Benchmark Analyses of the π^0 Production off the ⁴He Target

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Exclusive pion-induced Drell-Yan process at J-PARC



Measurement of Exclusive π^0 Electroproduction at JLab



CLAS Collab.PRL 109, 112001 (2012); PRC90,025205(2014)



CLAS Collaboration







PRL123,032502(2019)

PRL119,202004(2017)



C.Ji, H.-M.Choi, A.Lundeen, B.Bakker PRD99,116008(2019)

Salient Features

- No interference from Bethe-Heitler process
- BSA of exclusive coherent electroproduction of the π^0 off ⁴He has been measured.
- Data appear consistent with our benchmark BSA prediction for 0⁻⁺ meson production off the scalar target.
- General formulation of hadronic amplitudes in Meson Production off the Scalar Target (0⁺⁺ vs. 0⁻⁺)
- Comparison/Contrast with the leading twist GPD formulation.



5-Fold Differential Cross Section

for unpolarized target and without recoil polarization

The following notation of the coincidence cross section will be used in our calculations. Further details can be found in D. Drechsel and L. Tiator, J. Phys. G 18 (1992) 449-497. (scanned version) (click here for a larger image)

$$\frac{d\sigma}{d\Omega_f dE_f d\Omega} = \Gamma \frac{d\sigma_v}{d\Omega}, \quad \Gamma = \frac{\alpha}{2\pi^2} \frac{E_f}{E_i} \frac{k_\gamma}{Q^2} \frac{1}{1-\varepsilon}$$
$$\frac{d\sigma_v}{d\Omega} = \frac{d\sigma_T}{d\Omega} + \varepsilon \frac{d\sigma_L}{d\Omega} + \sqrt{2\varepsilon(1+\varepsilon)} \frac{d\sigma_{LT}}{d\Omega} \cos\phi$$
$$+ \varepsilon \frac{d\sigma_{TT}}{d\Omega} \cos 2\phi + h\sqrt{2\varepsilon(1-\varepsilon)} \frac{d\sigma_{LT'}}{d\Omega} \sin\phi$$

https://maid.kph.uni-mainz.de/maid2007/cross.html Mainz analysis interactive database (MAID)

ed from the
$$q \cdot J = 0$$
 relation, we get the following inv
 $|\mathcal{M}|^2 = \begin{pmatrix} e_2^2 \\ e_2^2 \end{pmatrix}^{2\mathcal{H}_{\mu\nu}} = J^{\dagger}_{\mu}J_{\nu}$
the leptonic tensor including the electron brack polarization λ is
 $\mathcal{L}^{\mu\nu} = q^2 \begin{bmatrix} g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu}\mu + k'^{\mu}k'^{\nu}) \\ g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ g^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\nu} + k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\mu} + k'^{\mu}k'^{\mu}k'^{\nu}) \\ \mathcal{H}^{\mu\nu} = g^{\mu\nu} + \frac{2}{q^2} (k^{\mu}k'^{\mu}k'^{\mu} + k'^{\mu}k'$

 $\mu^{\mu\nu} = q^{\mu\nu} + \frac{2}{2}(k^{\mu}k'^{\nu} + k'^{\mu}k^{\nu})$. Here, $\mathcal{L}^{\mu\nu}$ and $\mathcal{H}_{\mu\nu}$ are c

$$J^{\mu}_{PS} = F_{PS} \epsilon^{\mu\nu\alpha\beta} q_{\nu} \bar{P}_{\alpha} \Delta_{\beta}$$

$$\begin{aligned} \mathcal{H}_{\mu\nu} &= J^{\dagger}_{\mu} J_{\nu} \\ &= |F_{PS}|^{2} \epsilon_{\mu\alpha\beta\gamma} \epsilon_{\nu\alpha'\beta'\gamma'} q^{\alpha} \bar{P}^{\beta} \Delta^{\gamma} q^{\alpha'} \bar{P}^{\beta'} \Delta^{\gamma'} \\ &= \mathcal{H}_{\nu\mu} \\ \hline \epsilon^{\mu\nu\alpha\beta} k_{\alpha} k'_{\beta} \mathcal{H}_{\mu\nu} = 0 \end{aligned}$$

$$\frac{d\sigma_{h=+1}^{PS} - d\sigma_{h=-1}^{PS}}{d\sigma_{h=+1}^{PS} + d\sigma_{h=-1}^{PS}} = 0$$



$$d\sigma^{PS} = d\sigma_T^{PS} + d\sigma_{TT}^{PS}\epsilon\cos 2\phi$$
$$= d\sigma_T^{PS}(1 - \epsilon\cos 2\phi)$$

 $d\sigma_T^{PS} = \kappa \frac{e^4 |F_{PS}(Q^2, t, x)|^2 \sin^2 \theta}{4M^2 x^4 (1 - \epsilon)} \left(4M^2 x^2 + Q^2 \right) \left[x^2 \left(t^2 - 4m^2 M^2 \right) + Q^4 + 2Q^2 tx \right]$

Equivalently general benchmark prediction is feasible to the reverse processes envisioned in JPARC experiments, e.g. $\pi^- + {}^4\text{He} \rightarrow {}^4\text{He}{}^+\text{e}{}^-$

Although the Hydrogen-4 is a highly unstable isotope, it may be still identifiable from its decay to tritium through neutron emission.

If the lepton spin polarizations can be measured, then the prediction of the zero spin asymmetry of the lepton pair may be testable.

Conclusion and Outlook

- The general formulation with the Rosenbluth-type separation of the differential cross section for the exclusive π^0 production off ⁴He is provided.
- In particular, its BSA is predicted to be absent from the general symmetry ground and the data support this prediction.
- Equivalently general formulation is feasible for the reverse processes envisioned in JPARC experiments, e.g.

 π^- + ⁴He \rightarrow ⁴H e⁺e⁻