**Strong QCD from Hadron Structure Experiments 2019** 

November 05 - 09, 2019, Newport News, VA

Experimental access to baryon-to-meson transition distribution amplitudes: A new window to the 3 dimensional nucleon structure

## JUSTUS-LIEBIG-UNIVERSITÄT GIESSEN



# Stefan Diehl

Justus Liebig University Giessen University of Connecticut

# Outline

- Introduction to TDAs and comparison to GPDs
- Theoretical modelling and features of the TDA based desciption
- Experimental access to TDAs at JLAB (CLAS and HALL C)
- Perspectives for the experimental access to TDAs with PANDA at FAIR and at JPARC
- Summary and Outlook

Theory collaborators: B. Pire, L. Szymanowski, K. Semenov-Tian-Shansky



## **Physics motivation**

# GPDs

- → Light-cone matrix elements of non-local bilinear quark and gluon operators
- ➔ Describe hadronic structural information in terms of quark and gluon degrees of freedom
- Spin-dependent 2D transverse coordinate space
   + 1D longitudinal momentum space images of the nucleon
  - ➔ Tool to study the nature and origin of the nucleon spin
  - ➔ Impact parameter space: spatial femto-photographs of the hadron structure in the transverse plane



of Z

## **Physics motivation**

## **Baryon to meson TDAs**

- → Light-cone matrix elements of non-local three quark operators
- → Encoded physical picture close to GPDs
- Probe partonic correlations between states of different baryonic charge
  - Access to non-minimal Fock components of baryon light-cone wave functions



- → <u>Impact parameter space</u>: Femto-photography of hadrons from a new perspective
  - ➔ Spatial imaging of the structure of the pion cloud inside the nucleon



## **Theoretical modelling of TDAs**

**GPDs:** Description in the skewness  $\xi = 0$  limit

- **TDAs:** The soft pion theorems fixe the  $\xi \rightarrow 1$  limit in terms of nucleon DAs and thus provides the overall magnitude of TDAs
  - → A factorized Ansatz with input at  $\xi$  = 1 designed in

J.P. Lansberg, B. Pire, K. Semenov and L. Szymanowski (2012)

➔ Provides the unpolarized cross section for hard leptoproduction of a pion off nucleon:

$$\frac{d^{5}\sigma}{dE'd\Omega_{e'}d\Omega_{\pi}} = \Gamma \times \frac{\Lambda\left(s, m^{2}, M^{2}\right)}{128\pi^{2}s\left(s - M^{2}\right)} \times \sum_{s_{1}, s_{2}} \left\{ \frac{1}{2} \left( \left| \mathcal{M}_{s_{1}s_{2}}^{1} \right|^{2} + \left| \mathcal{M}_{s_{1}s_{2}}^{-1} \right|^{2} \right) + \ldots \right\} = \Gamma \times \left( \frac{d^{2}\sigma_{T}}{d\Omega_{\pi}} + \ldots \right)$$

$$N\gamma^{*} \to \pi N \text{ helicity amplitudes}$$

### → Polartized terms:

Need dominant leading twist transverse amplitude

+ next-to-leading twist subdominant longitudinal amplitude (twist-4 nucleon DAs or twist-4 nucleon-to-pion TDAs)

## **Experimental acessible features of the TDA based mechanism**

### **Characteristic features of the TDA-based mechanism:**

- → Dominance of the transverse polarisation of the virtual photon leads to a suppression of the longitudinal cross section σ<sub>L</sub> at large Q<sup>2</sup> by at least a factor 1/Q<sup>2</sup>
- The transverse cross section σ<sub>T</sub> shows a charakteristic 1/Q<sup>8</sup> scaling behaviour for fixed x<sub>B</sub>

### More distinguishing features become accesible with a polarized target:

- $\rightarrow$  Transverse target single spin asymmetry ~ *Im* part of the amplitude
- → TDA approach predicts a non vanishing and Q<sup>2</sup>-independent TSA
- → Two component TDA model predicts 10-15 % TSA for  $\gamma^* N \rightarrow \pi N$

## Backward $\pi$ electroproduction with CLAS at JLAB





- CLAS (e16 + e1f run period)
- 5.5, 5.75 GeV longitudinally polarized electron beam
- unpolarized hydrogen target
- Electron ID based on electromagnetic calorimeter and Cherenkov counters
- **π ID** based on a maximum likelyhood
   particle selection from TOF based
   β vs p correlation

# Hard exclusive $\pi^+$ electroproduction cross section in the backward direction



#### K. Park, M. Guidal et al. PL B 780, 340 (2018)

Variable	Number of bins	Range	Bin size
W	1	2.0-2.4 GeV	400 MeV
Q <sup>2</sup>	6	$1.6-4.5 \text{ GeV}^2$	varying
$\Delta_T^2$	1	0–0.5 GeV <sup>2</sup>	$0.5 \text{ GeV}^2$
$arphi^*_\pi$	9	$0^{\circ}$ –360 $^{\circ}$	$40^{\circ}$



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## Hard exclusive $\pi^+$ electroproduction



K. Park, M. Guidal et al. PL B 780, 340 (2019)

$$\frac{d\sigma}{d\Omega_{\pi}^{*}} = A + B\cos\varphi_{\pi}^{*} + C\cos 2\varphi_{\pi}^{*}$$
$$A = \sigma_{T} + \epsilon\sigma_{L} \quad B = \sqrt{2\epsilon(1+\epsilon)}\sigma_{LT}$$
$$C = \epsilon\sigma_{TT}$$

#### TDA model calculations for $\sigma_{U}$ :

- → Results depend on the input for the nucleon distribution aplitude
- dark blue (COZ): V.L. Chernyak et al., Z. Phys. C 42, 583 (1989)
- light blue (KS): I.D. King, C.T. Sachrajda, Nucl. Phys. B 279, 785 (1987)
- red: NNLO calculation: A. Lenz et al., Phys. Rev. D 79, 093007 (2009)

Hard exclusive  $\pi^+$  electroproduction beam spin asymmetry



**<u>Cross section</u>** (longitudinally pol. beam and unpol. target):

$$d\sigma = d\sigma_0 (1 + A_{UU}^{\cos(2\phi)} \cos(2\phi) + A_{UU}^{\cos(\phi)} \cos(\phi) + h A_{LU}^{\sin(\phi)} \sin(\phi))$$

$$BSA = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-} = \frac{A_{LU}^{\sin\phi} \sin\phi}{1 + A_{UU}^{\cos\phi} \cos\phi + A_{UU}^{\cos(2\phi)} \cos(2\phi)}$$

## Kinematic coverage and exclusivity cuts



## Beam spin asymmetry

$$BSA_i = \frac{1}{P_e} \cdot \frac{N_i^+ - N_i^-}{N_i^+ + N_i^-} \qquad \begin{array}{c} \mathsf{P}_e = \mathsf{75~\%: average e} \text{-} \text{ beam} \\ \text{polarisation} \end{array}$$

### Integrated over all kinematic variables in forward / backward region:



## **BSA for different -t bins**



χ² / ndf

 $\chi^2$  / ndf

p0

p0

df 15.57 / 11 0.09353 ± 0.005818

¢ [°]

ື (°]

df 3.817 / 11 -0.03933 ± 0.01594

# -t dependence of $A_{LU}^{\sin(\phi)}$



# -t dependence of $A_{LU}^{\sin(\phi)}$



# $\mathbf{Q}^{2}$ and $\mathbf{x}_{\mathbf{B}}$ dependence of $A_{LU}^{\sin(\phi)}$



BSA is a subleading twist effect both in the forward and backward regimes

## **Perspectives for CLAS12**

• CLAS 12 can map out froward and backward kinematics for

**RG-A:**  $ep \rightarrow e' p \pi^0$   $ep \rightarrow e' n \pi^+$   $ep \rightarrow e' p \omega$ **RG-B:**  $en \rightarrow e' p \pi^-$ 

➔ A significant amount of data has already been recorded and is currently in the analysis phase

**Example:** Kinematic coverage of  $ep \rightarrow e' \pi^+(n)$ 



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## Backward $\boldsymbol{\omega}$ production at JLAB Hall C

**Analysis by:** W. B. Li, G. Huber *et al.* (Jefferson Lab  $F_{\pi}$  collaboration) Phys. Rev. Lett. 123, 182501 (2019)

JLAB Hall C: 2.6 - 5.2 GeV electron beam on a liquid hydrogen target

 Recoil protons and scattered electrons detected with the hall C high precision particle spectrometers



## Backward $\omega$ production at JLAB Hall C

$$2\pi \frac{d^2\sigma}{dtd\phi} = \frac{d\sigma_{\rm T}}{dt} + \epsilon \frac{d\sigma_{\rm L}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{\rm LT}}{dt} \cos\phi + \epsilon \frac{d\sigma_{\rm TT}}{dt} \cos 2\phi$$

→ Full Rosenbluth separation to extract  $\sigma_{T}$  and  $\sigma_{L}$ 



## Backward $\omega$ production at JLAB Hall C

### Combined (CLAS and $F_{\pi-2}$ data for $\gamma^* p \rightarrow \omega p$ )

TDA-based predictions vs the Regge-based J.M. Laget's (JML18) model





- Antiproton beam
- cluster jet / pellet target
   (pp, pA)
- $E_{ar{p}} \leq 15 \,\, {
  m GeV}$  $W^2 \leq 30 \,\, {
  m GeV}^2$

 $L = 2 \cdot 10^{32} \text{ cm}^{-2} \text{s}^{-1}$ 



TDAs also occur in factorized description of:

 $W^2$ 

 $p_{\bar{N}}$ 

 $p_{\mathcal{M}}$ 

$$\bar{N} + N \rightarrow \gamma^{*}(q) + \pi \rightarrow \ell^{+} + \ell^{-} + \pi;$$

$$\bar{N} + N \rightarrow J/\psi + \pi \rightarrow \ell^{+} + \ell^{-} + \pi;$$
Theory: J.P. Lansberg et al. (2012)  
B. Pire, L. Szymanowski,  
K. Semenov-Tian-Shansky (2013)  
**Two regimes:** forward and backward  

$$\int_{P(q)}^{P(q)} \int_{P(q)}^{P(q)} \int_{P(q)}^{P(q)}$$



 $p_N$ 

CF

MNTD

21

 $W^2$ 

CF'

MNTDA

 $p_{\mathcal{M}}$ 

 $\bar{p}p \rightarrow \pi^0 \gamma^* \rightarrow \pi^0 \ell^+ \ell^-$  cross section estimates for **PANDA** 

Integrated cross section for different inputs of the nucleon DAs:



B. Pire, L. Szymanowski, K. Semenov-Tian-Shansky (2013)

• Cross section of  $\overline{p}$  n  $\rightarrow \pi^{-}$  l<sup>+</sup> l<sup>-</sup> is larger by a factor 2 (neutron target required)

### Several feasability studies:

M. C. Mora Espi, M. Zambrana, F. Maas (PANDA collaboration) (2015) B. Ramstein, E. Atomssa (PANDA collaboration) PRD 95 (2017)

#### Feasibility study for the measurement of $\pi N$ TDAs at $\overline{P}ANDA$ in $\bar{p}p \rightarrow J/\psi\pi^0$

B. Singh,<sup>1</sup> W. Erni,<sup>2</sup> B. Krusche,<sup>2</sup> M. Steinacher,<sup>2</sup> N. Walford,<sup>2</sup> H. Liu,<sup>3</sup> Z. Liu,<sup>3</sup> B. Liu,<sup>3</sup> X. Shen,<sup>3</sup> C. Wang,<sup>3</sup> J. Zhao,<sup>3</sup> M. Albrecht,<sup>4</sup> T. Erlen,<sup>4</sup> M. Fink,<sup>4</sup> F.H. Heinsius,<sup>4</sup> T. Held,<sup>4</sup> T. Holtmann,<sup>4</sup> S. Jasper,<sup>4</sup> I. Keshk,<sup>4</sup> H. Koch,<sup>4</sup> B. Kopf,<sup>4</sup> M. Kuhlmann,<sup>4</sup> M. Kümmel,<sup>4</sup> S. Leiber,<sup>4</sup> M. Mikirtychyants,<sup>4</sup> P. Musiol,<sup>4</sup> A. Mustafa,<sup>4</sup> M. Pelizäus,<sup>4</sup> J. Pychy,<sup>4</sup> M. Richter,<sup>4</sup> C. Schnier,<sup>4</sup> T. Schröder,<sup>4</sup> C. Sowa,<sup>4</sup> M. Steinke,<sup>4</sup> T. Triffterer,<sup>4</sup> U. Wiedner,<sup>4</sup> M. Ball,<sup>5</sup> R. Beck,<sup>5</sup> C. Hammann,<sup>5</sup> B. Ketzer,<sup>5</sup> K. Biguenko,<sup>23</sup> K.T. Brinkmann,<sup>23</sup> V. Di Pietro,<sup>23</sup> S. Diehl,<sup>23</sup> V. Dormenev,<sup>23</sup> P. Drexler,<sup>23</sup> M. Düren,<sup>23</sup> E. Etzelmüller,<sup>23</sup>

#### Event generator based on: B. Pire, L. Szymanowski, K. Semenov-Tian-Shansky (2013)

### **Study of:** $\overline{p} p \rightarrow J/\Psi \pi^0$ (signal)

and different background sources ( $\overline{p} p \rightarrow \pi^+ \pi^- \pi^0, \overline{p} p \rightarrow J/\Psi \pi^0 \pi^0, ...$ )

- Simulations for s = 5 GeV<sup>2</sup>, 10 GeV<sup>2</sup>, 12.2 GeV<sup>2</sup>, 16.9 GeV<sup>2</sup>, 24.3 GeV<sup>2</sup>
- <u>Assumption</u>: 2 fb<sup>-1</sup> of integrated luminosity (~ 5 month of high lumi. data taking)





- Signal and background count rates for 2 fb<sup>-1</sup> (~ 5 months in High Luminosity mode)
   Worst case scenario at p<sub>p</sub> = 5.5 GeV/c: S/B at least factor 10.
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## Baryon to Meson TDA measuremets at JPARC

**J-PARC:** intense pion beams with  $P_{\pi} = 10 - 20 \text{ GeV}$ 



- Charmonium production in assoziation with a nucelon can be used to access TDAs
- Theory: B. Pire, L. Szymanowski and K. Semenov-Tian-Shansky, PRD 95 (2017)



# **Summary and Outlook**

- Nucleon-to-meson TDAs provide new information about correlations of partons inside hadrons
- In the impact parameter space, TDAs provide a spatial imaging of the structure of the pion cloud inside the nucleon
- JLAB 6 GeV data provided first hints for the validity of the TDA based description
- The BSA of the hard exclusive π<sup>+</sup> production shows a clear sign change from forward to backward angles, which may indicate a transition from the GPD to the TDA regime.
- JLAB 12 GeV data will provide more detailed measuremnts in the TDA regime.
- PANDA will allow a check of the universality of TDAs.





