

# Two-photon transitions of charmonia

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# Outline

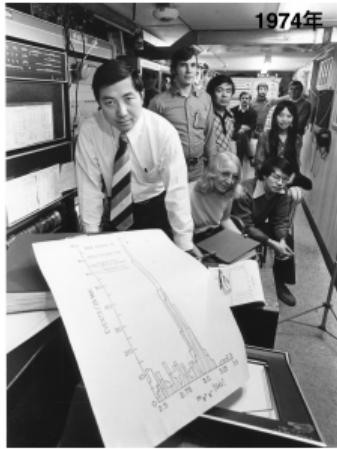
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- ▶ Physics of two-photon transition of charmonium
- ▶ Light front dynamics and basis light-front quantization
- ▶ Numerical results: two-photon width and transition form factors
- ▶ Summary

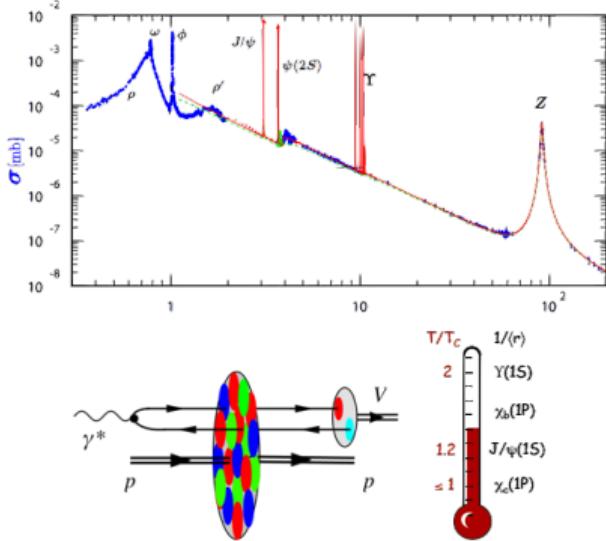
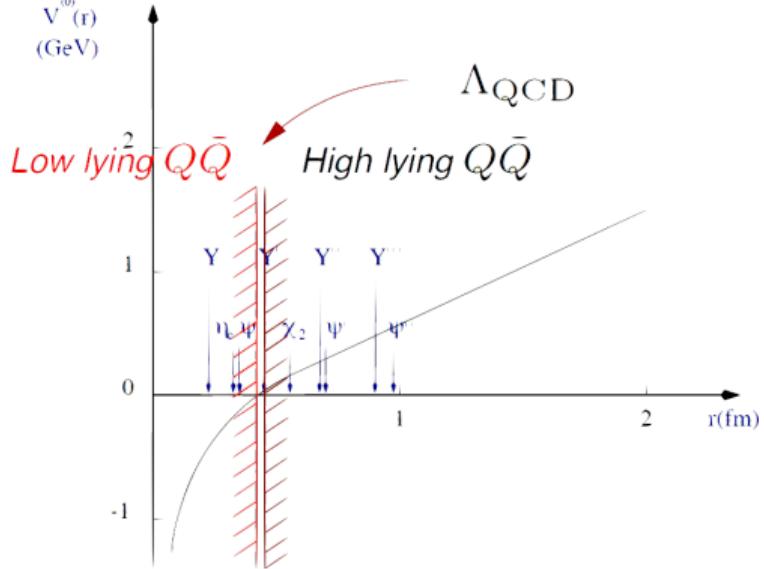
Based on: YL, M. Li (李枚键) and J.P. Vary, Phys. Rev. D letter **105** (2022); arXiv:2111.14178 [hep-ph]

Wave functions available on Mendeley Data, doi: 10.17632/cjs4ykv8cv.2

# Charmonium: "a golden system to study strong interactions"



November Revolution



Tremendous applications in ee, ep, eA, pp, AA collisions

- Theoretically a hard problem: multiscale, multi-physics

$$\Lambda_{\text{QCD}} \lesssim \alpha_s^2 m_c < \alpha_s m_c < m_c$$

- Physically a simple system: nonrelativistic ( $v_c \ll 1$ ), perturbative ( $\alpha_s \ll 1$ )

$$v_c^2 \sim 0.3, \alpha_s(m_c) \sim (0.3 - 0.6) (?)$$

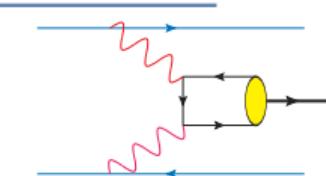
- Potential model, pNRQCD, Lattice QCD

[Brambilla '14]

# Two-photon transitions of charmonia

A clean & important probe to hadron structures [Reviews: Berger '87]

- ▶ Decay  $H \rightarrow \gamma + \gamma$ : golden channel for hadron identification
  - ▶ Selection rules:  $P, C, \text{angular momentum, gauge symmetry, ...}$
- ▶ Exclusive photoproductions  $\gamma + \gamma \rightarrow H$ : channel of discovery
  - ▶  $p\text{QCD factorization at large } Q^2,$



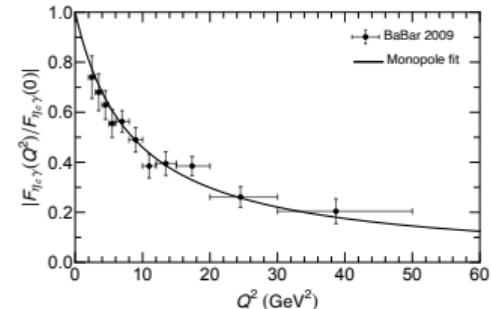
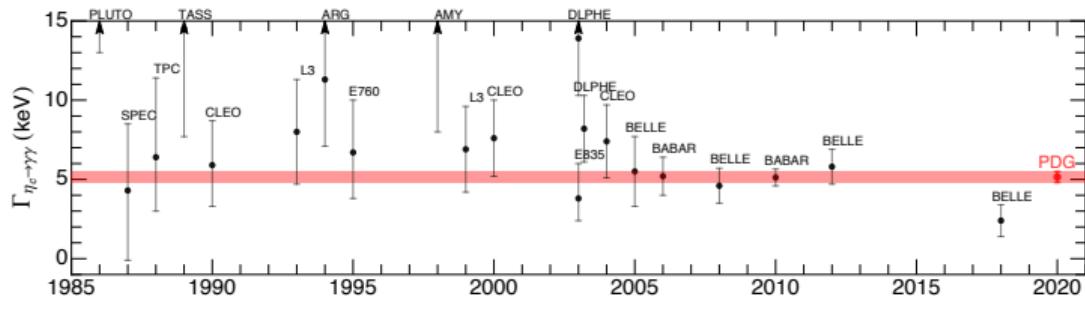
[Lepage '80 & Chernyak '84]

$$F_{\gamma\eta_c}(Q^2) = \int_0^1 dx T_H(x, Q^2) \phi_{\eta_c}(x; \mu),$$

## Experimental measurements

[Review of particle physics 2020]

- ▶ Diphoton width: extensive measurements for  $\eta_c, \eta'_c, \chi_{c0}, \chi_{c2}, X(3872)$ ;
- ▶ Transition form factors:  $F_{\eta_c\gamma}(Q^2)$  by BABAR 2010;  $F_{\chi_c\gamma}(Q^2)$  by Belle 2017 with limited statistics



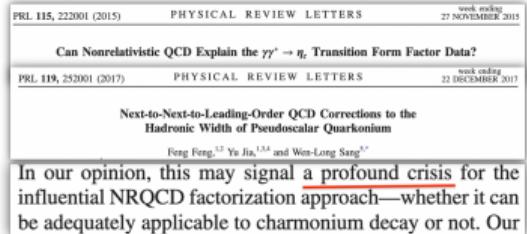
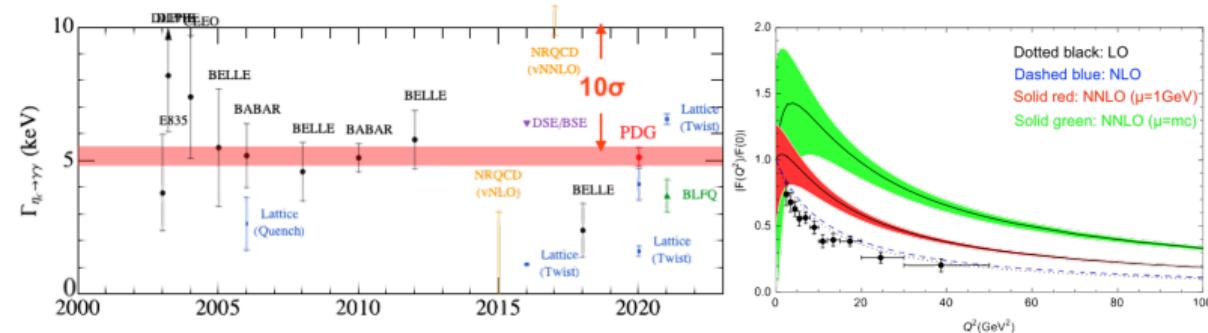
# A crisis in theories for charmonium?

## Status of theoretical predictions

- ▶ Potential model: large relativistic corrections [e.g. Babiarz '19]
- ▶ NRQCD:  $10\sigma$  discrepancy at NNLO -- a crisis for NRQCD? [Feng PRL '15&'17]
- ▶ Lattice QCD: challenge of representing  $\gamma^*$  on the lattice [e.g., Liu '20]

## Why charmonium is so challenging?

- ▶  $\alpha_s \sim (0.3 - 0.6)$  is not that small -- non-perturbative effects
- ▶  $v_c^2 \sim 0.3$  is not that small -- relativistic effects
- ▶  $am_c \sim 0.5$  is not that small -- high order  $O(a^2)$  effects



In our opinion, this may signal a profound crisis for the influential NRQCD factorization approach—whether it can be adequately applicable to charmonium decay or not. Our

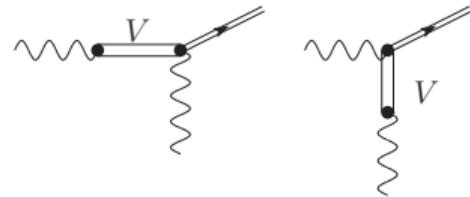
# T $\gamma$ P: from vector meson dominance to light-cone dominance

[Berger '87]

$Q^2 \rightarrow 0$

- ▶ Low  $Q^2$ : vector meson dominance [Sakurai '63, Novikov '78]

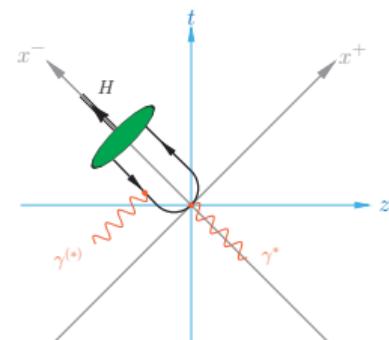
$$i\mathcal{M} \sim (Q^2 + M_V^2)^{-1} \underbrace{\psi(\vec{r} = 0)}_{\text{wave function at origin}}$$



- ▶ Large  $Q^2$ : light-cone dominance [Lepage '80 & Chernyak '84]

$$i\mathcal{M} = \epsilon_\mu \epsilon_\nu^* \int d^4x e^{iq \cdot z} \langle 0 | J^\mu(z) J^\nu(0) | P \rangle \sim \int dx T_H(x, Q^2) \underbrace{\phi_P(x; \mu)}_{\text{light-cone distribution amplitude}}$$

Large- $Q^2$  limit:  $z^2 \sim 1/Q^2 \rightarrow 0$  (the light cone) [Gribov '83, Nandi '07 & Li '09]



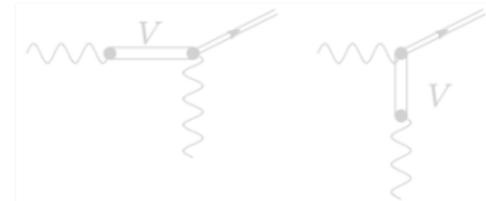
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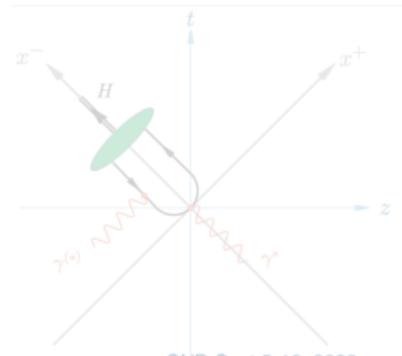


- ▶ Intermediate  $Q^2$

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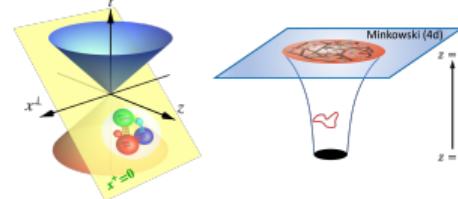
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- ▶ Intermediate  $Q^2$ : get off the light cone ( $z^2 = 0$ )  $\rightarrow$  light front ( $z^+ = 0$ )

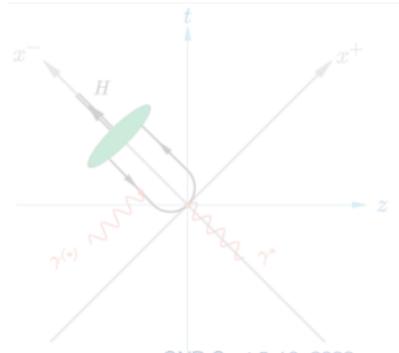
$$i\mathcal{M} \sim \int_0^1 dx \int \frac{d^2 k_\perp}{16\pi^3} T_H(x, \vec{k}_\perp; Q^2) \underbrace{\psi_P(x, \vec{k}_\perp; \mu)}_{\text{light-front wave function}}$$



- ▶ Large  $Q^2$ : light-cone dominance [Lepage '80 & Chernyak '84]

$$i\mathcal{M} = \epsilon_\mu \epsilon_\nu^* \int d^4x e^{iq \cdot z} \langle 0 | J^\mu(z) J^\nu(0) | P \rangle \sim \int dx T_H(x, Q^2) \underbrace{\phi_P(x; \mu)}_{\text{light-cone distribution amplitude}}$$

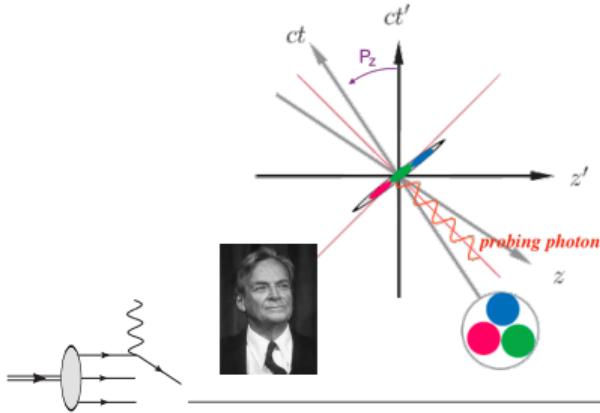
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# Physics on the light front

[Reviews: Brodsky '98, Burkardt '02, Bakker '14, Ji '21]

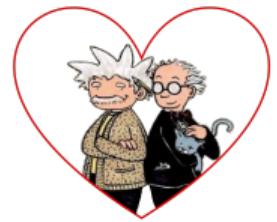
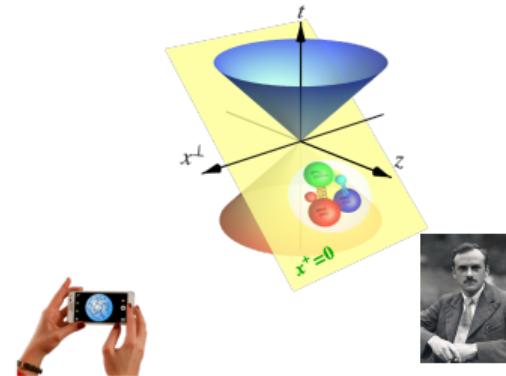
infinite momentum frame  
 $P_z \rightarrow \infty$  (Fubini, '65)



$$\begin{aligned}x^{\pm} &= x^0 \pm x^3 \\p^{\pm} &= p^0 \pm p^3 \\\underline{\mathcal{M}}^2 &= P^+ \underline{P}^- - \vec{P}_{\perp}^2\end{aligned}$$

$$\begin{aligned}i \frac{\partial}{\partial x^+} |\psi(x^+)\rangle &= \frac{1}{2} \underline{P}^- |\psi(x^+)\rangle \\&\Downarrow \\&\underline{\mathcal{M}}^2 |\psi_h(P, j, m_j)\rangle = M_h^2 |\psi_h(P, j, m_j)\rangle\end{aligned}$$

light front quantization  
 $x^+ = 0$  (Dirac, '49)

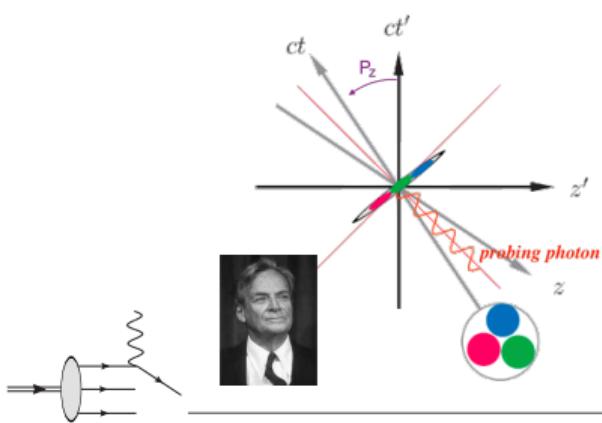


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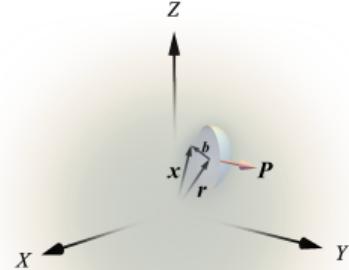
infinite momentum frame

$$P_z \rightarrow \infty \text{ (Fubini, '65)}$$



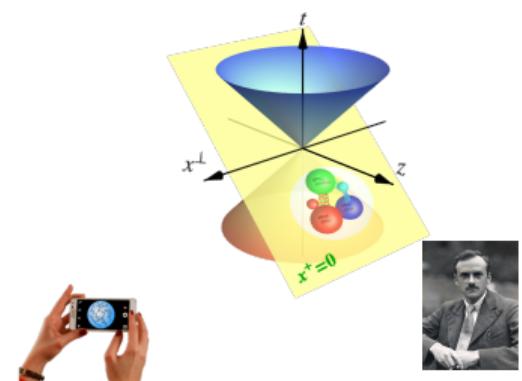
localized wavepacket

$$R_z \rightarrow 0 \text{ (2206.12903)}$$



light front quantization

$$x^+ = 0 \text{ (Dirac, '49)}$$



$$x^\pm = x^0 \pm x^3$$

$$p^\pm = p^0 \pm p^3$$

$$\underline{\mathcal{M}}^2 = P^+ \underline{P}^- - \vec{P}_\perp^2$$

$$i \frac{\partial}{\partial x^+} |\psi(x^+)\rangle = \frac{1}{2} \underline{P}^- |\psi(x^+)\rangle$$

⇓

$$\underline{\mathcal{M}}^2 |\psi_h(P, j, m_j)\rangle = M_h^2 |\psi_h(P, j, m_j)\rangle$$



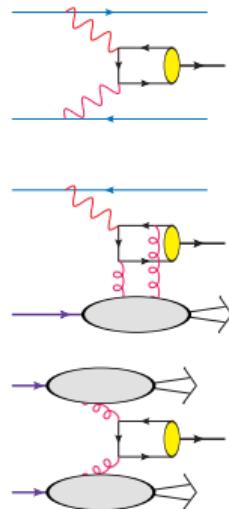
	<b>instant form</b>	<b>front form</b>	<b>point form</b>
time variable	$t = x^0$	$x^+ \triangleq x^0 + x^3$	$\tau \triangleq \sqrt{t^2 - \vec{x}^2 - a^2}$
quantization surface			
Hamiltonian	$H = P^0$	$P^- \triangleq P^0 - P^3$	$P^\mu$
kinematical	$\vec{P}, \vec{J}$	$\vec{P}^\perp, P^+, \vec{E}^\perp, E^+, J_z$	$\vec{J}, \vec{K}$
dynamical	$\vec{K}, P^0$	$\vec{F}^\perp, P^-$	$\vec{P}, P^0$
dispersion relation	$p^0 = \sqrt{\vec{p}^2 + m^2}$	$p^- = (\vec{p}_\perp^2 + m^2)/p^+$	$p^\mu = mv^\mu \ (v^2 = 1)$

$$P^\pm \triangleq P^0 \pm P^3, \vec{P}^\perp \triangleq (P^1, P^2), x^\pm \triangleq x^0 \pm x^3, \vec{x}^\perp \triangleq (x^1, x^2), E^i = M^{+i}, E^+ = M^{+-}, F^i = M^{-i}, \\ K^i = M^{0i}, J^i = \frac{1}{2}\epsilon^{ijk}M^{jk}.$$

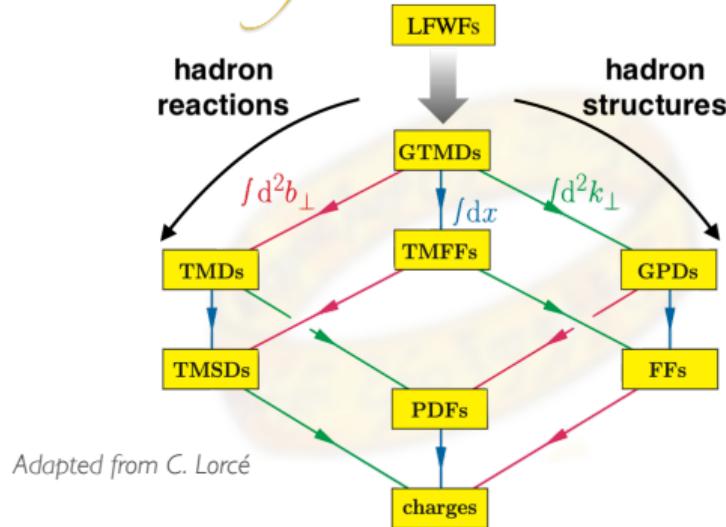
# Light-front wave functions (LFWFs)

[Reviews: Brodsky '98, Diehl '03, Lorcé '11]

LFWFs are frame independent and can directly access the partonic information of hadrons



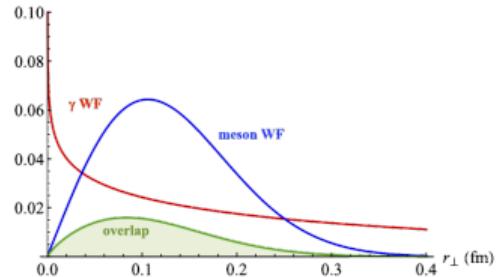
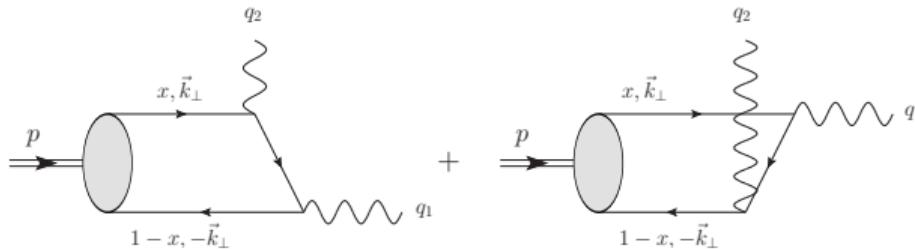
*One wavefunction to rule them all!*



*Hadron Physics without LFWFs is like Biology without DNA!*

— Stanley J. Brodsky

# LFWF representation of two-photon transitions



- ▶ The amplitude can be accessed in light-cone perturbation theory and also through hadronic matrix elements,  
[Lepage '80, Feldmann '97, Kroll '10, Babiarz '19]

$$\epsilon_\mu^*(q_1, \lambda_1) \epsilon_\nu^*(q_2, \lambda_2) e_\alpha(p, \lambda) \mathcal{M}^{\mu\nu\alpha} = \epsilon_\nu^*(q_2, \lambda_2) \langle \gamma^*(q_1, \lambda_1) | J^\nu(0) | H(p, \lambda) \rangle.$$

It is convenient to adopt a frame in which  $q_1^- = q_2^+ = 0$ , i.e. with manifest light-cone dominance.

- ▶ Example: LFWF representation of a pseudoscalar meson ( $0^{-+}$ ),

$$F_{P\gamma}(Q^2) = e_f^2 2\sqrt{2N_C} \int \frac{dx}{2\sqrt{x(1-x)}} \int \frac{d^2 k_\perp}{(2\pi)^3} \frac{\psi_{\uparrow\downarrow-\downarrow\uparrow/p}^*(x, \vec{k}_\perp)}{k_\perp^2 + m_f^2 + x(1-x)Q^2} + \dots$$

Intuitively, this is the **overlap** of the photon wave function with the meson wave function.  
[Beuf '16, Lappi '20]

# Light-front dynamics $\neq$ NR dynamics w. relativistic corrections

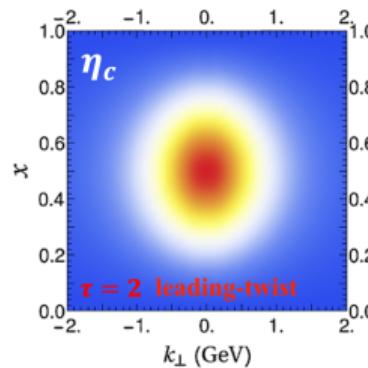
N.B. wave functions from non-relativistic dynamics with relativistic corrections in general are different from wave functions from relativistic dynamics, e.g. LFD.

- ▶ Parities in NRQM:  $P = (-1)^{L+1}, C = (-1)^{L+S}$  are approximations since  $L$  is not a good quantum number

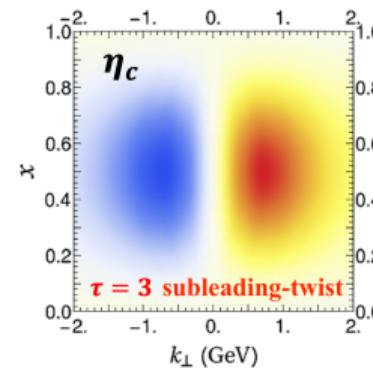
NRQM + relativistic correction: spin-orbital coupling  $\rightarrow$  partial wave mixing subject to parities

- ▶ Parities in LFD:  $m_P = (-1)^{L_z+S+1}, C = (-1)^{L_z+S+\ell}$  are **exact**

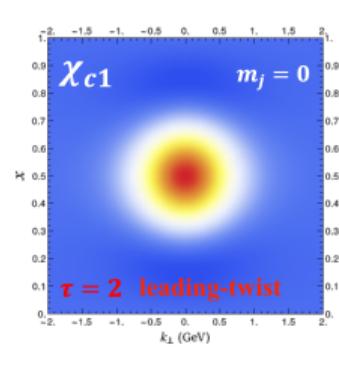
There exist leading-twist wavefunctions that are absent in NRQM (including relativistic corrections) due to parities



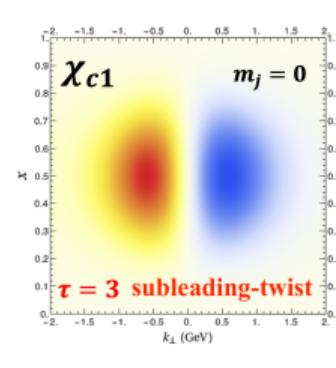
(a)  $\psi_{\uparrow\downarrow-\downarrow\uparrow}(\vec{k}_\perp, x)$



(b)  $\psi_{\uparrow\downarrow}(\vec{k}_\perp, x) = \psi_{\uparrow\uparrow}^*(\vec{k}_\perp, x)$



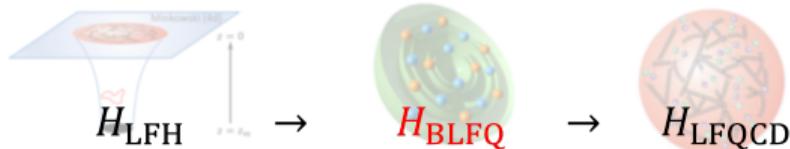
(c)  $\psi_{\uparrow\downarrow-\downarrow\uparrow}(\vec{k}_\perp, x)$



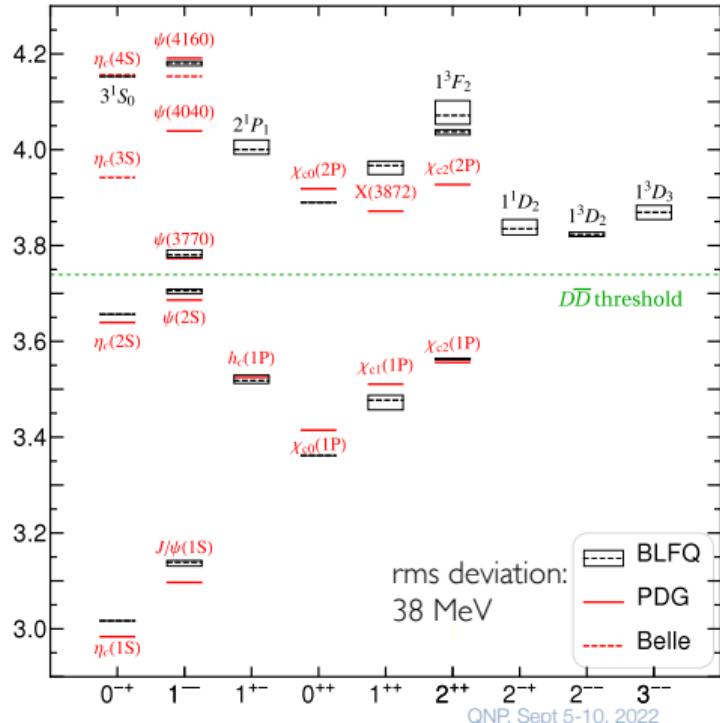
(d)  $\psi_{\uparrow\downarrow}(\vec{k}_\perp, x) = \psi_{\uparrow\uparrow}^*(\vec{k}_\perp, x)$

# Basis light-front quantization

[Vary et al. PRC '09; YL, Maris, Zhao, Vary, PLB '16]



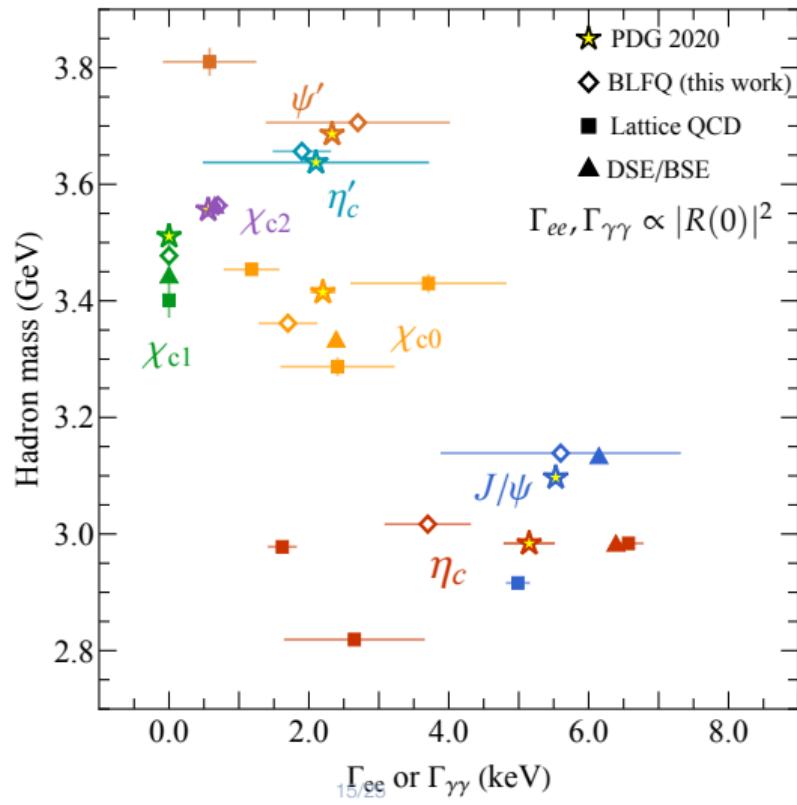
- ▶ Holographic light-front QCD confinement plus one-gluon-exchange interaction [Review: Brodsky '14]
- ▶ Solved in basis function approach,  $\Lambda_{\text{UV}} \approx b\sqrt{N_{\text{max}}}$ .
- ▶ Two free parameters  $m_c, \kappa$  fitted to the mass spectrum. Posterior rms deviation:  $\lesssim 40$  MeV
- ▶ Application to a variety of systems:
  - ▶  $c\bar{c}, b\bar{b}$ : YL, PLB '16 & PRD '17
  - ▶  $b\bar{c}, b\bar{q}, c\bar{q}$ : Tang, PRD '18 & EPJC '20
  - ▶  $q\bar{q}$ : Jia, PRC '19; Qian, PRC '20; YL, '21
  - ▶ Baryons: Mondal, PRD '20; Xu, '21
- ▶ Access to a variety of observables
  - ▶ Form factors: YL, PRD '18; Mondal, PRD '20
  - ▶ (Semi-)leptonic decay: Li, PRD '18 & '19; Tang, PRD '21
  - ▶ PDFs/GPDs: Lan, PRL '19 & PRD '20; Adhikari, PRC '18 & '21
  - ▶ Diffractive production: Chen, PLB '17 & PRC '18



# Numerical results: overview

Parameter-free predictions:

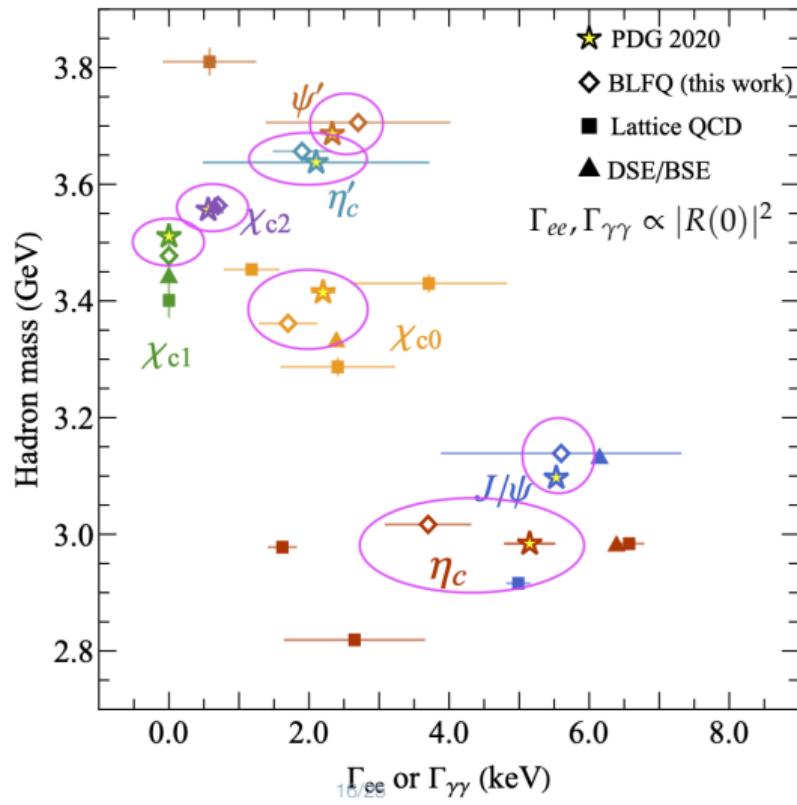
[Lattice: Dudek '06, Chen '16, Chen '20, Meng '21, Zou '21; DSE: Chen '16]



# Numerical results: overview

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[Lattice: Dudek '06, Chen '16, Chen '20, Meng '21, Zou '21; DSE: Chen '16]



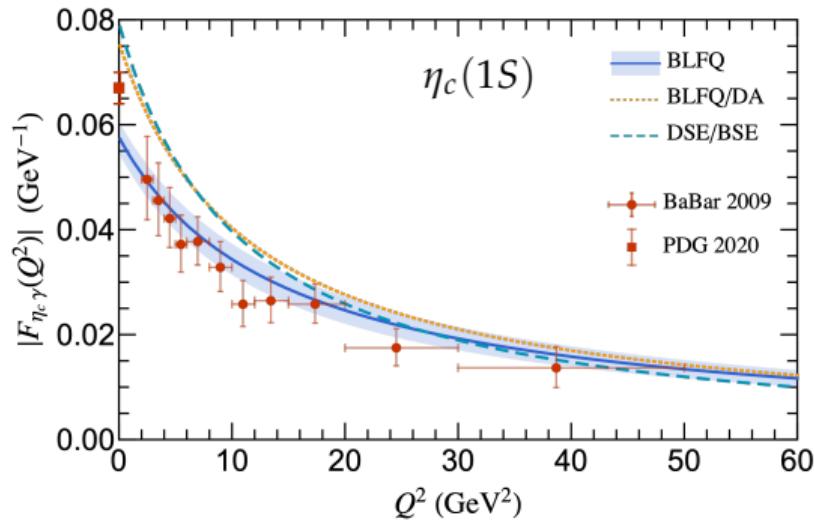
# Transition form factor: $\eta_c$

$$\mathcal{M}^{\mu\nu} = 4\pi\alpha_{\text{em}}\epsilon^{\mu\nu\rho\sigma}q_{1\rho}q_{2\sigma}F_{P\gamma\gamma}(q_1^2, q_2^2), \quad F_{P\gamma}(Q^2) \equiv F_{P\gamma\gamma}(Q^2, 0) = F_{P\gamma\gamma}(0, Q^2)$$

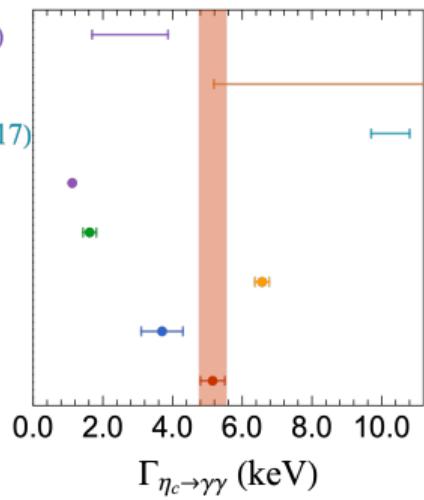
Diphoton width:  $\Gamma_{\gamma\gamma} = \frac{\pi}{4}\alpha_{\text{em}}^2 M_P^3 |F_{P\gamma\gamma}(0, 0)|^2$ .

[Lepage '81, Babiarz '19, Hoferichter '20]

$$F_{P\gamma}(Q^2) = e_f^2 2\sqrt{2N_C} \int \frac{dx}{2\sqrt{x(1-x)}} \int \frac{d^2k_\perp}{(2\pi)^3} \frac{\psi_{\uparrow\downarrow-\downarrow\uparrow/P}^*(x, \vec{k}_\perp)}{k_\perp^2 + m_f^2 + x(1-x)Q^2}$$



NRQM/LF (Babiarz 2019)  
 NRQM (Babiarz 2019)  
 NNLO NRQCD (Feng 2017)  
 Lattice (Chen 2016)  
 Lattice (Chen 2020)  
 Lattice (Meng 2021)  
 BLFQ (this work)  
 PDG 2020

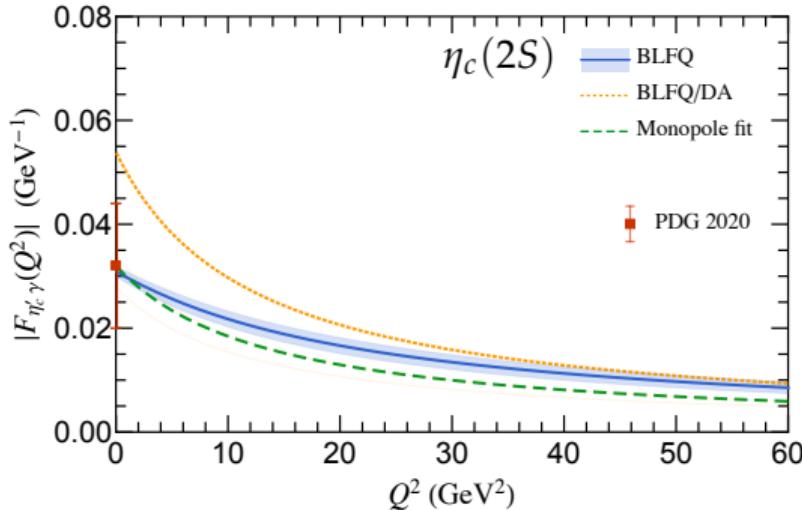


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- ▶ No experimental measurement yet.
- ▶ A monopole fit using  $\Lambda^2 = M_{\psi'}^2$  is included for comparison.
- ▶ Note that a VMD prediction requires the off-shell coupling  $g_V(Q^2) = V_{PV\gamma}(Q^2)$ : [Lakhina '06]

$$F_{P\gamma}^{(\text{VMD})}(Q^2) = \sum_V \frac{e_f^2 f_V}{1 + \frac{M_p}{M_V}} \left[ \frac{g_V(0)}{M_V^2 + Q^2} + \frac{g_V(Q^2)}{M_V^2} \right]$$

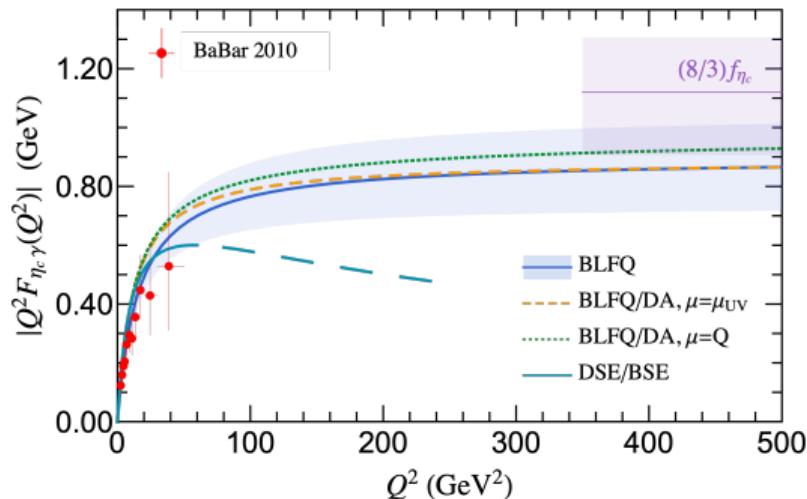
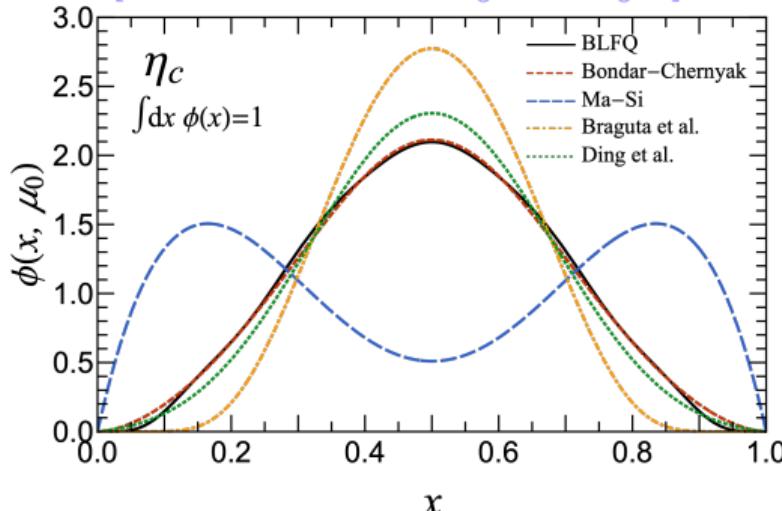
# Light cone distribution amplitude (LCDA)

At large- $Q^2$ , viz.  $Q^2 + \langle m_f^2/x(1-x) \rangle \gg \langle k_\perp^2/x(1-x) \rangle$ ,

$$F_{P\gamma}(Q^2) \approx e_f^2 f_P \int_0^1 dx \frac{\phi_P(x, \mu)}{x(1-x)Q^2 + m_f^2} \xrightarrow{Q \rightarrow \infty} \frac{6e_f^2 f_P}{Q^2}.$$

- ▶ LCDA plays a pivotal role in hard exclusive charmonium production. [See, e.g., Braguta '12]
- ▶ Our LCDA agrees with the Bondar-Chernyak model. Both fit the BABAR **normalized** TFF well.

[LCDAs: Ma '04, Bondar '05, Braguta '07, Ding '16]



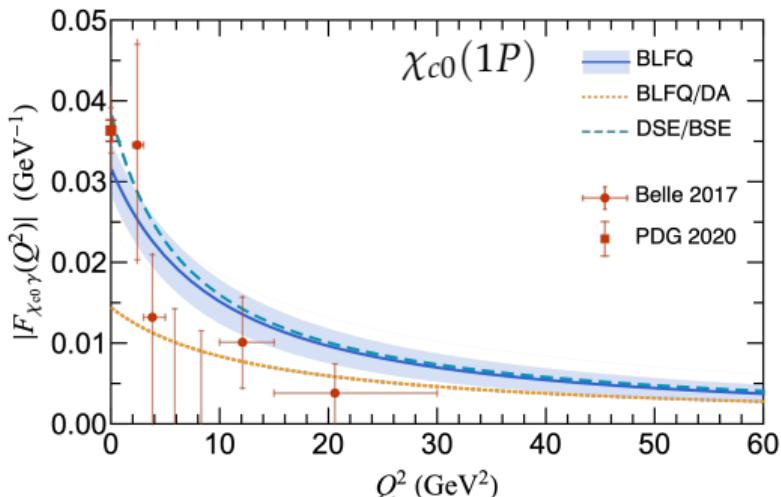
# Transition form factor: $\chi_{c0}$

[Babiarz '20, Hoferichter '20; DSE/BSE: Chen '17]

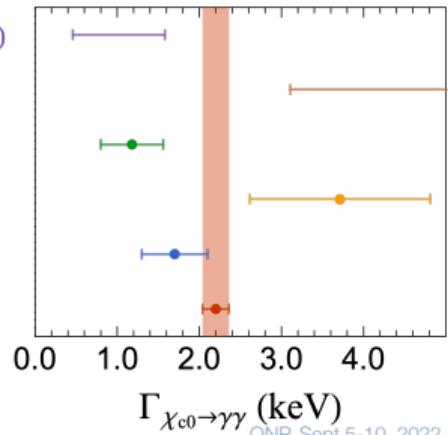
$$\mathcal{M}_{S \rightarrow \gamma\gamma}^{\mu\nu} = 4\pi\alpha_{\text{em}} \left\{ [(q_1 \cdot q_2)g^{\mu\nu} - q_2^\mu q_1^\nu] F_1^S(q_1^2, q_2^2) + \frac{1}{M_S^2} [q_1^2 q_2^2 g^{\mu\nu} + (q_1 \cdot q_2) q_1^\mu q_2^\nu - q_1^2 q_2^\mu q_2^\nu - q_2^2 q_1^\mu q_1^\nu] F_2^S(q_1^2, q_2^2) \right\}$$

Single-tag TFF:  $F_{S\gamma}(q^2) = F_1^S(q^2, 0) = F_1^S(0, q^2)$ . Width  $\Gamma_{\gamma\gamma} = \frac{\pi\alpha_{\text{em}}^2}{4} M_S^3 |F_{S\gamma}(0)|^2$ . Belle provides the first measurement of the TFF, albeit with limited statistics.

[Belle, PRD 2017]



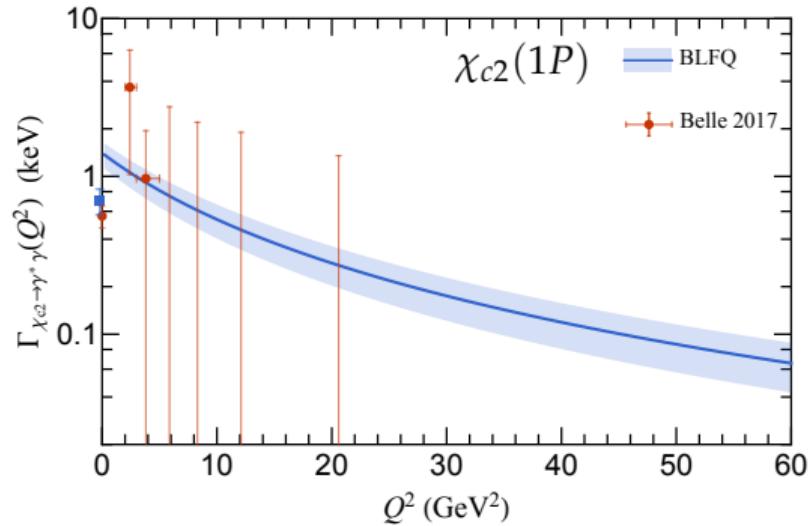
NRQM/LF (Babiarz 2019)  
 NRQM (Babiarz 2019)  
 Lattice (Chen 2020)  
 Lattice (Zou 2021)  
 BLFQ (this work)  
 PDG 2020



$$\Gamma_{T \rightarrow \gamma\gamma} = \frac{\pi \alpha_{\text{em}}^2}{5M_T} \left( |\mathcal{M}_{++;0}|^2 + |\mathcal{M}_{+-;2}|^2 \right).$$

Belle provides the first measurement of the  $Q^2$  dependent width, albeit with limited statistics.

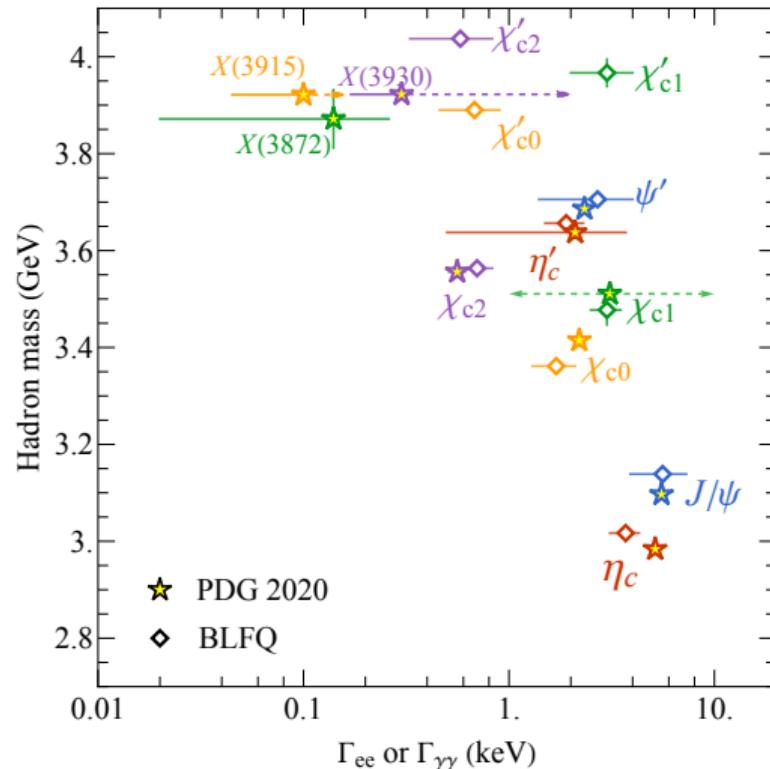
[Belle, PRD 2017]



# The $c\bar{c}$ contents of the X's

Belle recently measured the reduced diphoton width of  $X(3872)$ . The value is much smaller than our prediction of a pure  $c\bar{c} \chi_{c1}(2P)$ , indicating large portion of non- $c\bar{c}$  content (model dependence?)

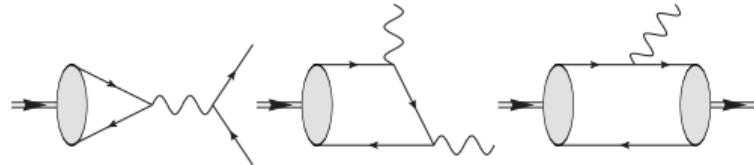
[Belle, PRL 2021]



# Radiative transitions

Leptonic and radiative transitions probe the fundamental structure of the hadrons:

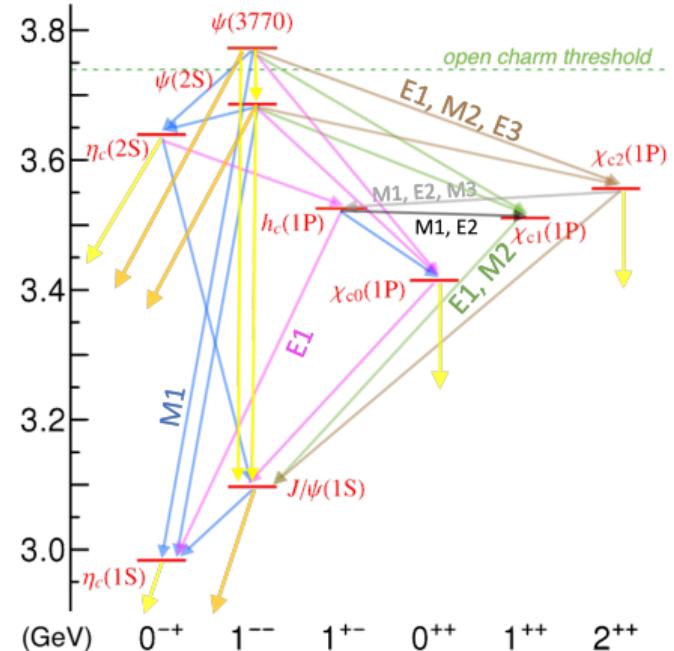
[Review: Barnes & Yuan, Int. J. Mod. Phys. A 2009]



	decay width (keV)	$\Gamma_{ee}$	$\Gamma_{\gamma\gamma}$	
$\eta_c$	PDG	-	5.15(35)	
	BLFQ	-	3.7(6)	$\Gamma_{\eta_c\gamma}$
$J/\psi$	PDG	5.53(10)	-	1.6(4)
	BLFQ	5.7(1.9)	-	2.6(1) $\Gamma_{J/\psi\gamma}$
$\chi_{c0}$	PDG	-	2.1(1.6)	-
	BLFQ	-	1.9(4)	- $15(1) \times 10^3$
$\chi_{c1}$	PDG	-	-	-
	BLFQ	-	-	- 288(16)
⋮				

[YL, PRD '17; Li, PRD '18; Chen, in progress]

[PDG, PTEP '20 + '21 (update)]

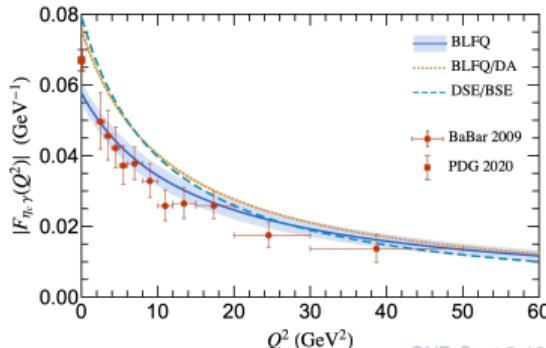
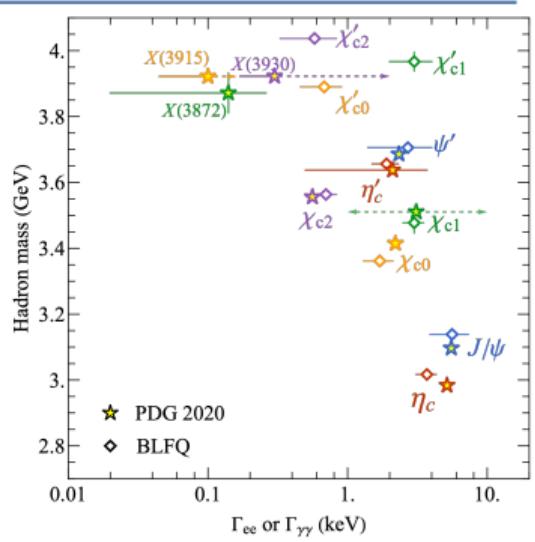


# Summary

- ▶ Light-front Hamiltonian formalism provides unique tools to access the hadronic observables
  - ▶ Light-cone dominance
  - ▶ Collinear factorization and  $k_T$  factorization
- ▶ We computed the two-photon width and transition form factors of charmonia within the basis light-front quantization approach.
  - ▶ Excellent agreements with the available experimental measurements.
  - ▶ No parameters are dialed to obtain these results.
  - ▶ Reveal relativistic nature of charmonium system
- ▶ The obtained wave functions await further experimental measurements and further applications.

Based on: YL, M. Li (李枚键) and J.P. Vary, Phys. Rev. D **L105** (2022); arXiv:2111.14178 [hep-ph]

LFWFs available on Mendeley Data, doi: 10.17632/cjs4ykv8cv.2



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Thank you for your attention.

