

Far-Forward and Far-Backward Detectors at the Electron-Ion Collider


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## The Electron-Ion Collider (EIC)

- Two interaction regions (IRs) for possible detector locations.
- Only one IR (IP6) part of the project scope.



## EIC Interaction Region



- Hadron and electron beams collide with a 25 mrad crossing angle (in IP6).
- Design enables maximum luminosity reach of the EIC ( $10^{34} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ ).
- Complicated layout + limited space makes full detector integration a challenge!


## EIC Far-Forward and Far-Backward Regions



- In addition to the central detector $\rightarrow$ detectors integrated into the beamline on both the hadron-going (far-forward) and electron-going (far-backward) direction.

The far-forward and backward systems are crucial for delivery of full EIC physics program!

[^0]Far-Backward Detectors

## Luminosity Monitor

- Process of elastic bremsstrahlung, ep $\rightarrow e_{\gamma} p, e_{A u} \rightarrow e_{\gamma} \mathrm{Au}$
- Large cross section peaked for photons at small angles
- Two methods for $\gamma$ detection: direct detector and $e^{ \pm}$ spectrometer:


The cross section is precisely known from QED

## Luminosity Monitor

- Must make measurement in challenging environment.
- High synchrotron radiation, high bremsstrahlung rates ( $\sim 10 \mathrm{GHz}$ ), etc.
- Need $\sim 1 \%$ for absolute luminosity measurement, $\sim 10^{-4}$ for relative luminosity measurement.

Figure: Direct photon detector


- Simple concept, approximate measurement
- More $\gamma$ incident in every bunch crossing
- Online machine performance
- Precise measurement for physics results
- More complex implementation
- Not sensitive to synchrotron radiation


## Luminosity Monitor

Figure: Towards central detector


Figure: Towards the tunnel


## Low-Q² Tagger

- Two detectors along the outgoing electron beam
- Same $Q^{2}$ is reached at different energies and angles:

Figure: $Q^{2}$ vs. energy and scattering angle


Figure: Towards central detector


Figure: Towards the tunnel


## Low-Q² Tagger

- Production rates are dominated by large bremsstrahlung cross section
- The rates give normalization to spectrum of reconstructed $Q^{2}$
- Clean photoproduction signal can be taken over a limited region in $Q^{2}$
- Bremsstrahlung electron are important to calibrate the luminosity measurement


## Low-Q² Tagger

- Reconstructed virtuality $Q^{2}$ is compared to generated true event $Q^{2}$
- Machine learning connects detected track with original scattered electron
- The $Q^{2}$ is given by electron energy and scattering angle
- Beam effects (vertex spread, angular divergence) are included in the simulation
- Beam angular divergence limits the resolution at $Q^{2}<10^{-3} \mathrm{GeV}^{2}$


Far-Forward Detectors

## Far-Forward Physics at the EIC


e+He3 spectator tagging ${ }^{2}$
$\left.\vec{\nabla}_{e}^{\prime}\right)$

coherent/incoherent
$\mathrm{J} / \psi$ production in $\mathrm{e}+\mathrm{A}^{3}$


Quasi-elastic electron scattering ${ }^{4}$

e+d DIS spectator tagging ${ }^{5}$

[1] Z. Tu, AJ, et al., Phys. Lett. B 811, 135877 (2020) [2] I. Friscic, D. Nguyen, J. R. Pybus, AJ, et al., Phys. Lett. B, 823, 136726 (2021)
[3] W. Chang, E.C. Aschenauer, M. D. Baker, AJ, J.H. Lee, Z. Tu, Z. Yin, and L. Zheng, Phys. Rev. D 104, 114030 (2021)
[4] F. Hauenstein, AJ, J. R. Pybus, A. Kiral, M. D. Baker, Y. Furletova, O. Hen, D. W. Higinbotham, C. Hyde, V. Morozov, D. Romanov, and L. B. Weinstein, Phys. Rev.
C 105, 034001 (2022)
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## Rare isotopes**




Many examples of detailed impact studies with full detector simulations! (non-exhaustive)

## Far-Forward Physics at the EIC

$>$ Physics channels require tagging of charged hadrons (protons, pions) or neutral particles (neutrons, photons) at very-forward rapidities ( $\eta>4.5$ ).
$>$ Different final states $\rightarrow$ tailored detector subsystems.
$>$ Various collision systems and energies (h: 41, 100-275 GeV, e: 5-18 $\mathrm{GeV} ; \mathrm{e}+\mathrm{p}, \mathrm{e}+\mathrm{d}, \mathrm{e}+\mathrm{Au}, \mathrm{etc}$.).
$>$ Placing of far-forward detectors uniquely challenging due to integration with accelerator.
>Details studied in EIC Yellow Report and Conceptual Design Report, and the detector design proposals.

# The Far-Forward Detectors 

BO Silicon Tracker and Preshower


Far-Forward Detector Subsystems

## B0 Detectors



## B0 Detectors

$>$ Charged particle reconstruction and photon tagging.
> MAPS for tracking + timing layer (e.g. LGADs).
$>$ Photon detection (tagging or full reco).

Credit to Ron Lassiter


This is the opening where the detector planes will be inserted

Preliminary Parameters:
$229.5 \mathrm{~cm} \times 121.1 \mathrm{~cm} \times 195 \mathrm{~cm}$ (Actual length will be shorter)

## B0 Detectors

( $5.5<\boldsymbol{\theta}<20.0 \mathrm{mrad}$ )


DD4HEP Simulation
$>$ Higher granularity silicon (e.g. MAPS) required.
$>$ Tagging photons important in differentiating between coherent and incoherent heavy-nuclear scattering, and for reconstructing $\pi^{0} \rightarrow \gamma \gamma$.
$>\begin{aligned} & \text { Space is a major concern here - an EMCAL is highly } \\ & \text { preferred, but may only have space for a preshower. }\end{aligned}$
Space is a major concern here - an EMCAL is highly
preferred, but may only have space for a preshower.
$>$ Still under evaluation.


## B0 Detectors in CAD

Credit to Ron Lassiter
Blue lines represent where element locations are along beamline

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## Roman Pots @ the EIC

Protons<br>$\mathrm{E}=275 \mathrm{GeV}$<br>$0<\boldsymbol{\theta}<5 \mathrm{mrad}$

## Full GEANT4 simulation.

Roman "Pots" @ the EIC


- Silicon detectors placed directly into machine vacuum!
- Allows maximal geometric coverage!
- Technology: AC-LGAD sensor provides both fine pixilation ( $\sim 140$ um spatial resolution), and fast timing ( $\sim 35 \mathrm{ps}$ ).

Roman "Pots" @ the EIC
25.6 cm

$\sigma(z)$ is the Gaussian width of the beam, $\beta(z)$ is the RMS transverse beam size.
$\varepsilon$ is the beam emittance.

$$
\sigma(z)=\sqrt{\varepsilon \cdot \beta(z)}
$$



Low-pT cutoff determined by beam optics.
$>$ The safe distance is $\sim 10 \sigma$ from the beam center.
$>1 \sigma \sim 1 \mathrm{~mm}$
$>$ These optics choices change with energy, but can also be changed within a single energy to maximize either acceptance at the RP, or the luminosity.

## Off-Momentum Detectors



Off-momentum detectors implemented as horizontal "Roman Pots" style sensors.

- Same technology choice(s) as for the Roman Pots.
- Need to also study use of OMD on other side for tagging negative pions.

```
Protons
```

Protons
123.75 < E < 151.25 GeV
123.75 < E < 151.25 GeV
(45%<\frac{\mp@subsup{p}{z,proton}{}}{\mp@subsup{p}{z,\mathrm{ beam }}{}}<55%)
(45%<\frac{\mp@subsup{p}{z,proton}{}}{\mp@subsup{p}{z,\mathrm{ beam }}{}}<55%)
0<\boldsymbol{0}<5\textrm{mrad}

```
0<\boldsymbol{0}<5\textrm{mrad}
```


## Preliminary CAD drawings of RP and OMD Supports and Magnet Cryostats



## Summary of Detector Performance (Trackers)



- Includes realistic considerations for pixel sizes and materials
- More work needed on support structure and associated impacts.
- Roman Pots and Off-Momentum detectors suffer from additional smearing due to improper transfer matrix reconstruction.
- This problem is close to being solved!


## Digression: particle beams

- Angular divergence
- Angular "spread" of the beam away from the central trajectory.
- Gives some small initial transverse momentum to the beam particles.

- Crab cavity rotation
- Can perform rotations of the beam bunches in 2D.
- Used to account for the luminosity drop due to the crossing angle allows for head-on collisions to still take place.



## Summary of Detector Performance (Trackers)



- All beam effects included!
- Angular divergence.
- Crossing angle.
- Crab rotation/vertex smearing.


## Beam effects the dominant source of momentum

 smearing!
## Zero-Degree Calorimeter



## Zero-Degree Calorimeter

Credit to Shima Shimizu


Credit to Ron Lassiter

## Summary and Takeaways

- All FF detector acceptances and detector performance well-understood with currently available information.
- Numerous impact studies done!
- Primary technology choices done for all subsystems, with exception to systems dependent on evolving engineering design (e.g. B0 calorimetry).
- More realistic engineering considerations need to be added to simulations as design of IR vacuum system and magnets progresses toward CD-2/3a.
- Lots of work is currently on-going!
- Very busy fall ahead of us!


## Working Group Information

## - Far-Backward

- Conveners: Igor Korover [korover@mit.edu](mailto:korover@mit.edu), Nick Zachariou [nick.zachariou@york.ac.uk](mailto:nick.zachariou@york.ac.uk), Krzyzstof Piotrzkowski [krzysztof.piotrzkowski@cern.ch](mailto:krzysztof.piotrzkowski@cern.ch), Jaroslav Adam [jaroslavadam299@gmail.com](mailto:jaroslavadam299@gmail.com)
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- Wiki: https://wiki.bnl.gov/EPIC/index.php?title=FarForward


## Backup

# The Far-Forward Detectors 



Digression: Machine Optics
275 GeV DVCS Proton Acceptance



High Divergence: smaller $\beta^{*}$ at IP, but bigger $\beta(z=30 m)$-> higher lumi., larger beam at RP

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275 GeV DVCS Proton Acceptance




High Divergence: smaller $\beta^{*}$ at IP, but bigger $\beta(z=30 m)$-> higher lumi., larger beam at RP

High Acceptance: larger $\beta^{*}$ at IP, smaller $\beta(z=30 m)$->
lower lumi., smaller beam at RP

## Digression: Machine Optics

275 GeV DVCS Proton Acceptance



High Acceptance: larger $\beta^{*}$ at IP, smaller $\beta(z=30 m)$->
lower lumi., smaller beam at RP

So how does the FF system perform for measurements (non-exhaustive)?

## Off-Momentum Detectors

## Off-Momentum Detectors

- Off-momentum protons $\rightarrow$ smaller magnetic rigidity $\rightarrow$ greater bending in dipole fields.


## Luminosity Monitor

- Must make measurement in challenging environment.
- High synchrotron radiation, high bremsstrahlung rates ( $\sim 10 \mathrm{GHz}$ ), etc.
- Need $\sim 1 \%$ for absolute luminosity measurement, $\sim 10^{-4}$ for relative luminosity measurement.
- Can make direct photon measurement, or indirect via pair conversion in exit window, where $\mathrm{e}^{+} \mathrm{e}^{-}$pair is steered toward two calorimeters opposite a dipole magnet.
- Direct photon calorimeter includes moveable SR filters/monitors (F1 and F2), and has configurations for high (PCALf) and low (PCALc) luminosity running.



## Luminosity Monitor

- Conversion layer is part of beam layout
- Need for precise knowledge of conversion probability
- Heat load from synchrotron radiation is incident on the layer
- Several considerations for the design:

II: thin converter in vacuum


I: baseline design, converter holds the vacuum


III: vacuum up to detectors


## Exit window for luminosity monitor

- Part of outgoing electron beam pipe
- Conversion layer for bremsstrahlung photons
- Tilt angle vs. electron (and photon) beam axis against synchrotron radiation



## Low-Q² Taggers

- Two taggers for reconstructing electrons from low- $\mathrm{Q}^{2}\left(<10^{-1} \mathrm{GeV}^{2}\right)$ reactions.
- Combination of EM calorimetry for energy reconstruction, and silicon layers (High Resolution Hodoscope - HIHS) for position and angular resolution.



## Performance for low-Q2 tagger

- Tagger 1 and 2 are placed closer (further) from the IP
- Overlap in Q2 acceptance (<0.1 GeV^2)
- Complementary in electron energy (higher energies reach Tagger 2)
- Consistent for Pythia6 and quasi-real photoproduction (QR)

(a) Tagger 1

(b) Tagger 2




## Machine Optics: Roman Pots



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## Momentum Resolution - Timing

For exclusive reactions measured with the Roman Pots we need good timing to resolve the position of the interaction within the proton bunch. But what should the timing be?


RMS hadron bunch length $\sim 10 \mathrm{~cm}$.


Looking along the beam with no crabbing.
${ }^{\sim} 1.25 \mathrm{~mm}$


What the RP sees.

- Because of the rotation, the Roman Pots see the bunch crossing smeared in x .
- Vertex smearing $=12.5 \mathrm{mrad}$ (half the crossing angle) $* 10 \mathrm{~cm}=1.25 \mathrm{~mm}$
- If the effective vertex smearing was for a 1 cm bunch, we would have .125 mm vertex smearing.
- The simulations were done with these two extrema and the results compared.
$>$ From these comparisons, reducing the effective vertex smearing to that of the 1 cm bunch length reduces the momentum smearing to negligible from this contribution.
- This can be achieved with timing of $\sim 35$ ps ( $1 \mathrm{~cm} /$ speed of light).



## Momentum Resolution - Comparison

- The various contributions add in quadrature (this was checked empirically, measuring each effect independently).

$$
\Delta p_{t, \text { total }}=\sqrt{\underbrace{\left(\Delta p_{t, A D}\right)^{2}}_{\begin{array}{l}
\text { Angular } \\
\text { divergence }
\end{array}}+\underbrace{\left(\Delta p_{t, C C}\right)^{2}}_{\begin{array}{l}
\text { Primary vertex } \\
\text { smearing from crab } \\
\text { cavity rotation. }
\end{array}}+\underbrace{\left(\Delta p_{t, p x l}\right)^{2}}_{\begin{array}{l}
\text { Smearing from } \\
\text { finite pixel size. }
\end{array}}}
$$


- Beam angular divergence
- Beam property, can't correct for it - sets the lower bound of smearing.
- Subject to change (i.e. get better) - beam parameters not yet set in stone
- Vertex smearing from crab rotation
- Correctable with good timing ( $\sim 35$ ps)
- Finite pixel size on sensor


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