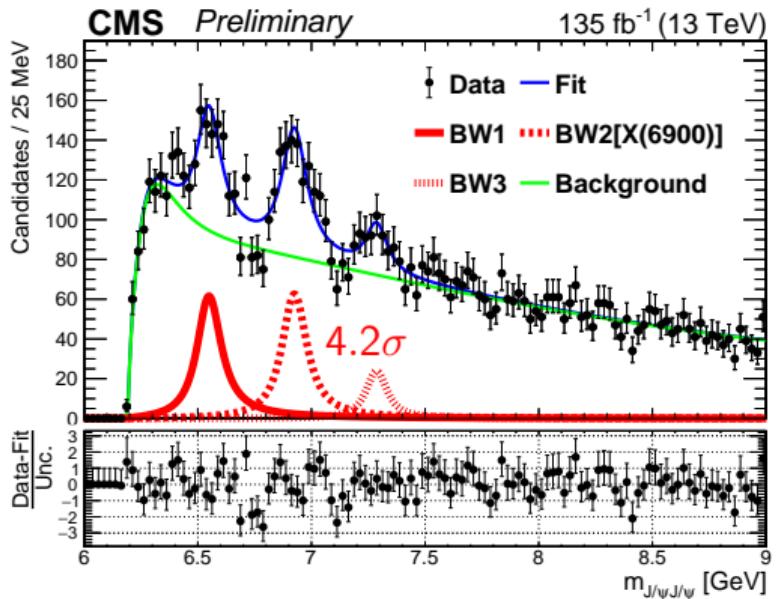


LHCb:

- ramping up after major Upgrade I
- ×5 statistics in Run 3(2023-2025) @13.6 TeV + Run 4(2029-2032) @14 TeV

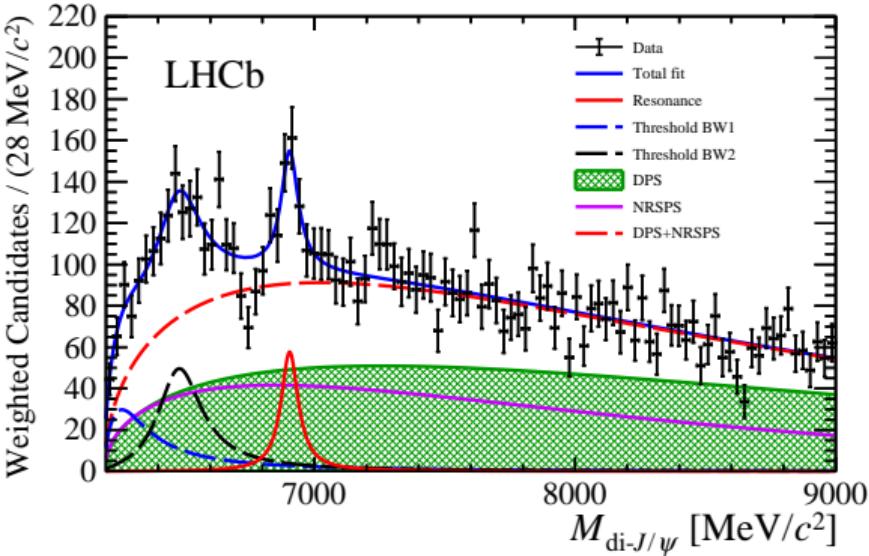
# Thank you for your attention

# CMS confirms $T_{\psi\psi}$ structures



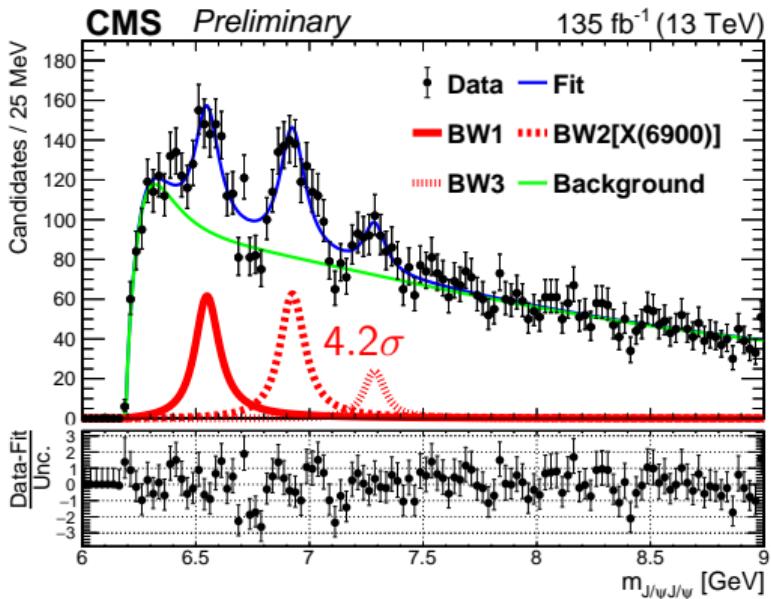
[CMS-PAS-BPH-21-003]

- Clear dips is present that makes the incoherent fit struggles
- Third state is significant



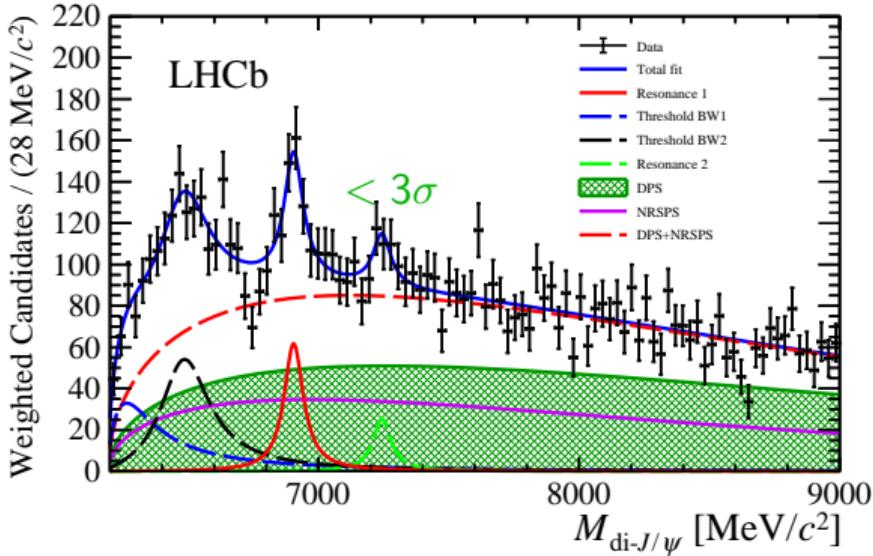
[LHCb, Sci.Bull. 65 (2020) 23, 1983-1993]

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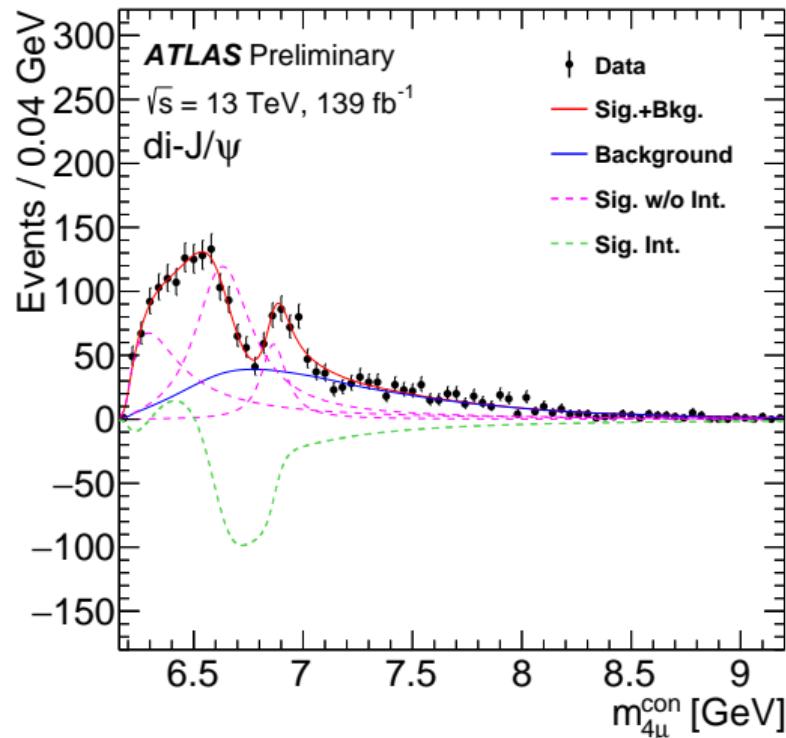
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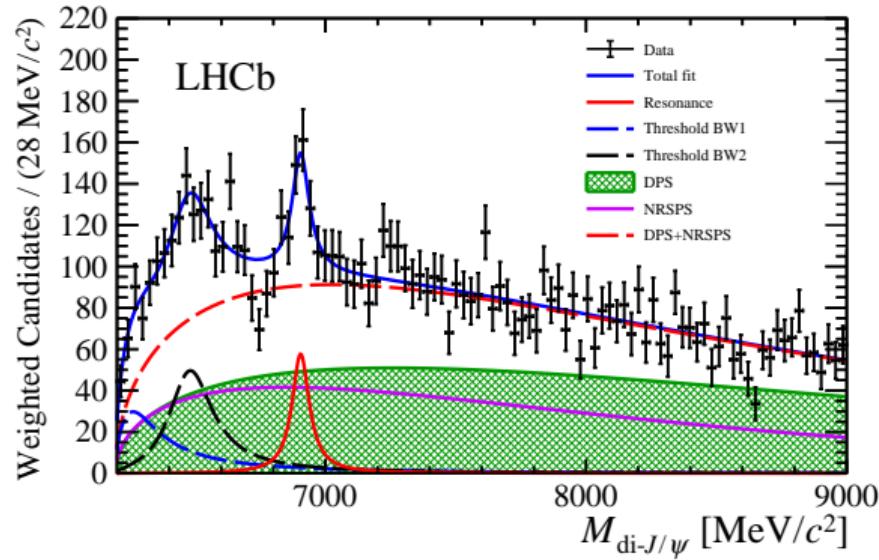


[LHCb, Sci.Bull. 65 (2020) 23, 1983-1993]

# ATLAS also finds structures in $J/\psi J/\psi$

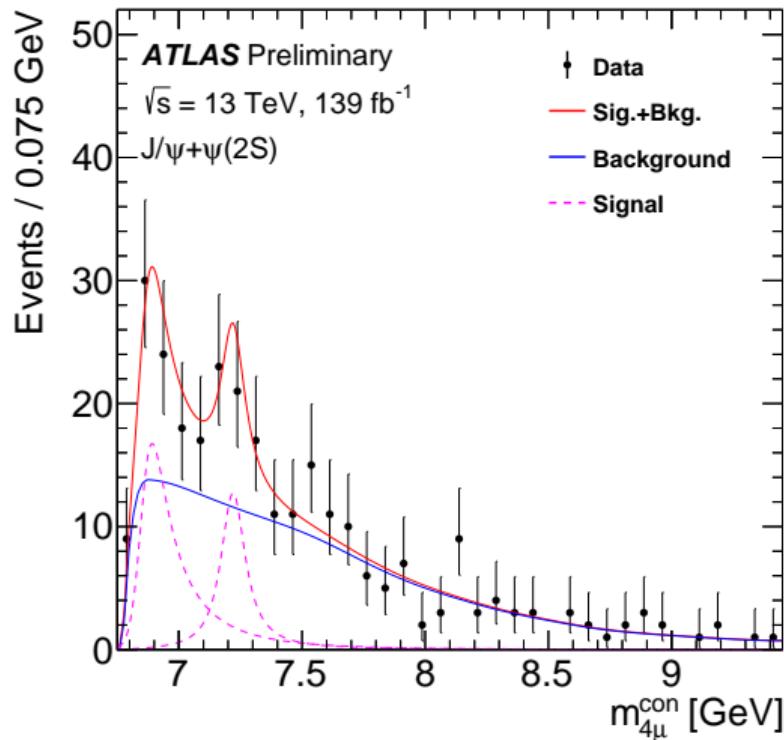


[ATLAS-CONF-2022-040]



[LHCb, Sci.Bull. 65 (2020) 23, 1983-1993]

# ATLAS also finds structures in $J/\psi J/\psi$ and $\psi' J/\psi$

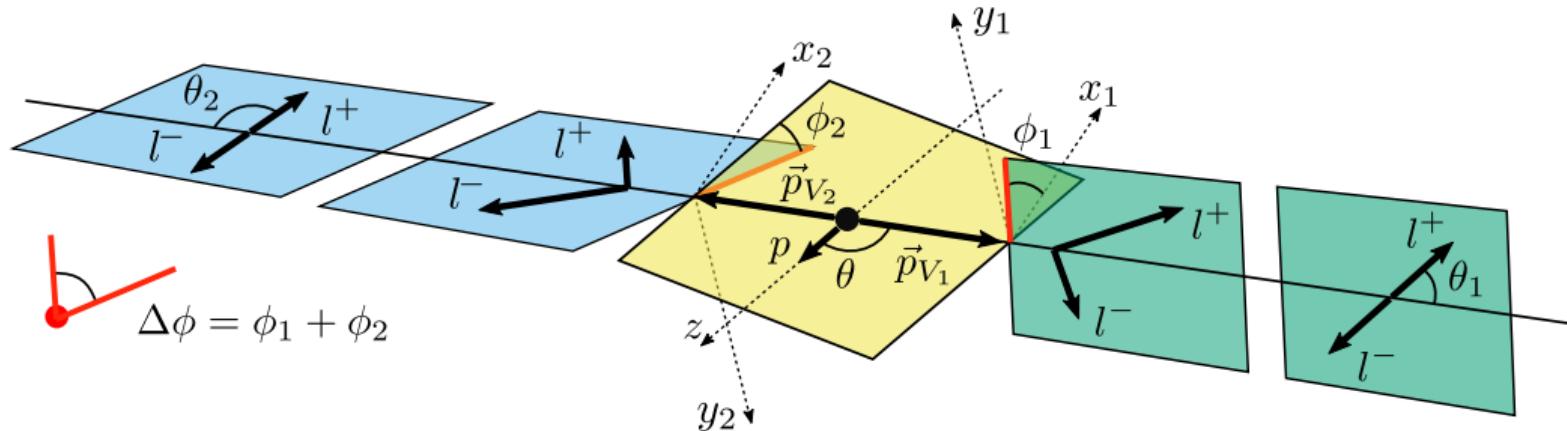


- Hints for the near-threshold structure
- Resonances in  $\psi' J/\psi$  might produce structures in  $J/\psi J/\psi$  as partial-reconstructed decays,  
 $\psi' \rightarrow J/\psi + \text{ neutrals}$

[ATLAS-CONF-2022-040]

# Four-body decay angles

[MM, L. An, R. McNulty, arXiv:2007.05501]



- $\theta$  is the polar angle of  $(J/\psi)_1$  with respect to the polarization direction
- $(\theta_1, \phi_1)$  are the spherical angles of  $\mu^+$  in the  $(J/\psi)_1$  helicity frame
- $(\theta_2, \phi_2)$  are the spherical angles of  $\mu^+$  in the  $(J/\psi)_2$  helicity frame

No polarization  $\Rightarrow$  no  $z$   $\Rightarrow$  no decay plane (pink)  $\Rightarrow$  only  $\phi = \phi_1 + \phi_2$  matters.

# Matrix of helicity couplings

[MM, L. An, R. McNulty, arXiv:2007.05501]

$$H_{\lambda_1, \lambda_2} = \begin{pmatrix} h_{1,1} & h_{1,0} & h_{1,-1} \\ h_{0,1} & h_{0,0} & h_{0,-1} \\ h_{-1,1} & h_{-1,0} & h_{-1,-1} \end{pmatrix}$$

The same-color elements are connected by symmetries.

The symmetry relates the couplings

$$H_{\lambda_1, \lambda_2} = (-1)^J \eta_X H_{-\lambda_1, -\lambda_2}, \quad H_{\lambda_1, \lambda_2} = (-1)^J H_{\lambda_2, \lambda_1}.$$

# Four categories of possible helicity matrices

[arXiv:2007.05501]

group	$\eta_X(-1)^J, (-1)^J$	$J^P$	symmetry
I	+, +	$0^+$ , $2^+, 4^+, 6^+$	symmetric, S
II	- , +	$0^-$ , $2^-, 4^-, 6^-$	symmetric, S
III	+, -	$1^-, 3^-, 5^-, 7^-$	antisymmetric, A
IV	- , -	$1^+$ , $3^+, 5^+, 7^+$	antisymmetric, A

$$\begin{array}{c}
 H^{(I)} \qquad \qquad \qquad H^{(II)} \qquad \qquad \qquad H^{(III)} \qquad \qquad \qquad H^{(IV)} \\
 \left( \begin{matrix} b & a & c \\ a & d & a \\ c & a & b \end{matrix} \right)_S \quad \left( \begin{matrix} b & a & \\ a & -a & -b \\ & -a & \end{matrix} \right)_S \quad \left( \begin{matrix} & a & \\ -a & & -a \\ & a & \end{matrix} \right)_A \quad \left( \begin{matrix} & a & c \\ -a & & a \\ -c & -a & \end{matrix} \right)_A
 \end{array}$$

$a, b, c, d$  are still unknown coefficients, complex in general.

# Four categories of possible helicity matrices

[arXiv:2007.05501]

group	$\eta_X(-1)^J, (-1)^J$	$J^P$	symmetry
I	+, +	$0^+$ , $2^+, 4^+, 6^+$	symmetric, S
II	- , +	$0^-$ , $2^-, 4^-, 6^-$	symmetric, S
III	+, -	$1^-, 3^-, 5^-, 7^-$	antisymmetric, A
IV	- , -	$1^+$ , $3^+, 5^+, 7^+$	antisymmetric, A

$$\begin{array}{ccc}
 \boxed{0^+} & \boxed{0^-} & \boxed{1^+} \\
 \left( \begin{matrix} b & & \\ & d & \\ & & b \end{matrix} \right)_S & \left( \begin{matrix} b & & \\ & & -b \end{matrix} \right)_S & \left( \begin{matrix} & a & \\ -a & & a \\ & -a & \end{matrix} \right)_A
 \end{array}$$

$a, b, c, d$  are still unknown coefficients, complex in general.

# Rederivation of the Landau-Yang theorem

[arXiv:2007.05501]

"A massive particle with spin 1 cannot decay into two photons", wikipedia

Photons do not carry the longitudinal polarization  $\Rightarrow H_{0,i} = H_{i,0} = 0$

$H^{(I)}$

$$\begin{pmatrix} b & a & c \\ a & d & a \\ c & a & b \end{pmatrix}_S$$

$$\begin{pmatrix} b \\ d \\ b \end{pmatrix}_S$$

$0^+$

$H^{(II)}$

$$\begin{pmatrix} b & a & \\ a & -a & \\ & a & -b \end{pmatrix}_S$$

$$\begin{pmatrix} b \\ \\ -b \end{pmatrix}_S$$

$0^-$

$H^{(III)}$

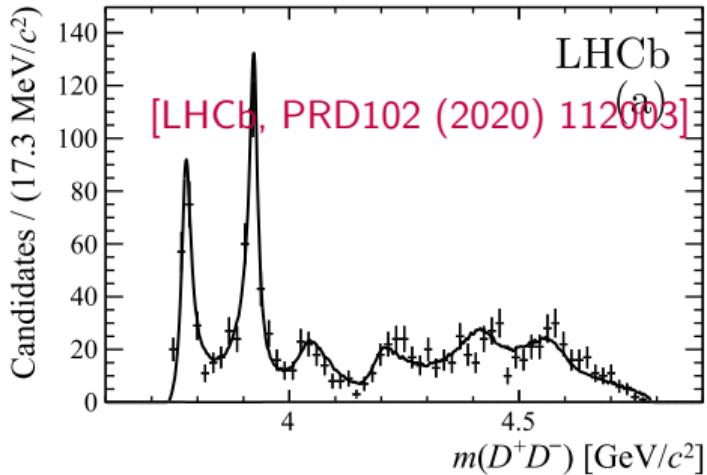
$$\begin{pmatrix} & a & \\ -a & a & \\ & a & -a \end{pmatrix}_A$$

$$\begin{pmatrix} & a & c \\ -a & a & a \\ -c & -a & \end{pmatrix}_A$$

$1^+$

No decay to two photons for  $1^+$ , and group- $III$ :  $1^-, 3^-, 5^-, \dots$

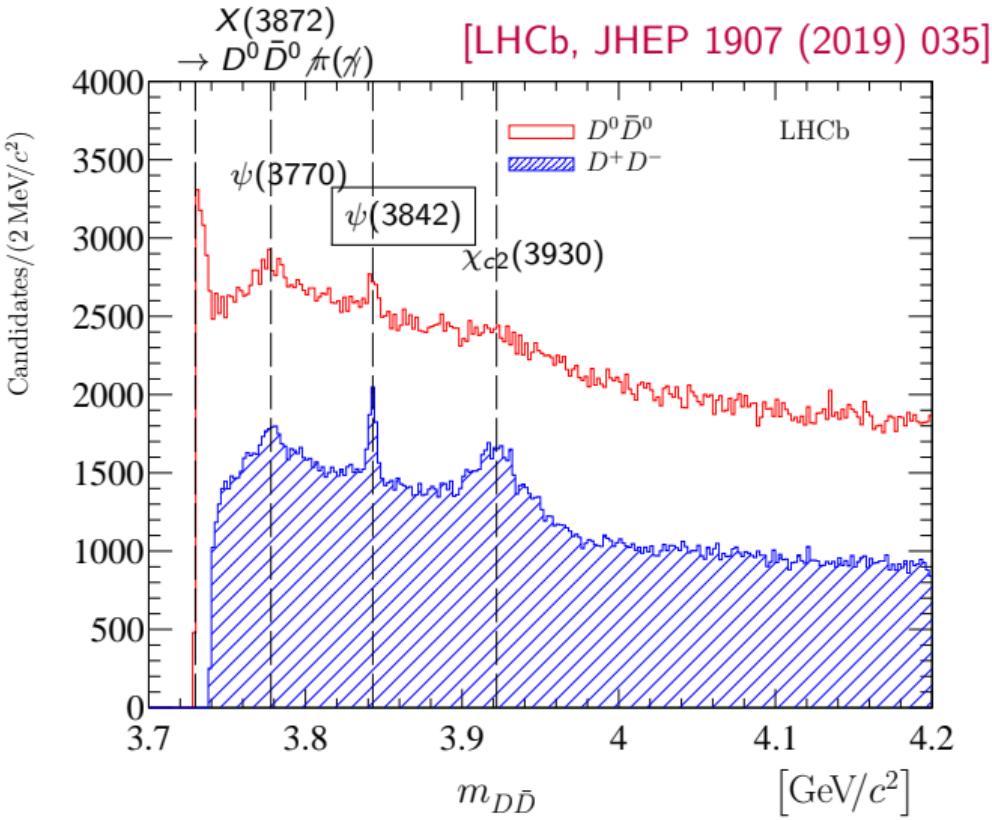
# $D\bar{D}$ spectroscopy



Natural parity chanmonia above  $D^+D^-$  threshold:

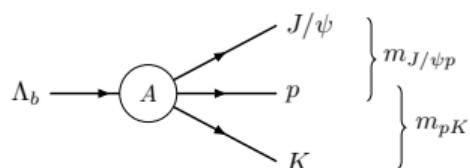
- $\psi(3770)$ ,  $\chi_{c0}(3930)$ ,  $\chi_{c2}(3930)$ ,  
 $\psi(4040)$ ,  $\psi(4160)$ ,  $\psi(4415)$

Compare to inclusive  $DD$  spectra:



# Dalitz Plot Decomposition (DPD)

Update to the angular analysis formalism



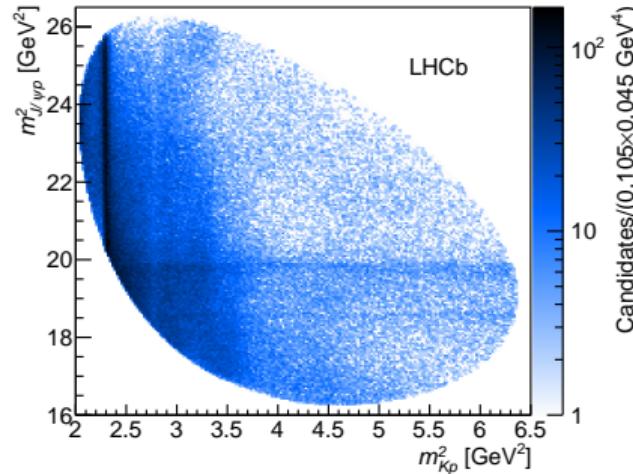
Spin in 3-body decays  
[MM et al. PRD (2019)]  
[Mengzhen Wang et al.,  
2012.03699 (2020)]

$$A_{\lambda_0, \lambda_1, \lambda_2} = \underbrace{A_{\lambda_0, \lambda_1, \lambda_2}^{(12)}}_{0 \rightarrow X, 1} + \underbrace{A_{\lambda_0, \lambda_1, \lambda_2}^{(23)}}_{X \rightarrow 2, 3} + \underbrace{A_{\lambda_0, \lambda_1, \lambda_2}^{(31)}}_{0 \rightarrow X, 2}$$

## Used in the past

$$A_{\lambda_0, \lambda_1, \lambda_2}^{(23)} = \underbrace{\square(\phi_i, \theta_i)}_{0 \rightarrow X, 1} \times \underbrace{\square(\phi'_i, \theta'_i)}_{X \rightarrow 2, 3} \times \underbrace{\square(\phi''_i, \theta''_i)}_{\text{spin align.}}$$

- unphysical inhomogeneity
- spin 1/2:  $A(\pi) \neq A(-\pi)$
- range of  $\phi$  matters  $[-\pi, \pi]$  vs  $[0, 2\pi]$



## Proposed in DPD

$$A_{\lambda_0, \lambda_1, \lambda_2} = \sum_{\nu} D_{\lambda_0, \nu}^{1/2*}(\alpha, \beta, \gamma) \underbrace{O_{\lambda_1, \lambda_2}^{\nu}(m_{12}^2, m_{23}^2)}_{O^{(12)} + O^{(23)} + O^{(31)}}$$

$$O_{\nu, \lambda_1, \lambda_2}^{(23)}(m_{12}^2, m_{23}^2) = \underbrace{\square}_{0 \rightarrow X, 1} \times \underbrace{\square}_{X \rightarrow 2, 3} \times \underbrace{\square}_{\text{spin align.}}$$

- correct  $\phi$  dependence by construction