

Towards a combined analysis of inclusive/exclusive electroproduction

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in collaboration with



Mainz Institute for Nuclear Physics



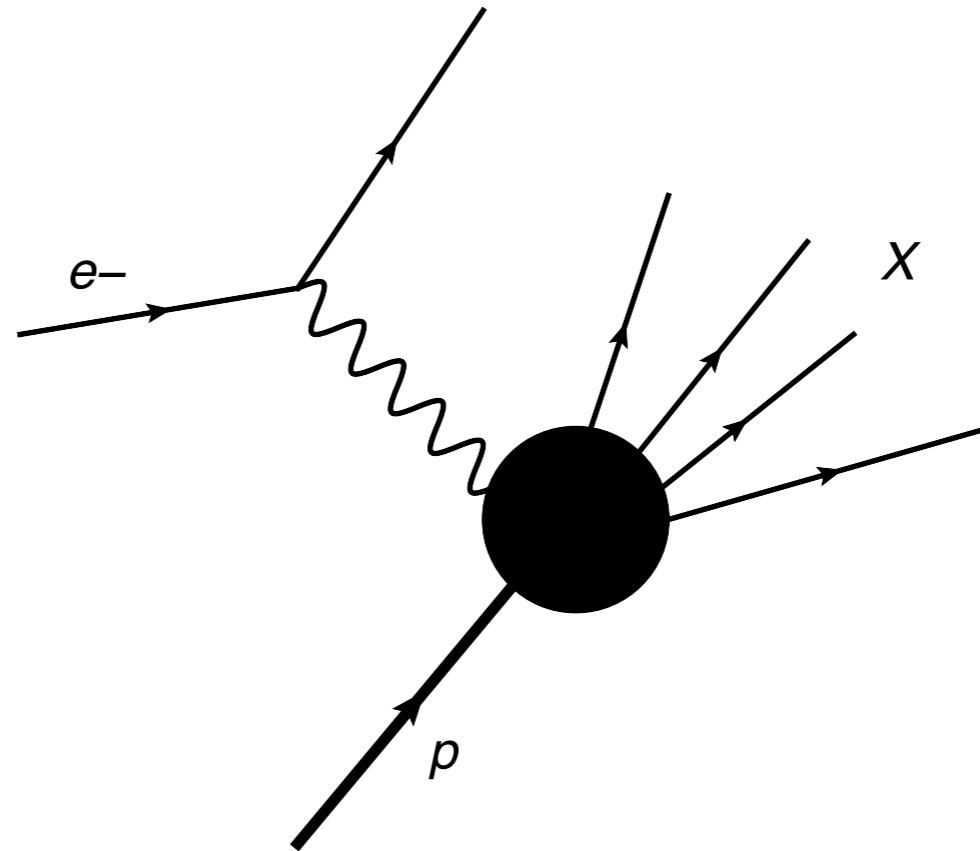
Phys. Rev. C100 (2019) 035201

arXiv:1904.08016 [hep-ph]

STRONG QCD FROM HADRON STRUCTURE EXPERIMENTS 2019

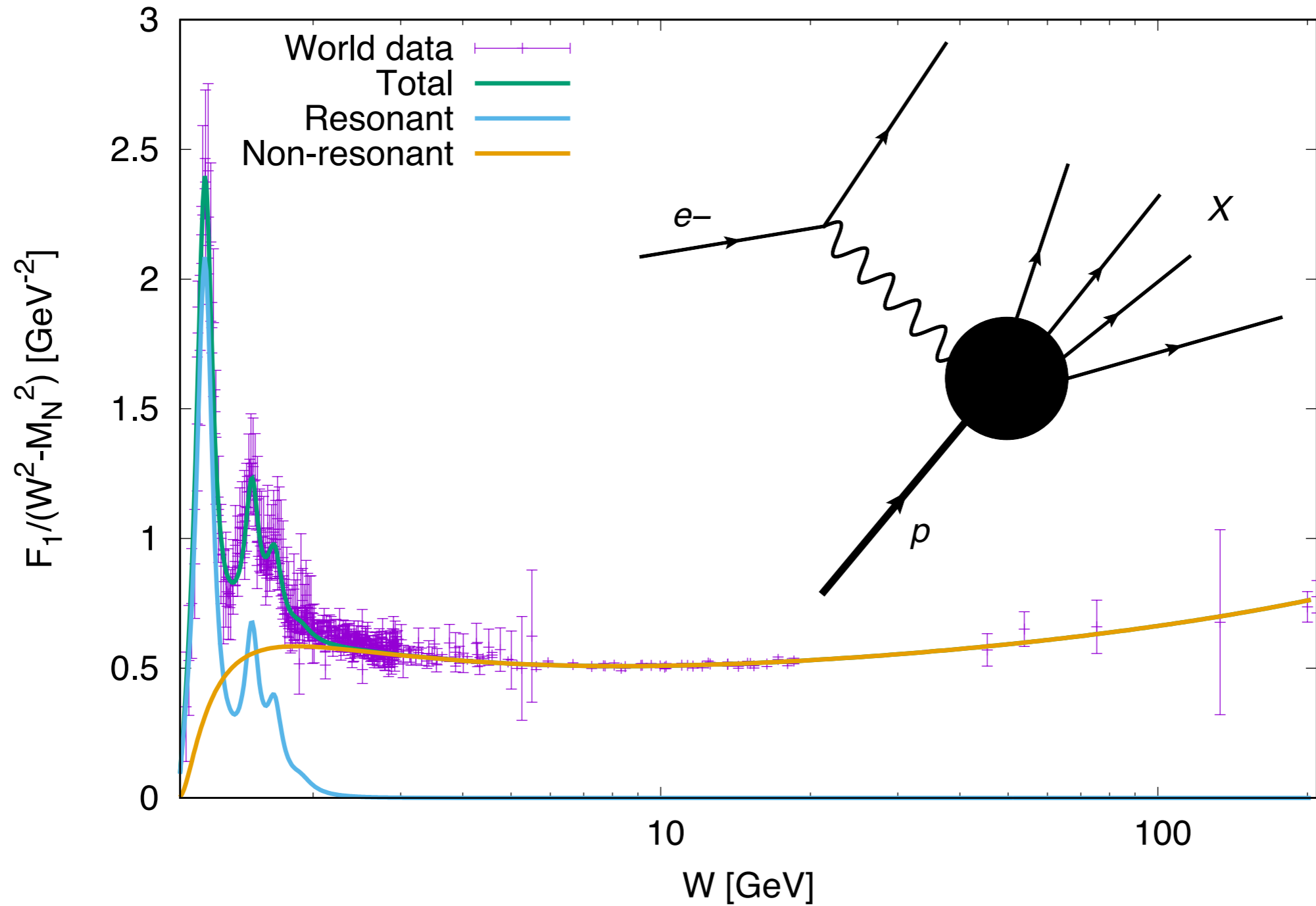
2019-11-08

Nucleon excitation structure



Nucleon excitation structure

Photoproduction



Inclusive and exclusive reactions: can we use one for the other?

Results of resonant contributions to inclusive observables

Non-resonant piece: low and high energies communicate

Results of combined description (preliminary)

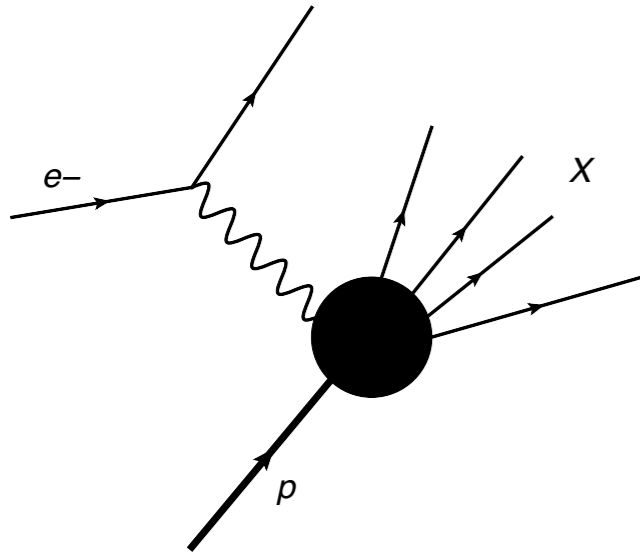
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Inclusive electron scattering



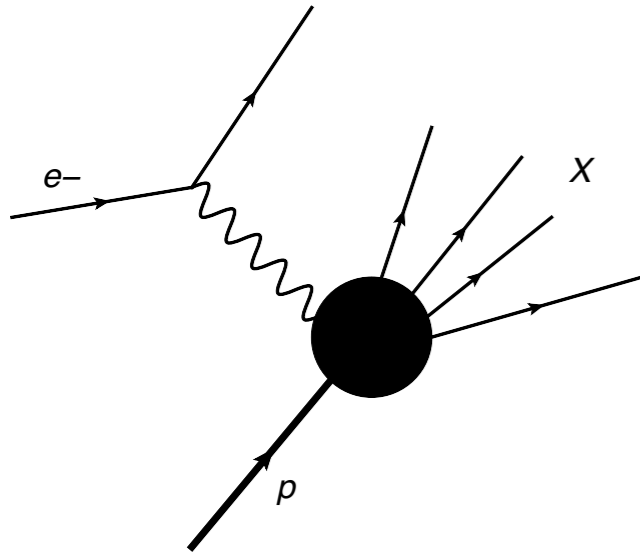
$$F_1 \propto \sigma_T(W, Q^2)$$

$$F_2 \propto \sigma_T(W, Q^2) + \sigma_L(W, Q^2)$$

$$\sigma_U(W, Q^2) = \sigma_T(W, Q^2) + \epsilon_T \sigma_L(W, Q^2)$$

- Gives access to structure functions
- Here: unpolarized structure functions

Inclusive electron scattering

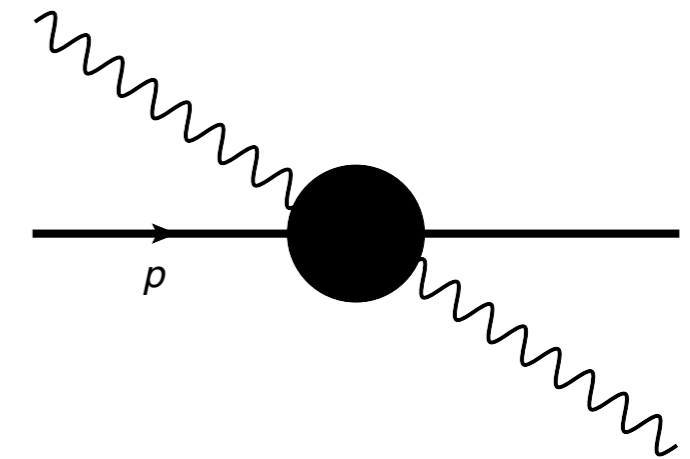


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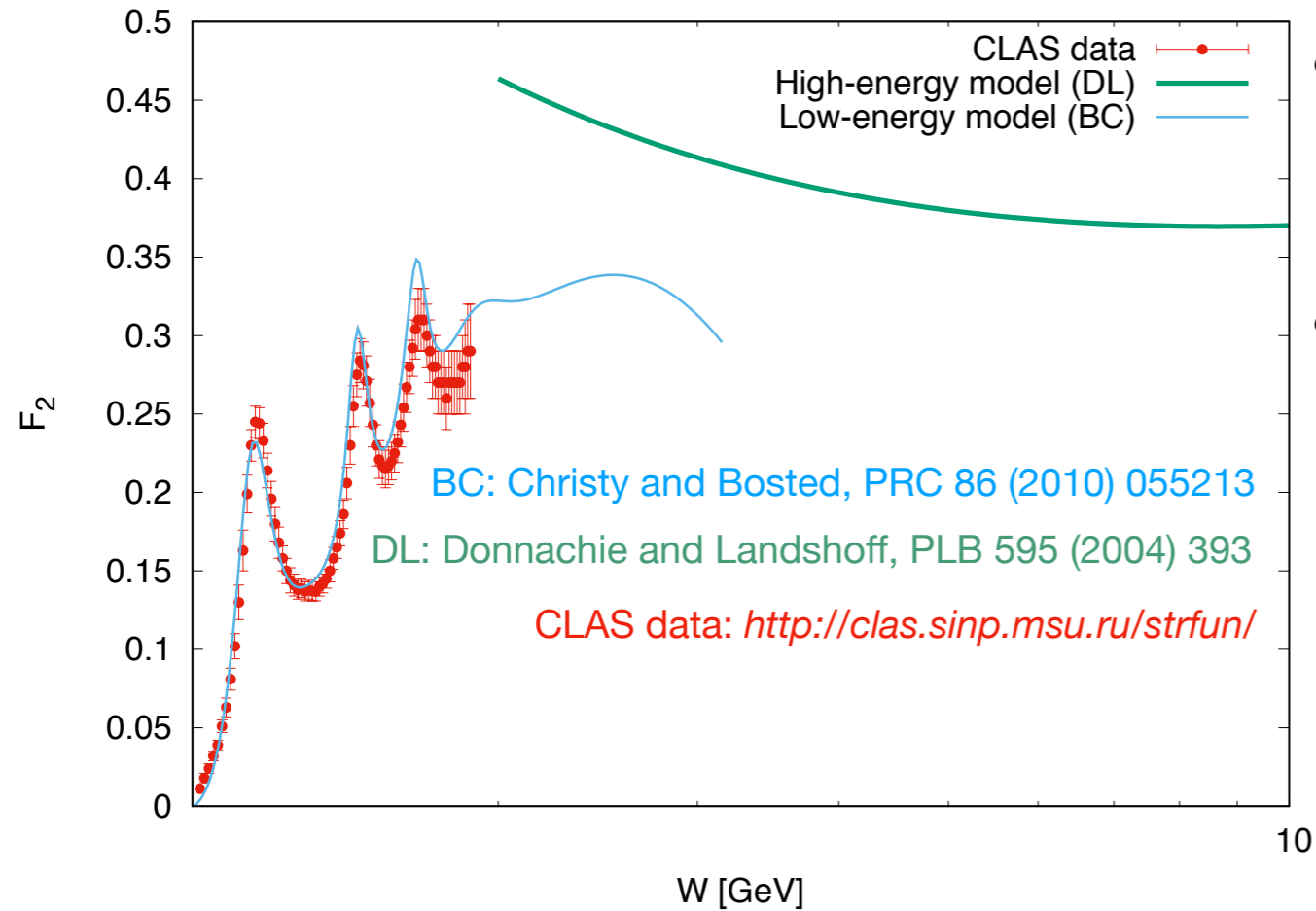
$$\Im T_1(\nu, Q^2) = \frac{\pi\alpha}{M} F_1(\nu, Q^2)$$

$$\Im T_2(\nu, Q^2) = \frac{\pi\alpha}{\nu} F_2(\nu, Q^2)$$

- Related to forward virtual Compton scattering (VVCS) amplitudes T_1 and T_2
- Information about hadronic contribution to **muonic hydrogen Lamb shift!**

Inclusive data and models on the market

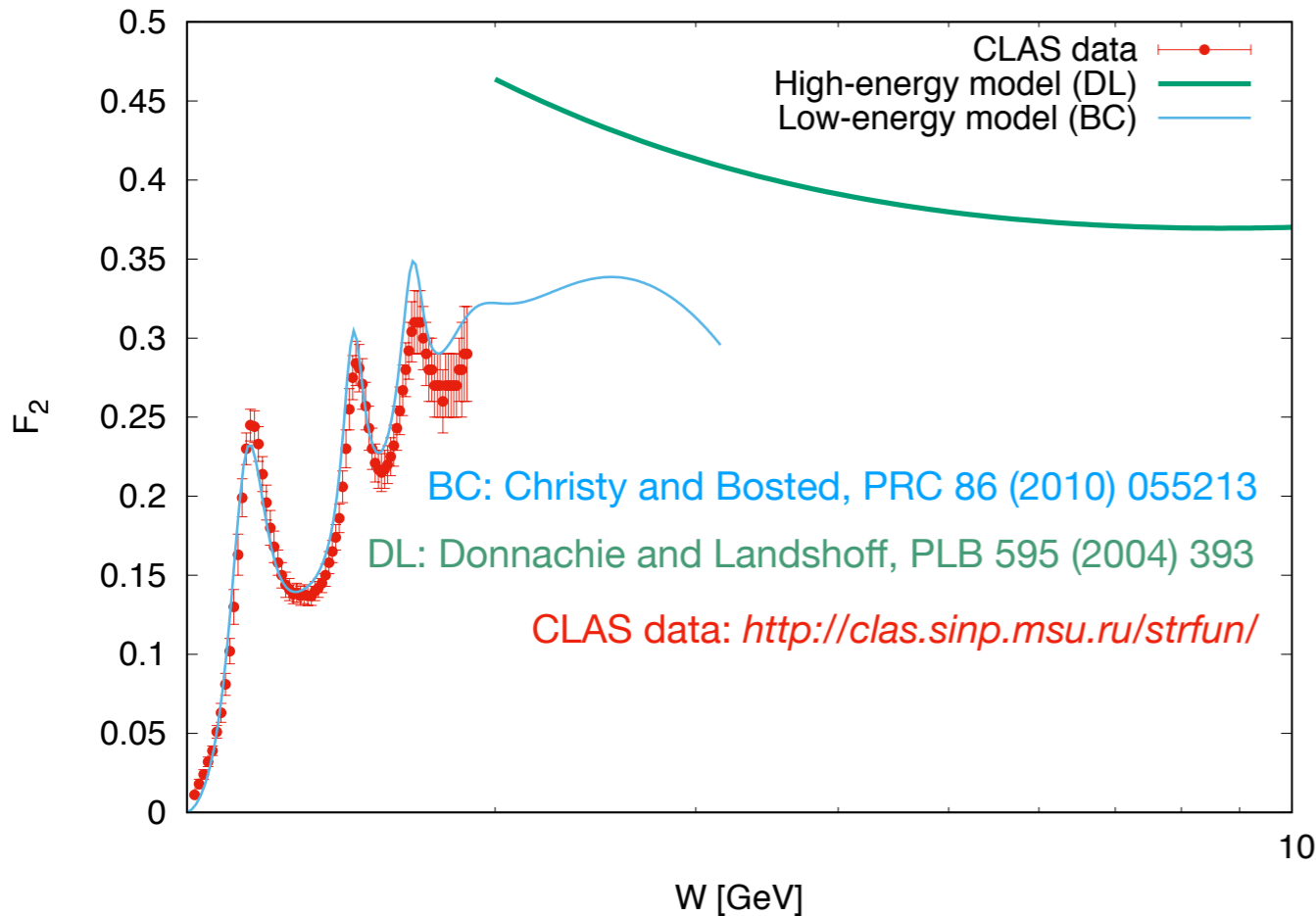
$$Q^2=1.0 \text{ GeV}^2$$



- **At low energies: precise CLAS data**
Well covered up to 3rd resonance regime:
towards parton distributions at large x
- CLAS12 is to reach
 $0.05 \text{ GeV}^2 < Q^2 < 12 \text{ GeV}^2$, W up to 4 GeV

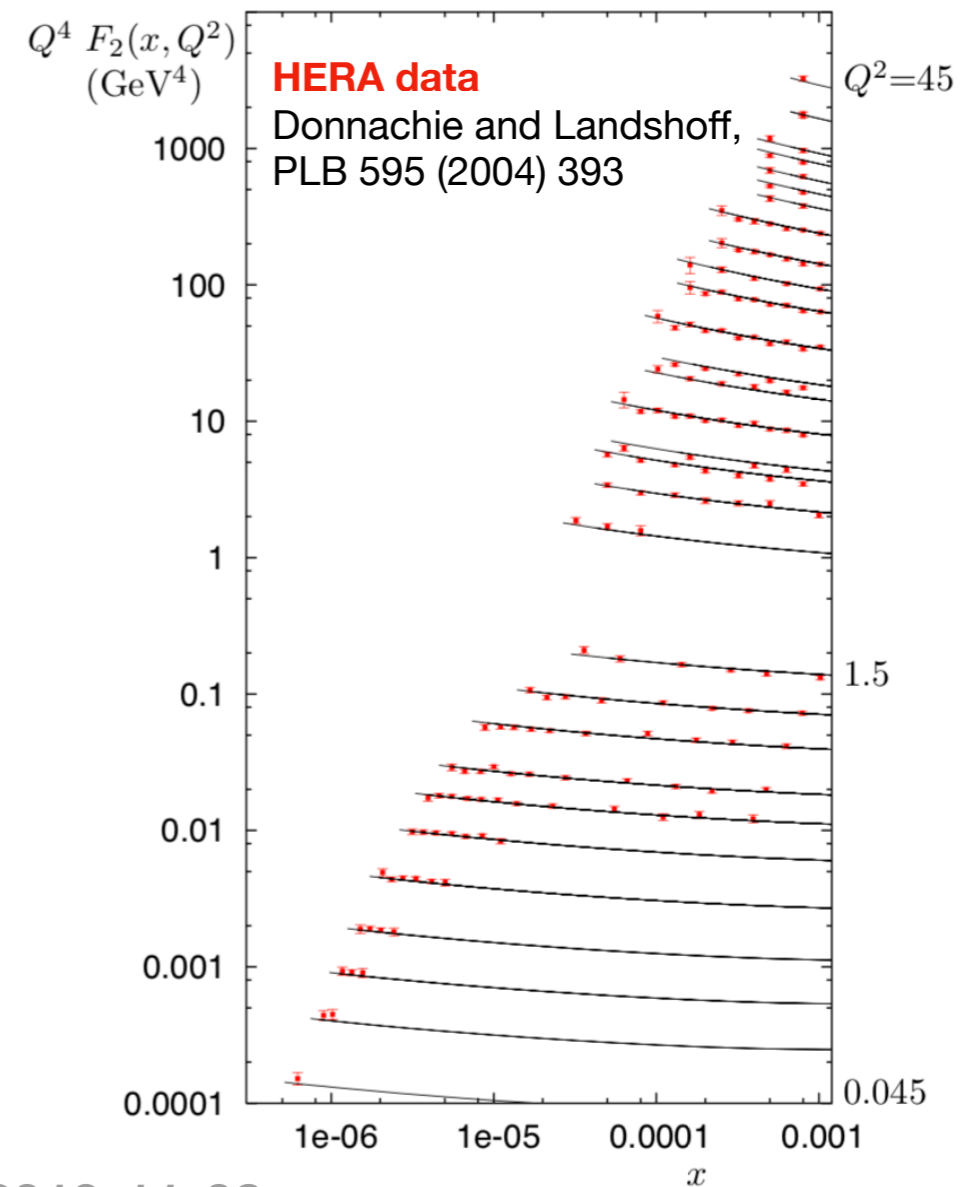
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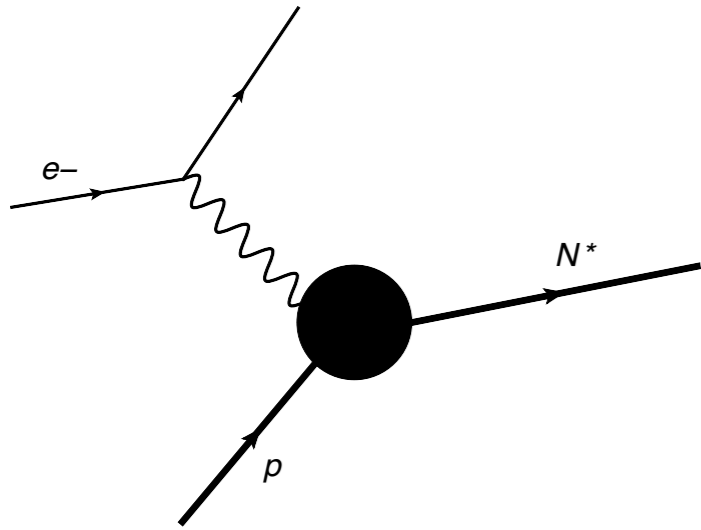


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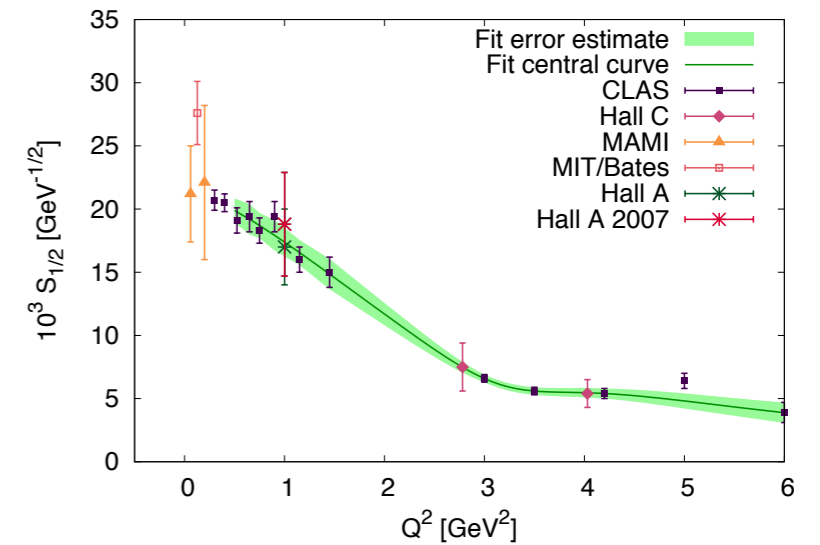
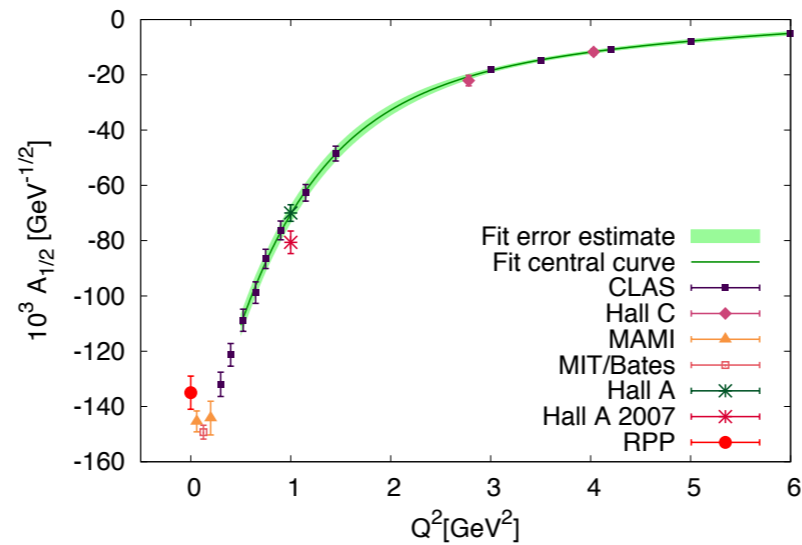
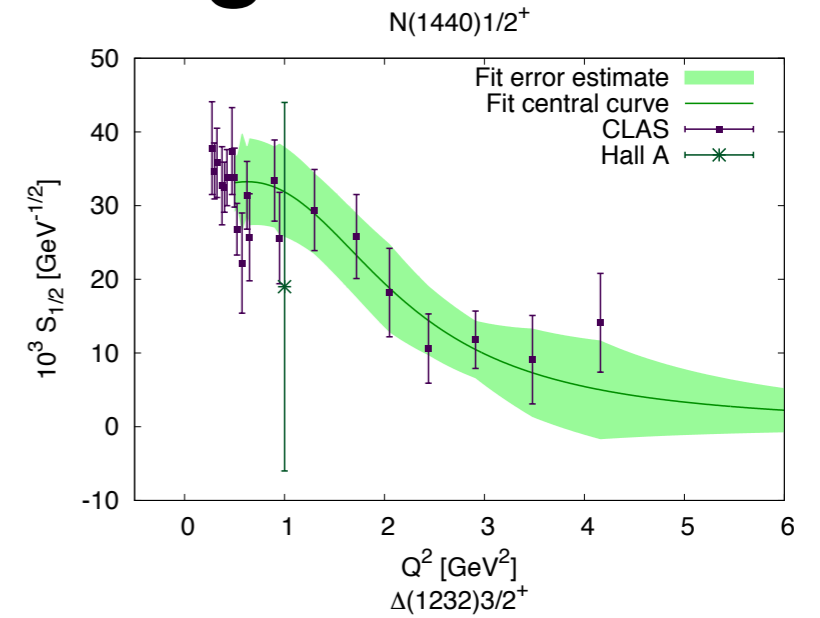
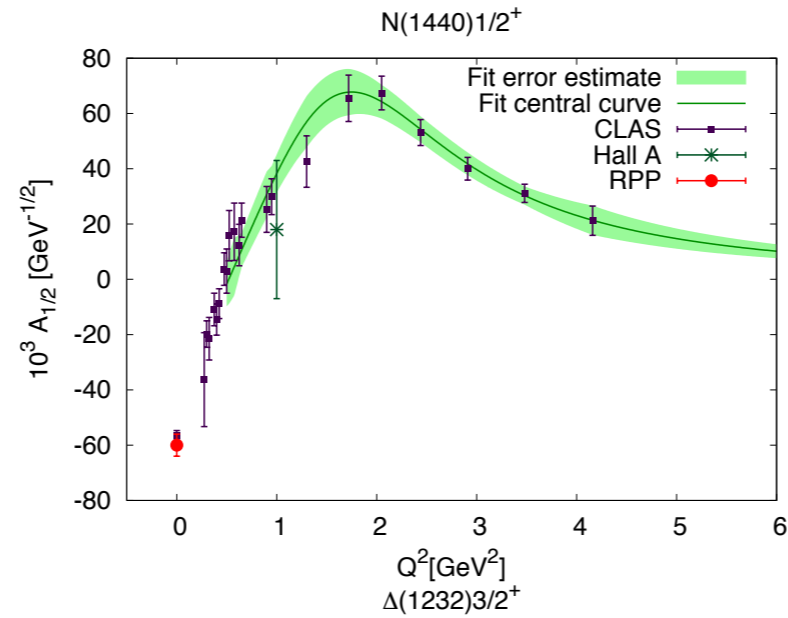
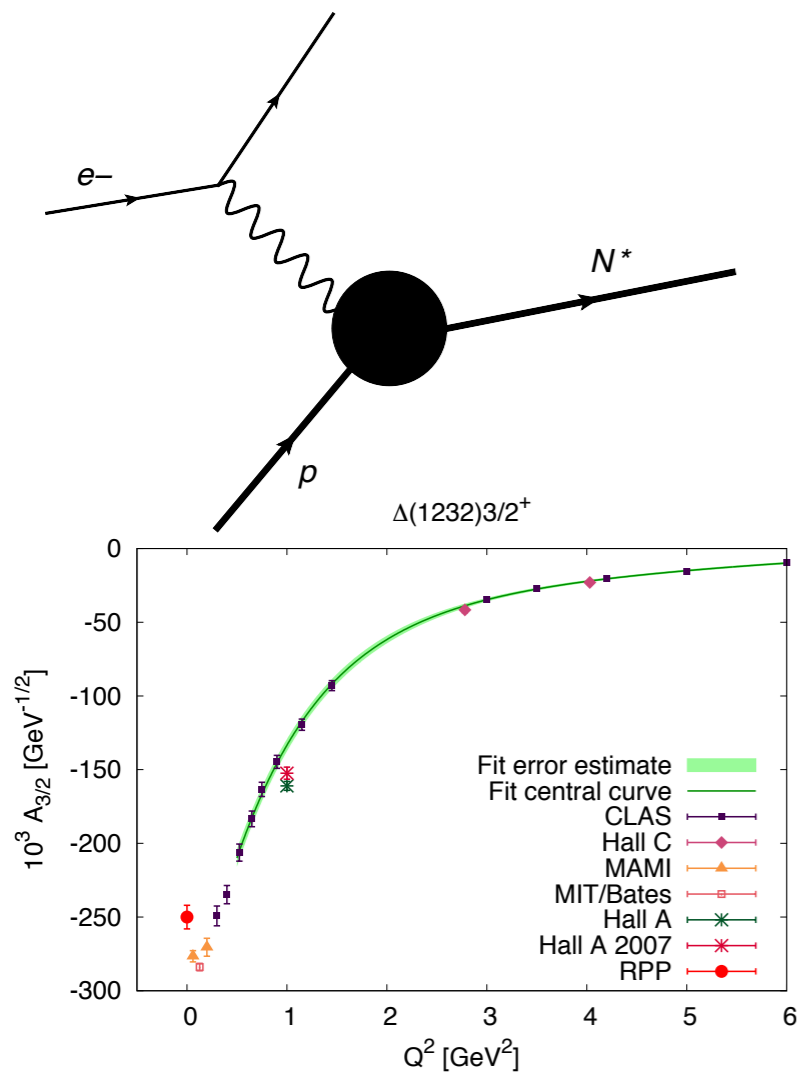
- **At high energies: HERA and ZEUS data**
Extremely good Q^2 coverage at low x
- **Goal:** combining high and low-energy models
- Description of observables integrated over x :
Cottingham formula, subtraction function in VVCS,
Lamb shift, ...
- Tests on quark-hadron duality



Exclusive electron scattering



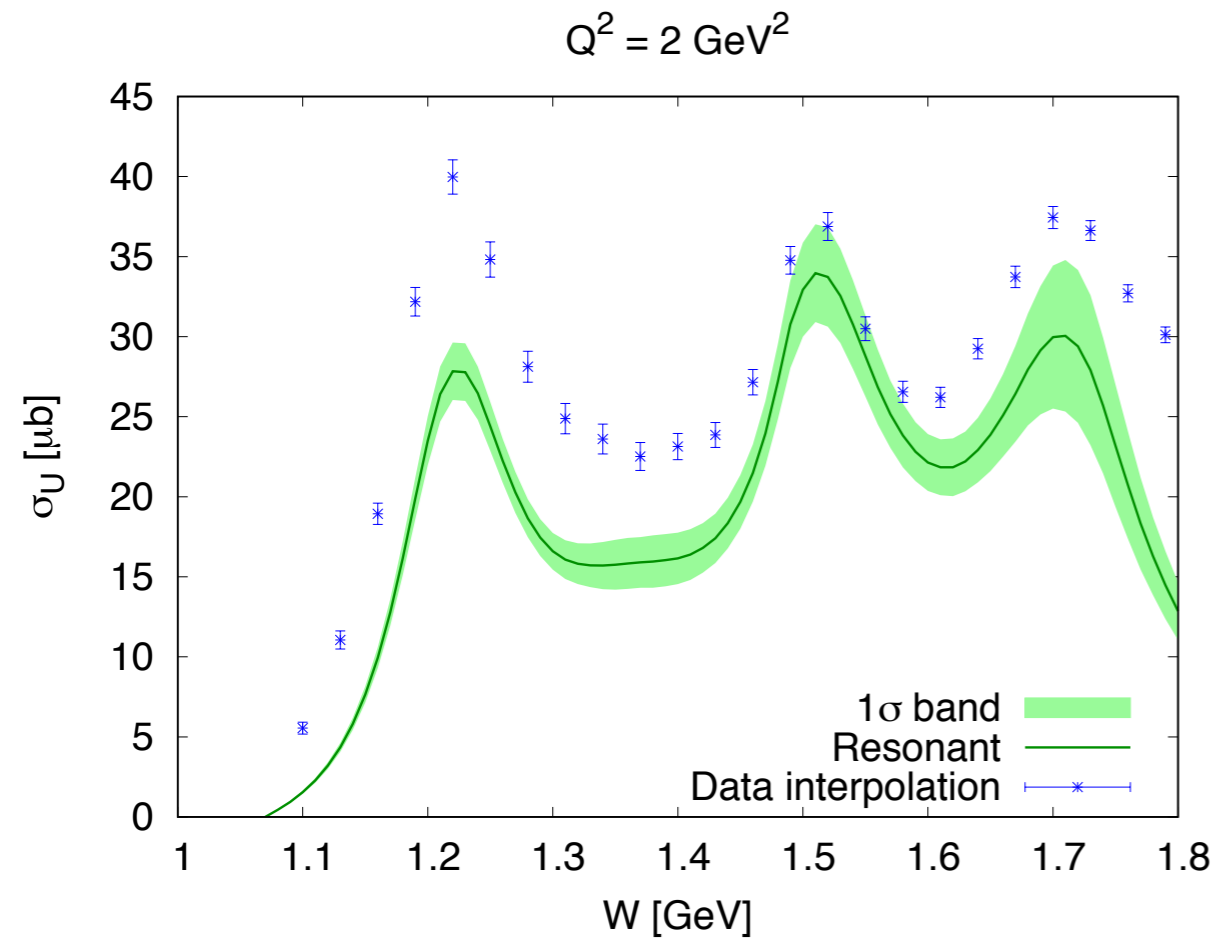
Exclusive electron scattering



- Longitudinal and transverse electrocouplings
- Allows us to determine each of the resonant contributions separately
- Good world data on some of the main resonances; CLAS data up to 3rd resonance region (and soon full resonance region coverage with CLAS12)

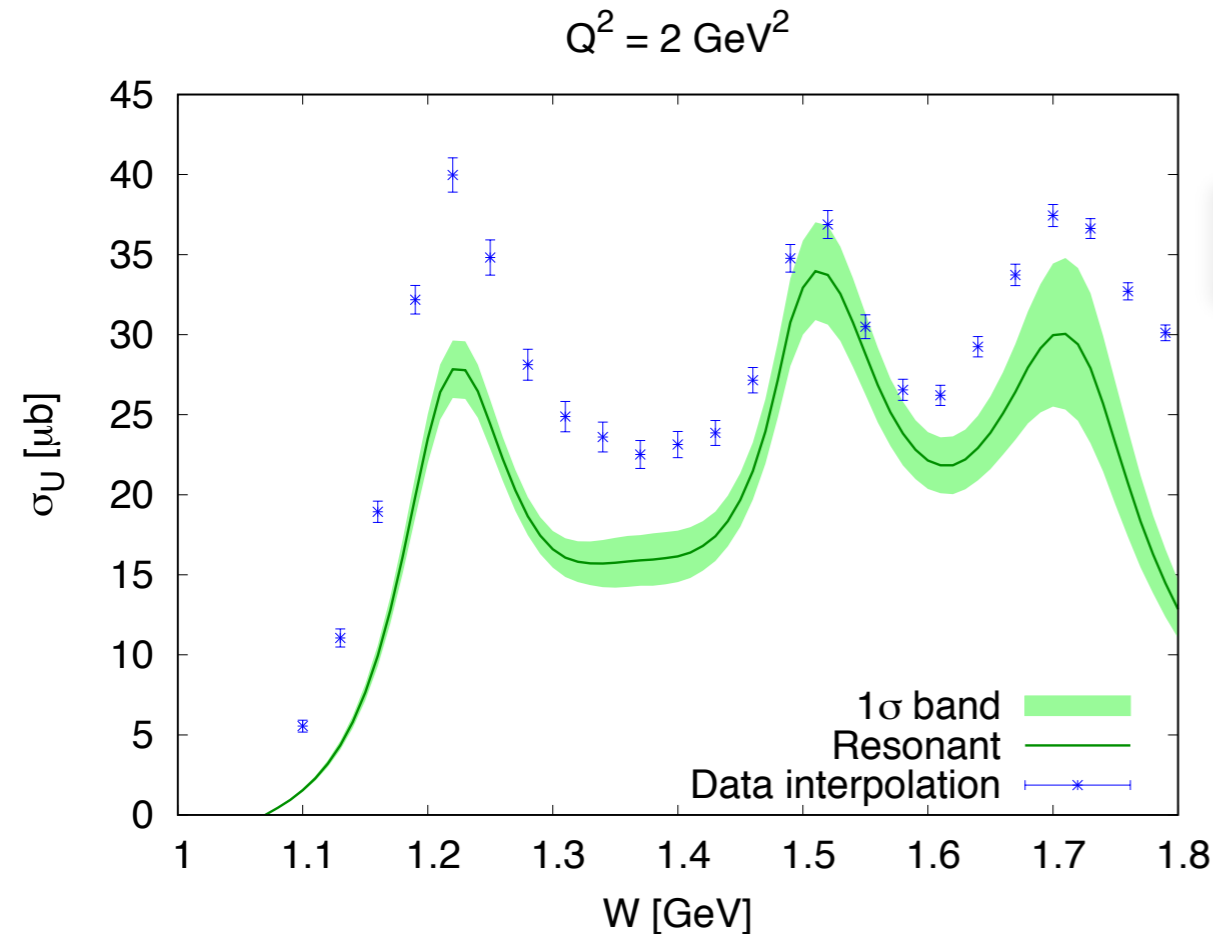
From exclusive to inclusive electron scattering

$$\sigma_{T,L}(W, Q^2) = \sigma_{T,L}^R(W, Q^2) + \sigma_{T,L}^{NR}(W, Q^2)$$



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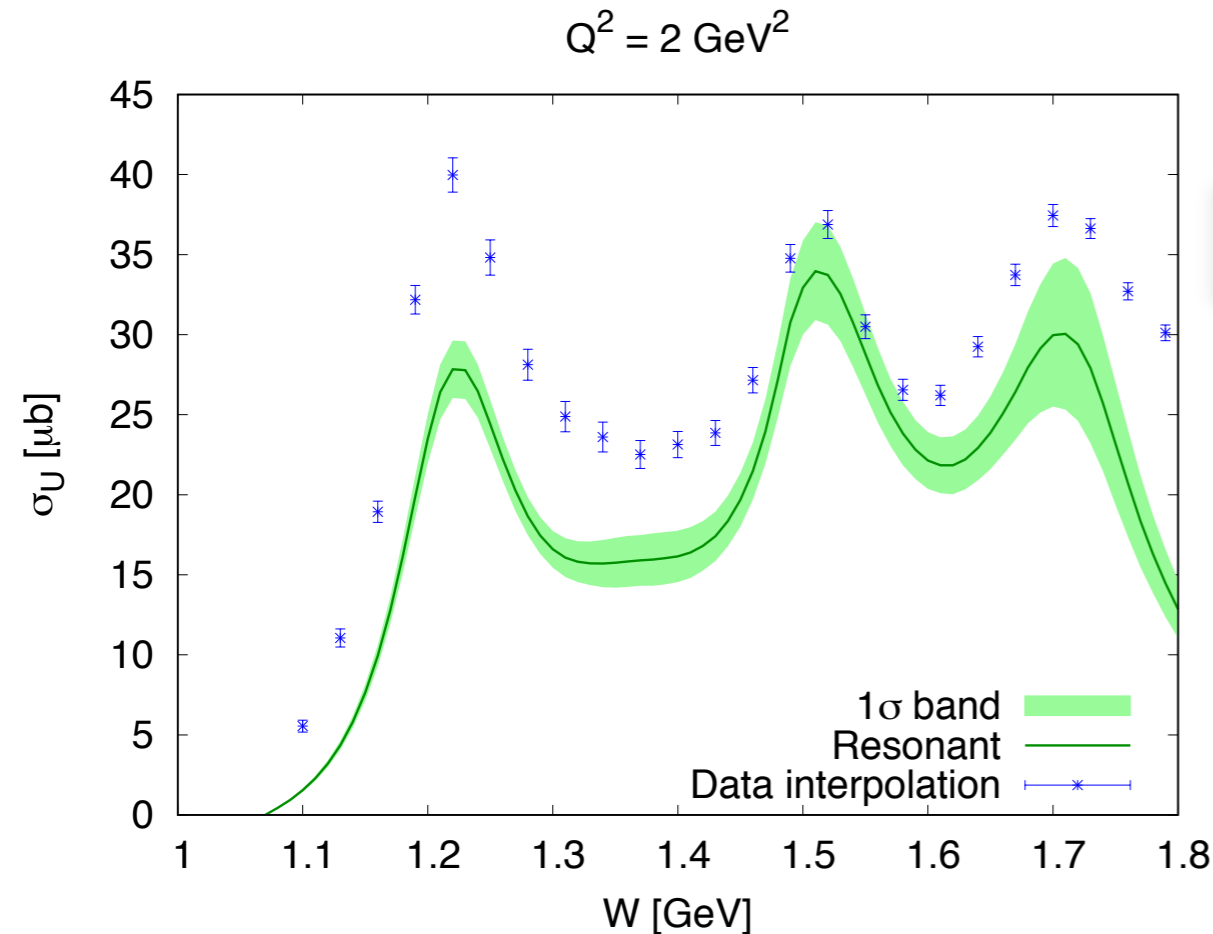
Breit-Wigner resonance model

Moiseev et al., PRC 86 (2012) 035203

$$\sigma_{T,L}^R(W, Q^2) = \frac{\pi}{q_\gamma^2} \sum_{N^*, \Delta^*} (2J_r + 1) \frac{M_r^2 \Gamma_{\text{tot}}(W) \Gamma_\gamma^{T,L}(M_r)}{(M_r^2 - W^2)^2 + M_r^2 \Gamma_{\text{tot}}^2(W)}$$

From exclusive to inclusive electron scattering

$$\sigma_{T,L}(W, Q^2) = \sigma_{T,L}^R(W, Q^2) + \sigma_{T,L}^{NR}(W, Q^2)$$



Breit-Wigner resonance model

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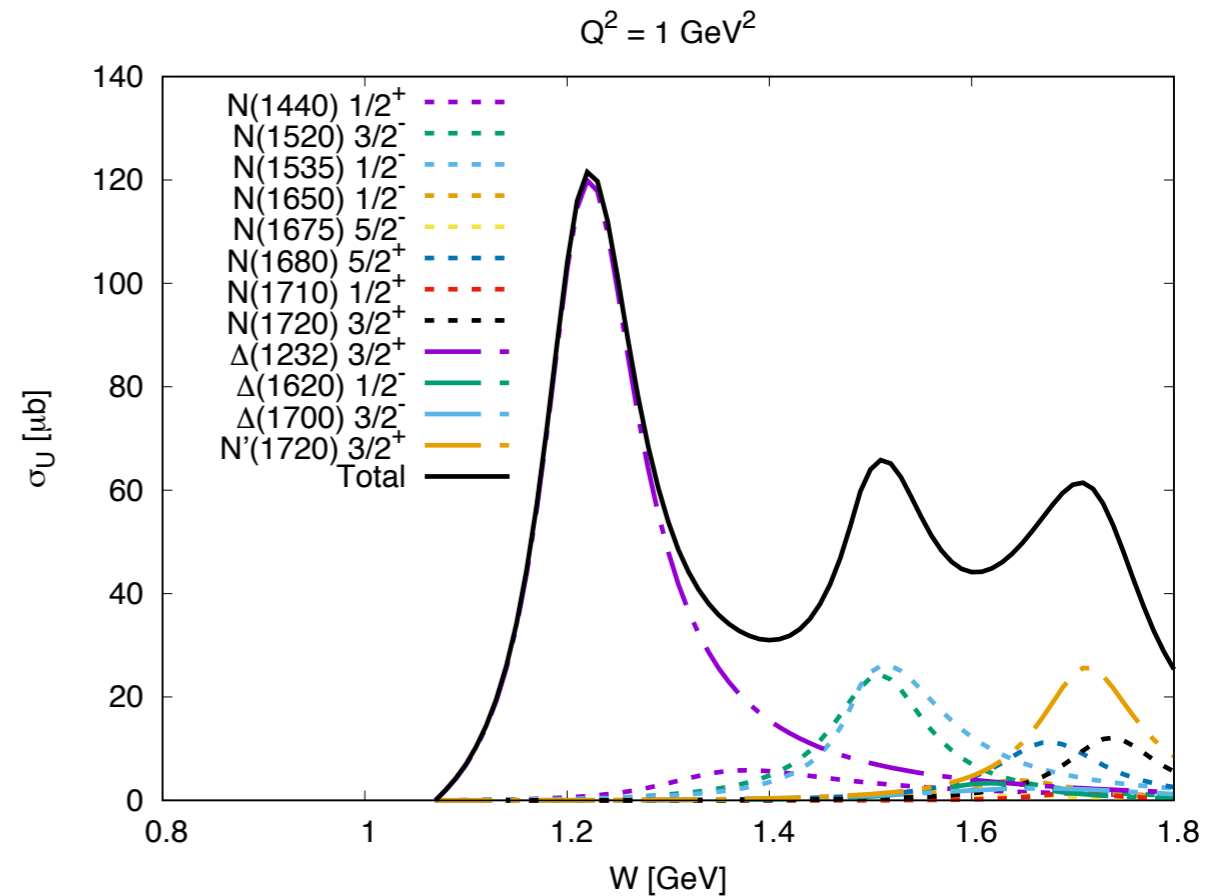
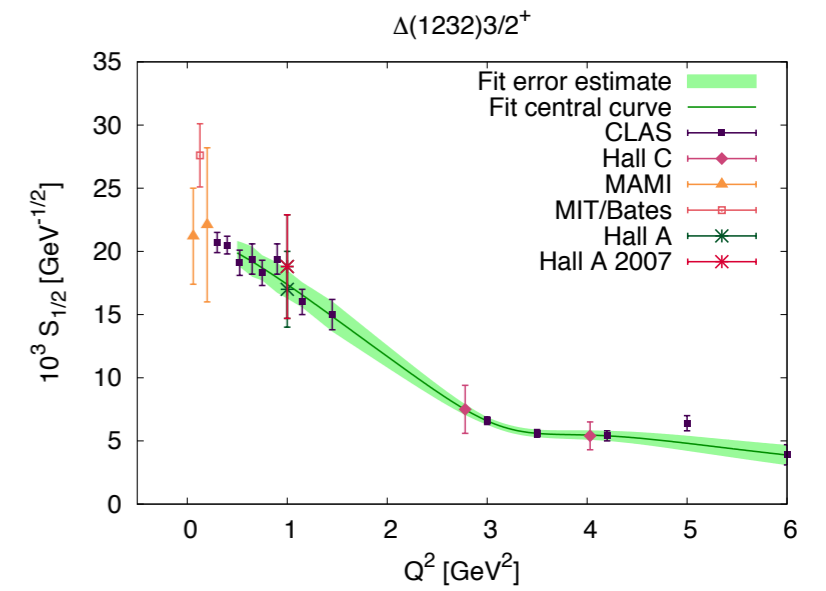
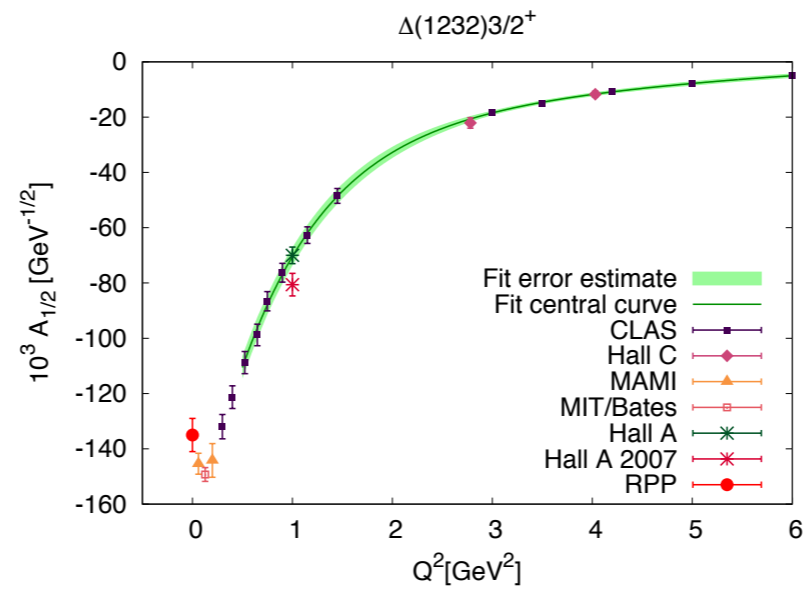
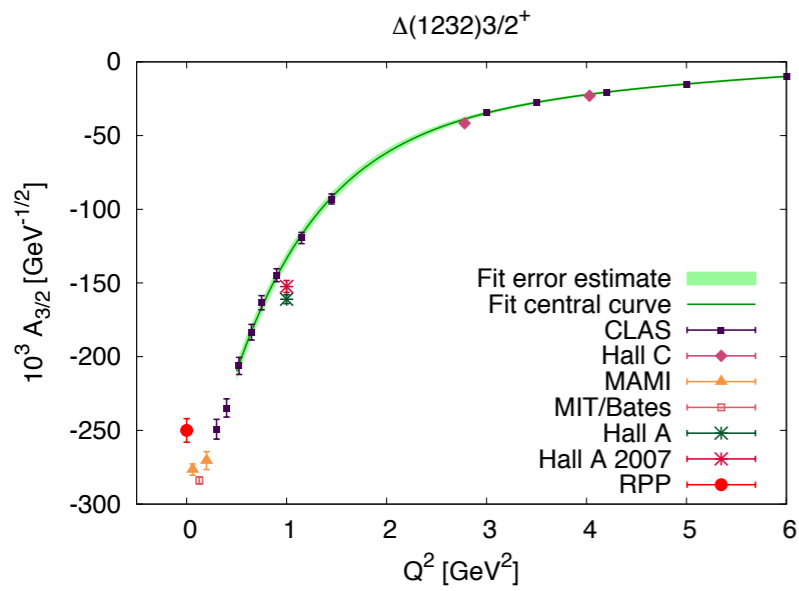
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$$\Gamma_\gamma^T(M_r, Q^2) \sim \left| A_{1/2}(Q^2) \right|^2 + \left| A_{3/2}(Q^2) \right|^2$$

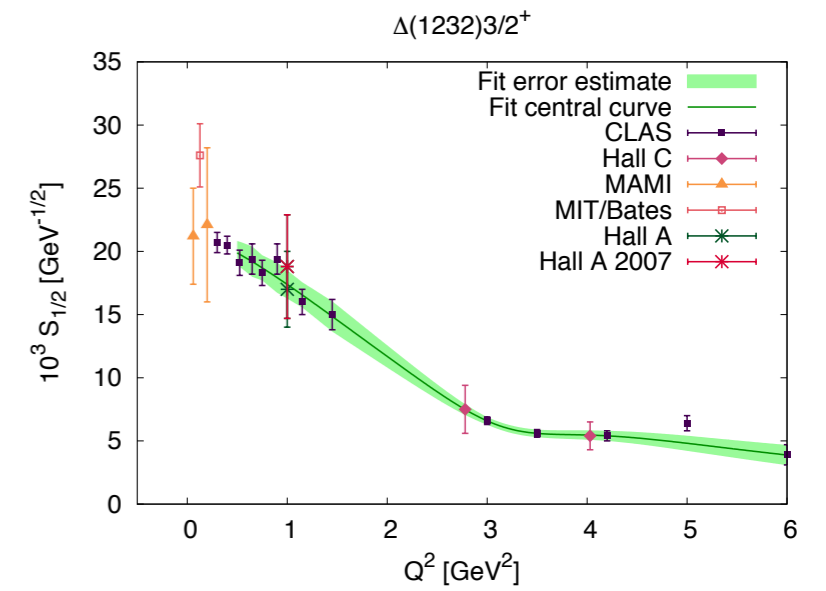
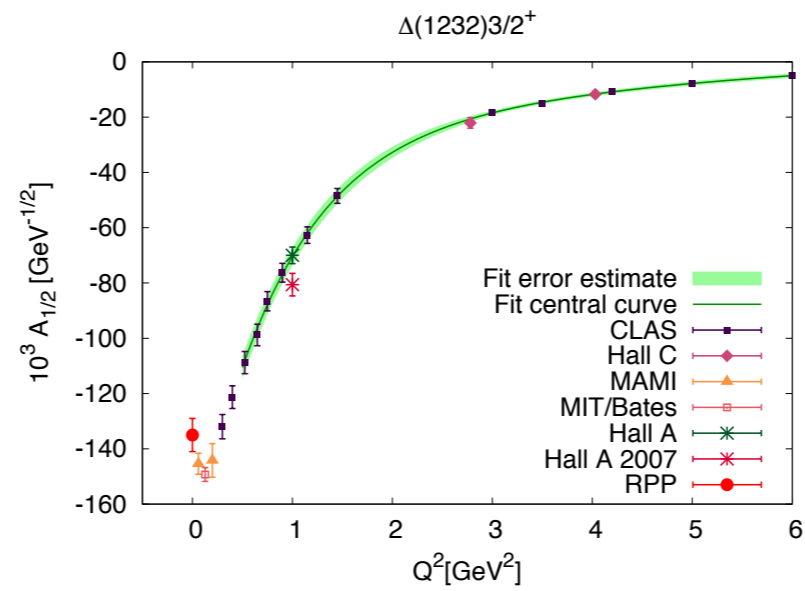
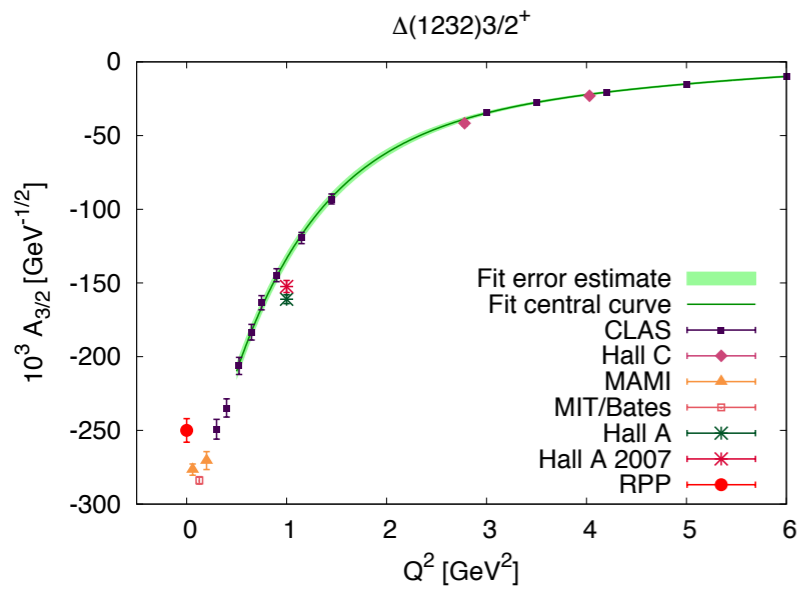
$$\Gamma_\gamma^L(M_r, Q^2) \sim \left| S_{1/2}(Q^2) \right|^2$$

Electrocouplings from data

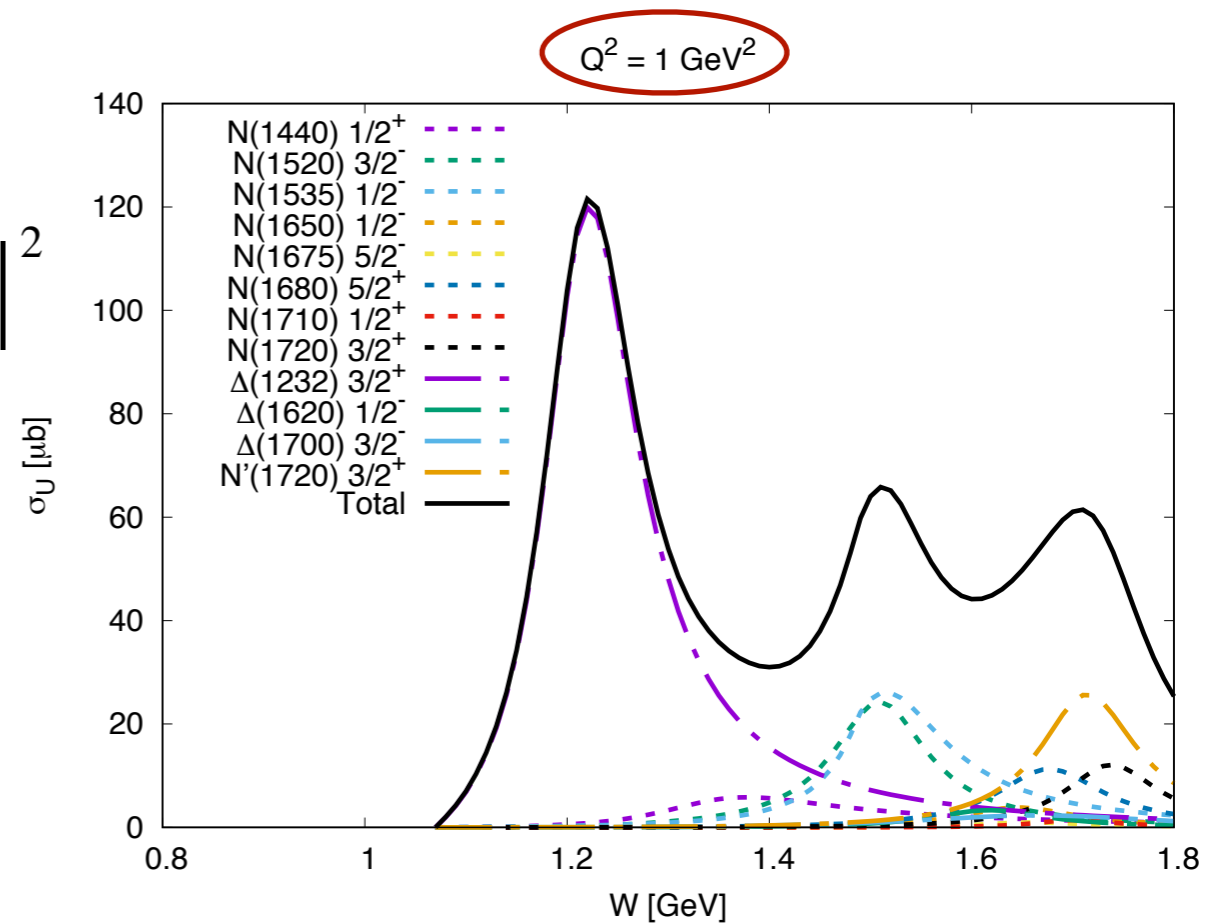
How electrocouplings enter inclusive data



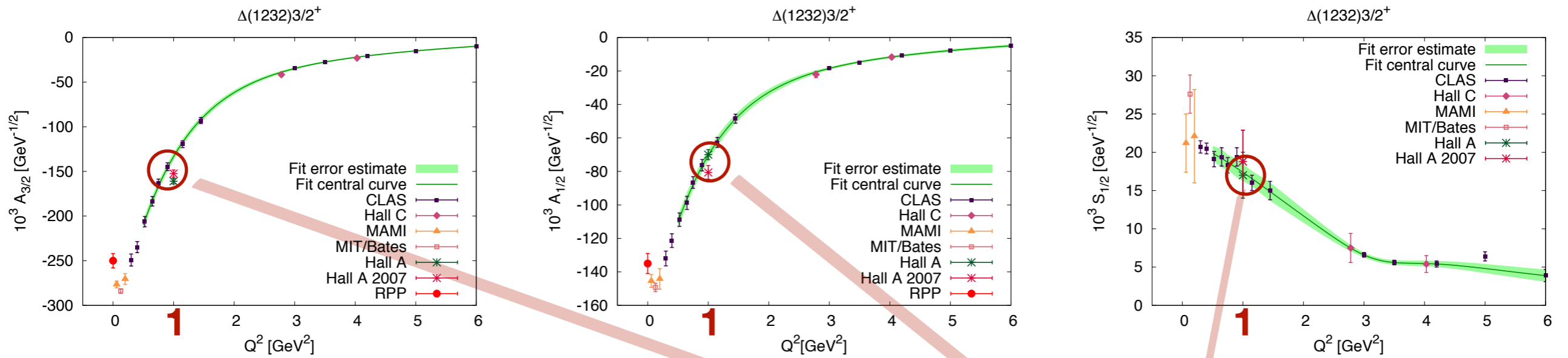
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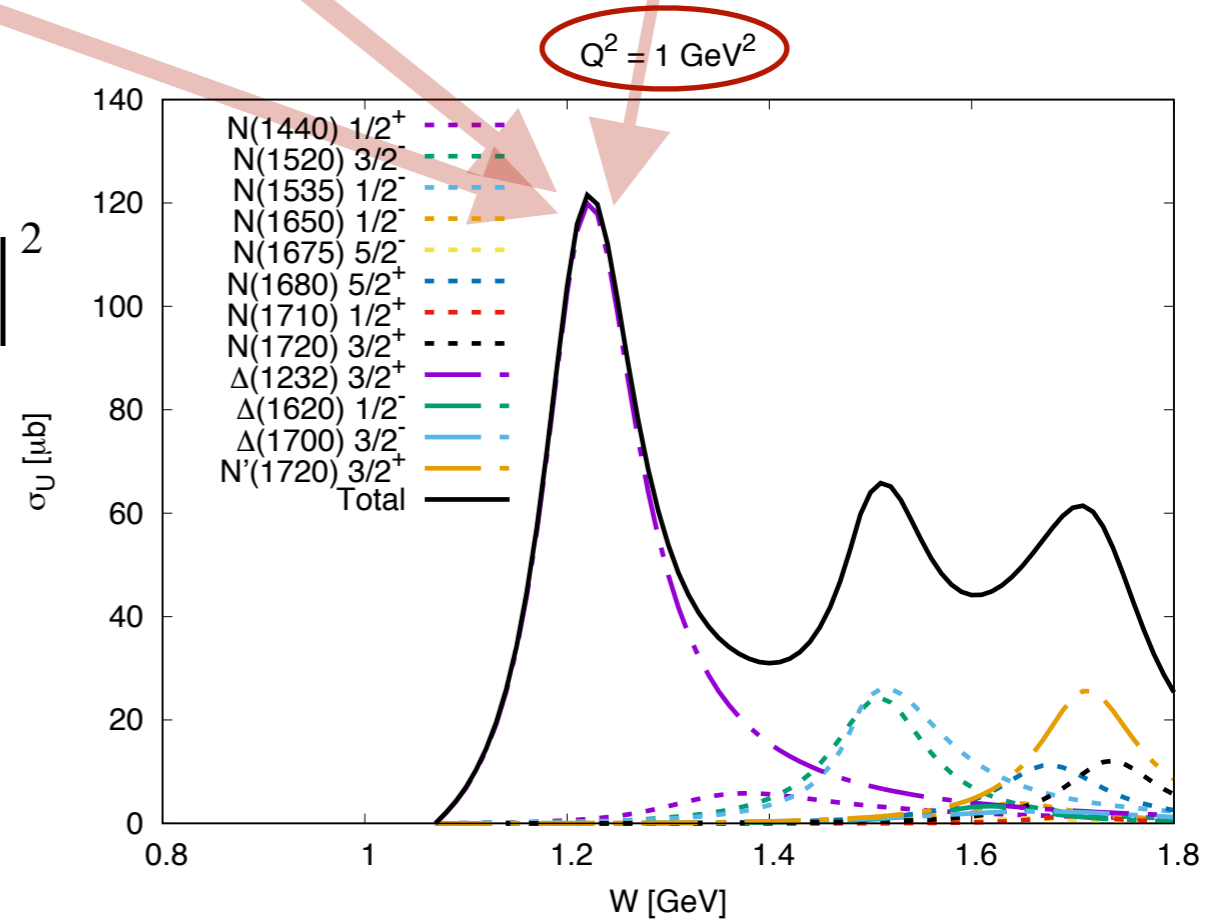
$$\sigma_{T,L}^R(W, Q^2) \sim \left| A_{1/2}(Q^2) \right|^2, \left| A_{3/2}(Q^2) \right|^2, \left| S_{1/2}(Q^2) \right|^2$$



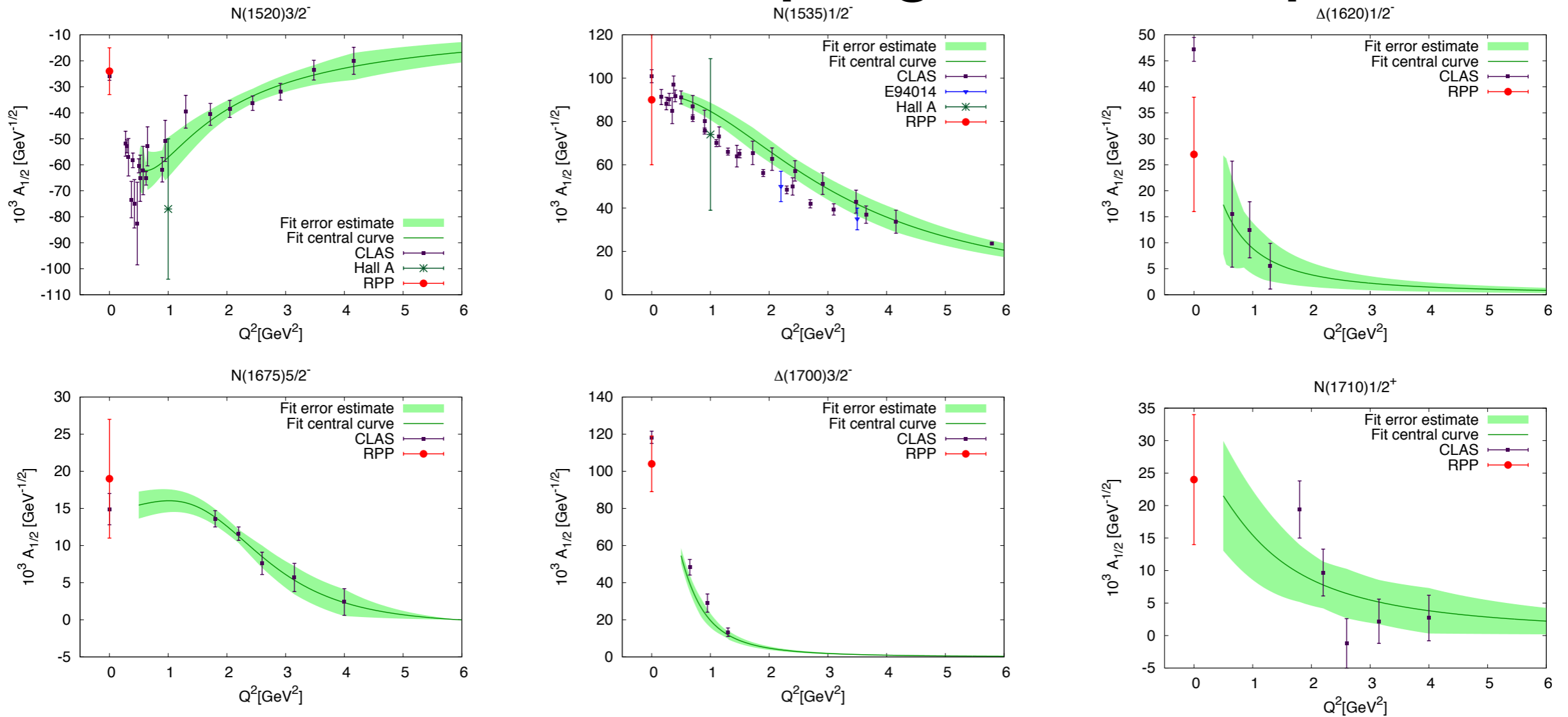
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Further electrocoupling data examples



https://userweb.jlab.org/~mokeev/resonance_electrocouplings/

- Interpolation/extrapolation functions: <https://userweb.jlab.org/~isupov/couplings/>
in good agreement with world data and preliminary CLAS12 results at higher Q^2
- Error bands estimated from data uncertainties
and scaled with coupling size in extrapolation region

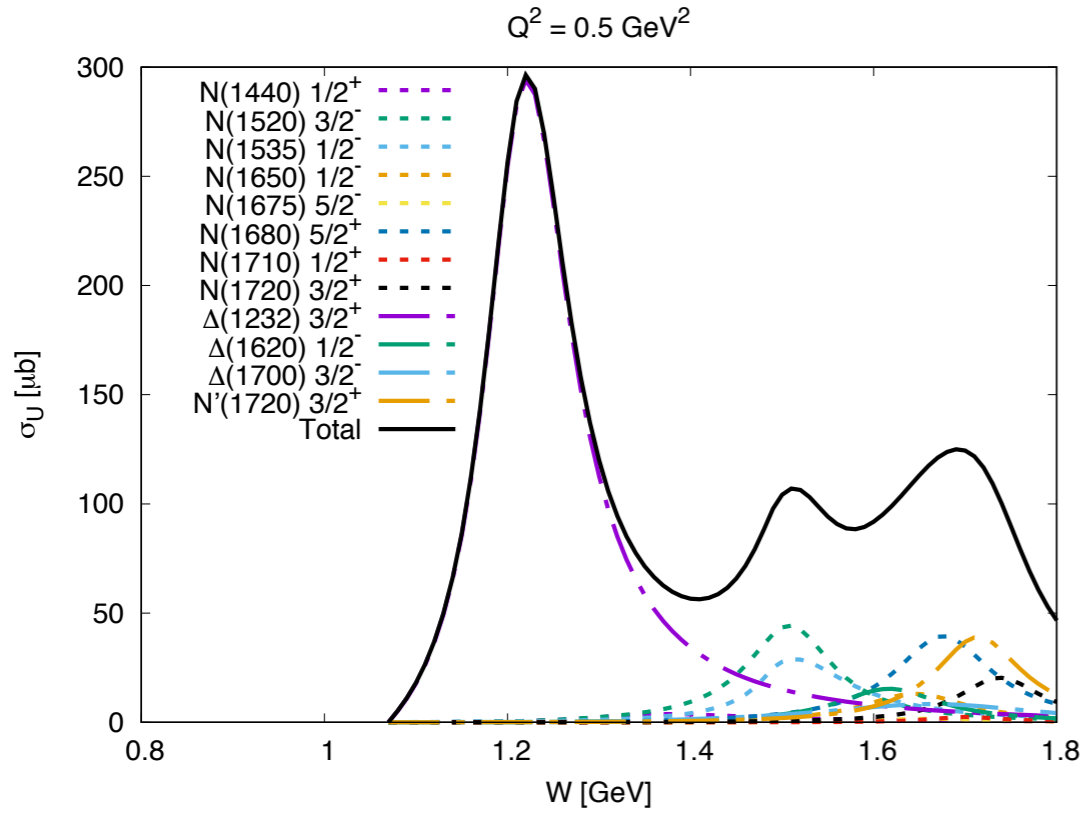
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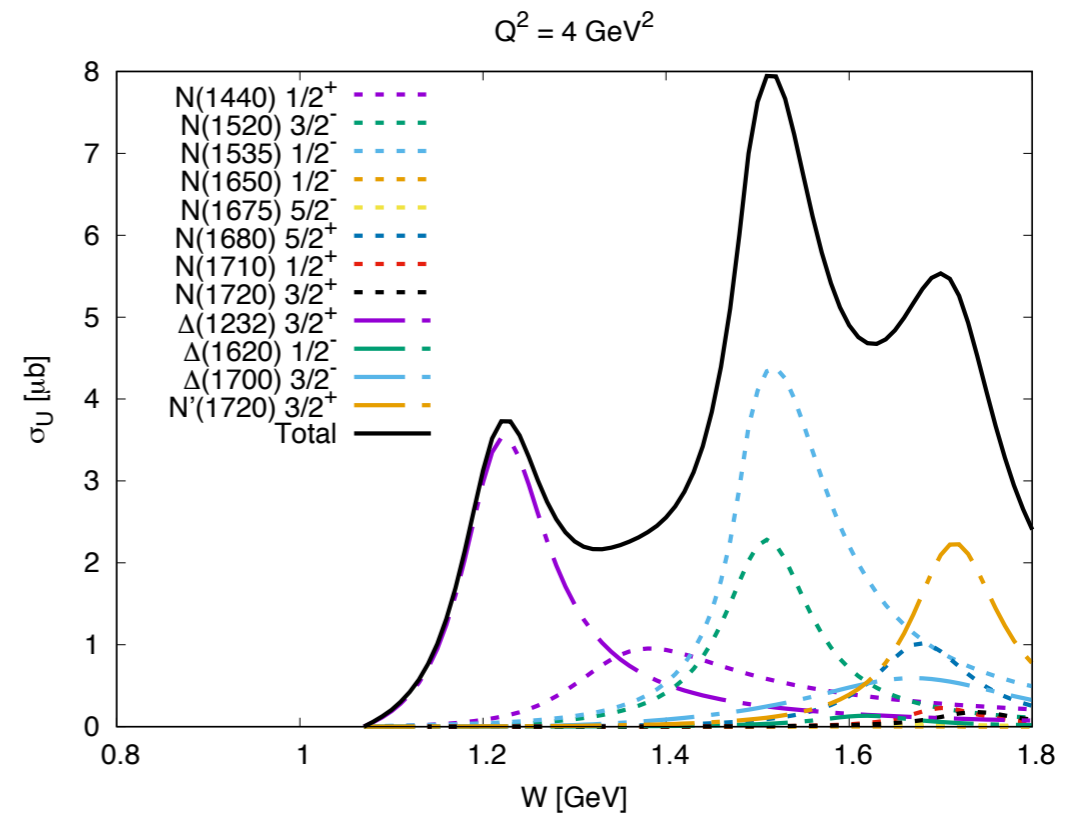
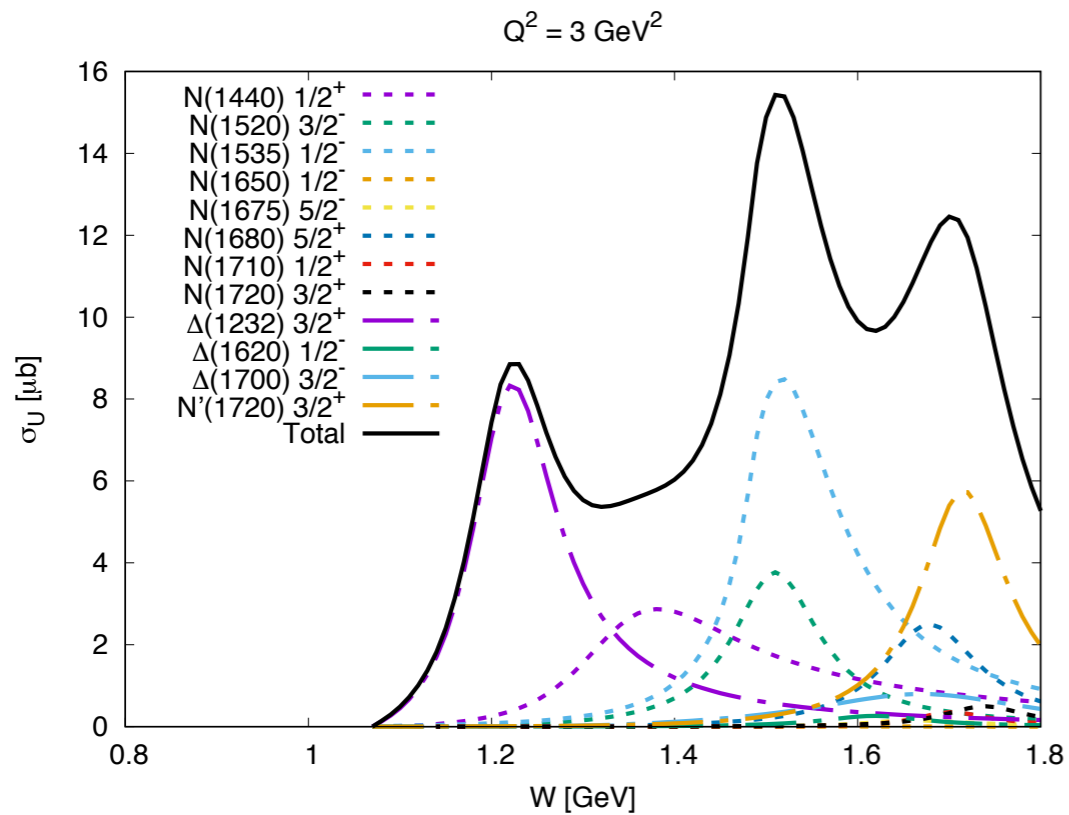
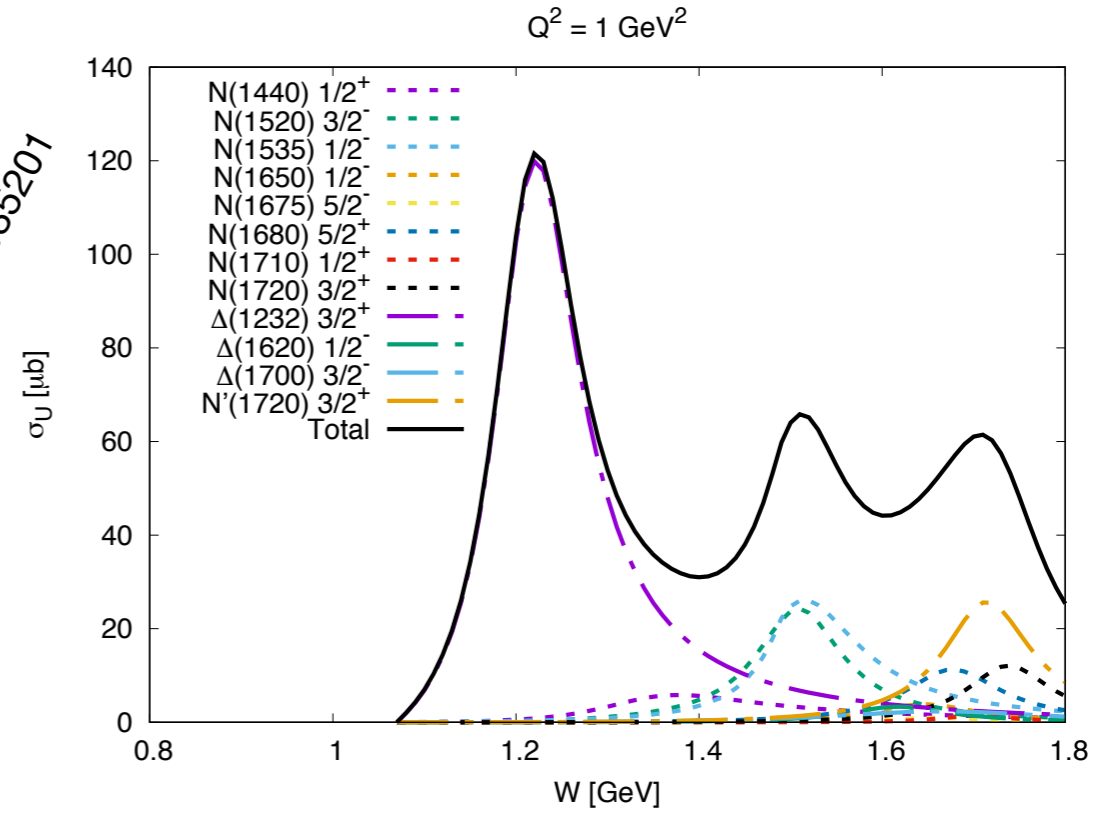
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Results of combined description (preliminary)

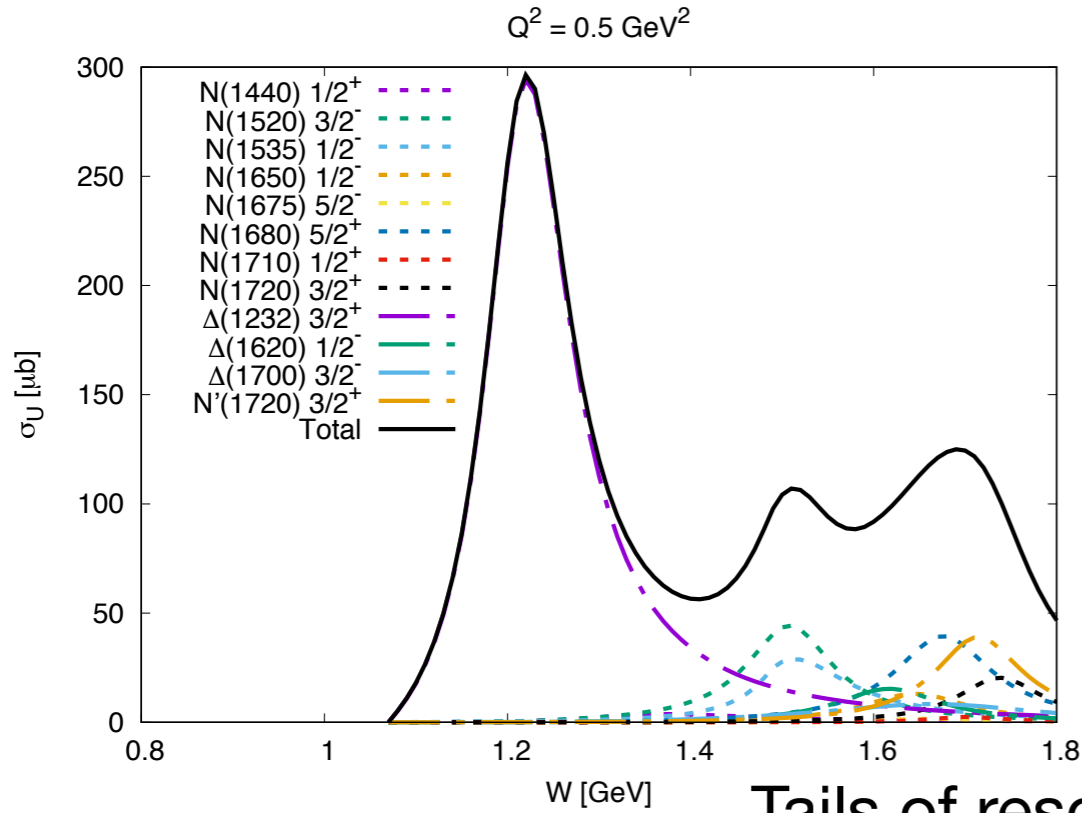
Resonant contributions at different Q^2



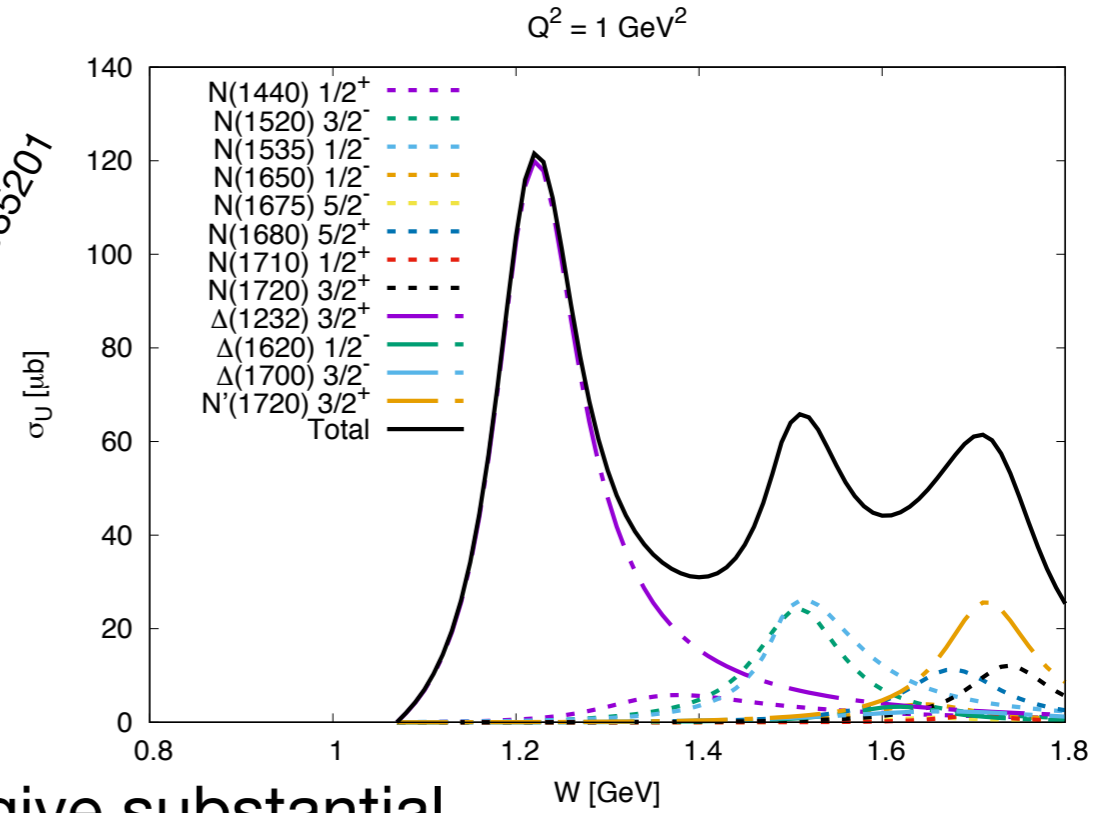
ANHB et al., PRC100 (2019) 035201



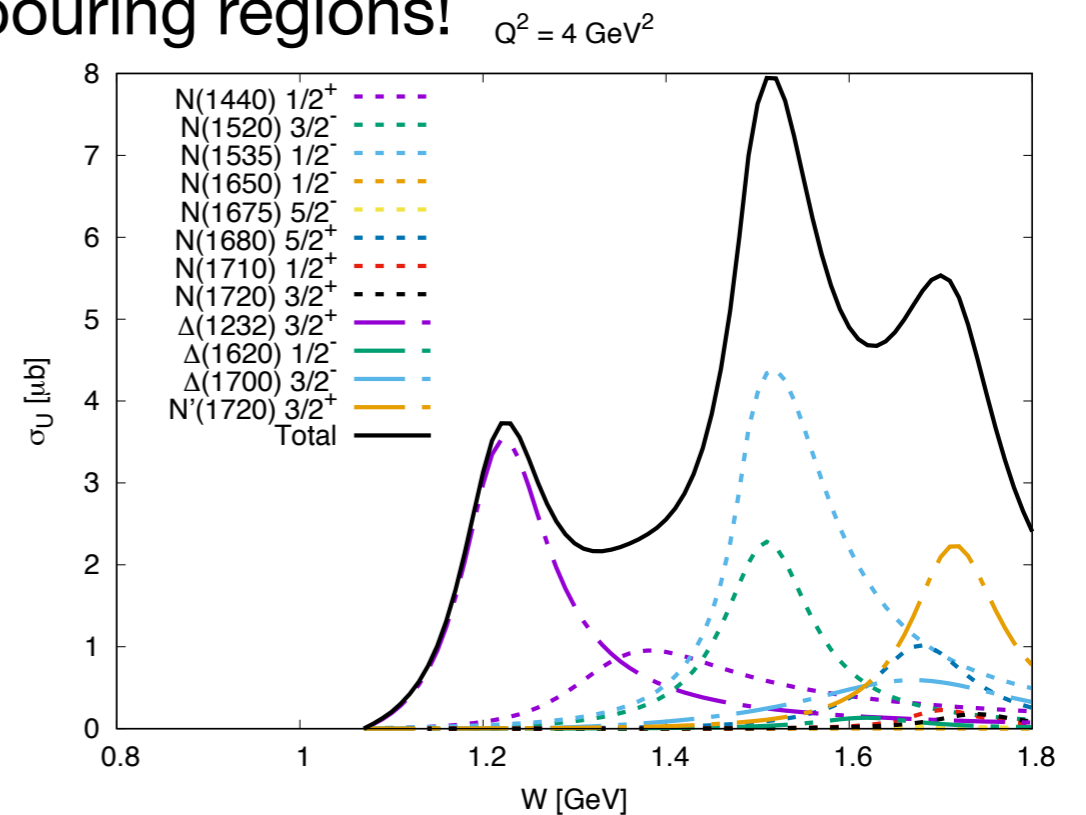
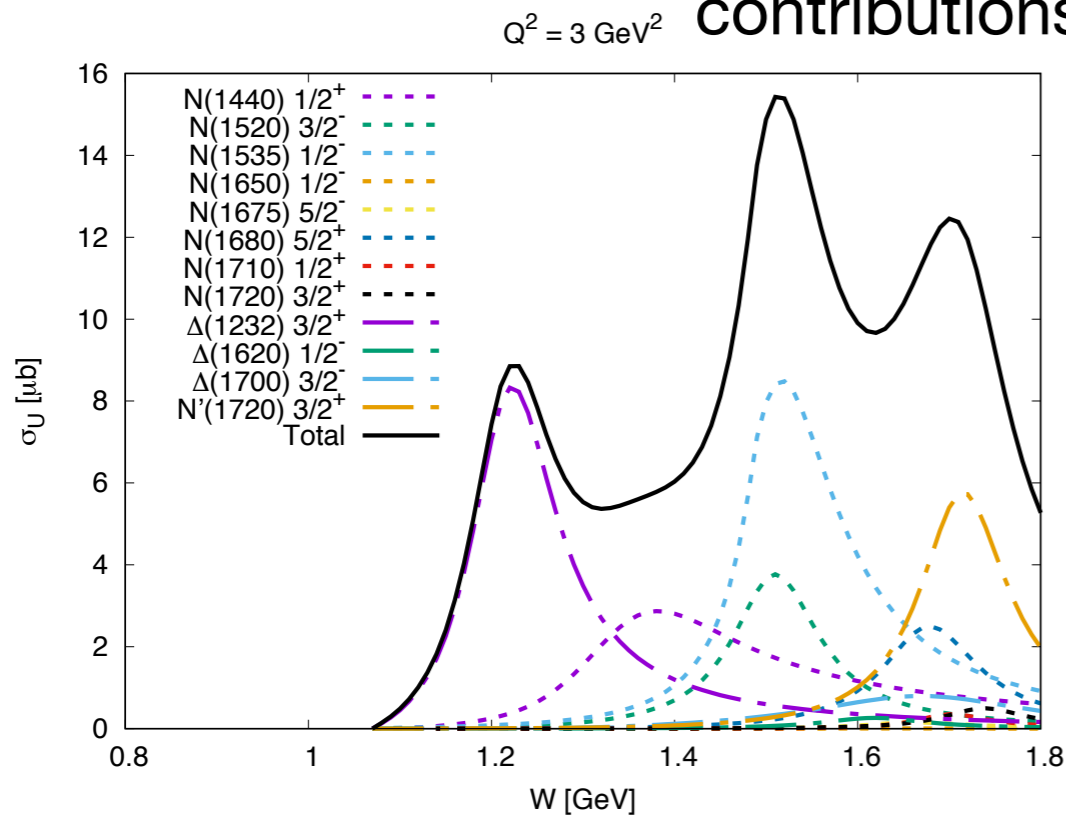
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ANHB et al., PRC100 (2019) 035201

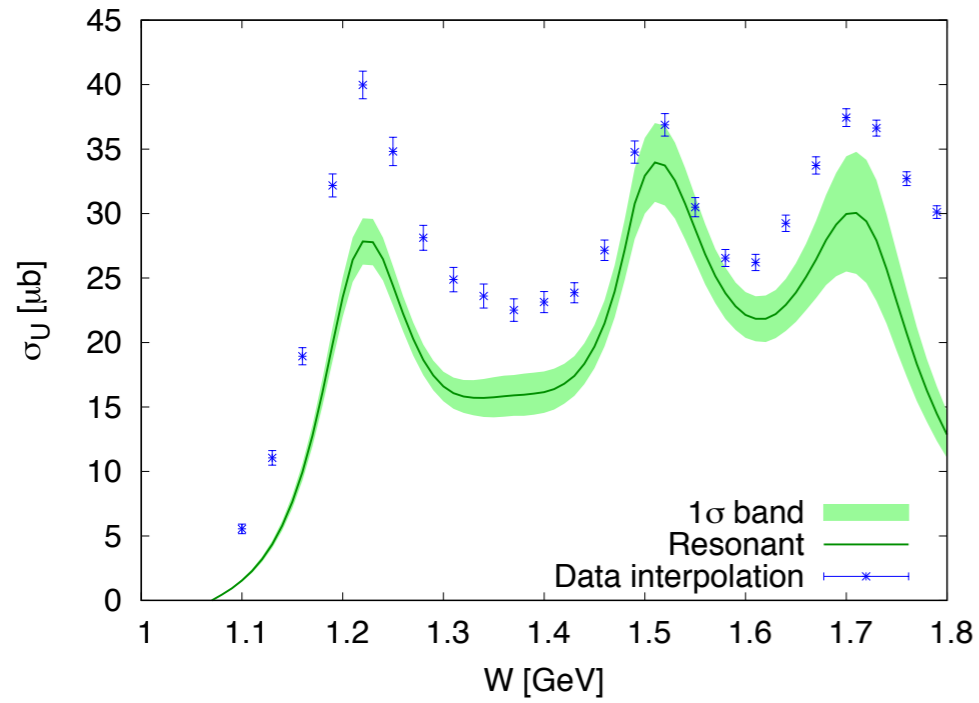


Tails of resonances give substantial contributions to neighbouring regions!

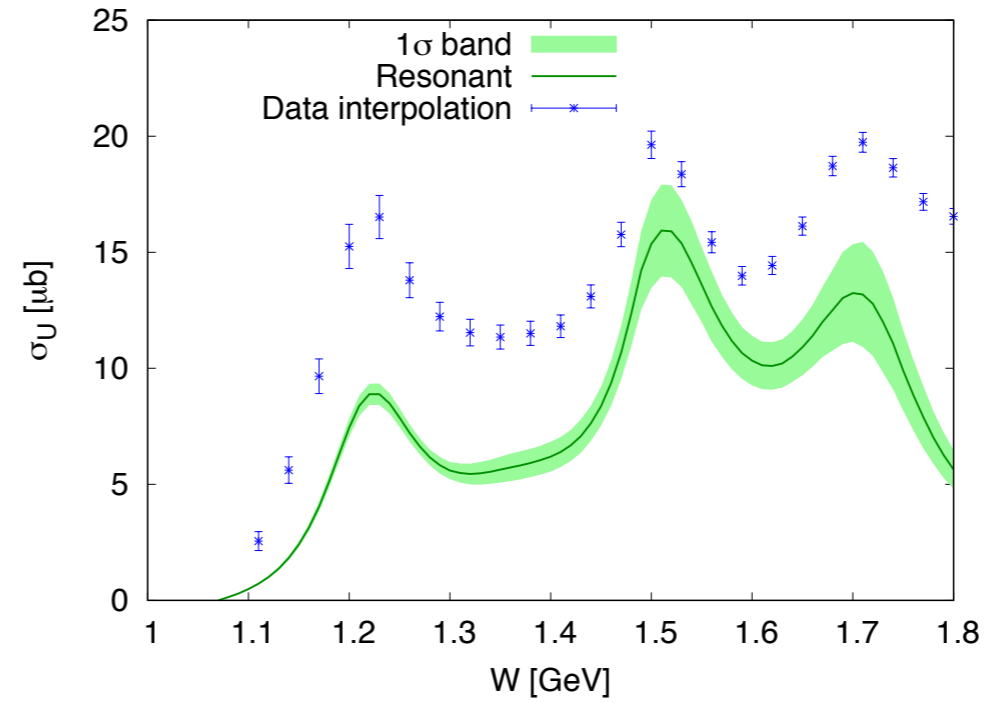


Inclusive unpolarized cross sections

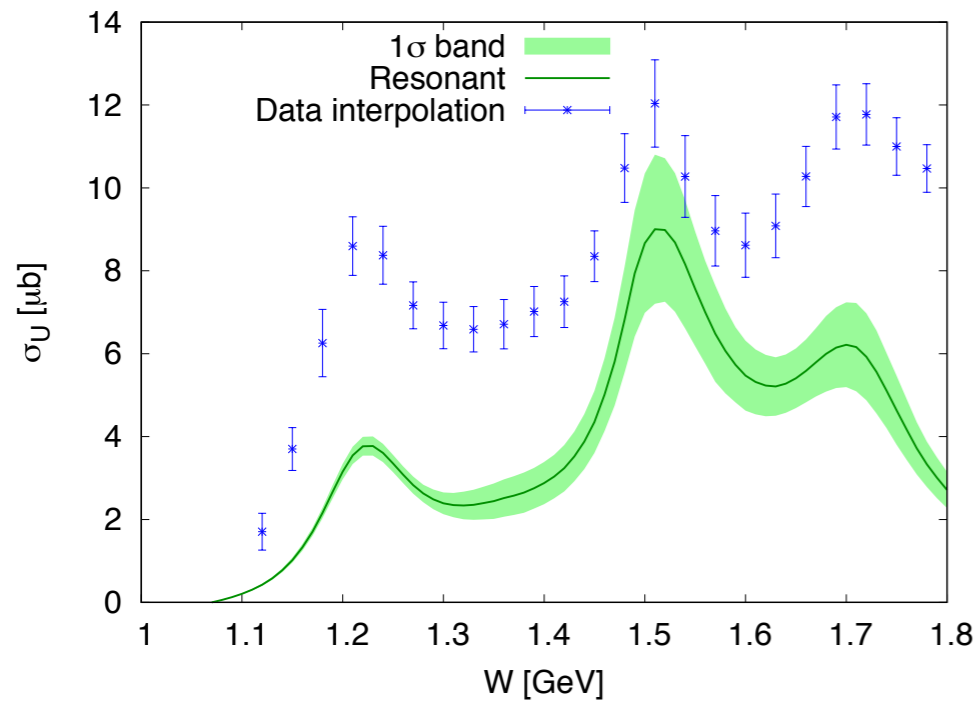
$Q^2 = 2 \text{ GeV}^2$



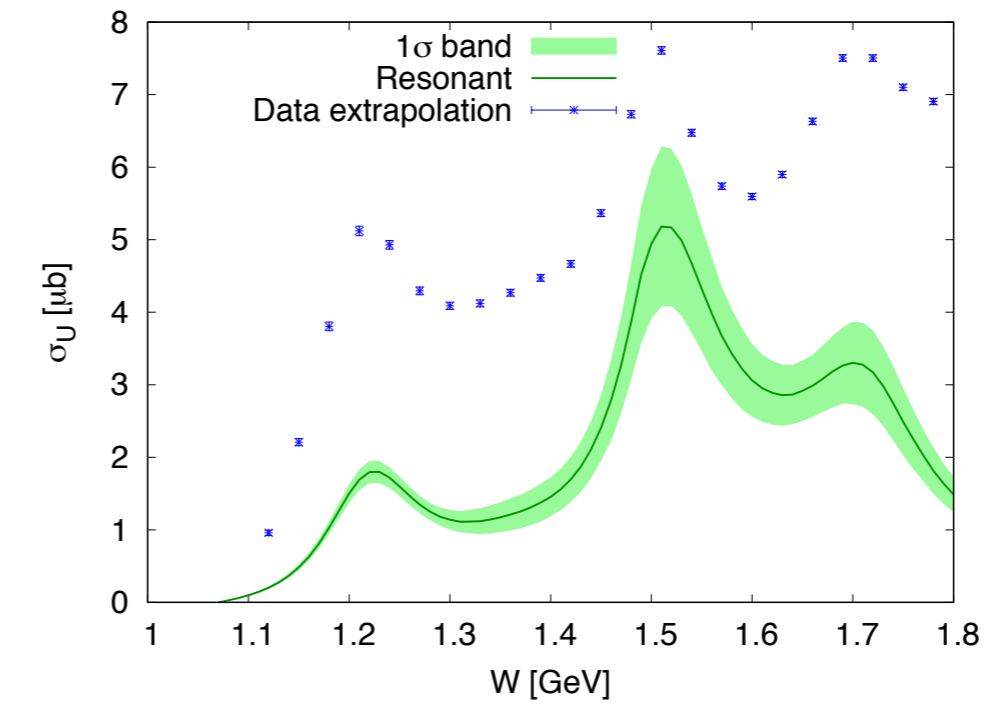
$Q^2 = 3 \text{ GeV}^2$



$Q^2 = 4 \text{ GeV}^2$



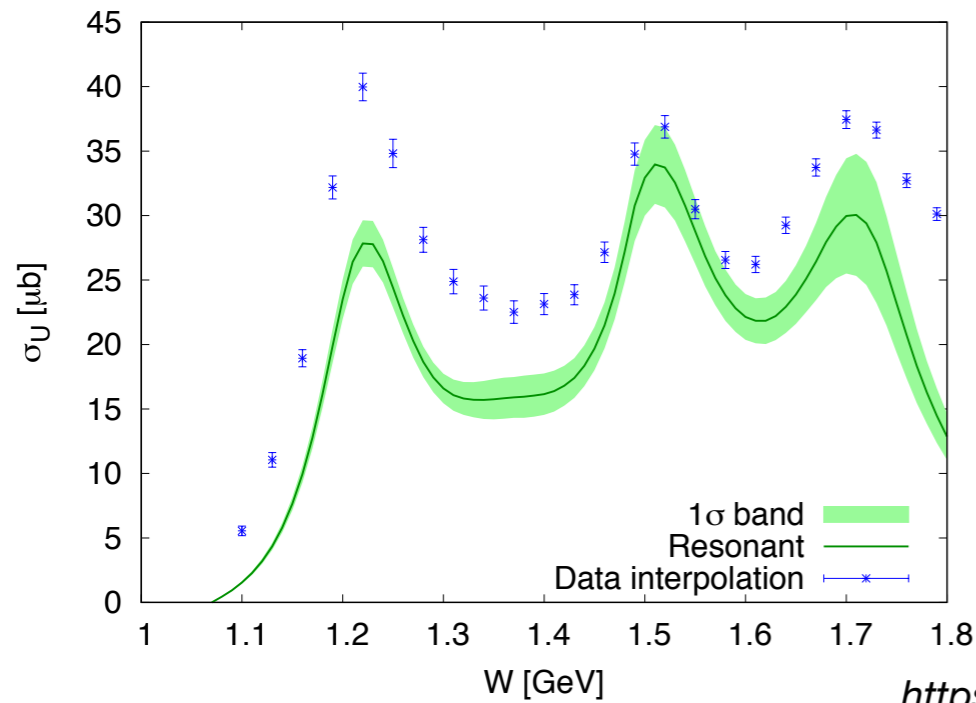
$Q^2 = 5 \text{ GeV}^2$



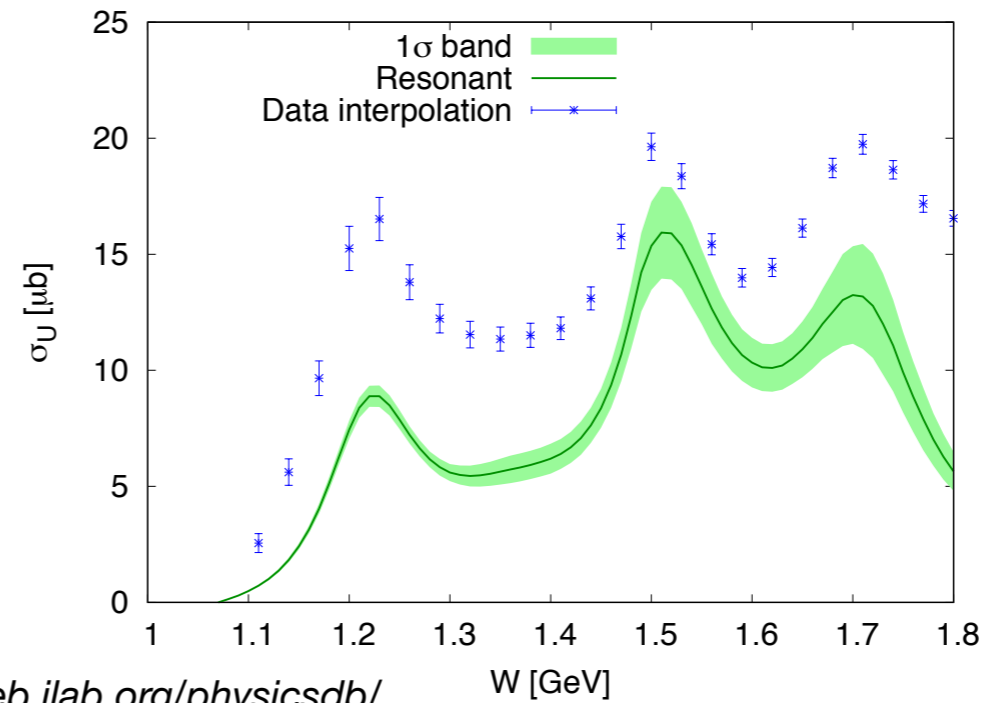
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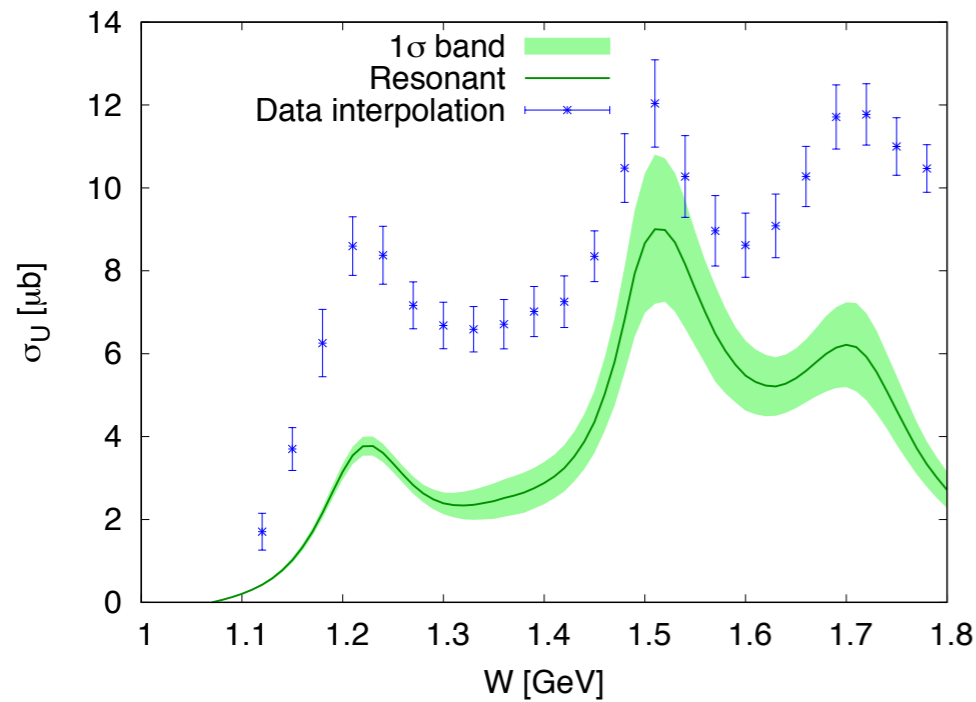
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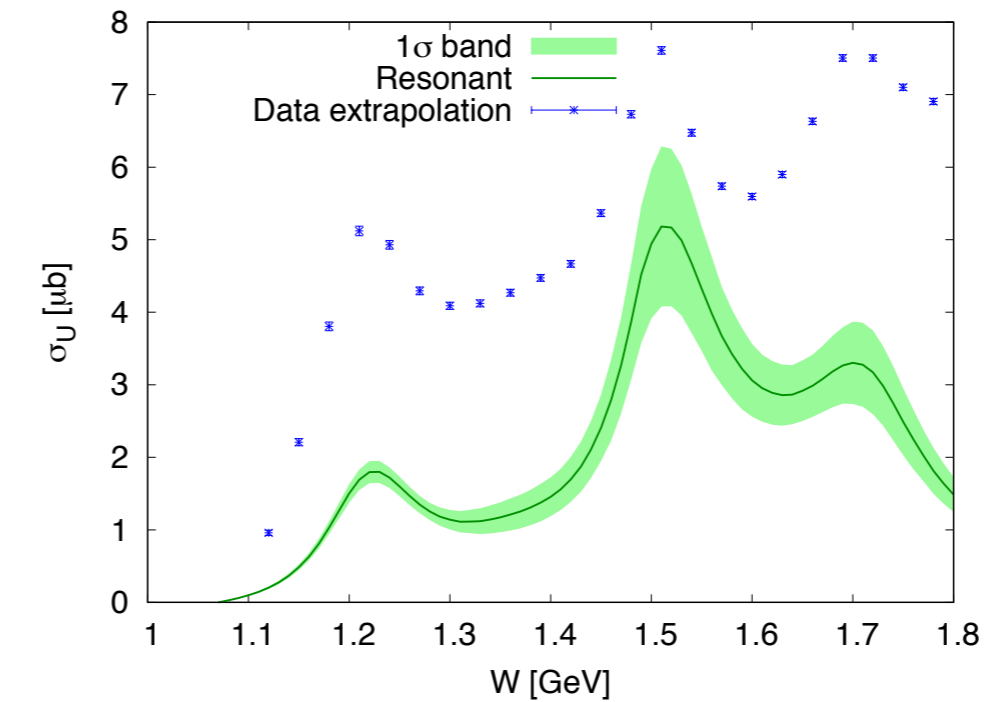
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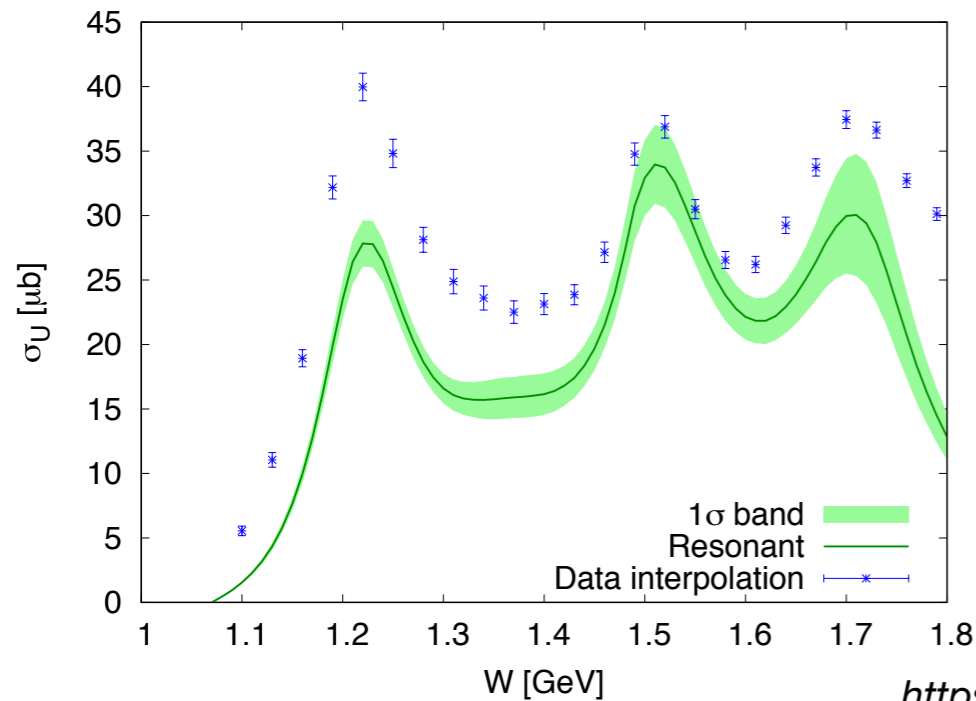
<https://clasweb.jlab.org/physicsdb/>
Golubenko et al., PPN 50 (2019) 587

Second resonance region decreases less with Q^2 :
intricate differences in Q^2 evolution of electrouplings

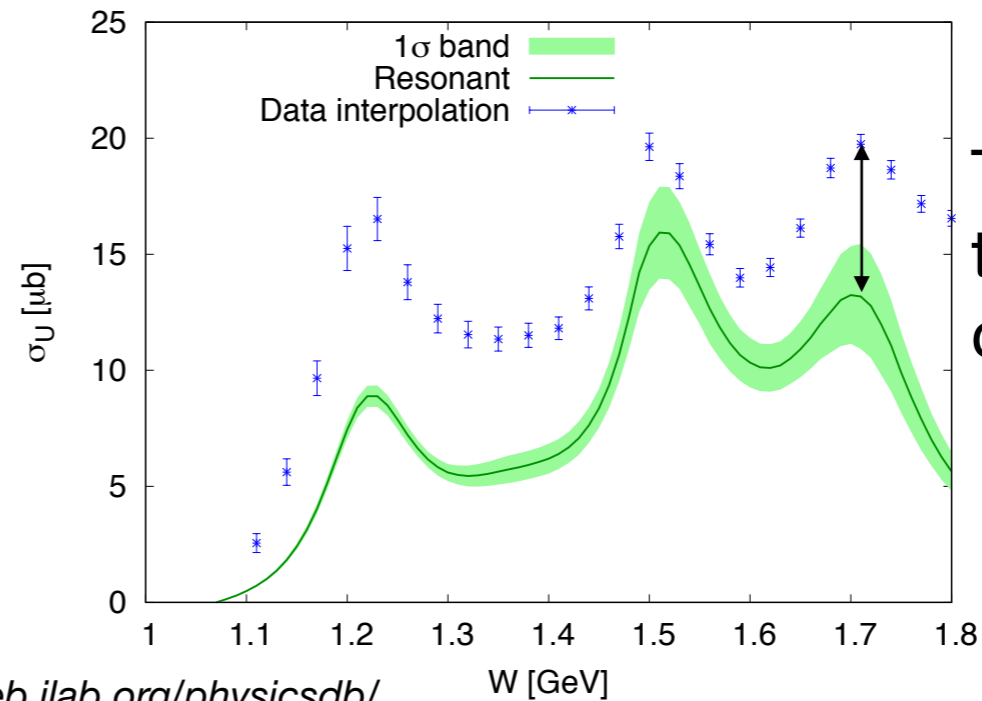
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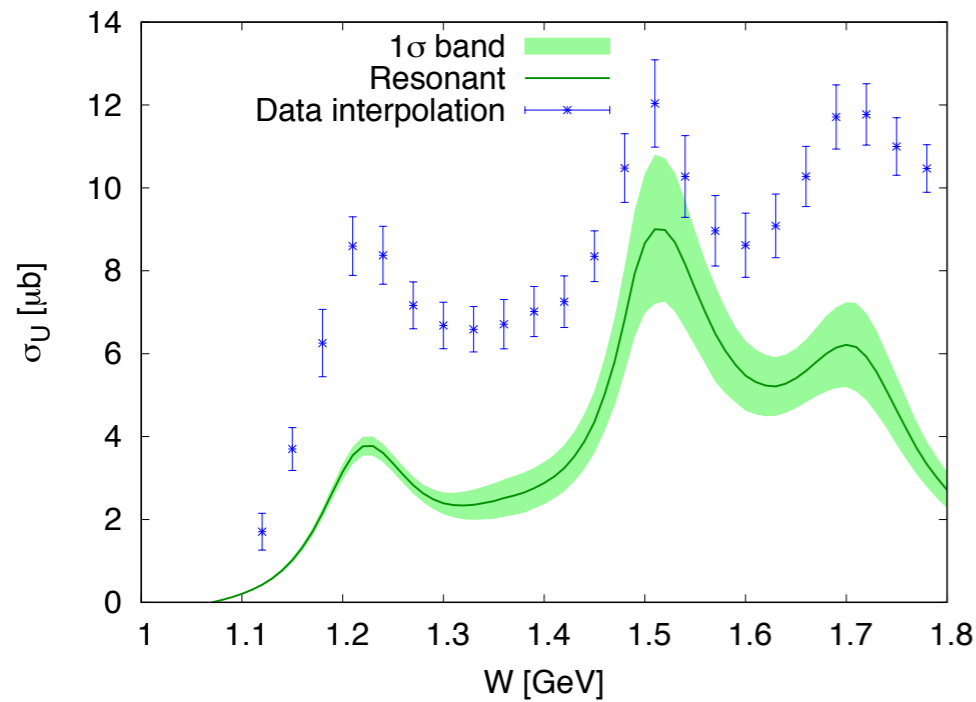
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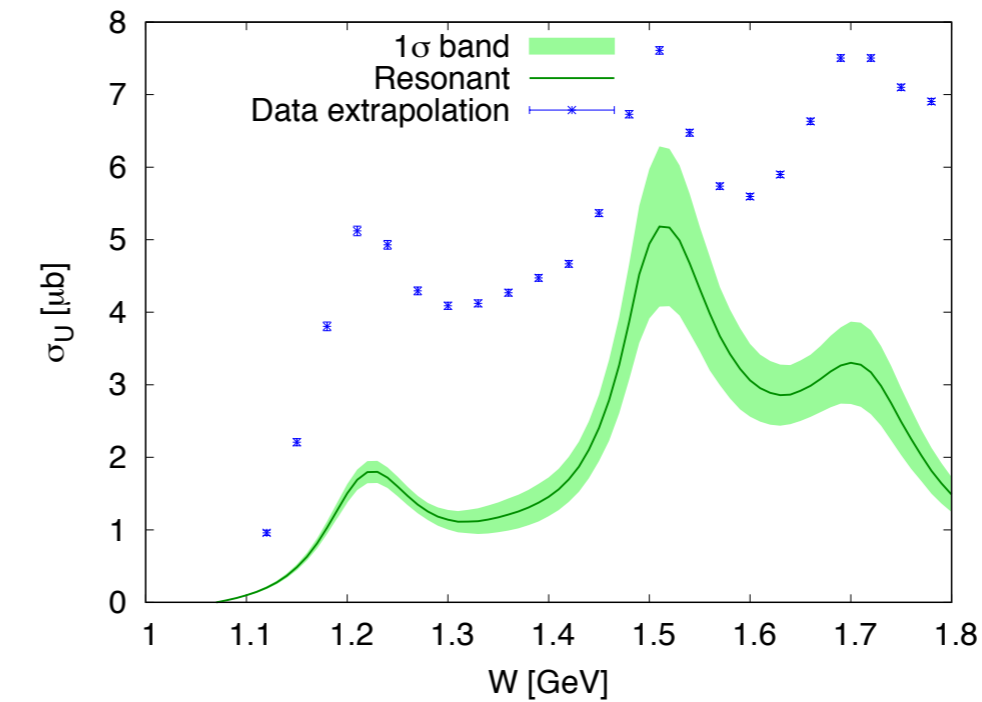
Time to figure out the non-resonant contributions!

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Golubenko et al., PPN 50 (2019) 587

$Q^2 = 4 \text{ GeV}^2$



$Q^2 = 5 \text{ GeV}^2$



ANHB et al., PRC100 (2019) 035201

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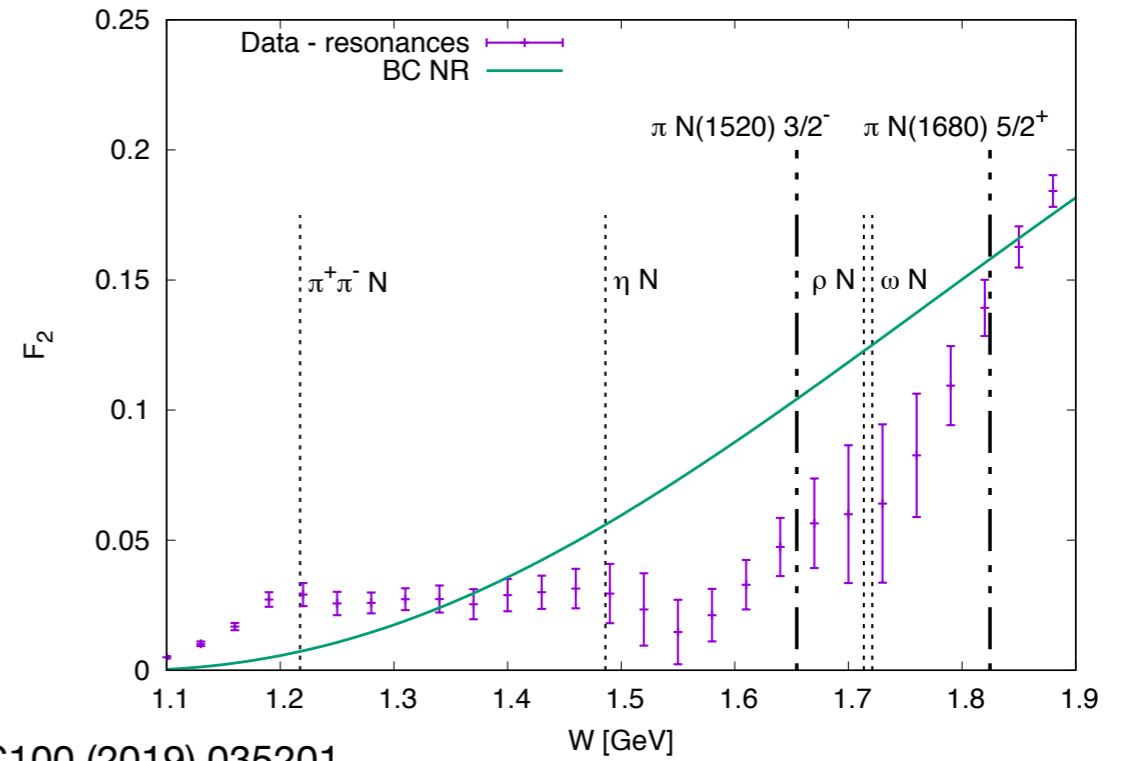
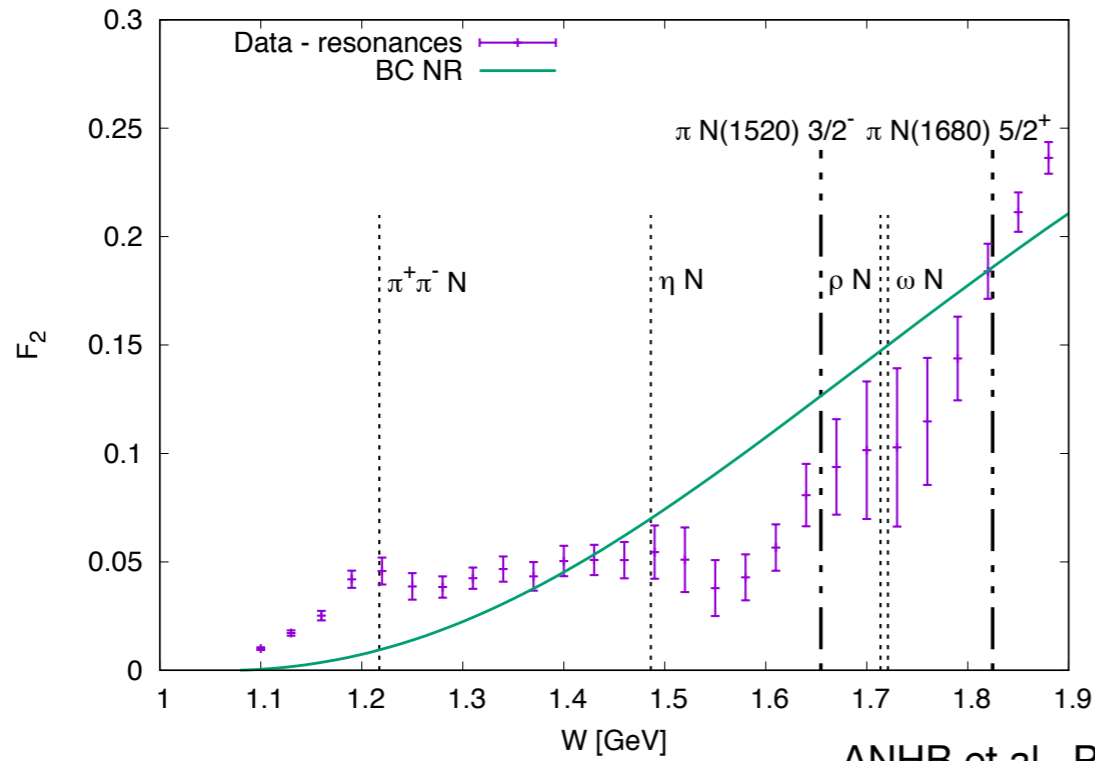
Non-resonant piece: low and high energies communicate

Results of combined description (preliminary)

First estimates of non-resonant contribution

$Q^2=1.5 \text{ GeV}^2$ BC: Christy and Bosted, PRC 86 (2010) 055213

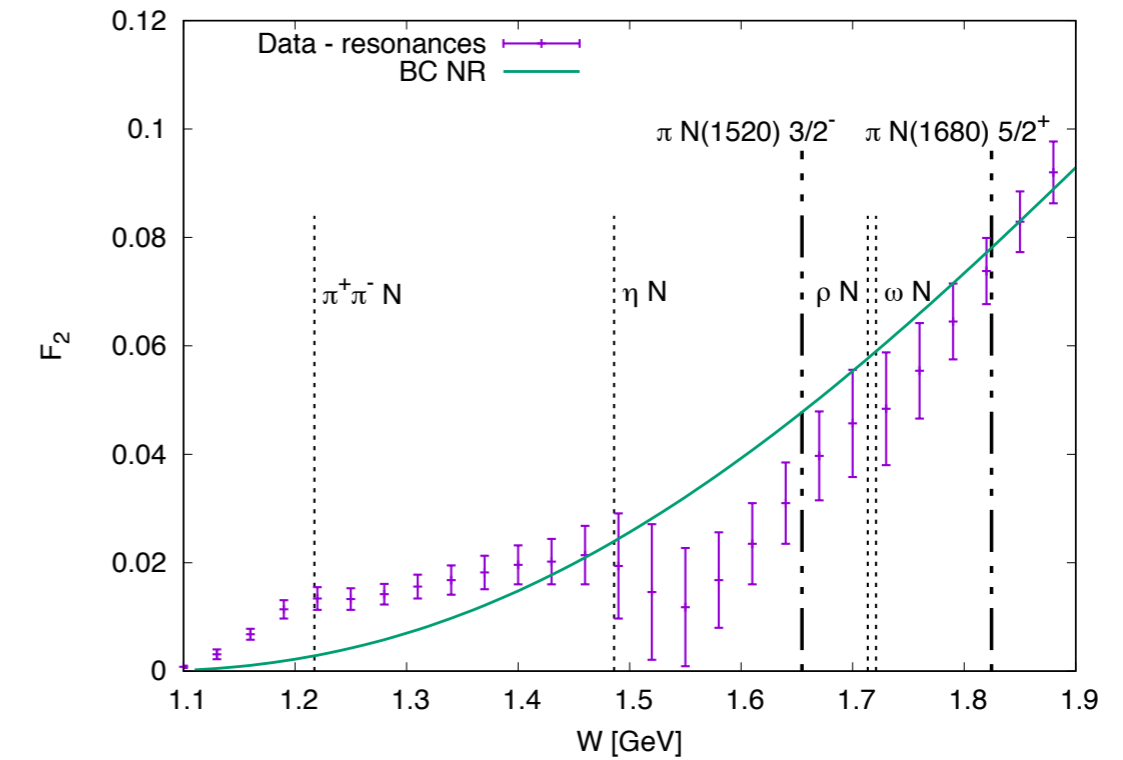
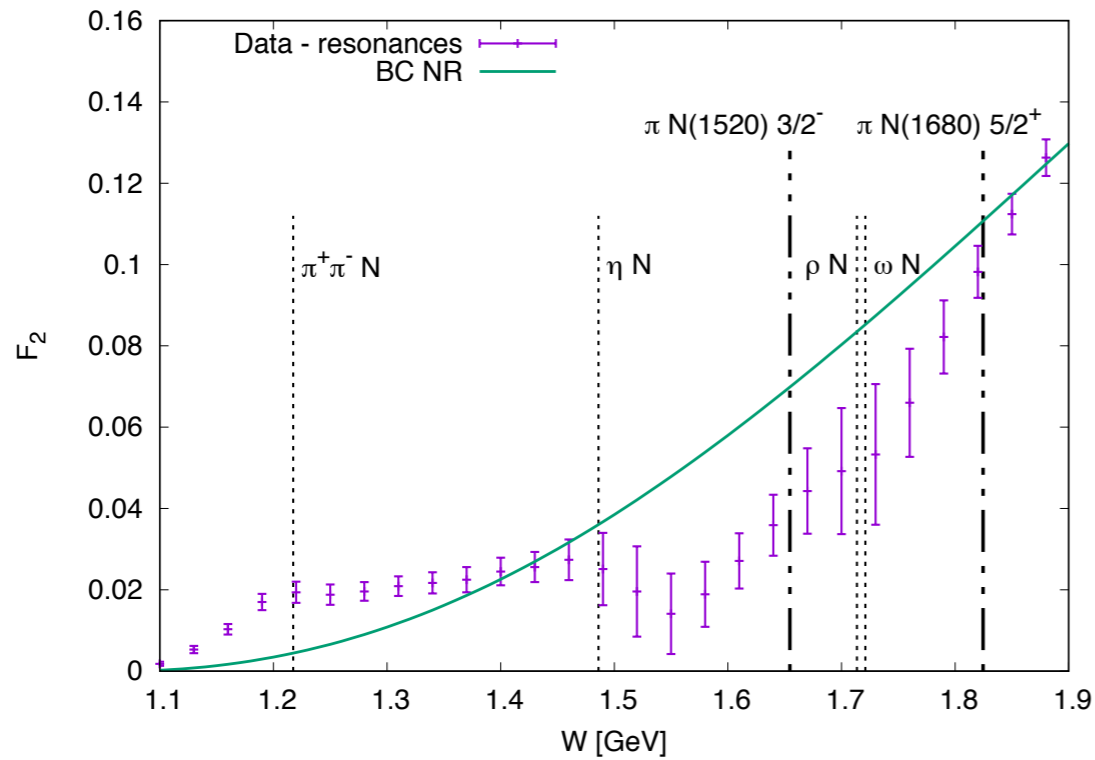
$Q^2=2.0 \text{ GeV}^2$



ANHB et al., PRC100 (2019) 035201

$Q^2=3.0 \text{ GeV}^2$

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Continuation to high energies

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Non-resonant
Regge piece

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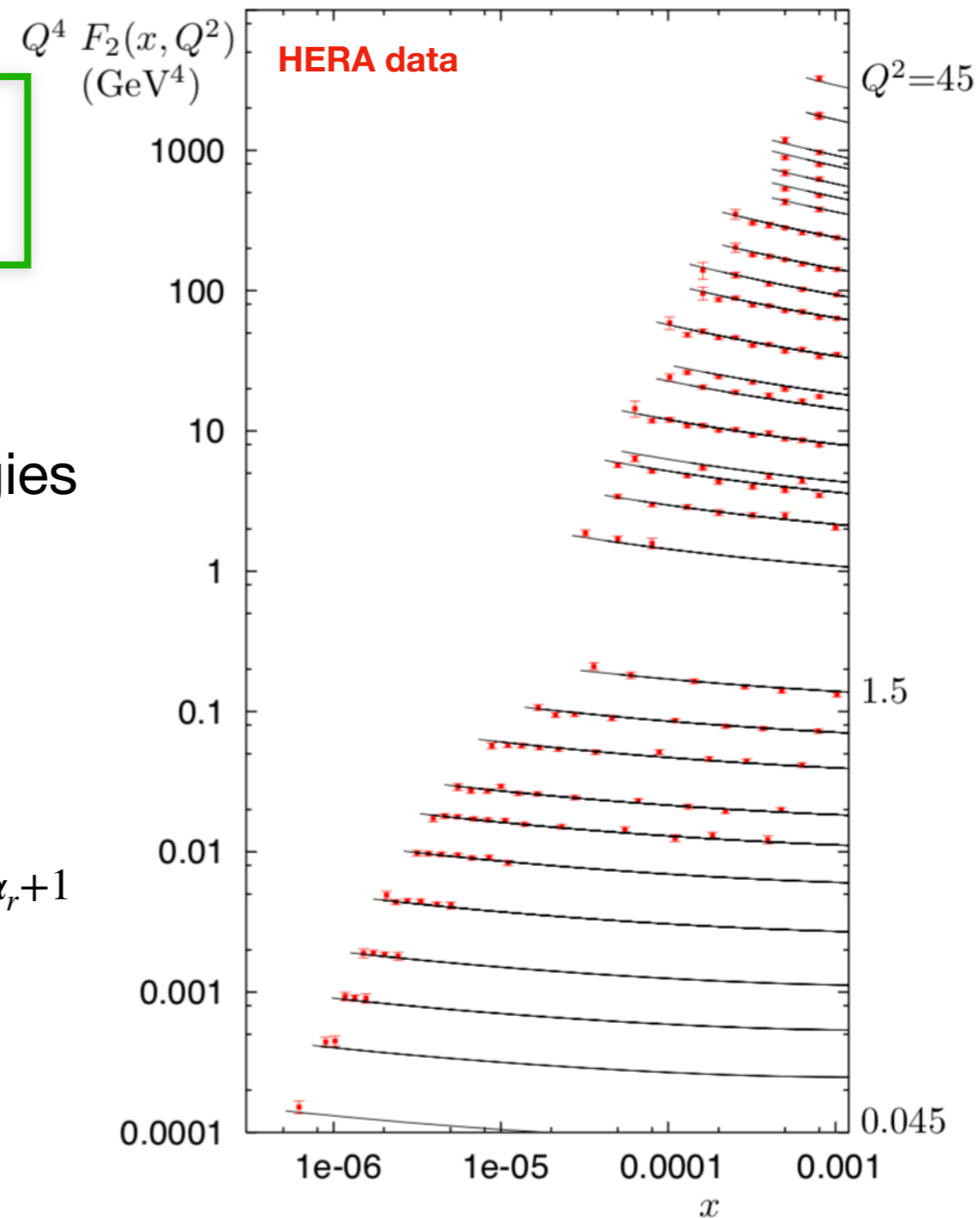
Non-resonant
Regge piece

Excellent fit to data at high energies

Donnachie and Landshoff, PLB 595 (2004) 393

$$F_2(x, Q^2) = f_h(Q^2)x^{-\alpha_h+1} + f_s(Q^2)x^{-\alpha_s+1} + f_r(Q^2)x^{-\alpha_r+1}$$

$$x = \frac{Q^2}{2M_N\nu}$$



Continuation to high energies

$$\sigma_{T,L}(W, Q^2) = \sigma_{T,L}^R(W, Q^2) + \sigma_{T,L}^{NR}(W, Q^2)$$

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Hard pomeron

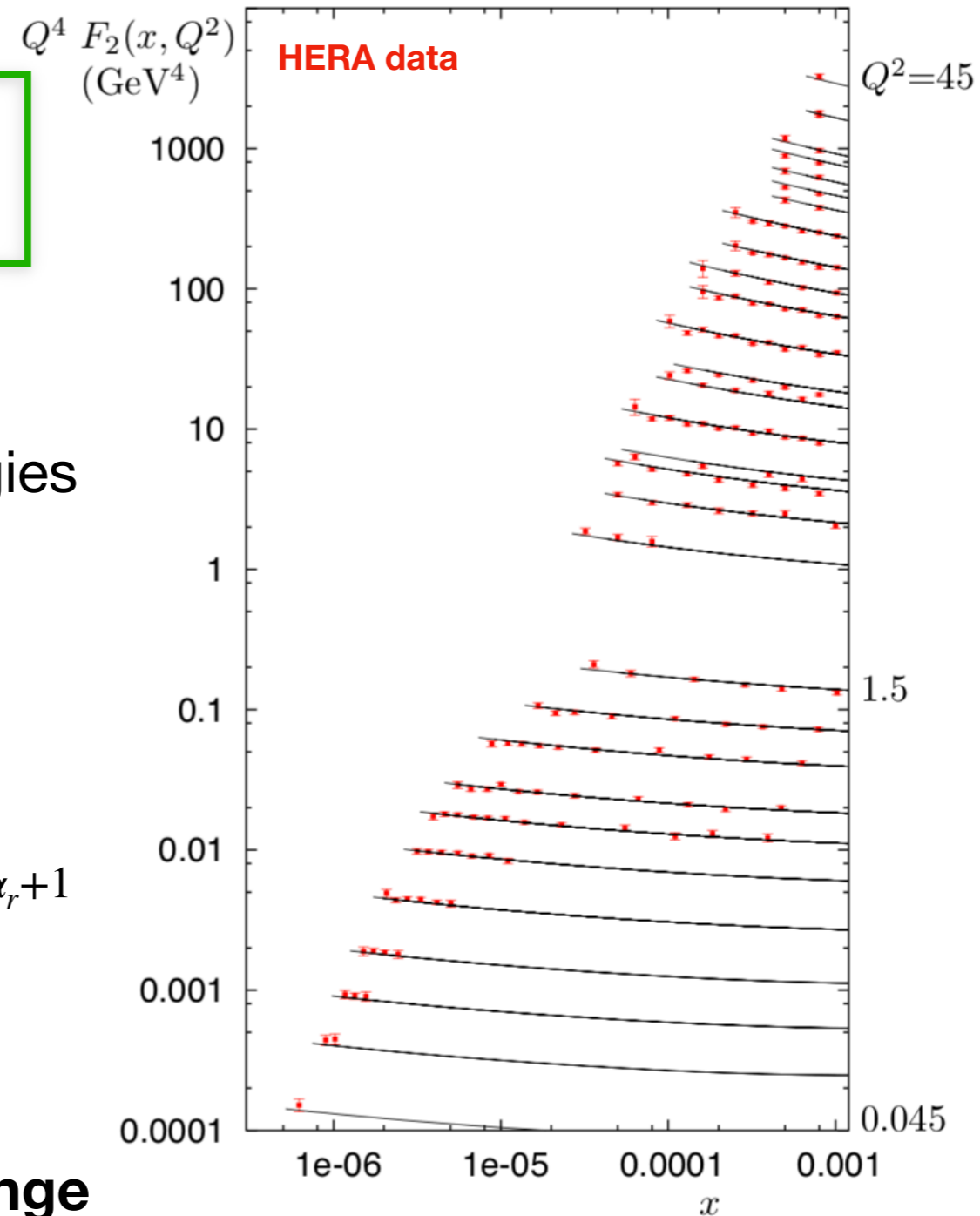
$$\alpha_h = 1.452$$

Soft pomeron

$$\alpha_s = 1.0667$$

Meson exchange

$$\alpha_r = 0.524$$



High and low energies: previous works

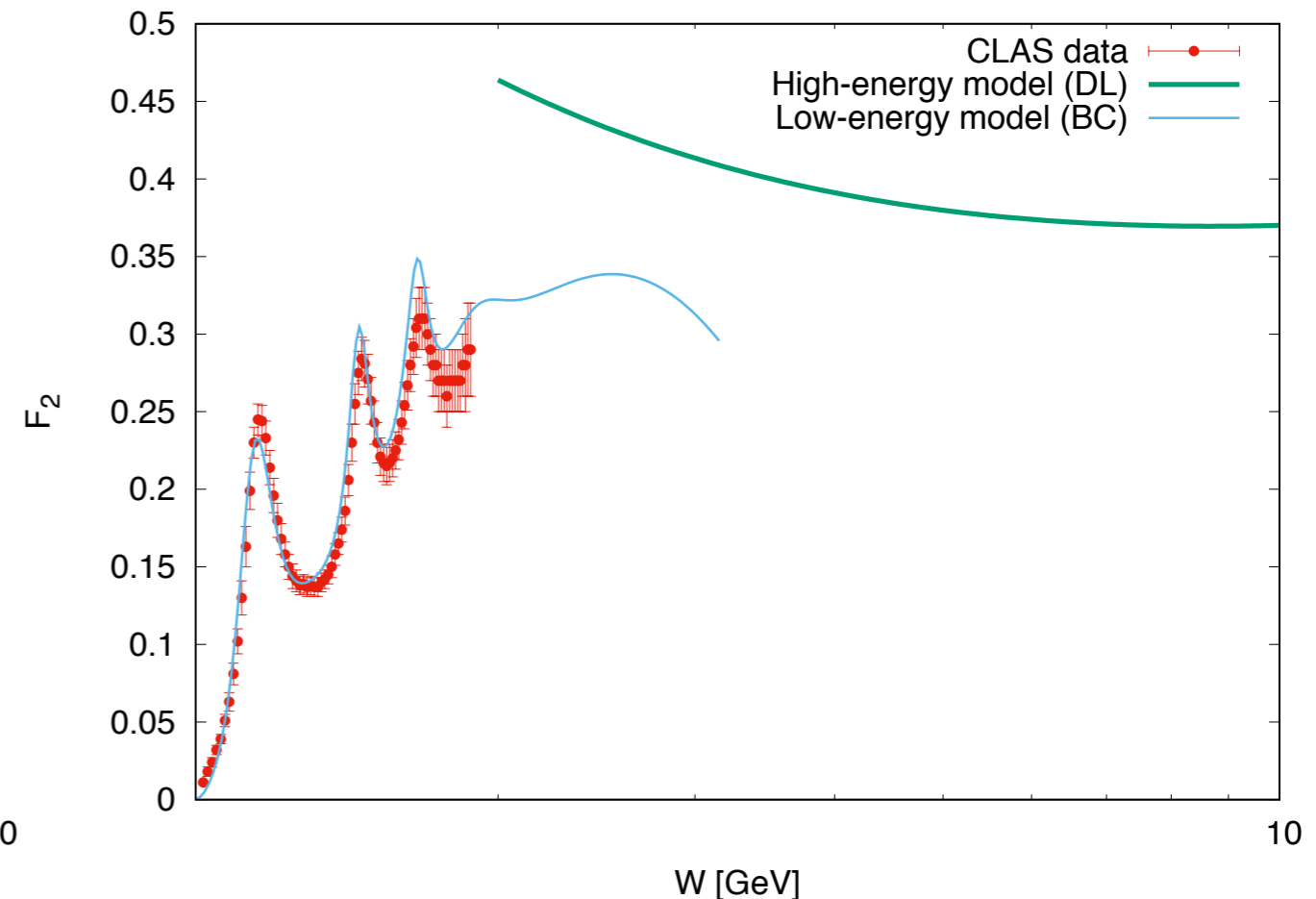
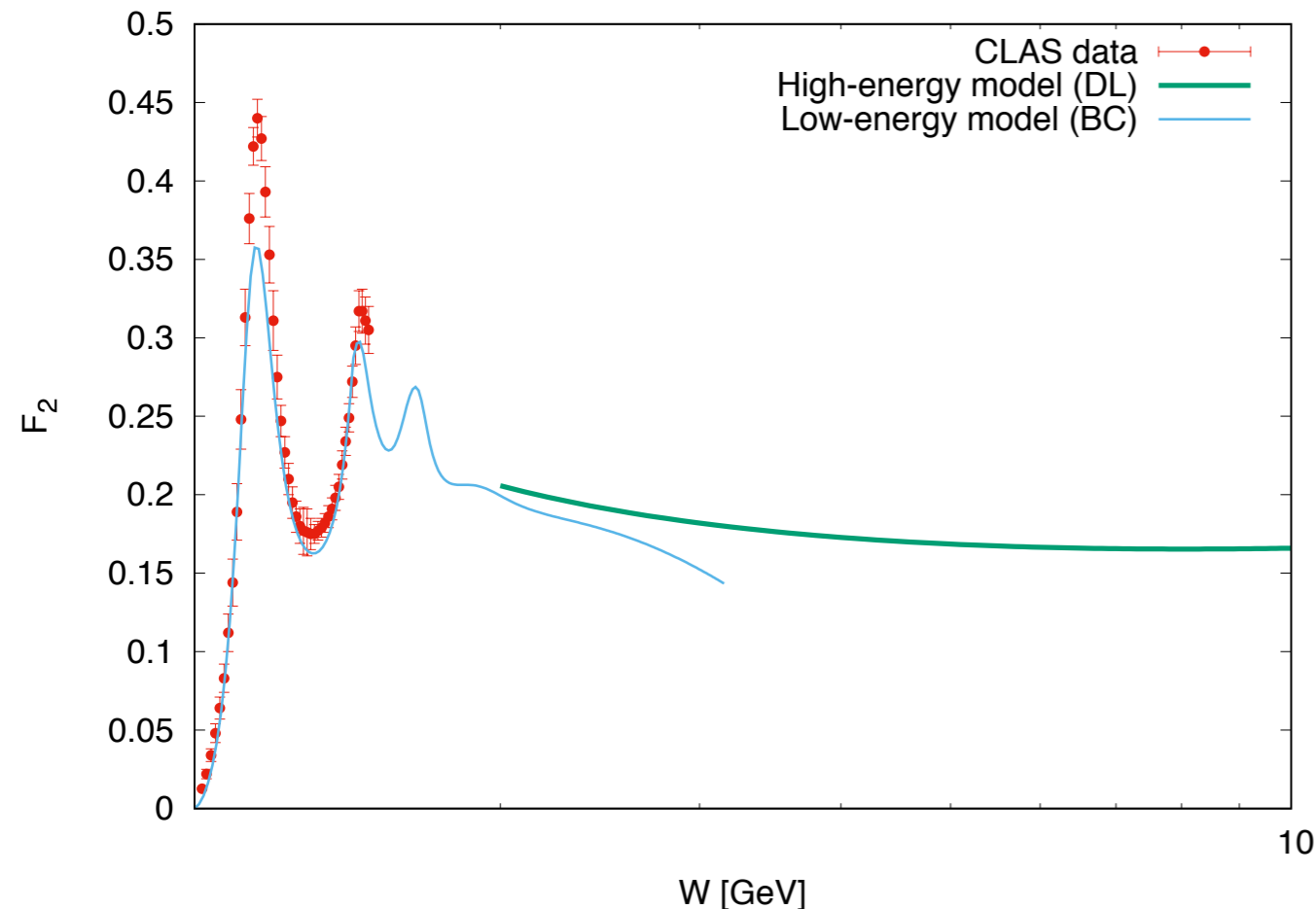
BC: Christy and Bosted, PRC 86 (2010) 055213

DL: Donnachie and Landshoff, PLB 595 (2004) 393

CLAS data: <http://clas.sinp.msu.ru/strfun/>

$Q^2=0.23 \text{ GeV}^2$

$Q^2=1.0 \text{ GeV}^2$



- Excellent description of data at low and high energies, separately
- At low Q^2 : high- and low-energy theories compatible in overlap region
- At slightly **higher Q^2** : **huge gap** between high and low-energy models
- Leads to difficulties in extracting observables integrated over x (or energies)

Proposal for updated parametrization

$$F_1(\nu, Q^2) = F_1^{\text{res}}(\nu, Q^2) + \sum_{i=0}^2 \gamma_{\alpha_i}(Q^2) (\nu - \nu_{\text{thr}})^{\alpha_i} \left(1 - \frac{\nu_{\text{thr}}}{\nu}\right)^{a_i(Q^2)} \left(1 + \frac{\nu_{\text{thr}}}{\nu}\right)^{b_i(Q^2)}$$

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- Automatically recovers **Regge behaviour** at high energies for $\alpha_i < 1$

$$F_1(\nu, Q^2) \longrightarrow \sum_{i=0}^2 \gamma_{\alpha_i}(Q^2) \nu^{\alpha_i}$$

- For the Pomerons we have the **requirement**

$$b_i(Q^2) = a_i(Q^2) + \alpha_i, \quad 1 \leq \alpha_i < 2$$

- Implements **threshold** and **resonant** behaviour

Proposal for updated parametrization

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- Automatically recovers **Regge behaviour** at high energies for $\alpha_i < 1$

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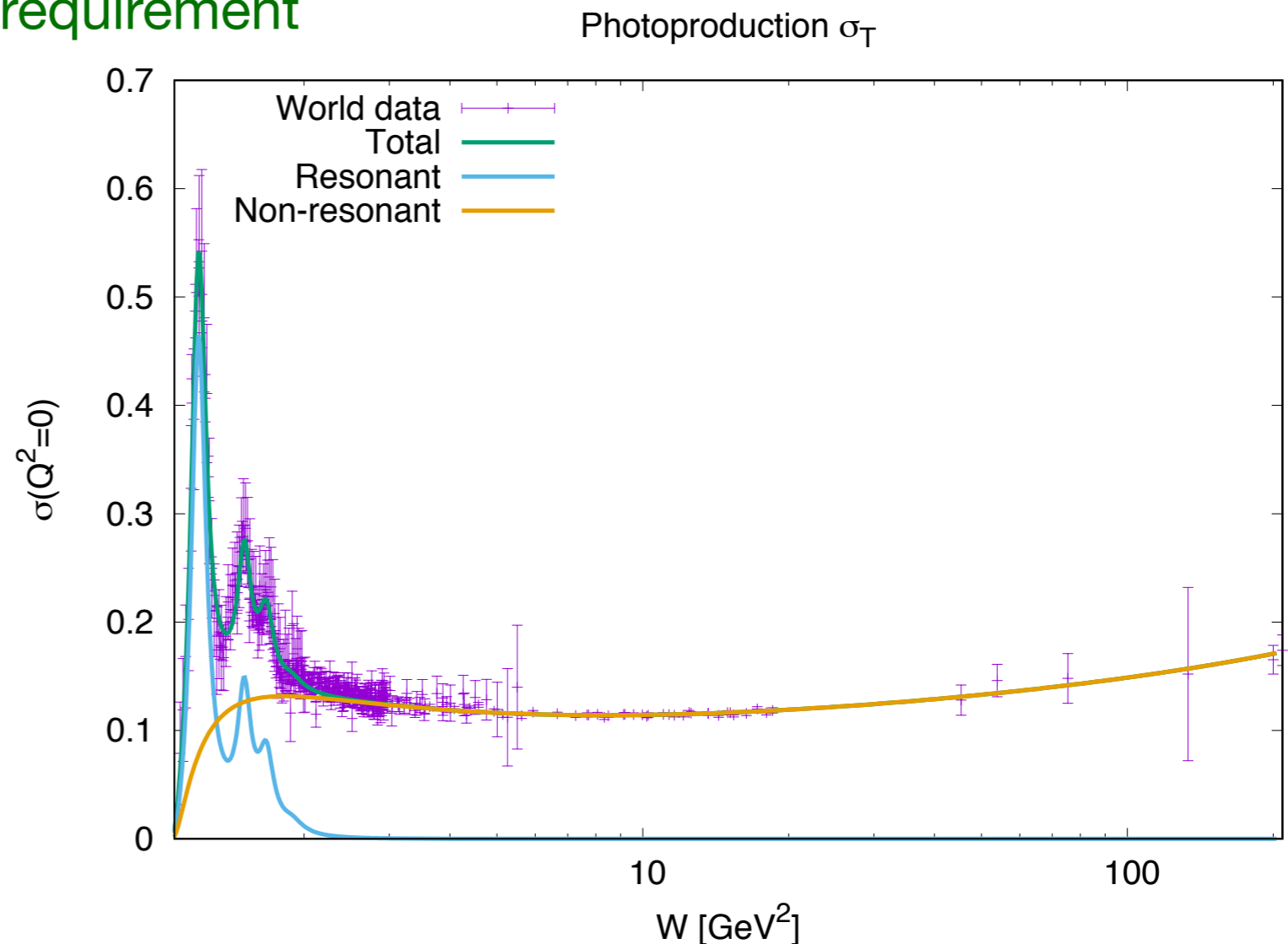
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Dispersion relation and Lamb shift

VVCS amplitude T_1 obeys once-subtracted dispersion relation

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Source of the largest uncertainty of hadronic contribution to Lamb shift
relevant to improve muonic spectroscopy in proton charge radius extraction

Birse and Govern, EPJ A 48 (2012) 120

Alarcón et al., EPJ C 74 (2014) 2852

Tomalak and Vanderhaeghen, EPJ C 76 (2016) 125

Fit constraints

$$F_1(\nu, Q^2) = F_1^{\text{res}}(\nu, Q^2) + \sum_{i=0}^2 \gamma_{\alpha_i}(Q^2) (\nu - \nu_{\text{thr}})^{\alpha_i} \left(1 - \frac{\nu_{\text{thr}}}{\nu}\right)^{a_i(Q^2)} \left(1 + \frac{\nu_{\text{thr}}}{\nu}\right)^{b_i(Q^2)}$$

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- $F_1 \sim$ imaginary part of T_1 :
additional constraints obtained by imposing for subtraction function

$$\bar{T}_1(\mathbf{0}, \mathbf{0}) = \mathbf{0}$$

and using the experimental value of the magnetic polarizability:

$$\beta_{M1} = \left. \frac{d}{dQ^2} \bar{T}_1(0, Q^2) \right|_{Q^2=0}$$

- Constrains the fit parameter space even further

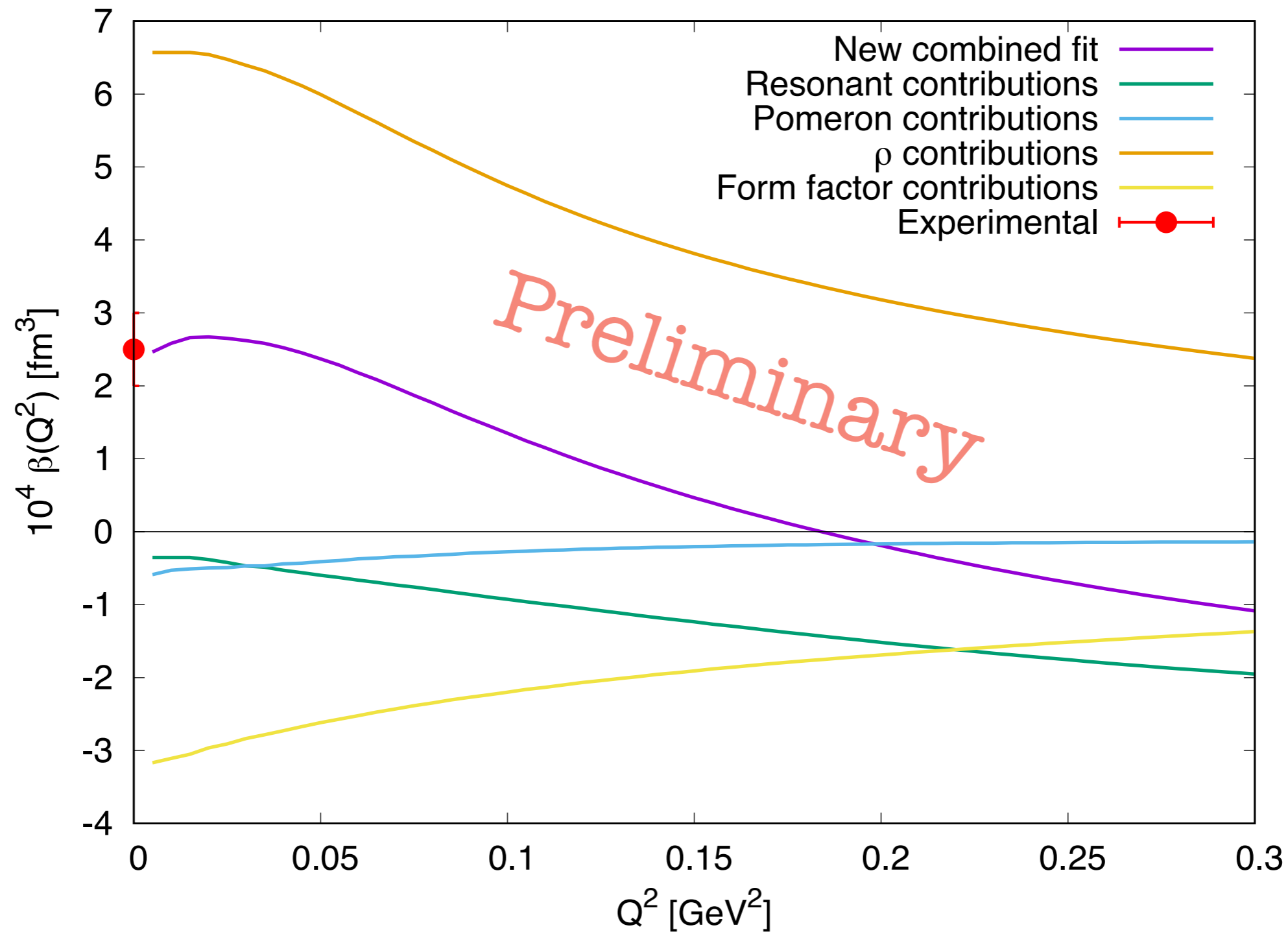
Inclusive and exclusive reactions: can we use one for the other?

Results of resonant contributions to inclusive observables

Non-resonant piece: low and high energies communicate

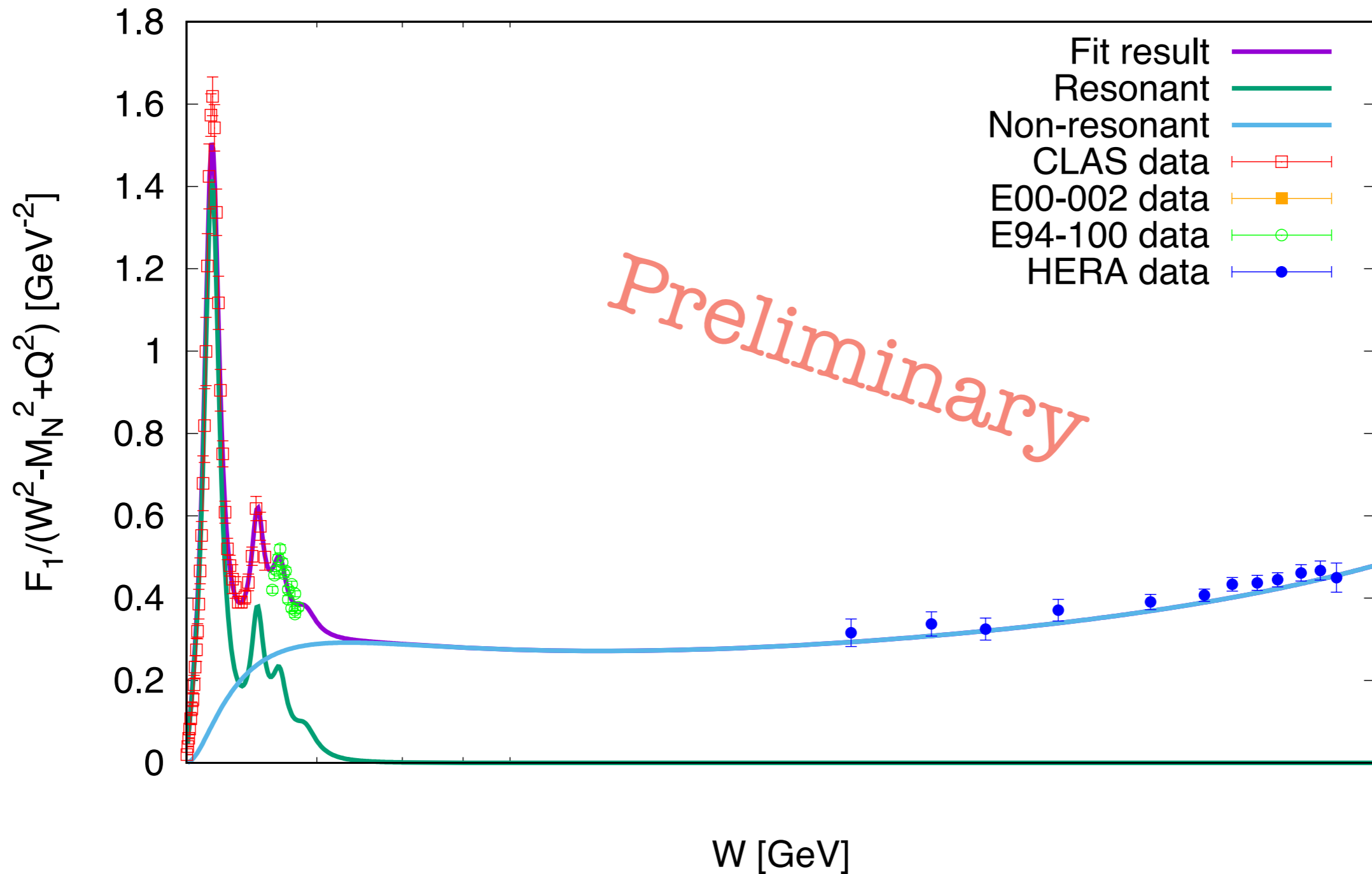
Results of combined description (preliminary)

Magnetic polarizability: our fit results



Structure functions at non-vanishing Q^2

$$0.25 \leq Q^2 < 0.3 \text{ GeV}^2$$



First combined model for low and high energies!

Summary

- New resonance model: **N* electrocouplings** from CLAS(12) allow to describe the **resonant contributions** to inclusive electron-scattering
- Intricate behaviour with W and Q^2 in the resonance regime
- First estimates for the non-resonant behaviour

Summary

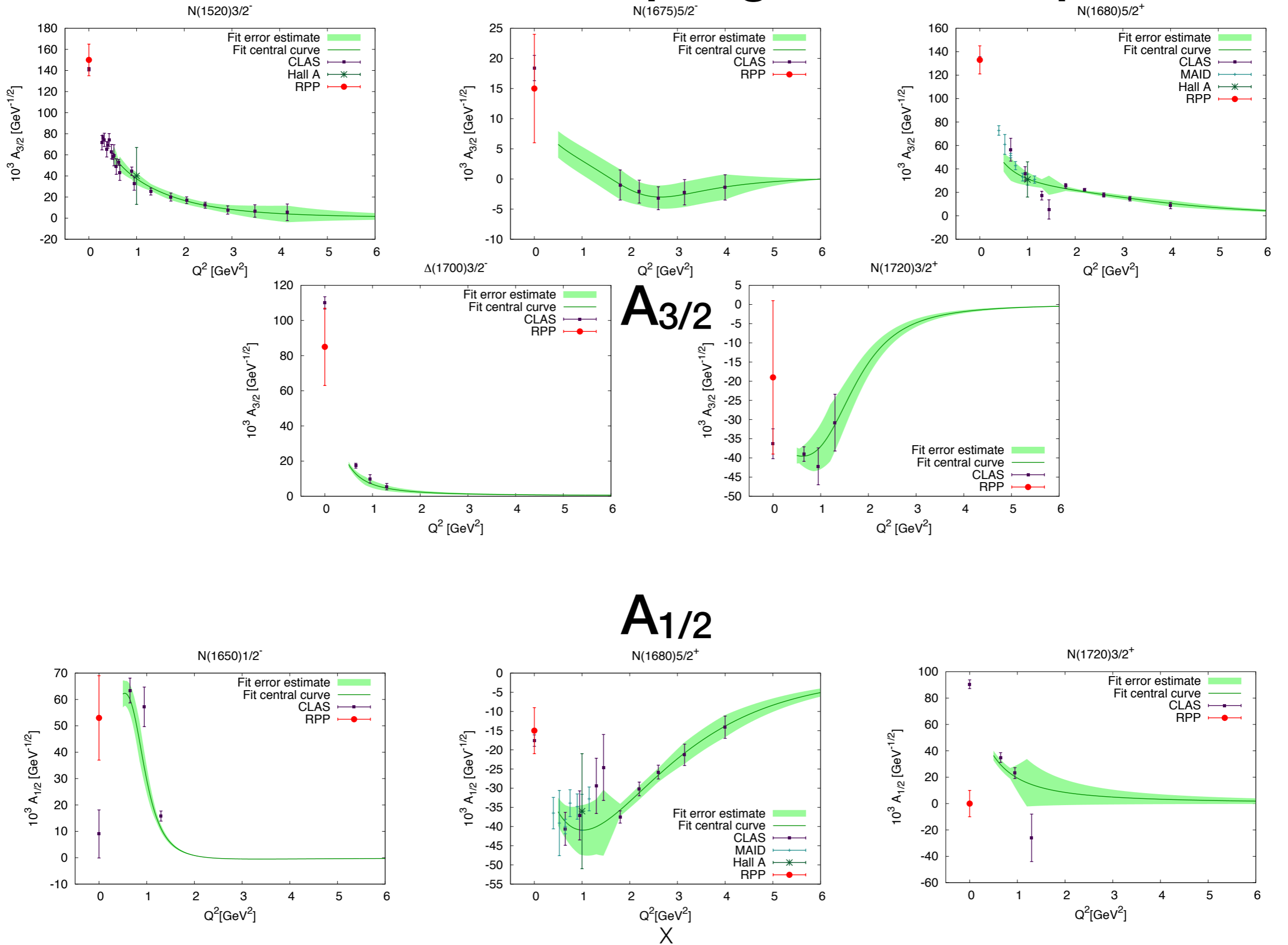
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Work in progress

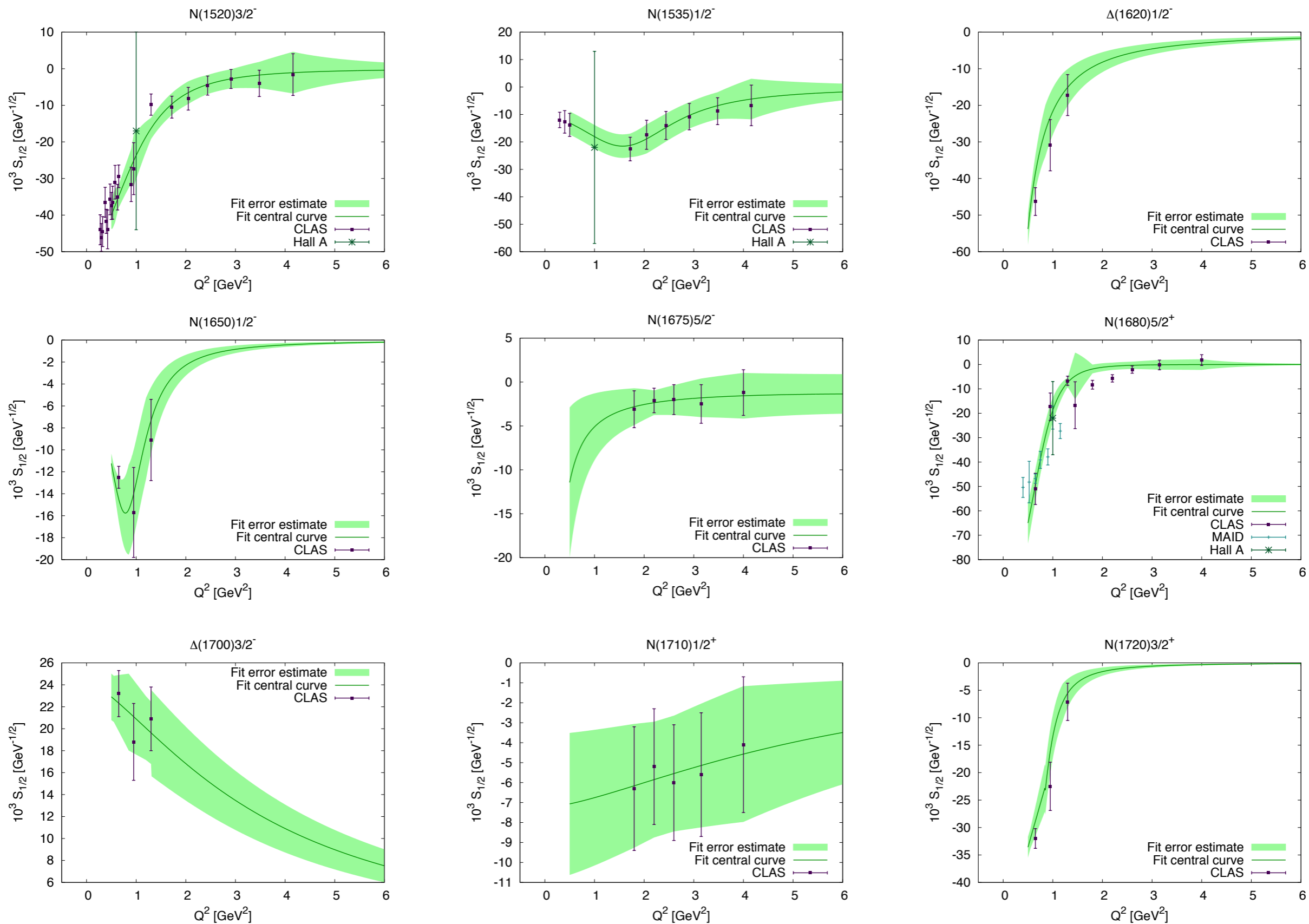
- Analytical transition between low and high x (or energies): important for VVCS subtraction function
- First phenomenological determination of subtraction contribution to Lamb shift
- Updates on inclusive and exclusive electron scattering CLAS12: coming soon!

Backup

Further electrocoupling data examples

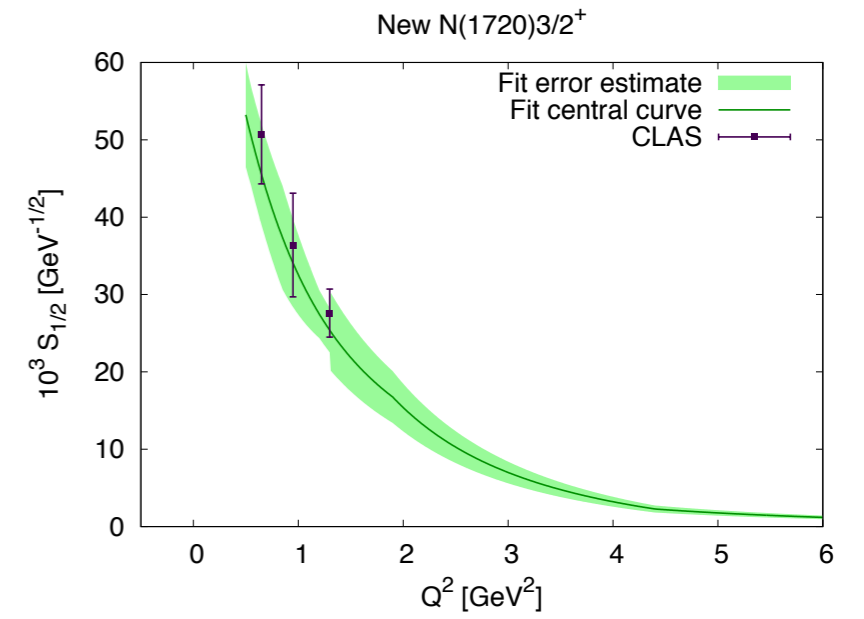
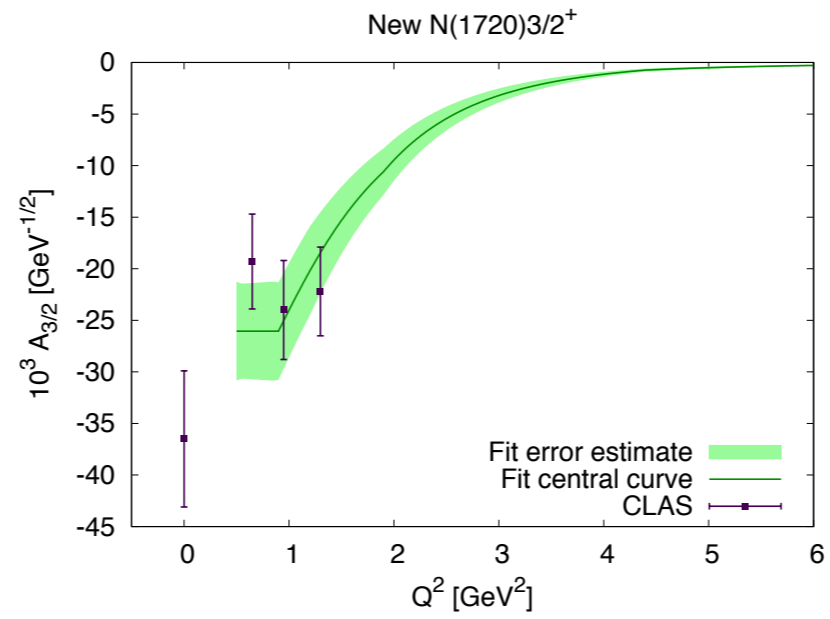
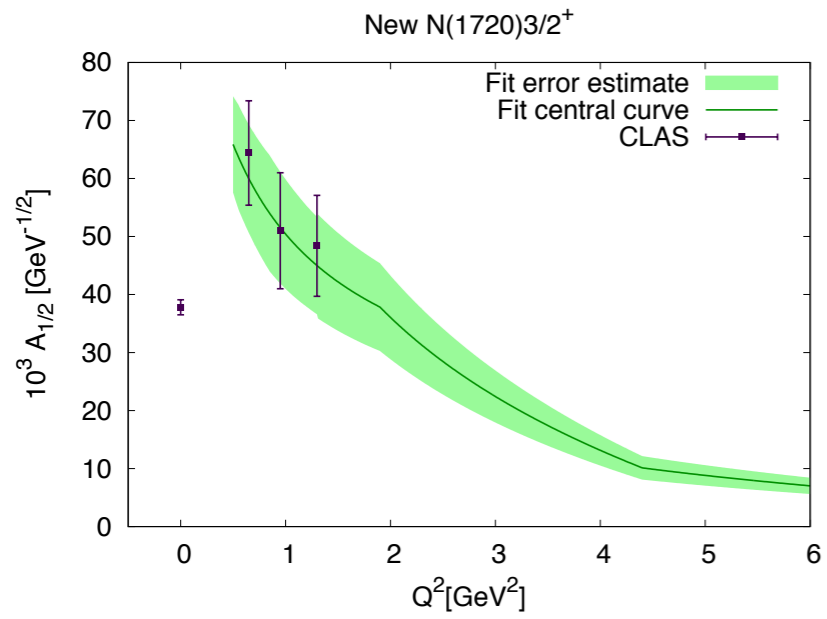


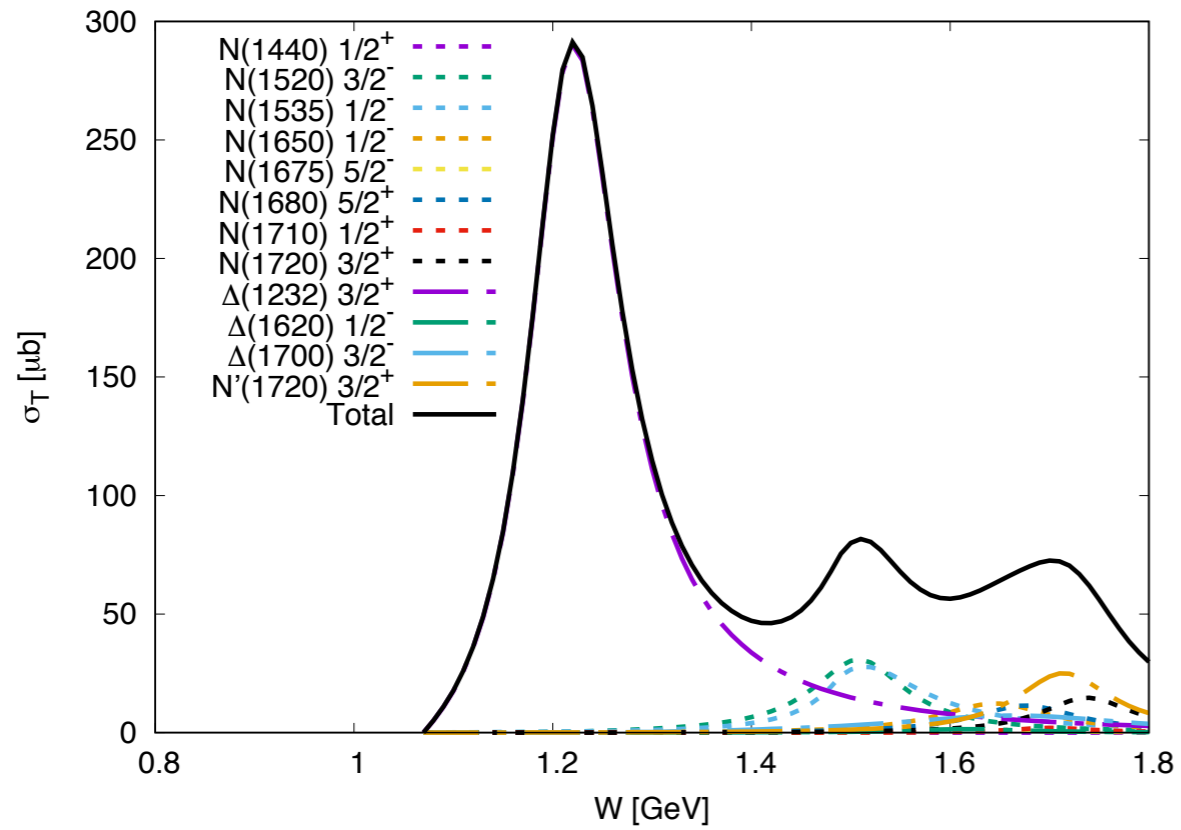
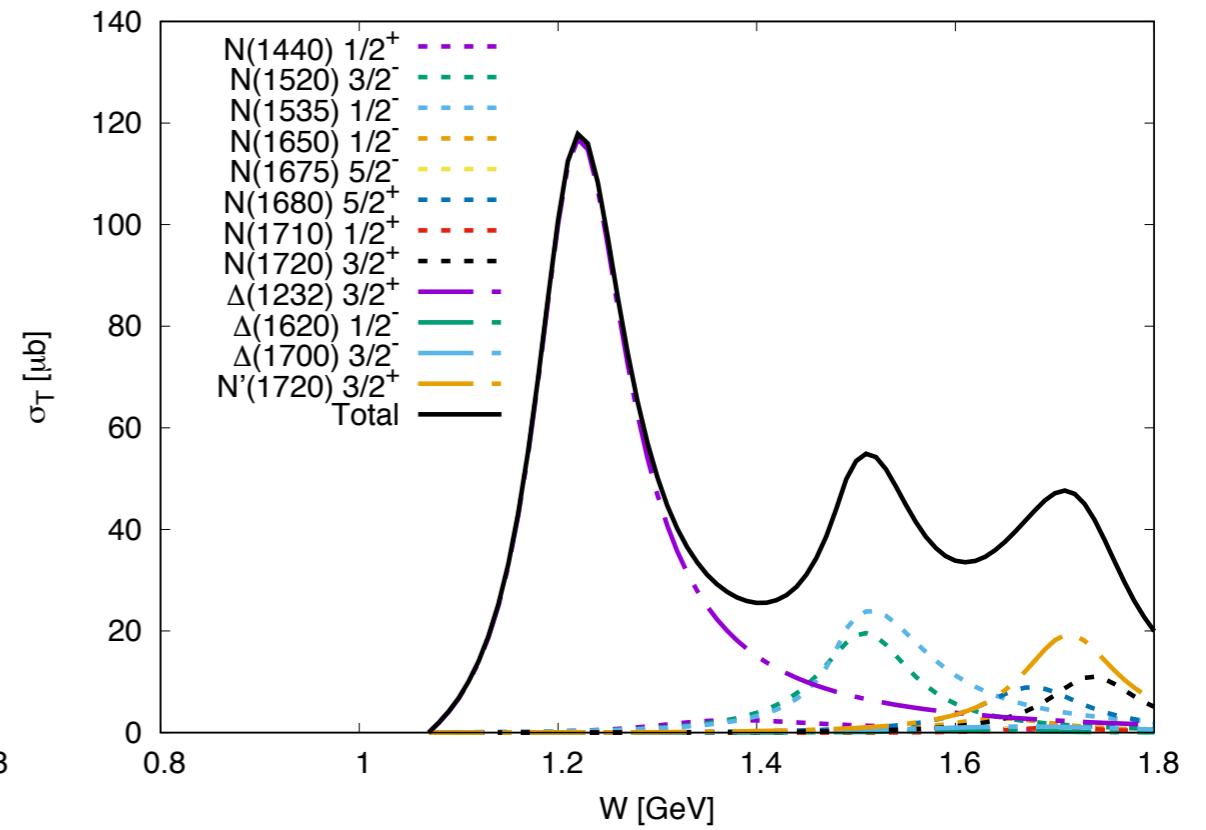
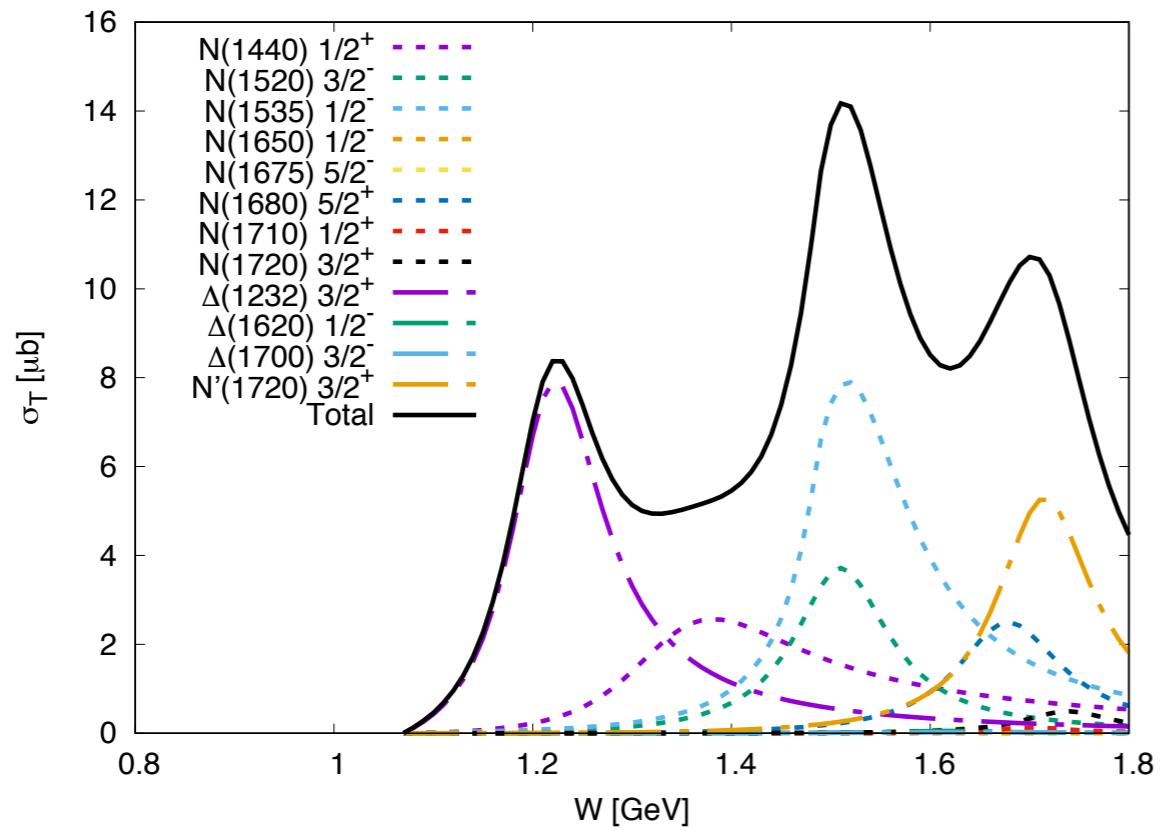
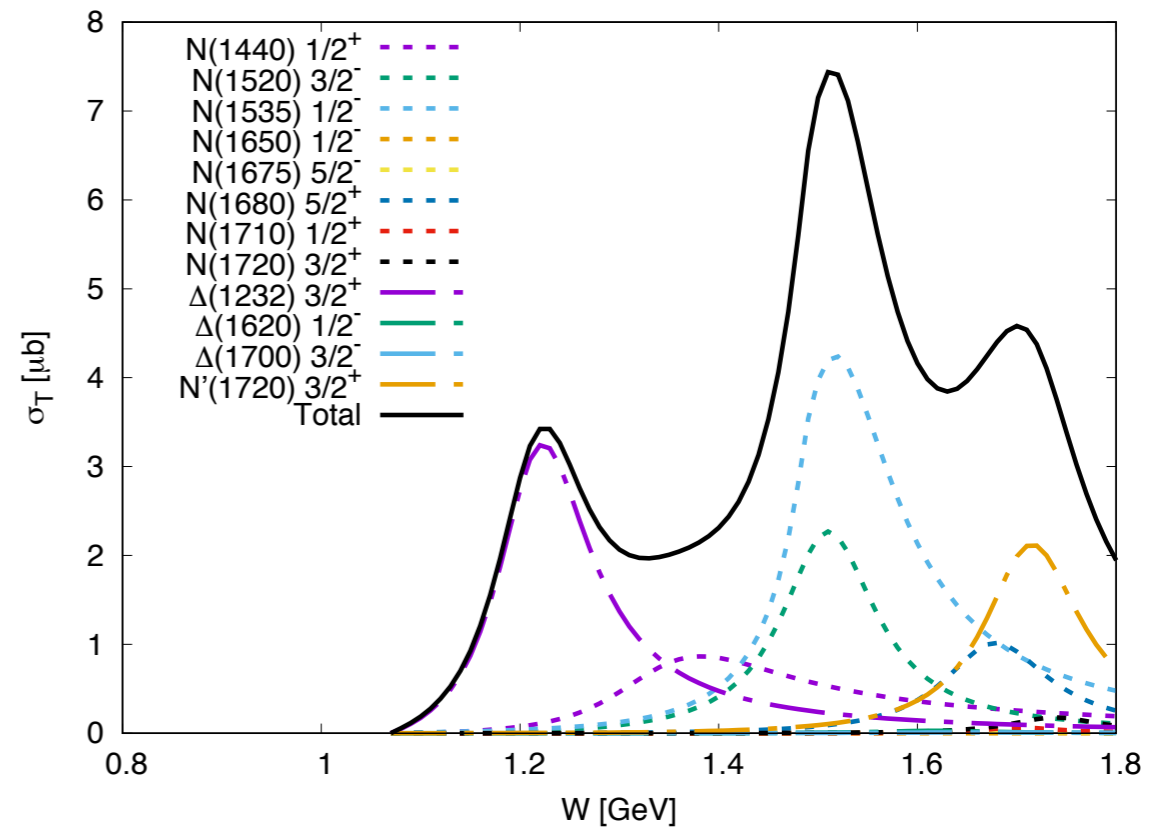
S_{1/2}

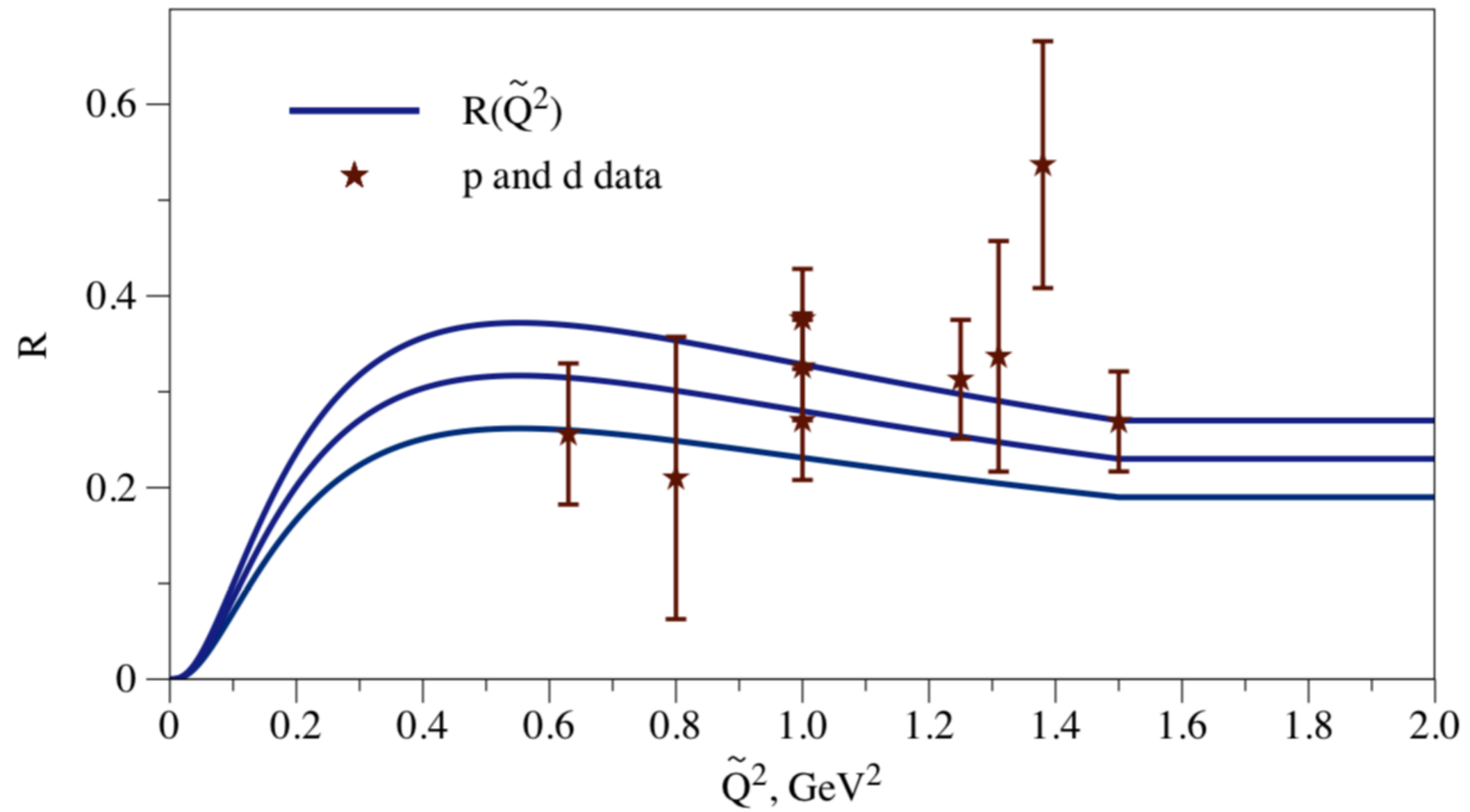


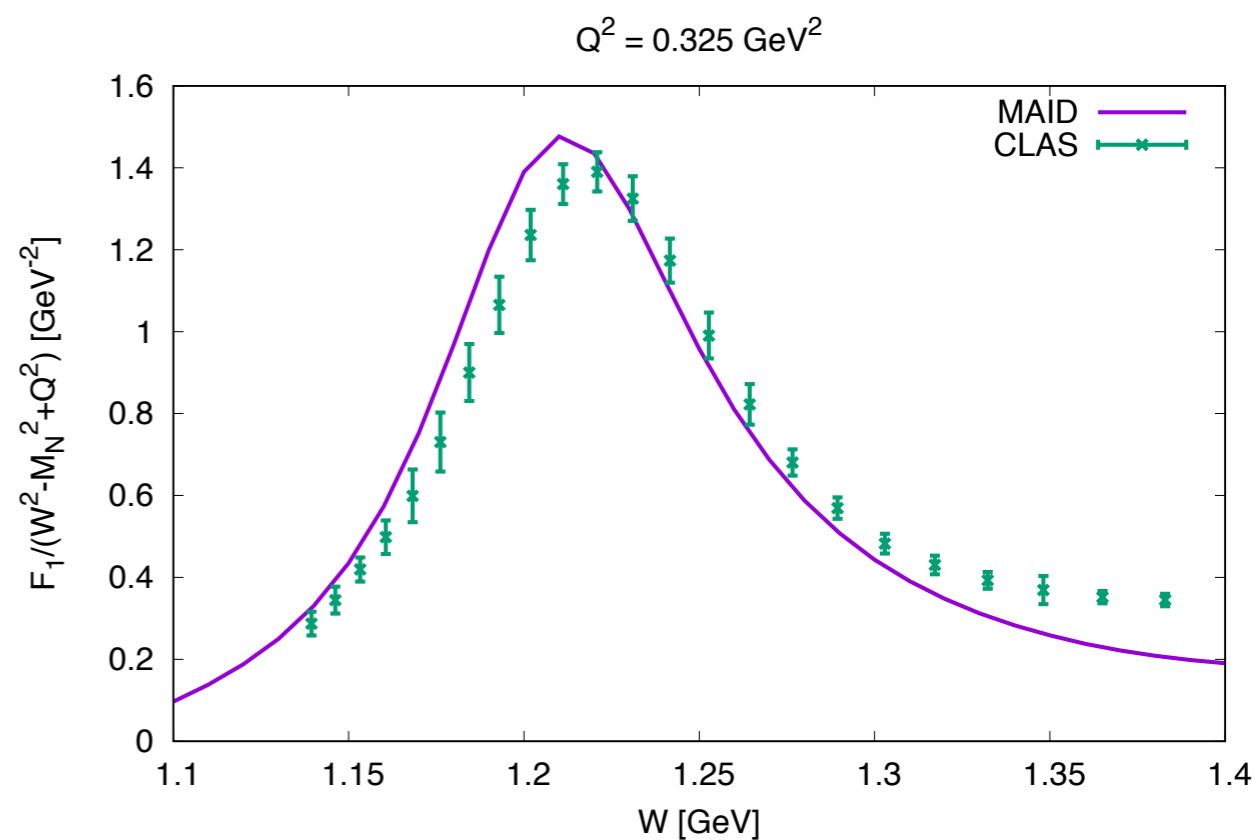
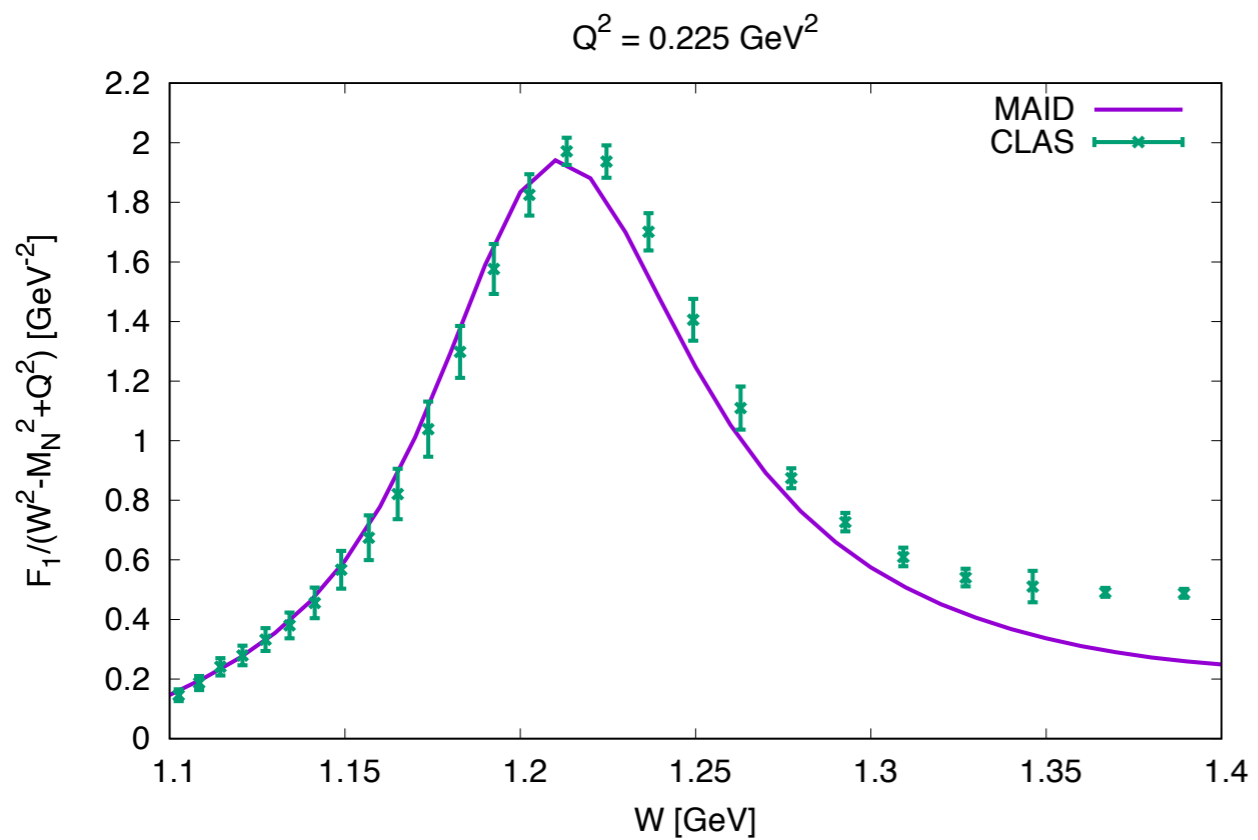
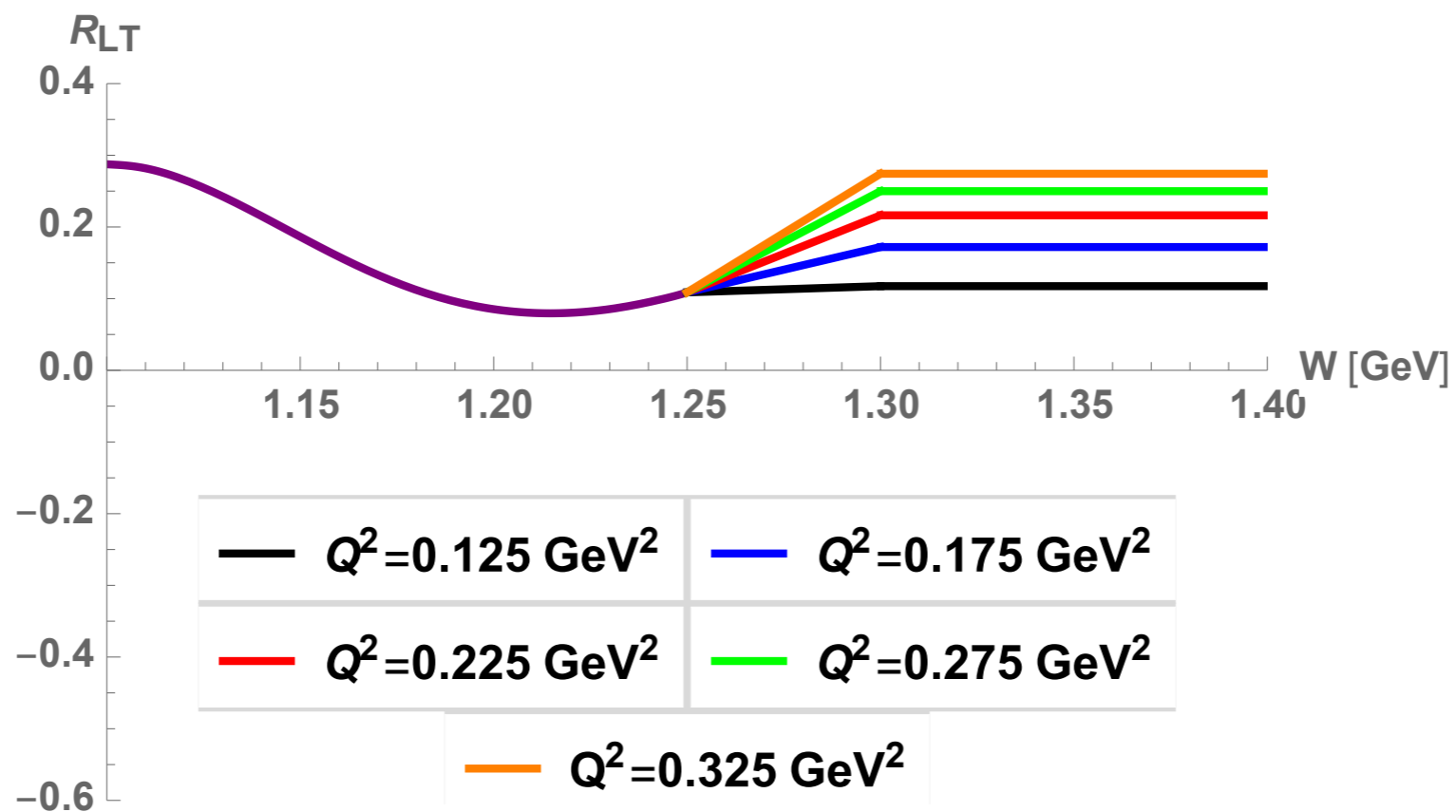
X

New resonance

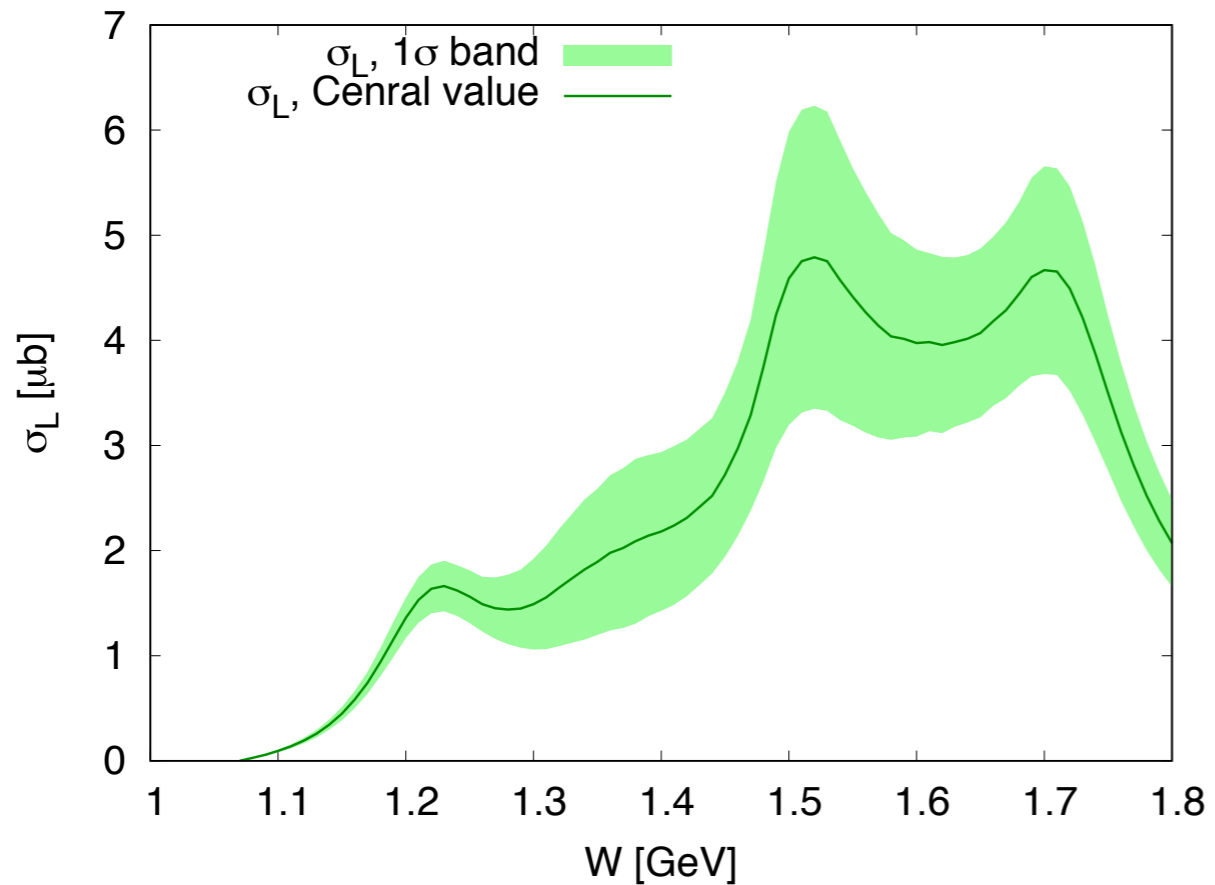
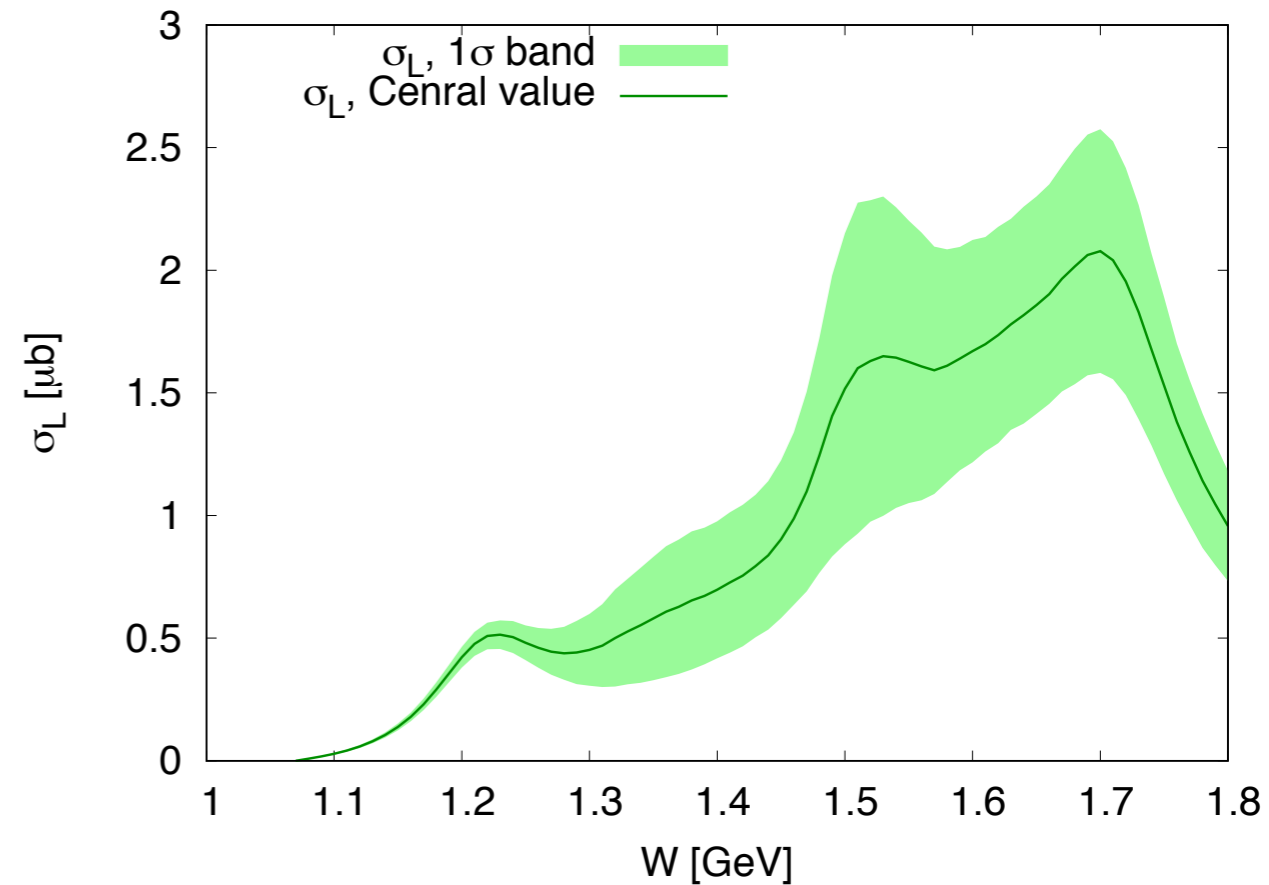
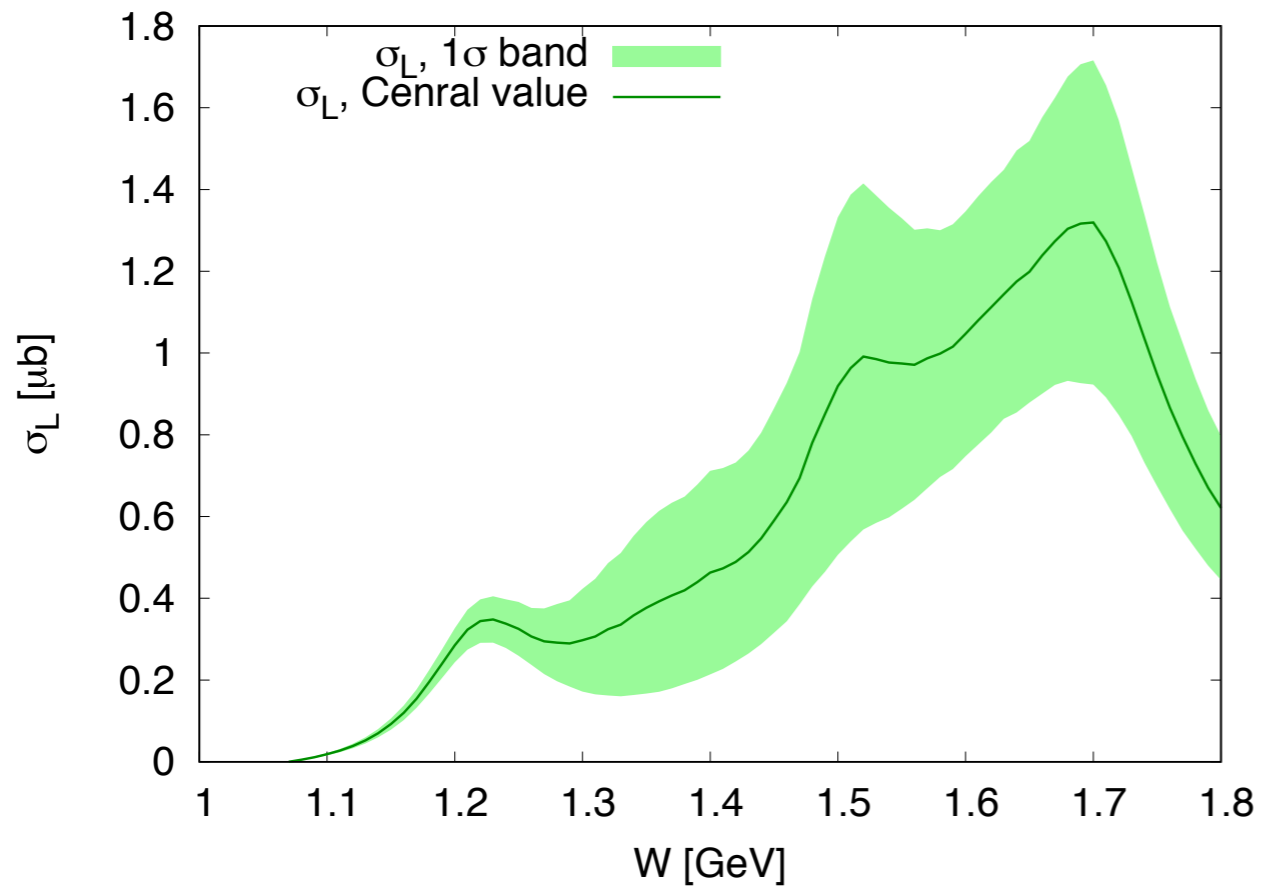
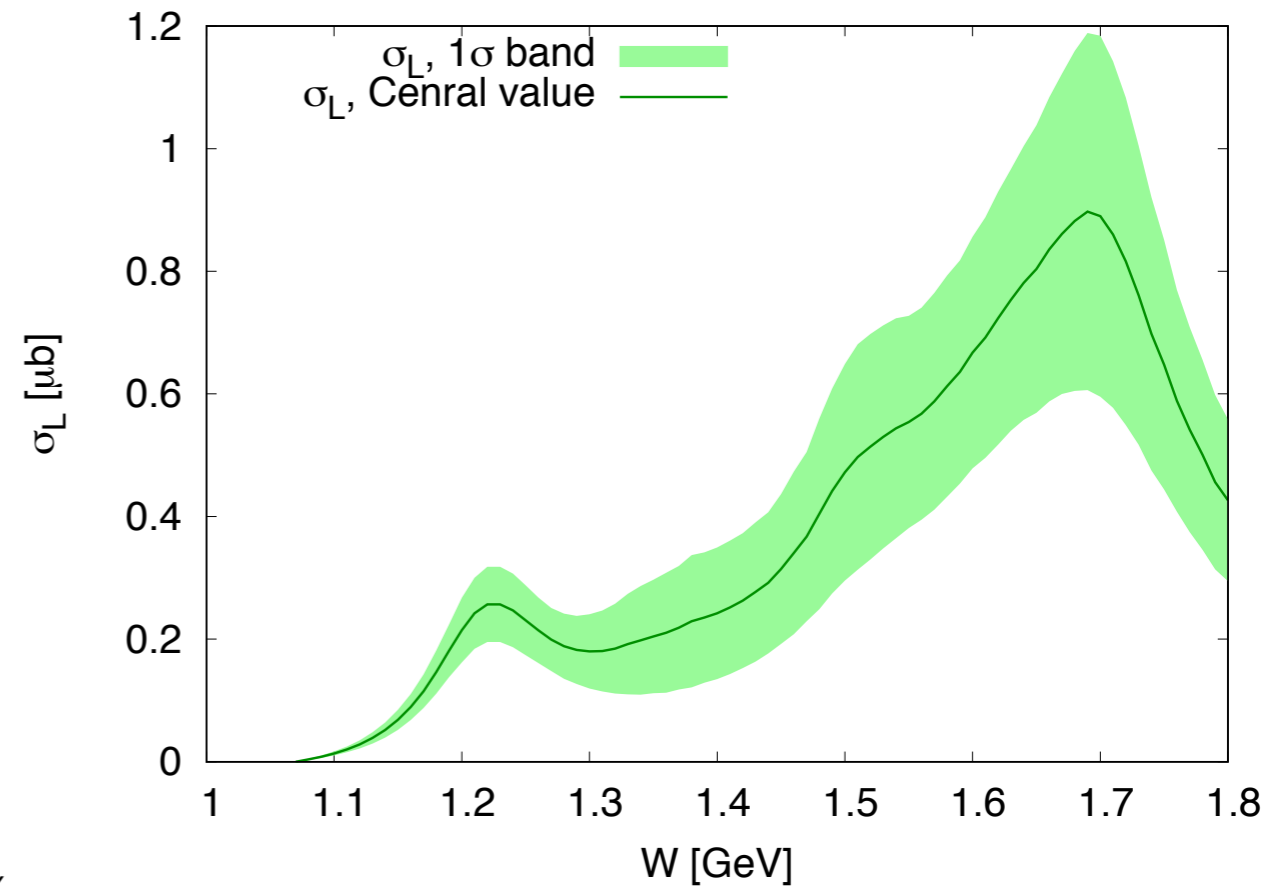


$Q^2 = 0.5 \text{ GeV}^2$  $Q^2 = 1 \text{ GeV}^2$  $Q^2 = 3 \text{ GeV}^2$  $Q^2 = 4 \text{ GeV}^2$ 





X

$Q^2 = 2 \text{ GeV}^2$  $Q^2 = 3 \text{ GeV}^2$  $Q^2 = 4 \text{ GeV}^2$  $Q^2 = 5 \text{ GeV}^2$ 

X

Subtraction function calculation

$$\bar{T}_1(0, Q^2) = T_1^{\text{NR}}(0, Q^2) + \frac{\alpha}{M} F_D^2(Q^2) + \frac{2\alpha}{M} \int_{\nu_0}^{\infty} d\nu \frac{F_1^{\text{R}}(\nu, Q^2)}{\nu}$$