

Implementation of DVMP in PARTONS

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Chiral-even GPDs

- Chiral-even GPDs parametrize the following off-forward matrix elements of quark operators at a light-like separation

D.Müller, D.Robaschik, B.Geyer, F.-M.Dittes, J.Hořejši, Fortsch. Phys. **42**, 101 (1994)

X.D.Ji, PRL **78**, 610 (1997); PRD **55**, 7114 (1997).

A.V.Radyushkin, PLB **380**, 417 & **385**, 333 (1996).

$$\begin{aligned} P^+ \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle p', \lambda' | \bar{\psi}(-\frac{z}{2}) \gamma^+ \psi(\frac{z}{2}) | p, \lambda \rangle \Big|_{z^+=0, \vec{z}_T=0} \\ = \bar{u}(p', \lambda') \left[H^q \gamma^+ + E^q \frac{i\sigma^{+\alpha} \Delta_\alpha}{2m} \right] u(p, \lambda) \end{aligned}$$

$$\begin{aligned} P^+ \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle p', \lambda' | \bar{\psi}(-\frac{z}{2}) \gamma^+ \gamma_5 \psi(\frac{z}{2}) | p, \lambda \rangle \Big|_{z^+=0, \vec{z}_T=0} \\ = \bar{u}(p', \lambda') \left[\tilde{H}^q \gamma^+ \gamma_5 + \tilde{E}^q \frac{\gamma_5 \Delta^+}{2m} \right] u(p, \lambda) \end{aligned}$$

- GPDs depend on three parameters

$$x, \quad \xi = \frac{p^+ - p'^+}{p^+ + p'^+}, \quad t = \Delta^2$$

where $\Delta = p' - p$ and $P^+ = (p' + p)^+ / 2$.

Chiral-even GPDs

- Chiral-even GPDs are accessible through Deeply Virtual Compton Scattering

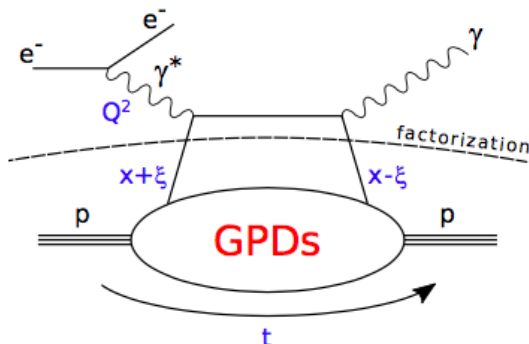


Fig: Eur. Phys. J. C78 (2018) 6, 478

- There are four chiral-odd GPDs $H_T, \tilde{H}_T, E_T, \tilde{E}_T$ at leading twist

M. Diehl, Eur. Phys. J. C **19**, 485 (2001)

$$\begin{aligned} & \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle p', \lambda' | \bar{\psi}(-\tfrac{1}{2}z) i\sigma^{+i} \psi(\tfrac{1}{2}z) | p, \lambda \rangle \Big|_{z^+=0, \mathbf{z}_T=0} \\ &= \frac{1}{2P^+} \bar{u}(p', \lambda') \left[H_T^q i\sigma^{+i} + \tilde{H}_T^q \frac{P^+ \Delta^i - \Delta^+ P^i}{m^2} \right. \\ & \quad \left. + E_T^q \frac{\gamma^+ \Delta^i - \Delta^+ \gamma^i}{2m} + \tilde{E}_T^q \frac{\gamma^+ P^i - P^+ \gamma^i}{m} \right] u(p, \lambda). \end{aligned}$$

where $i = 1, 2$ is the transversity index

Chiral-odd GPDs

- Accessible, if one assumes an effective handbag mechanism, through exclusive meson production processes

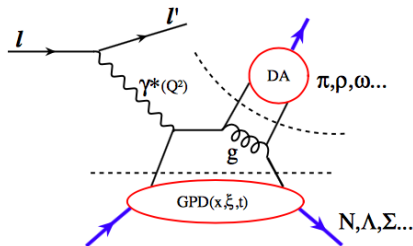


Fig: EPJ Web Conf. **137**, 01003 (2017)

Properties of GPDs

- In the forward limit $\Delta \rightarrow 0$, certain GPDs are related to PDFs

$$H^q(x, 0, 0) = f_1^q(x)$$

$$\tilde{H}^q(x, 0, 0) = g_1^q(x)$$

$$H_T^q(x, 0, 0) = h_1^q(x).$$

- Related to form factors via

$$\int_{-1}^1 \left\{ H, E, \tilde{H}, \tilde{E} \right\} (x, \xi, t) dx = F_1(t), F_2(t), G_A(t), G_P(t)$$

$$\int_{-1}^1 \left\{ H_T, \tilde{H}_T, E_T \right\} (x, \xi, t) dx = H_T(t), \tilde{H}_T(t), E_T(t)$$

$$\int_{-1}^1 \left\{ \tilde{E}_T \right\} (x, \xi, t) dx = 0$$

- Polynomiality

Goloskokov-Kroll Model

- Goloskokov-Kroll(GK) model for pseudoscalar meson production is based on an effective handbag mechanism

S. V. Goloskokov, P. Kroll, Eur. Phys. J. C **65**, 137 (2010) & Eur. Phys. J. A **47**, 112 (2011)

- For π^0 production, contributions from longitudinally polarized photons read

$$\mathcal{M}_{0+,0+} = \sqrt{1-\xi^2} \frac{e}{Q} [\langle \tilde{H} \rangle - \frac{\xi^2}{1-\xi^2} \langle \tilde{E} \rangle]$$

$$\mathcal{M}_{0-,0+} = \frac{e}{Q} \frac{\sqrt{-t'}}{2m} [\xi \langle \tilde{E} \rangle]$$

where $t' = t - t_0$ with $t_0 = -4M_N^2 \xi^2 / (1 - \xi^2)$.

first index = helicity of the produced meson,
second index = helicity of the final particle,
third index = helicity of the virtual photon,
fourth index = helicity of the initial particle.

Goloskokov-Kroll Model

- Generically, $\langle F \rangle$ represents a convolution of a GPD F with an appropriate subprocess amplitude

$$\langle F \rangle = \sum_{\lambda} \int_{-1}^1 dx \mathcal{H}_{0\lambda,0\lambda}(x, \xi, Q^2, t=0) F(x, \xi, t)$$

where λ denotes unobserved helicities of the partons.

- Subprocesses are calculated in the so-called modified perturbative approach: Transverse momenta of the quark and the anti-quark in the meson are kept and gluon radiations are taken into account through Sudakov factor S
- For π^0 electroproduction, GPDs appear in the form of

$$(e_u F^u - e_d F^d)/\sqrt{2}$$

- In impact space

$$\mathcal{H}_\pi = \int d\tau d^2\vec{b} \hat{\Psi}_\pi(\tau, -\vec{b}) \hat{\mathcal{F}}_\pi(\bar{x}, \xi, \tau, Q^2, \vec{b}) \alpha_s(\mu_R) \exp(-S(\tau, \vec{b}, Q^2))$$

where $\hat{\Psi}_\pi$ is the pion wave function and $\hat{\mathcal{F}}_\pi$ is the hard scattering kernel

- Contributions from transversely polarized photons are significant: Twist-3 effects
- For small $-t'$ and ξ , only H_T and $\bar{E}_T$ enter in twist-3 contributions

$$\mathcal{M}_{0-,++} = \sqrt{1-\xi^2} e \langle H_T \rangle$$

$$\mathcal{M}_{0+,\mu+} = -\frac{e}{4m} \sqrt{-t'} \langle \bar{E}_T \rangle$$

- Partial cross sections

$$\begin{aligned}\frac{d\sigma_L}{dt} &= \kappa[\mathcal{M}_{0+,0+}^2 + \mathcal{M}_{0-,0+}^2], \\ \frac{d\sigma_T}{dt} &= \frac{\kappa}{2}[\mathcal{M}_{0-,-+}^2 + \mathcal{M}_{0-,++}^2 + \mathcal{M}_{0+,-+}^2 + \mathcal{M}_{0+,++}^2], \\ \frac{d\sigma_{TT}}{dt} &= -\frac{\kappa}{2}\text{Re}[\mathcal{M}_{0-,++}^* \cdot \mathcal{M}_{0-,-+} + \mathcal{M}_{0+,++}^* \cdot \mathcal{M}_{0+,-+}], \\ \frac{d\sigma_{LT}}{dt} &= \frac{\kappa}{\sqrt{2}}\text{Re}[\mathcal{M}_{0-,0+}^* \cdot (\mathcal{M}_{0-,++} - \mathcal{M}_{0-,-+}) \\ &\quad + \mathcal{M}_{0+,0+}^* \cdot (\mathcal{M}_{0+,++} - \mathcal{M}_{0+,-+})].\end{aligned}$$

- GPDs are constructed from double distribution ansatz

$$F_i^a(\bar{x}, \xi, t) = \int_{-1}^1 d\rho \int_{-1+|\rho|}^{1-|\rho|} d\eta \delta(\rho + \xi\eta - \bar{x}) f_i^a(\rho, \eta, t)$$

where for valence-quark GPDs;

$$f_i(\rho, \eta, t) = \exp[(b_i - \alpha'_i \ln \rho)t] F_i^a(\rho, \xi = t = 0) \frac{3}{4} \frac{(1 - \rho)^2 - \eta^2}{(1 - \rho)^3} \Theta(\rho)$$

- Parameters of the forward limits

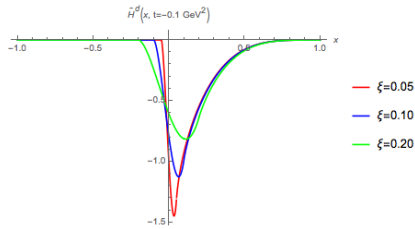
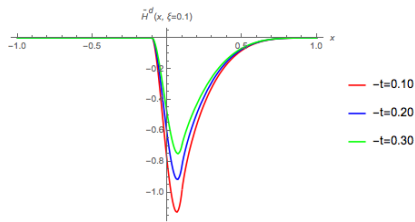
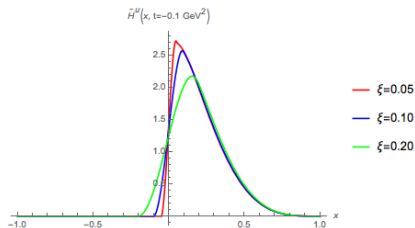
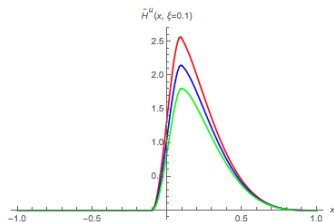
$\tilde{H}$: DSSV, [Phys. Rev. D **80**, 034030 \(2009\)](#)

H_T : ABM, [Phys. Rev. D **86**, 054009 \(2012\)](#) and DSSV, [Phys. Rev. D **80**, 034030 \(2009\)](#)

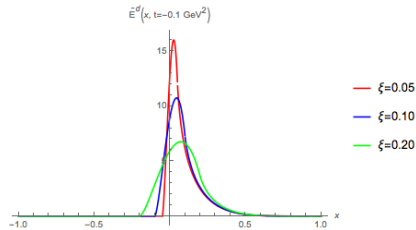
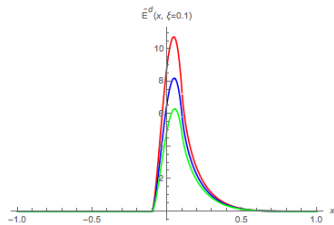
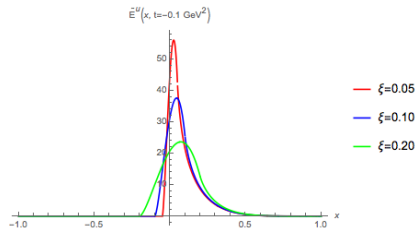
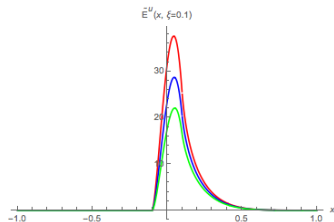
$\tilde{E}$: LHPC Collaboration, [Phys. Rev. D **77**, 094502 \(2008\)](#)

$\bar{E}_T$: QCDSF and UKQCD Collaborations, [Phys. Rev. Lett. **98**, 222001 \(2007\)](#)

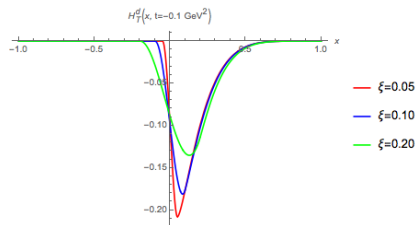
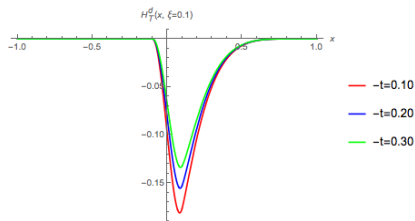
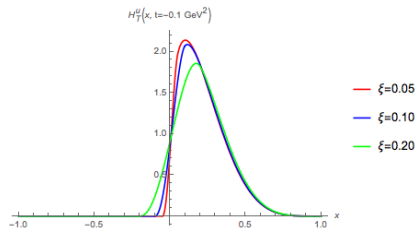
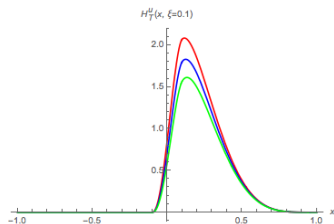
Goloskokov-Kroll Model - GPDs



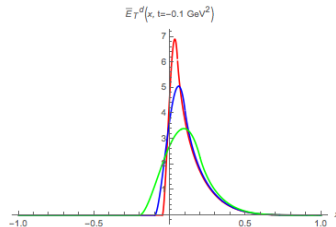
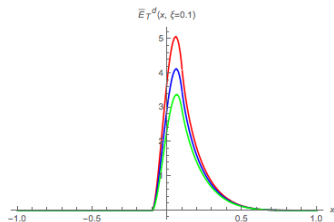
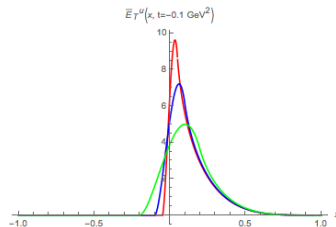
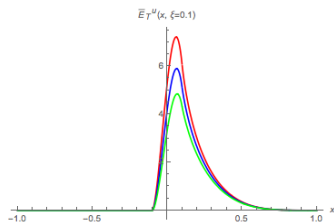
Goloskokov-Kroll Model - GPDs



Goloskokov-Kroll Model - GPDs



Goloskokov-Kroll Model - GPDs



- In π^0 electroproduction, $\bar{E}_T$ dominates the process since GPDs appear in the form of $(e_u \bar{E}_T^u - e_d \bar{E}_T^d)/\sqrt{2}$
- To compute longitudinally polarized photon amplitudes, we need: twist-2 meson wave function, kernel, running coupling, Sudakov factor and GPDs $\tilde{H}$ and $\tilde{E}$.
- To compute transversely polarized photon amplitudes, we need: twist-3 meson wave function, kernel, running coupling, Sudakov factor and GPDs H_T and $\bar{E}_T$.
- 3-dimensional integrals over $\bar{x}, \tau$ and b are performed in impact space
- In many different processes, such as vector meson or other pseudoscalar meson production, meson wavefunction as well as the Sudakov factor has the same structure

Goloskokov-Kroll Model - π^0

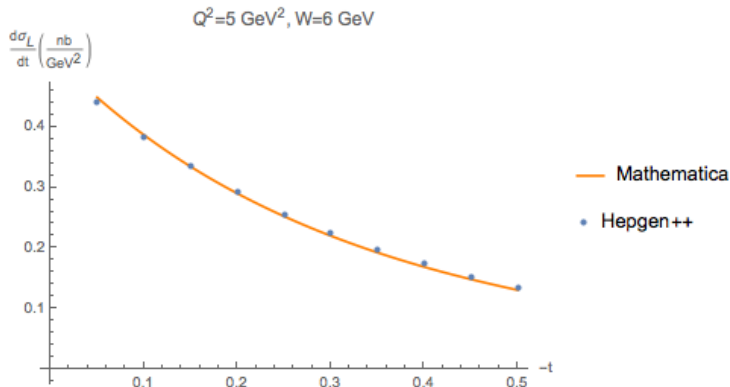
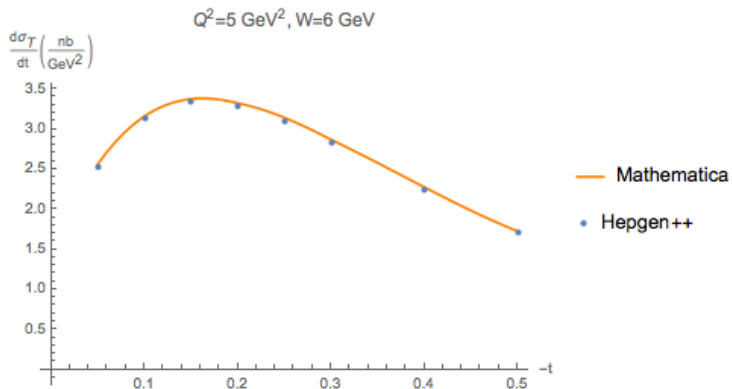
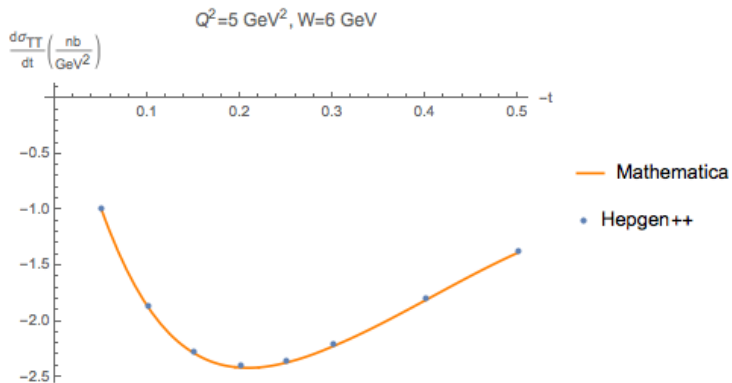


Figure: Comparison between our Mathematica results and *Hepgen* ++ (see C. Regali's PhD thesis, <https://freidok.uni-freiburg.de/data/11449>)

Goloskokov-Kroll Model - π^0



Goloskokov-Kroll Model - π^0



- Contrary to π^0 production, GPDs appear in the form of $(F^u - F^d)$ in π^+ production
- H_T dominates the process due to the relative sign of up and down quark GPDs
- One also needs to consider pion-pole terms in π^+ production



B. Berthou et al., Eur.Phys.J. C78 (2018) 6, 478

- PARtomic Tomography Of Nucleon Software (PARTONS) is a software framework dedicated to the phenomenology of GPDs
- Website: <http://partons.cea.fr>

- Bridge between models of GPDs and experimental data
- Written in C++: Inheritance and polymorphism
- XML interface for automated tasks
- Layered structure:
 - GPD Layer
 - CFF Layer
 - Process Layer
 - Observable Layer

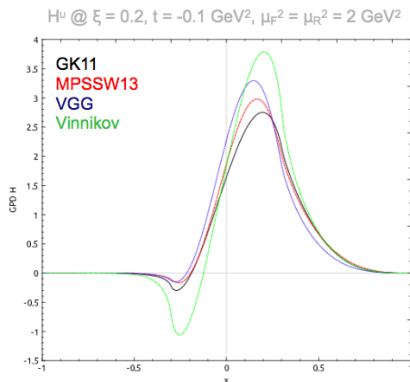
- Existing modules:

GPDs: GK, VGG, Vinnikov, MPSSW, MMS, HM

Evolution: Vinnikov's code

CFF (for DVCS): LO and NLO evaluation

Cross Section (for DVCS): VGG, BMJ, GV



- Recent applications with PARTONS:

"Border and skewness functions from a leading order fit to DVCS data",

H. Moutarde, P. Sznajder, J. Wagner, *Eur. Phys. J. C* **78**, no. 11, 890 (2018)

"Unbiased determination of Compton Form Factors",

H. Moutarde, P. Sznajder, J. Wagner, *Eur. Phys. J. C* **79**, no. 7, 614 (2019)

- Adding new modules as easy as possible (see the website for tutorials)

- PARTONS: A comprehensive framework dedicated to the phenomenology of GPDs
- DVCS process is already implemented in PARTONS
- Easy addition of other partonic processes
- Implementation of DVMP processes is in progress
- First model to be implemented: Goloskokov-Kroll model of π^0 and π^+ production