Generalized parton distributions and form factors of the proton

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Generalized parton distributions and form factors

- Generalized form factors (GPDs) are three-dimensional descriptions of partonic structure.
- GPDs contain the elastic form factors.
- Two-dimensional slices of GPDs are seen in **deeply virtual Compton scattering** (DVCS).
- GPDs & DVCS offer additional insight to form factors.
- I will present model calculations/predictions for proton GPDs.



Form factors and proton structure Point particle

$$\langle p's'|j^{\mu}(0)|ps\rangle = \bar{u}^{s'}(p')\gamma^{\mu}u^{s}(p)$$



$$\Delta = p' - p \qquad t = \Delta^2 = -Q^2$$

General spin-half particle

$$\langle p's'|j^{\mu}(0)|ps\rangle = \bar{u}^{s'}(p') \left[\gamma^{\mu}F_1(t) + \frac{i\sigma^{\mu\Delta}}{2m_p}F_2(t)\right] u^s(p)$$



- Form factors appear in the electromagnetic current.
- They tell us about substructure—how is the proton different from a point particle?
- Seen in reactions where proton gets a momentum kick.
- How kick is redistributed gives hint at structure—seen in t dependence.
 - Stay together easily with large kick? Slow t falloff
 - Trouble staying together? Fast t falloff

Electric and magnetic form factors

Sachs form factors Sum rules $G_E(t) = F_1(t) + \frac{t}{4m_p^2}F_2(t)$ $G_E(0) = F_1(0) = 1$ $G_M(t) = F_1(t) + F_2(t)$ $G_M(0) = F_1(0) + F_2(0)$ $\mu_p = 1 + \kappa_p$

• 3D Fourier transforms of $G_E(t)$ and $G_M(t)$ characterize electric and magnetic response of charged objects (electrons) to proton; e.g.,

$$\rho_E(\mathbf{r}) = \int \frac{d^3 \mathbf{k}}{2E_k (2\pi)^3} G_E(t = -\mathbf{k}^2) e^{-i(\mathbf{k} \cdot \mathbf{r})}$$

- These are not literal spatial densities (see e.g., Miller, PRC99 (2019) 035202)
- Charge and magnetization radii:

$$\langle r_E^2 \rangle = 6 \frac{\mathrm{d}G_E(t)}{\mathrm{d}t} \bigg|_{t=0} \qquad \langle r_M^2 \rangle = 6 \frac{\mathrm{d}G_M(t)}{\mathrm{d}t} \bigg|_{t=0}$$

 \ldots electrons respond as if proton had these radii.

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Flavor-separated structure

- Proton & neutron have valence *uud* and *udd* structure.
- Assume charge symmetry: neutron is a proton with $u \leftrightarrow d$ swap.

$$\begin{split} F_{ip}(t) &= \frac{2}{3}F_{ip}^u(t) - \frac{1}{3}F_{ip}^d(t) \\ F_{in}(t) &= -\frac{1}{3}F_{ip}^u(t) + \frac{2}{3}F_{ip}^d(t) \end{split}$$

(where i = 1, 2) and thus

 $F_{ip}^u(t) = 2F_{ip}(t) + F_{in}(t)$ $F_{ip}^d(t) = F_{ip}(t) + 2F_{in}(t)$

- Neutron measurements can tell us more about the proton.
- Different forms for $F_{ip}^u(t)$ and $F_{ip}^d(t)$ hint at interesting isospin structure!



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Diquark correlations

- One way to get behavior seen in Cates *et al.* is **diquark correlations**.
- Proton is made of three **dressed quarks**.
 - Dressed quarks are quasi-particles.
 - Mass of dressed quark: 300-400 MeV (model-dependent).
 - They are amalgamations of the true (current) quarks.
- Two of the three quarks are often bound together into a diquark.
- Proton then looks like a two-body state: quark and diquark.



Scalar diquarks

- The lightest diquarks are scalar/isoscalar.
- Isoscalarity \Rightarrow down quark is in diquark.
- Asymptotic form for two-body form factor $\sim \frac{1}{Q^2}$ at large Q^2 .
- Quark outside diquark (up quark): form factor goes as $\frac{1}{Q^2}$.
- Quark within diquark: two layers of two-body behavior, $\frac{1}{Q^4}$
- Cates *et al.* result highly suggestive of isoscalar diquarks!



Figure from Cates et al., PRL106 (2011) 252003

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Scalar & axial-vector diquarks



	Scalar diquarks	Scalar+axial	Experiment
κ_p	0.57	1.49	1.79

(Pion cloud improves model; Cloët et al., PRC90 (2014) 045202)

- More than just scalar diquarks may exist.
- Axial-vector/isovector diquarks are next-lightest diquark.
- Axial-vector diquarks needed to get large proton magnetic moment.
- **Question**: how could we see evidence of axial diquarks in observables? (a la Cates *et al.* for scalar diquarks).

Model calculations (on left):

- Nambu–Jona-Lasinio model
- No pion cloud included
- Q^2 falloff is slower than reality
- Introducing dipole with vector meson pole improves Q^2 falloff
- See arXiv:1907.08256 & Cloët &al. (2014) for details

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Proton GPDs

Multidimensional structure



- More proton structure information if a new particle is **created** and **detected** in final state.
- Two possible avenues for exploration:
 - **4** Hard exclusive reactions: proton remains intact.
 - **Semi-inclusive deeply inelastic scattering** (SIDIS): proton is (generally) destroyed, sum over all final states that include created/detected particle.
- **Option 1** is closest analogue to form factors (elastic scattering).
- **Deeply virtual Compton scattering** (DVCS): the created/detected particle is a photon.

DVCS and GPDs



Deeply virtual Compton scattering $\mathcal{H}(\xi, t)$



Generalized parton distribution $H(x,\xi,t)$

- **Deeply virtual Compton scattering** (DVCS) is the hard electro-production of a photon.
- Generalized parton distributions (GPDs) encode the QCD structure seen in DVCS.
- Presence of a loop means one GPD variable is integrated over.
- Integrated quantities seen in experiment: Compton form factors

$$\mathcal{H}(\xi,t) = \int \mathrm{d}x \, \left[\frac{1}{\xi - x - i0} \mp \frac{1}{\xi + x - i0} \right] H(x,\xi,t)$$

The GPD/DVCS variables



- $\frac{x+\xi}{1+\xi}$ and $\frac{x-\xi}{1-\xi}$ are initial and final light cone momentum fractions of struck quark.
- t is the invariant momentum transfer to the target.
- Q^2 is the invariant momentum transfer from the electron.
- $t \neq -Q^2$, in contrast to elastic scattering.
- Q^2 acts as a **resolution scale**, like in deeply inelastic scattering (DIS). (I will use $Q^2 = 4$ GeV² for remainder of talk.)
- t tells us about structure seen from redistribution of momentum kick.
- t is the "form factor variable" rather than Q^2 .

GPDs and form factors

• Helicity-independent GPDs defined using a non-local operator

$$\frac{1}{2} \int \frac{dz}{2\pi} e^{-iP \cdot nzx} \langle p' | \bar{q} \left(\frac{nz}{2}\right) \not n q \left(-\frac{nz}{2}\right) | p \rangle = \bar{u}(p') \left[\not n H^q(x,\xi,t) + \frac{i\sigma^{n\Delta}}{2m_N} E^q(x,\xi,t) \right] u(p)$$

...in the light cone gauge. Other gauges require a Wilson line.

• Breakdown analogous to form factors:

$$\langle p's'|j^{\mu}(0)|ps\rangle = \bar{u}^{s'}(p') \left[\gamma^{\mu}F_1(t) + \frac{i\sigma^{\mu\Delta}}{2m_p}F_2(t)\right] u^s(p)$$

• These are even related through integration

$$\int dx H^{q}(x,\xi,t;Q^{2}) = F_{1}^{q}(t) \qquad \int dx E^{q}(x,\xi,t;Q^{2}) = F_{2}^{q}(t)$$

• GPDs are, in a sense, extensions of form factors





Orange is up; blue is down.

- Forward limit: $H^q(x, 0, 0) = q(x)$
- Related to (flavor-separated) **Dirac form** factor:

$$\int \mathrm{d}x \, H^q(x,0,t) = F_1^q(t)$$

• Additionally related to gravitational form factor:

$$\int \mathrm{d}x \, x H^q(x,0,t) = A^q(t)$$

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describing distribution of energy

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• Show hybrid of behaviors of form factors and parton distribution functions (PDFs)

Non-skewed proton GPD: $E^q(x, \xi = 0, t)$



Orange is up; blue is down.

- No forward limit.
- Related to (flavor-separated) **Pauli form** factor:

$$\int \mathrm{d}x \, E^q(x,0,t) = F_2^q(t)$$

- Up and down quarks have opposite signs—contribute constructively to κ_p
- Additionally related to gravitational form factor:

$$\frac{1}{2}\int \mathrm{d}x\,x\Big(H^q(x,0,t)+E^q(x,0,t)\Big)=J^q(t)$$

describing distribution of total angular momentum

A normalized, flavor-separated super-ratio

- How much faster is down quark falloff for different *x* slices?
- Normalize out the t = 0 value:

$$S = \frac{H^d(x,0,t)}{H^d(x,0,0)} \bigg/ \frac{H^u(x,0,t)}{H^u(x,0,0)}$$

- In scalar diquark model, down quarks are within diquark.
 - The d/u ratio is small to begin with at large x
 - At small x, d/u ratio decreases as diquark dies off with t
- We see new behavior offered by this multidimensional description!



Scalar diquarks only

Super-ratio: a model comparison



Scalar diquarks only

Scalar and axial diquarks

- Noticeable qualitative difference!
- This is because down quark can be *outside* an axial diquark.



Skewed proton GPD: $H^q(x, \xi = 0.5, t)$



Orange is up; blue is down.

• Still related to (flavor-separated) **Dirac form** factor:

$$\int \mathrm{d}x \, H^q(x,\xi,t) = F_1^q(t)$$

(integration kills ξ dependence!)

• Related to *two* gravitational form factors:

$$\int \mathrm{d}x \, x H^q(x,\xi,t) = A^q(t) + \xi^2 C^q(t)$$

• $C^{q}(t)$ describes how shear and pressure are distributed within the proton (see Polyakov & Schweitzer).

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• Being able to study ξ dependence offers an exciting new window into probing forces!

Flavor-separated gravitational form factors?



• Sum rules: A(0) = 1, B(0) = 0, C(0) < 0, after sum over quarks & gluons ... hence why $A^u(0) + A^d(0) < 1$ in plot above

- $A^u(t) > A^d(t)$ since there are two up quarks, but ...
- ... down quarks contribute more to shear/pressure??? (We need to understand this better within the model)
- Flavor-separated gravitational form factors could tell us so much more about the proton.
- But GPDs & GFFs are not *directly* measured—long term project.
- What could we do more directly?

DVCS and Bethe-Heitler processes





Bethe-Heitler process

- DVCS interferes with **Bethe-Heitler** process: photon emission by electron
- This is a boon rather than a bane: interference between diagrams seen in beam-spin asymmetry!
- Interference term seen in $\sin \phi$ modulation of **beam spin asymmetry**:

$$A_{LU}(\phi) = \frac{\sigma^+(\phi) - \sigma^-(\phi)}{\sigma^+(\phi) + \sigma^-(\phi)} \propto \operatorname{Im}\left[F_1 \mathcal{H} - \frac{t}{4m_p^2} F_2 \mathcal{E} + \xi(F_1 + F_2)\tilde{\mathcal{H}}\right] \sin\phi + \cos\phi \text{ modulations} + \dots$$

Compton form factors

$$A_{LU}(\phi) \propto \operatorname{Im}\left[F_1 \mathcal{H} - \frac{t}{4m_p^2}F_2\mathcal{E} + \xi(F_1 + F_2)\tilde{\mathcal{H}}\right]\sin\phi + \dots$$

(diagram taken from Pisano et al., PRD91 (2015) 052014)

- $\tilde{\mathcal{H}}$ is a helicity-dependent Compton form factor.
- The imaginary parts of Compton form factors appear as:

$$\begin{aligned} \mathcal{H}(\xi,t) &= -\pi \sum_{q} e_{q} [H^{q}(\xi,\xi,t) - H^{q}(-\xi,\xi,t)] \\ \tilde{\mathcal{H}}(\xi,t) &= -\pi \sum_{q} e_{q} [\tilde{H}^{q}(\xi,\xi,t) + \tilde{H}^{q}(-\xi,\xi,t)] \\ \tilde{\mathcal{E}}(\xi,t) &= -\pi \sum_{q} e_{q} [\tilde{H}^{q}(\xi,\xi,t) + \tilde{H}^{q}(-\xi,\xi,t)] \\ \end{aligned}$$

leptonic plane

hadronic

- CFFs and beam spin asymmetry give **two-dimensional slices** of GPDs.
- Need precision determinations of form factors in order to extract CFFs.
- Question: is there anything about diquarks we can see in *just* imaginary CFFs?

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Model calculation: flavor-separated Compton form factors



Super-ratio for Compton form factors

- Let's see if this works for CFFs, which are measured.
- How much faster is down quark falloff for different *ξ* slices?
- This is a new lever-arm offered by virtual Compton scattering!
- Normalize out the t = 0 value:

$$S = \frac{\mathcal{H}^d(\xi, t)}{\mathcal{H}^d(\xi, 0)} \bigg/ \frac{\mathcal{H}^u(\xi, t)}{\mathcal{H}^u(\xi, 0)}$$

- High $x \Rightarrow$ high ξ in imaginary CFFs
- In scalar diquark model, down quarks are within diquark.
 - The d/u ratio is small to begin with at large ξ
 - At small ξ , d/u ratio decreases as diquark dies off with t



Scalar diquarks only

Super-ratio: Compton form factor case



Scalar diquarks only

Scalar and axial diquarks

- Noticeable qualitative difference, still!
- Again because down quark can be *outside* an axial diquark.

Conclusions and summary

- **Generalized parton distributions** (GPDs) are multi-dimensional descriptions of proton structure which contain the elastic form factors.
- GPDs can also be used to determine **gravitational form factors**.
- Compton form factors are two-dimensional slices of GPDs, which are measured in deeply virtual Compton scattering (DVCS).
- Compton form factors contain the potential to elaborate upon the quark-diquark structure of the proton.

Thank you for your time and attention!