

# Hyperon Form Factors

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on behalf of the BESIII collaboration

International Conference on Quarks and Nuclear Physics 2022  
Online  
2022-09-06



# The Strong Interaction

Non-perturbative effects manifest in:

- Spin
- Mass
- **Structure** ← **Our focus**

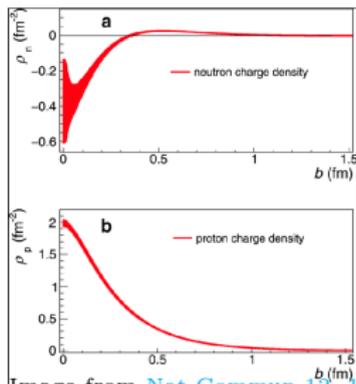
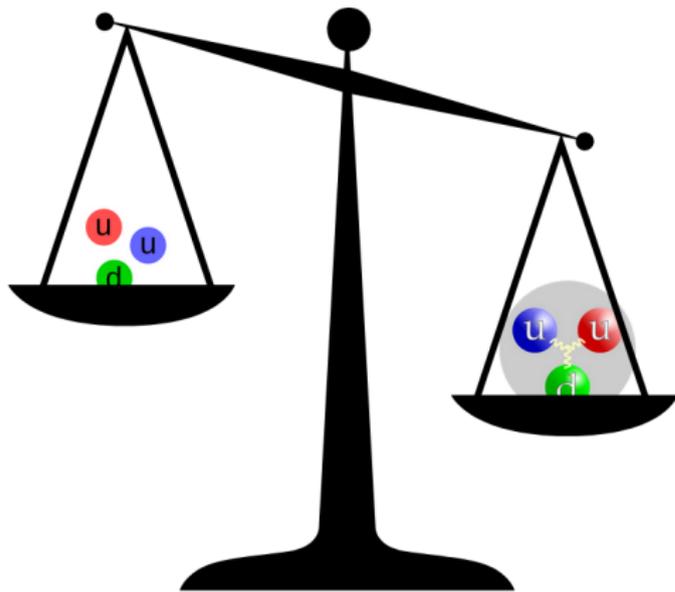


Image from [Nat Commun 12, 1759 \(2021\)](#).

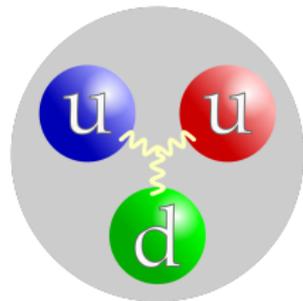
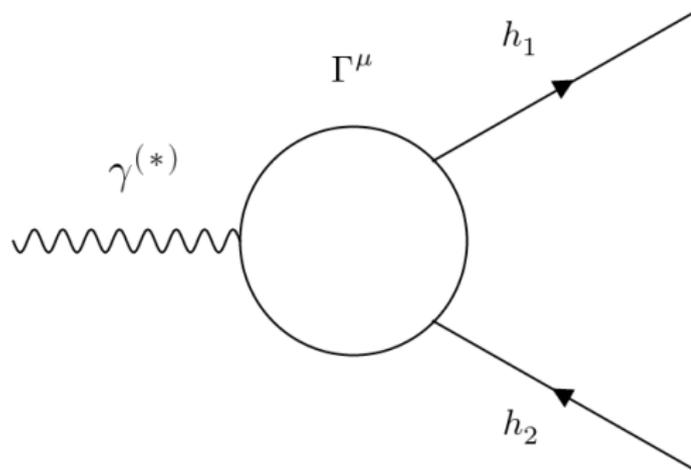


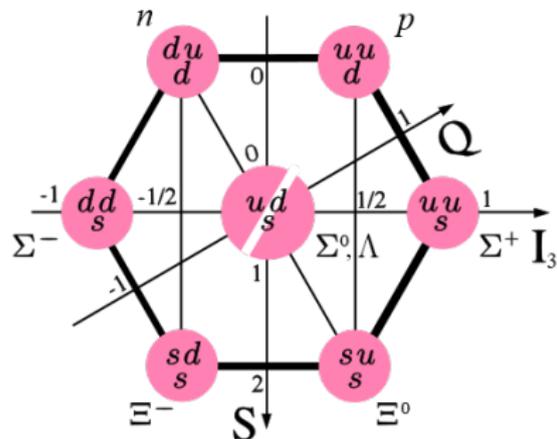
Image credit: [Jacek Rybak](#)

# Electromagnetic Form Factors

- Quantify deviation from pointlike nature
- Scalar functions of momentum transfer  $q^2$
- Elastic/transition ( $h_1 = h_2/h_1 \neq h_2$ )
- Spacelike/timelike ( $q^2 < 0/q^2 > 0$ )



# Hyperons



Hyperon	Mass [GeV/ $c^2$ ]	Decay (BF)
$\Lambda$	1.116	$p\pi^-$ (63.9%) $n\pi^0$ (35.8%)
$\Sigma^-$	1.197	$n\pi^-$ (99.8%)
$\Sigma^+$	1.189	$p\pi^0$ (51.6%) $n\pi^+$ (48.3%)
$\Xi^0$	1.315	$\Lambda\pi^0$ (99.5%)
$\Xi^-$	1.321	$\Lambda\pi^-$ (99.8%)
$\Omega$	1.672	$\Lambda K^-$ (67.8%) $\Xi^0\pi^-$ (23.6%) $\Xi^-\pi^0$ (8.6%)

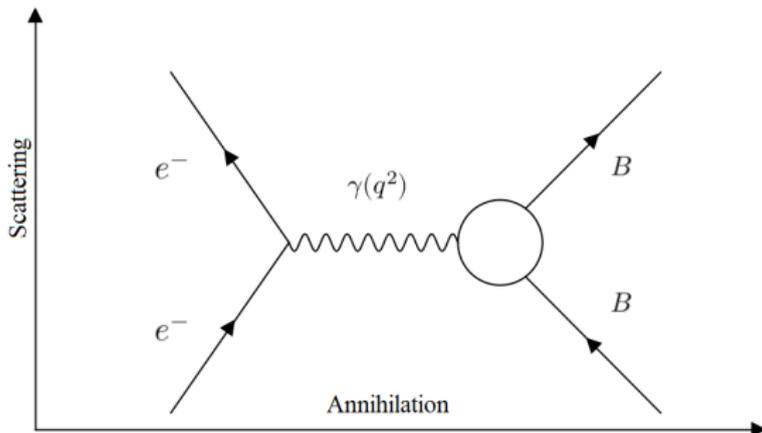
# One-Photon Exchange

**Spin 1/2 Baryons:** Two independent form factors

Matrix element:

$$\Gamma^\mu = \gamma^\mu F_1(q^2) + \frac{i\sigma^{\mu\nu}q_\nu}{2M} F_2(q^2)$$

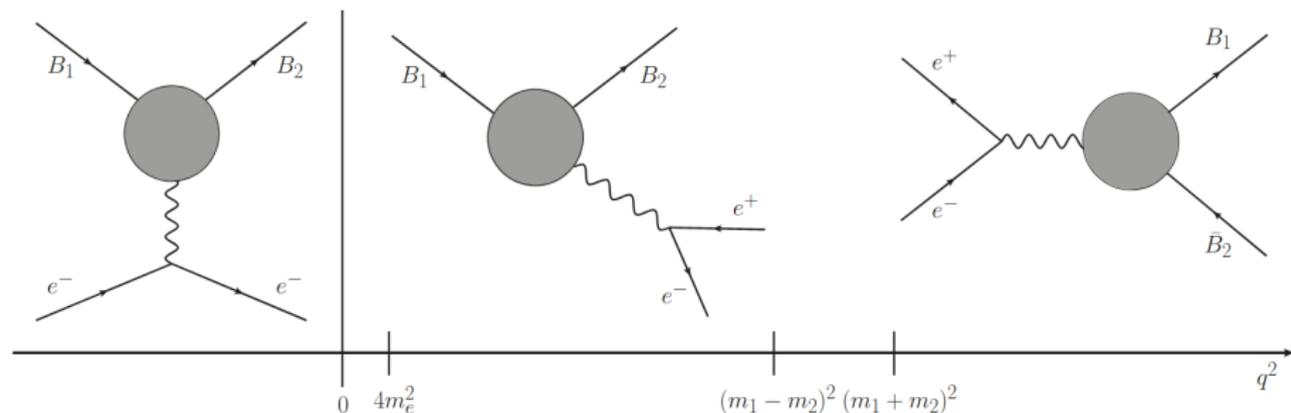
- Dirac and Pauli form factors  $F_1(q^2)$ ,  $F_2(q^2)$
- Sachs form factors  $G_E(q^2) = F_1 + \frac{q^2}{m_B^2} F_2$ ,  $G_M(q^2) = F_1 + F_2$   
 $G_E(0) = Q$ ,  $G_M(0) = \mu$



# Electromagnetic Form Factors

Spacelike ( $q^2 < 0$ )

Timelike ( $q^2 > 0$ )

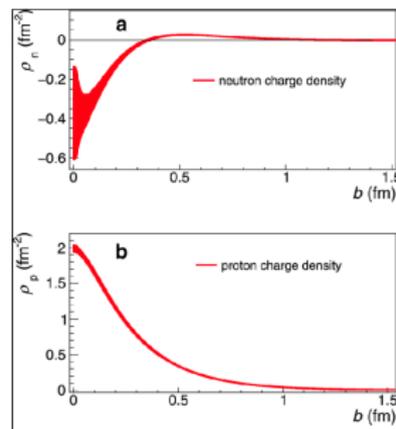


**BESIII contributes here**

Image credit: Perotti, E., PhD Thesis, Uppsala University (2020)

# Spacelike Form Factors

- Can be studied in elastic lepton scattering
- **Hyperons are unstable** → **Difficult!**
- Real-valued functions of  $q^2$ .
- Related to charge and magnetization densities



Nat. Commun. 12, 1759 (2021).

$$\langle r_E \rangle^2 = 6 \left. \frac{dG_E(q^2)}{dq^2} \right|_{q^2=0}$$
$$\langle r_M \rangle^2 = \frac{6}{G_M(0)} \left. \frac{dG_M(q^2)}{dq^2} \right|_{q^2=0}$$

# Timelike Form Factors

- Can be studied in lepton-antilepton annihilation  
**Hyperon FFs experimentally accessible!**
- Complex functions of  $q^2$ :  
 $G_M(q^2) = |G_M(q^2)|e^{i\Phi_M}$ ,  $G_E(q^2) = |G_E(q^2)|e^{i\Phi_E}$
- Observables  $R = |G_E/G_M|$ ,  $\Delta\Phi = \Phi_E - \Phi_M$

**Cross section for  $e^+e^- \rightarrow B\bar{B}$ :**

$$\frac{d\sigma(q^2)}{d\Omega} = \frac{\alpha^2\beta}{4q^2} \left( \frac{1}{\tau} |G_E|^2 \sin^2 \theta + |G_M|^2 (1 + \cos^2 \theta) \right)$$

**Effective form factor:**

$$|G_{eff}(q^2)|^2 = \frac{2\tau |G_M(q^2)|^2 + |G_E(q^2)|^2}{2\tau + 1} = \frac{\sigma_{Born}(q^2)}{\left(1 + \frac{1}{2\tau}\right) \left(\frac{4\pi\alpha^2\beta C}{3q^2}\right)}$$

# Phase Measurement

## What is the significance of $\Delta\Phi$ ?

- SL and TL FFs related by dispersion relations.

- As  $|q^2| \rightarrow \infty$ : **SL**  $\rightarrow$  **TL**

$$\implies \Delta\Phi = n \cdot \pi$$

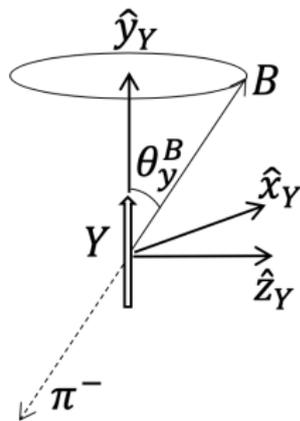
- Oscillations of  $\Delta\Phi$  reveal zero-crossings

Phys. Rev. D 104, 116016 (2021)

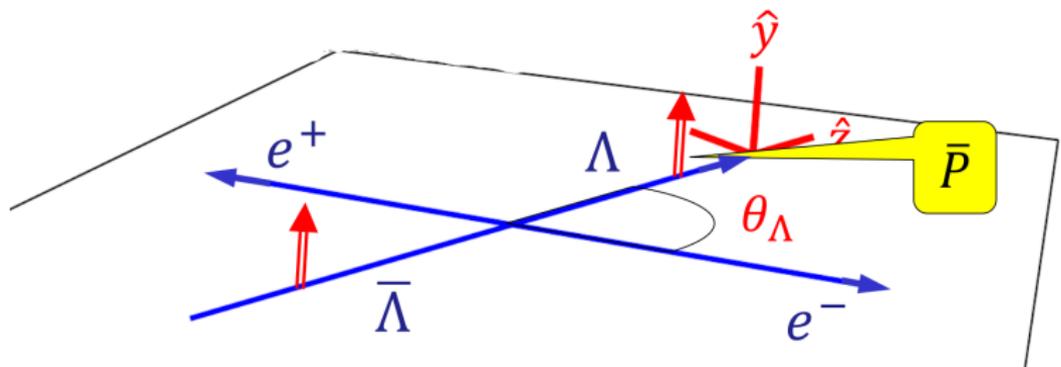
- Provides constraints for unmeasurable SL FFs

## How to measure it?

- If  $\sin \Delta\Phi \neq 0$   $B/\bar{B}$  can be polarized
- Experimental access to polarization in self-analyzing weak decays of hyperons!



# $B\bar{B}$ Production in $e^+e^-$ -Annihilation



Unpolarized  $e^+e^-$  beams  $\rightarrow$  Transverse polarization:

$$P_y(\cos\theta_\Lambda) = \frac{\sqrt{1-\eta^2} \cos\theta_\Lambda \sin\theta_\Lambda}{1+\eta \cos^2\theta_\Lambda} \sin(\Delta\Phi)$$

Angular distribution:  $\frac{d\Gamma}{d\Omega} \propto 1 + \eta \cos^2\theta_\Lambda$ ,  $-1 \leq \eta \leq 1$

$$R = \sqrt{\tau} \sqrt{\frac{1-\eta}{1+\eta}}$$

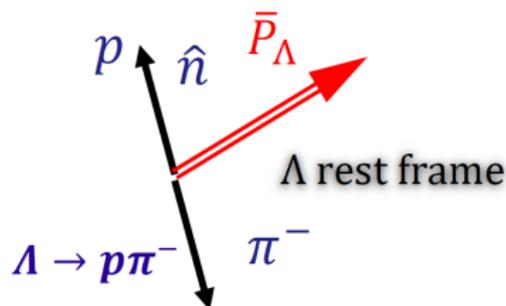
# Self-Analyzing Hyperon Decays

Angular distribution of daughter baryon reveals mother hyperon polarization

$$\frac{d\Gamma}{d\Omega} = \frac{1}{4\pi}(1 + \alpha_\Lambda \hat{n} \bar{P}_\Lambda)$$

Decay asymmetry parameter  $\alpha_\Lambda$

$CP$ -conservation:  $\alpha_\Lambda = -\alpha_{\bar{\Lambda}}$



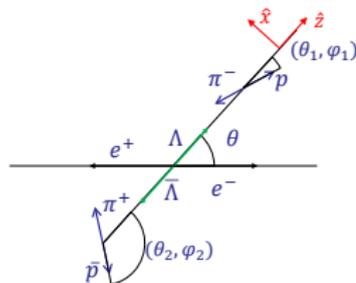
Precise measurements of  $\alpha_\Lambda$  from BESIII via  $J/\psi \rightarrow \Lambda \bar{\Lambda}^1$ ,  $J/\psi \rightarrow \Xi \bar{\Xi}^2$

BESIII, <sup>1</sup>Nature Phys. 15 (2019) 631, <sup>2</sup>Nature 606, 64–69 (2022)

# Formalism for the Process $e^+e^- \rightarrow \Lambda\bar{\Lambda} \rightarrow p\pi^-\bar{p}\pi^+$

Full reaction described by:  $\xi = (\theta, \theta_1, \phi_1, \theta_2, \phi_2) = (\theta, \Omega_1, \Omega_2)$

Fäldt, Kupsc, Phys. Lett. B 772 (2017) 16-20



$$d\Gamma \propto \mathcal{W}(\xi) = \mathcal{F}_0(\xi) + \eta \mathcal{F}_5(\xi)$$

$$- \alpha_{\Lambda}^2 \left( \mathcal{F}_1(\xi) + \sqrt{1 - \eta^2} \cos(\Delta\Phi) \mathcal{F}_2(\xi) + \eta \mathcal{F}_6(\xi) \right)$$

$$+ \alpha_{\Lambda} \sqrt{1 - \eta^2} \sin(\Delta\Phi) (\mathcal{F}_3(\xi) - \mathcal{F}_4(\xi))$$

Spin correlations

Polarization

- $\mathcal{F}_i$  are known functions of the angles
- Same method used by BESIII for precise  $CP$ -tests in  $J/\psi \rightarrow \Lambda\bar{\Lambda}, \Sigma\bar{\Sigma}, \Xi\bar{\Xi}$  (see talk by He Li on Wednesday)

When only  $\Lambda/\bar{\Lambda}$  is measured two angles  $\xi = (\theta, \theta_p)$  are sufficient.

$$\mathcal{W}_{\Lambda/\bar{\Lambda}}(\xi) = 1 + \eta \cos^2 \theta + \alpha_{\Lambda/\bar{\Lambda}} \sqrt{1 - \eta^2} \sin(\Delta\Phi) \sin \theta \cos \theta \cos \theta_p$$

# Beijing Electron-Positron Collider (BEPCII)

CMS Energy: 2-4.95 GeV

Design luminosity:  $10^{33} \text{ cm}^{-2} \text{ s}^{-1}$

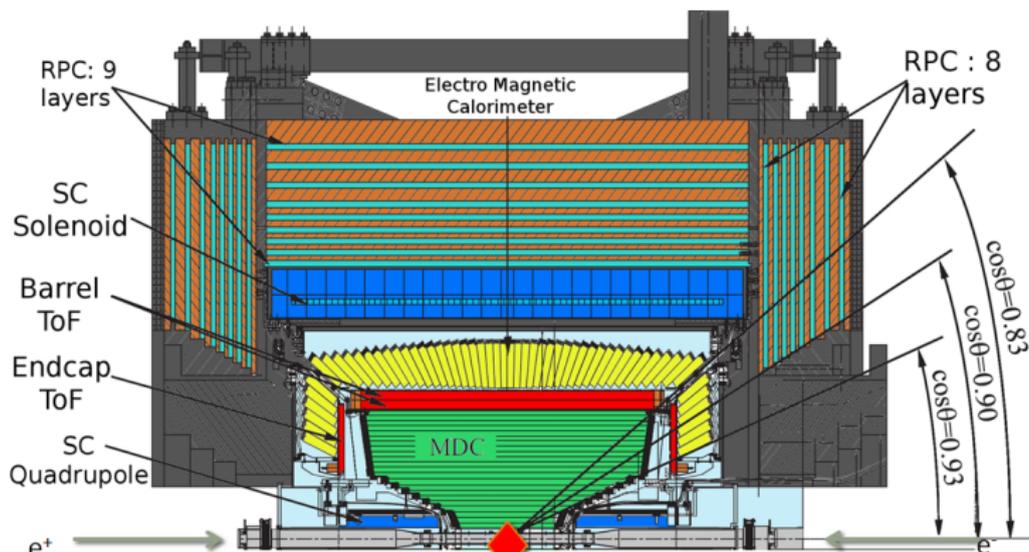
Linac

BEPCII



# Beijing Spectrometer (BESIII)

- Near  $4\pi$  coverage
- Helium-gas drift chamber
- CsI(Tl) crystal calorimeter
- Scan data:
  - $0.5 \text{ fb}^{-1}$  from 2.0 to 3.08 GeV
  - $20 \text{ fb}^{-1}$  from 3.51 to 4.6 GeV
- Plastic scintillator TOF-system (endcaps upgraded to MRPC in 2015)
- 1 T super-conducting solenoid
- RPC-based muon chamber



# $e^+e^- \rightarrow \Lambda\bar{\Lambda}, \Lambda \rightarrow p\pi^- + c.c.$ at 2.396 GeV

$\alpha_\Lambda = -\alpha_{\bar{\Lambda}} = 0.750$  fixed

555 events (2.5% bkg)

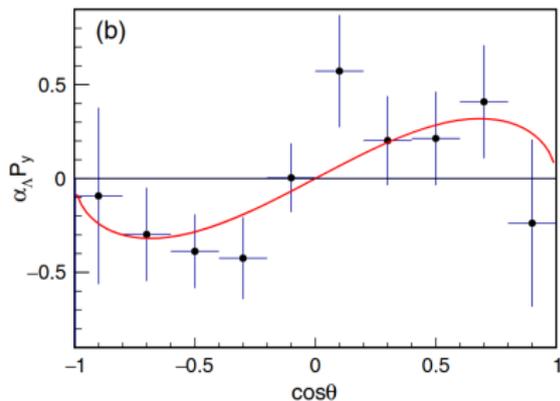
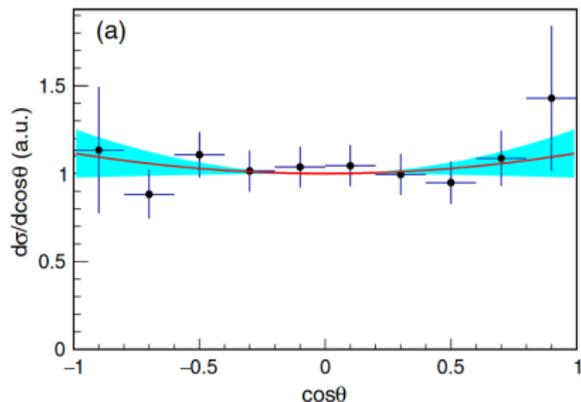
- Two-photon exchange contribution negligible
- **First measurement of relative phase for any baryon!!!**

$$R = 0.96 \pm 0.14 \pm 0.02$$

$$\Delta\Phi = 37 \pm 12 \pm 6^\circ$$

$$\sigma_{Born} = 118.7 \pm 5.3_{stat} \pm 5.1_{syst} \text{ pb}$$

BESIII



$e^+e^- \rightarrow \Lambda\bar{\Lambda}, \Lambda \rightarrow p\pi^- + c.c. \text{ at High } q^2$

$\mathcal{B}(\psi(3770) \rightarrow \Lambda\bar{\Lambda}) > 10$  times larger than previously assumed<sup>1</sup>

**Spin analysis at 3.773 GeV<sup>2</sup>:**

$\alpha_\Lambda = -\alpha_{\bar{\Lambda}} = 0.754$  fixed

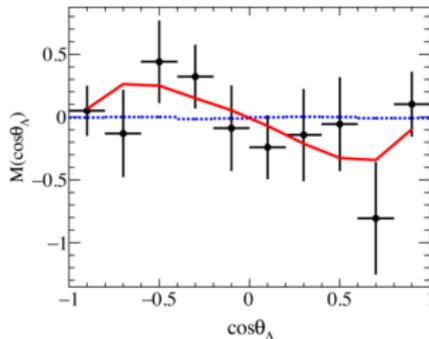
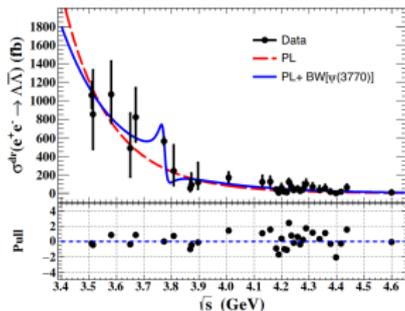
262 events (0.5% bkg)

$$\eta_\psi = 0.85_{-0.20}^{+0.12} \pm 0.02$$

$$R_\psi = 0.48_{-0.35}^{+0.21} \pm 0.03$$

$$\Delta\Phi_\psi = 71_{-46}^{+66} \pm 5^\circ$$

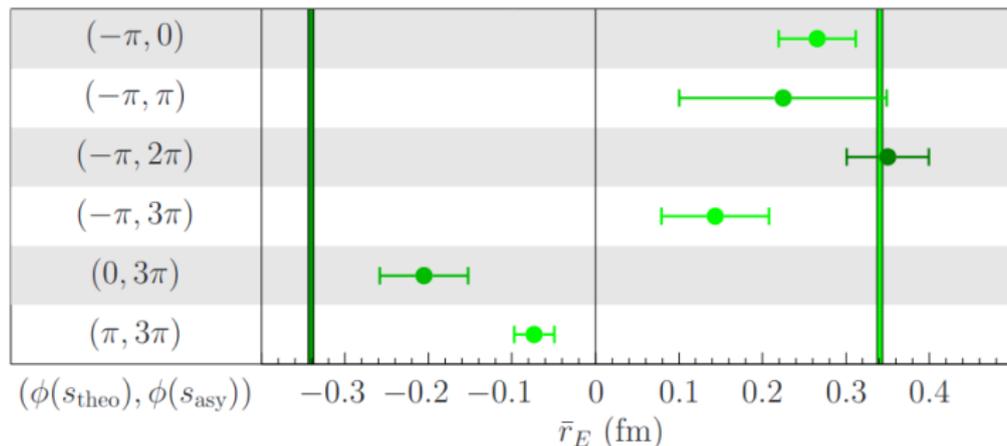
**BESIII**



<sup>1</sup>BESIII, Phys. Rev. D 104 (2021) 9, L091104 <sup>2</sup>BESIII, Phys. Rev. D 105 L011101 (2022)

# Interpretation

Phase measurement helps pin down charge  $\Lambda$  radius



Fit<sup>1</sup> to data from BESIII<sup>2</sup> and BaBar<sup>3</sup> in different scenarios

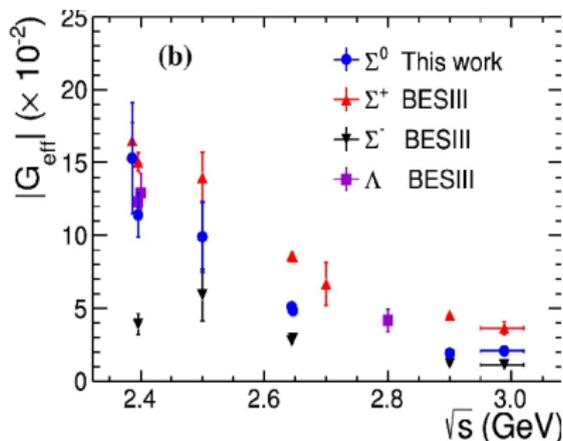
**Next step: What is the  $q^2$ -dependence of  $\Delta\Phi$ ?**

<sup>1</sup>Mangoni et al., Phys. Rev. D 104, 116016 (2021) <sup>2</sup>BESIII, Phys. Rev. Lett. 123 122003 (2019)

<sup>3</sup>BaBar, Phys. Rev. D 76, 092006 (2007)

# Effective FFs of Single Strange Hyperons

- Asymmetry in  $\Sigma$  isospin triplet FFs
- Relation between  $\Lambda/\Sigma$  FFs as expected from diquark correlations
- $\sigma_{Born}(e^+e^- \rightarrow \Sigma^+\bar{\Sigma}^-)/\sigma_{Born}(e^+e^- \rightarrow \Sigma^-\bar{\Sigma}^+) = 9.7 \pm 1.3$   
→ inconsistent with predictions from *e.g.*  
Phys. Rev. D 101, 014014 (2020)



<sup>1</sup>BESIII, Phys. Rev. D 97, 032013 (2018) <sup>2</sup>BESIII, Phys. Lett. B 814, 136110 (2021)

<sup>3</sup>BESIII, Phys. Lett. B 831, 137187 (2022)

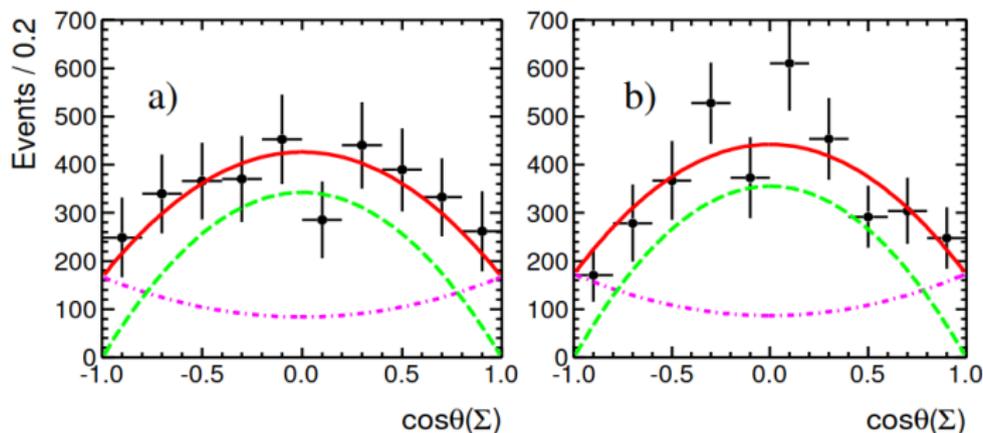
$$e^+e^- \rightarrow \Sigma^+\bar{\Sigma}^-$$

Form factor ratio  $R = |G_E/G_M|$  for  $\Sigma^+$

$\sqrt{s}$ [GeV]	$R$
2.3960	$1.83 \pm 0.26 \pm 0.24$
2.6454	$0.66 \pm 0.15 \pm 0.11$
2.9000	$1.06 \pm 0.36 \pm 0.09$



Example fit at  $\sqrt{s} = 2.396$  GeV:

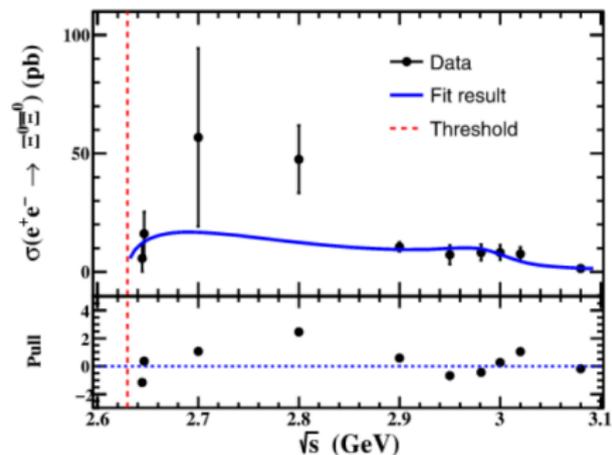


Solid line: Fit result, Dashed line:  $G_E$ , Dotted line:  $G_M$

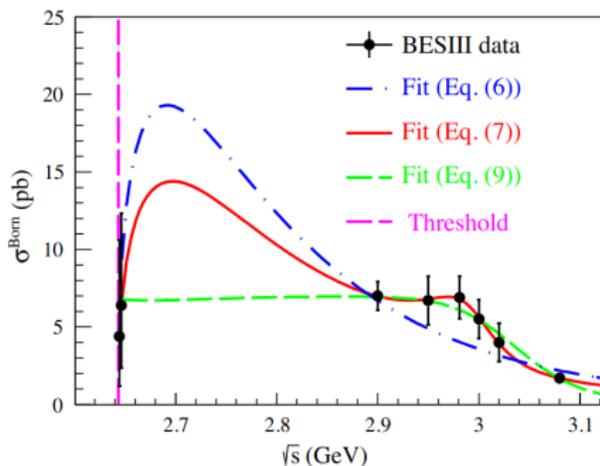
$$e^+e^- \rightarrow \Xi \bar{\Xi}$$

- $\sigma_{Born}$  and  $|G_{eff}|$  for  $e^+e^- \rightarrow \Xi^0\bar{\Xi}^0, \Xi^-\bar{\Xi}^+$  near threshold
- $\sim 2\sigma$  resonant structure at 3.0 GeV

$$e^+e^- \rightarrow \Xi^0\bar{\Xi}^0$$



$$e^+e^- \rightarrow \Xi^-\bar{\Xi}^+$$



BESIII

<sup>1</sup>BESIII: Phys. Rev. D 103, 012005 (2021) <sup>2</sup> BESIII, Phys. Lett. B 820, 136557 (2021)

# Summary & Outlook

- Hyperon EMFFs are important observables for understanding the structure of baryons
- Spin-polarisation analysis enables determination of relative phase  $\Delta\Phi$
- $\Delta\Phi$  helps pin down spatial structure
- Many new results from BESIII
  - **First phase measurement:**  $\Delta\Phi$  for  $\Lambda^0$
  - $R = |G_E/G_M|$  for  $\Lambda^0$  and  $\Sigma^+$
  - Effective form factors for  $\Lambda^0$ ,  $\Sigma^\pm$ ,  $\Sigma^0$ ,  $\Xi^0$ ,  $\Xi^-$
  - More data to be analysed!