

# Parton Distribution Functions from Lattice QCD

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with

Eloy Romero and Joe Karpie

# Introduction

- Goal: Compute hadron structure properties from QCD
  - Parton distribution functions (PDFs)
- Operator product: Mellin moments are local matrix elements that can be computed in Lattice QCD
  - Power divergent mixing limits us to few moments
- Few years ago X. Ji suggested an approach for obtaining PDFs from Lattice QCD
- First calculations already available
  - X. Ji, Phys.Rev.Lett. 110, (2013)*
  - Y.-Q. Ma J.-W. Qiu (2014) 1404.6860*
  - H.-W. Lin, J.-W. Chen, S. D. Cohen, and X. Ji, Phys.Rev. D91, 054510 (2015)*
  - C. Alexandrou, et al, Phys. Rev. D92, 014502 (2015)*
- A new approach for obtaining PDFs from LQCD introduced by A. Radyushkin
  - A. Radyushkin Phys.Lett. B767 (2017)*
- Hadronic tensor methods
  - K-F Liu et al Phys. Rev. Lett. 72 (1994) , Phys. Rev. D62 (2000) 074501*
  - Detmold and Lin 2005*
  - M. T. Hansen et al arXiv:1704.08993.*
  - UKQCD-QCDSF-CSSM Phys. Lett. B714 (2012), arXiv:1703.01153*
  - Ma and Qiu : [arXiv:1709.03018](https://arxiv.org/abs/1709.03018)*

# Pseudo-PDFs

Unpolarized PDFs proton:

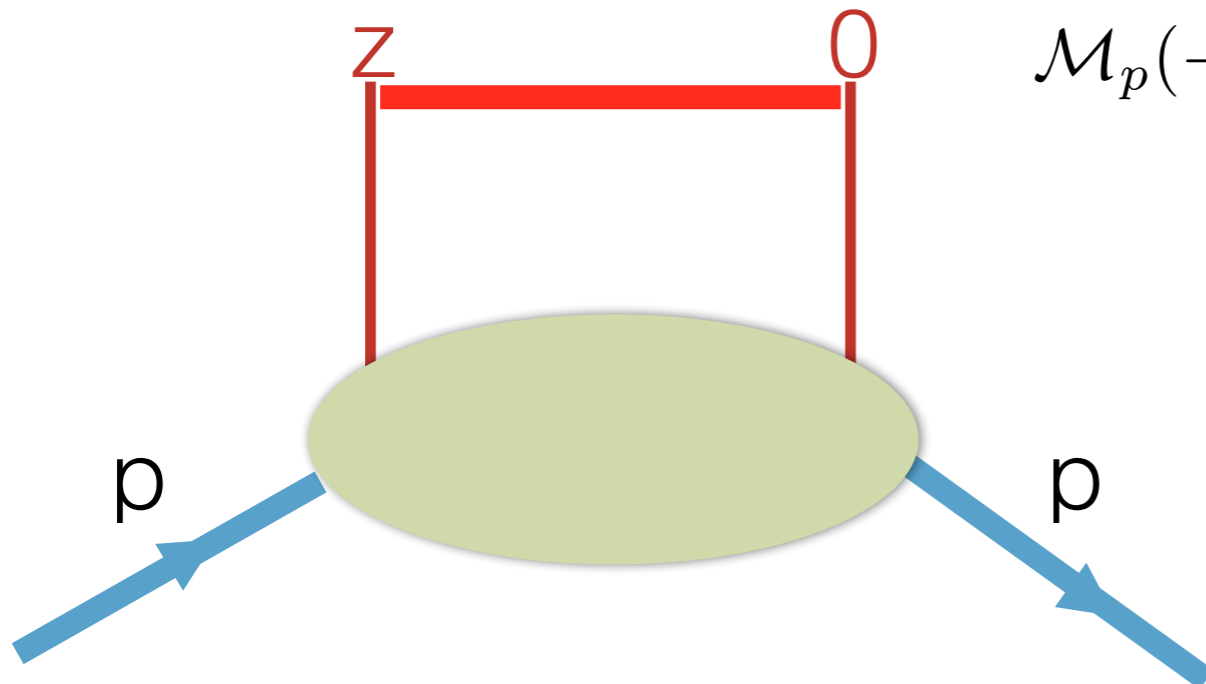
$$\mathcal{M}^\alpha(z, p) \equiv \langle p | \bar{\psi}(0) \gamma^\alpha \hat{E}(0, z; A) \psi(z) | p \rangle$$

$$\hat{E}(0, z; A) = \mathcal{P} \exp \left[ -ig \int_0^z dz'_\mu A_\alpha^\mu(z') T_\alpha \right]$$

$$\mathcal{M}^\alpha(z, p) = 2p^\alpha \mathcal{M}_p(-pz, -z^2) + z^\alpha \mathcal{M}_z(-pz, -z^2)$$

$$\mathcal{M}^+(z, p) = 2p^+ \mathcal{M}_p(-p_+ z_-, 0)$$

$$\mathcal{M}_p(-p_+ z_-, 0) = \int_{-1}^1 dx f(x) e^{-ixp_+ z_-}$$



$\mathcal{M}_p(-pz, -z^2)$  is a Lorentz invariant therefore computable in any frame

$\nu = -zp$  Ioffe time

**Goal:**

**Compute PDFs with Lattice QCD using space like separations  $z$**

$$\mathcal{M}_p(\nu, z^2) = \int_0^1 d\alpha \mathcal{C}(\alpha, z^2 \mu^2, \alpha_s(\mu)) \mathcal{Q}(\alpha\nu, \mu) + \mathcal{O}(z^2 \Lambda_{qcd}^2)$$

$\mathcal{Q}(\nu, \mu)$  is called the Ioffe time PDF

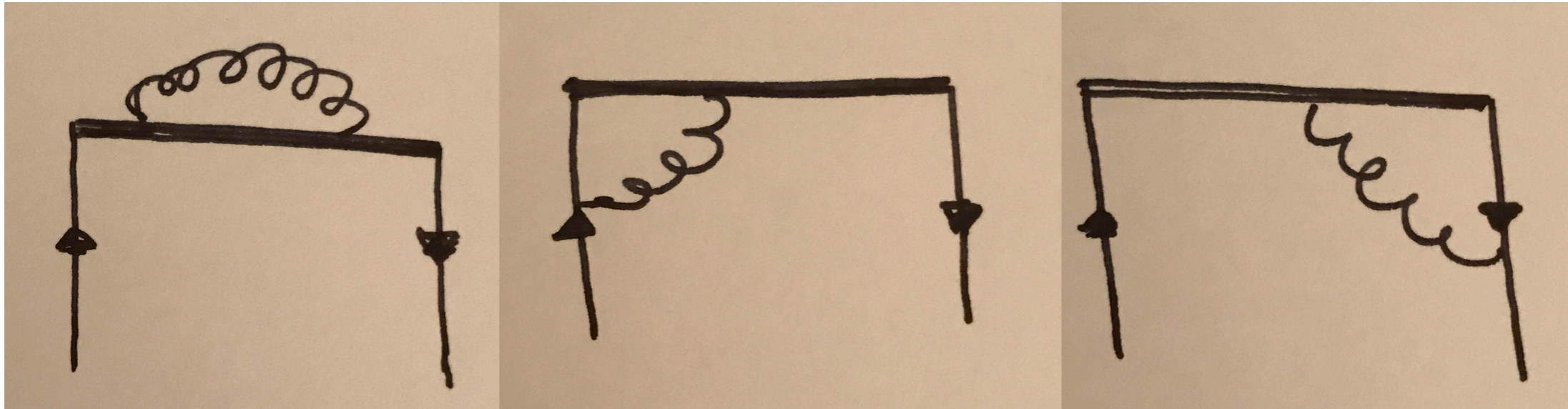
V. Braun, et. al Phys. Rev. D 51, 6036 (1995)

$$\mathcal{Q}(\nu, \mu) = \int_{-1}^1 dx e^{-ix\nu} f(x, \mu)$$

Matching to  $\overline{MS}$

Radyushkin Phys.Rev. D98 (2018) no.1, 014019  
Izubuchi et al. Phys.Rev. D98 (2018) no.5, 056004  
Zhang et al. Phys.Rev. D97 (2018) no.7, 074508

# UV behaviour



One loop calculation of the UV divergences results in

$$\mathcal{M}^0(z, P, a) \sim e^{-m|z|/a} \left( \frac{a^2}{z^2} \right)^{2\gamma_{end}}$$

after re-summation of one loop result resulting exponentiation

- J.G.M.Gatheral, Phys.Lett.133B,90(1983)
- J.Frenkel, J.C.Taylor, Nucl.Phys.B246,231(1984),
- G.P.Korchensky, A.V.Radyushkin, Nucl.Phys.B283,342(1987).

Multiplicatively renormalizable

Consider the ratio  $\mathfrak{M}(\nu, z_3^2) \equiv \frac{\mathcal{M}_p(\nu, z_3^2)}{\mathcal{M}_p(0, z_3^2)}$

UV divergences will cancel in this ratio resulting a renormalization group invariant (RGI) function

The lattice regulator can now be removed

$\mathfrak{M}^{cont}(\nu, z_3^2)$  Universal independent of the lattice

$\mathcal{M}_p(0, 0) = 1$  Isovector matrix element

$$\mathfrak{M}(\nu, z^2) = \int_0^1 d\alpha \mathfrak{E}(\alpha, z^2 \mu^2, \alpha_s(\mu)) \mathcal{Q}(\alpha\nu, \mu) + \sum_{k=1}^{\infty} \mathcal{B}_k(\nu) (z^2)^k$$

$$\mathcal{B}_k(\nu) (z^2)^k \sim \mathcal{O}(\Lambda_{qcd}^{2k})$$

Polynomial corrections to the Ioffe time PDF may be suppressed

B. U. Musch, *et al* Phys. Rev. D 83, 094507 (2011)

M. Anselmino *et al.* 10.1007/JHEP04(2014)005

A. Radyushkin Phys.Lett. B767 (2017)

Polynomial corrections will vanish in the  $z^2 = 0$  limit

## Factorization:

$$\mathfrak{M}(\nu, z^2) = \int_0^1 dx q(x, \mu) \mathcal{K}(x\nu, z^2 \mu^2) + \mathcal{O}(z^2 \Lambda_{qcd}^2)$$



**perturbative kernel (known to NLO)**

## Challenges:

- Compute matrix elements in lattice QCD
  - Optimal interpolating fields to isolate the ground state matrix element requires computation of many correlations functions. This task is dominated by dot-products which are typically not optimally implemented
- Solution of an inverse problem to obtain the PDF from lattice QCD data
  - Controlling the polynomial corrections is required for proper estimation of systematic errors. This in turn implies the need for very fine lattice spacings.



# Project Goals

- Create pseudo-data using the above model
  - Explore various options for parametrizing the transverse structure of the proton
- Develop methods that allow for reliable extraction of the PDF using pseudo-data
- Understand the systematics of the PDF extraction Karpie et. al JHEP 1904 (2019) 057
- Optimize the codes for computing the relevant matrix elements from lattice QCD
- Support goes to:
  - 1 graduate student to work on the modeling (J. Karpie)
  - 1 postdoc W&M to work on the LQCD code optimization (E. Romero)

# Matrix Element computation

*-Eloy Romero*

# Hadron correlation functions

- One of the LQCD goals is to predict low-energy hadron spectrum of nuclei
- That means measuring the two-point correlation functions of field operators with specific quantum numbers

$$C(t', t) = \langle \chi(t') \chi^\dagger(t) \rangle$$

- In Monte-Carlo calculations, the physical relevant signal in correlation functions falls exponentially and is rapidly dominated by statistical fluctuations
- Operators that create low-lying energy eigenstates quickly improve the quality of the data extracted exponentially
- Smearing methods project the operator on the low-energy states and make computationally feasible the evaluation of the correlation functions

$$\tilde{\chi}(t) = \square(t) \chi(t) \square(t), \quad \square(t) = V(t) V^\dagger(t), \quad C(t', t) \approx \langle \tilde{\chi}(t') \tilde{\chi}^\dagger(t) \rangle$$

- The computational cost is dominated by the creation and contraction of tensors with sizes of the smearing projection rank

# Higher-order Singular Value Decomposition

$$\tau = (U_1, U_2, \dots, U_d) \cdot s$$

where  $U_i$  are unitary matrices and  $s$  is the core tensor.

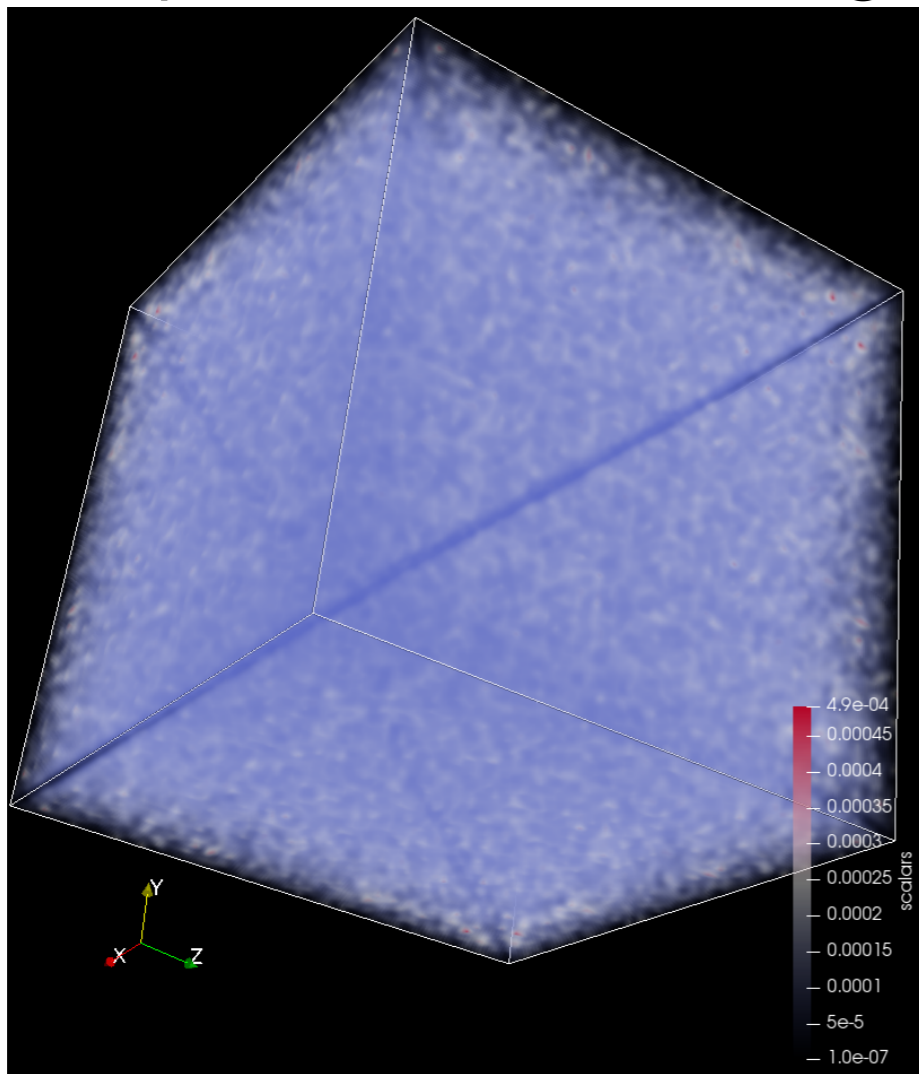
- HOSVD accelerates the contractions of tensors if near null rows, columns and fibers can be dropped from the core tensors

# Higher-order Singular Value Decomposition

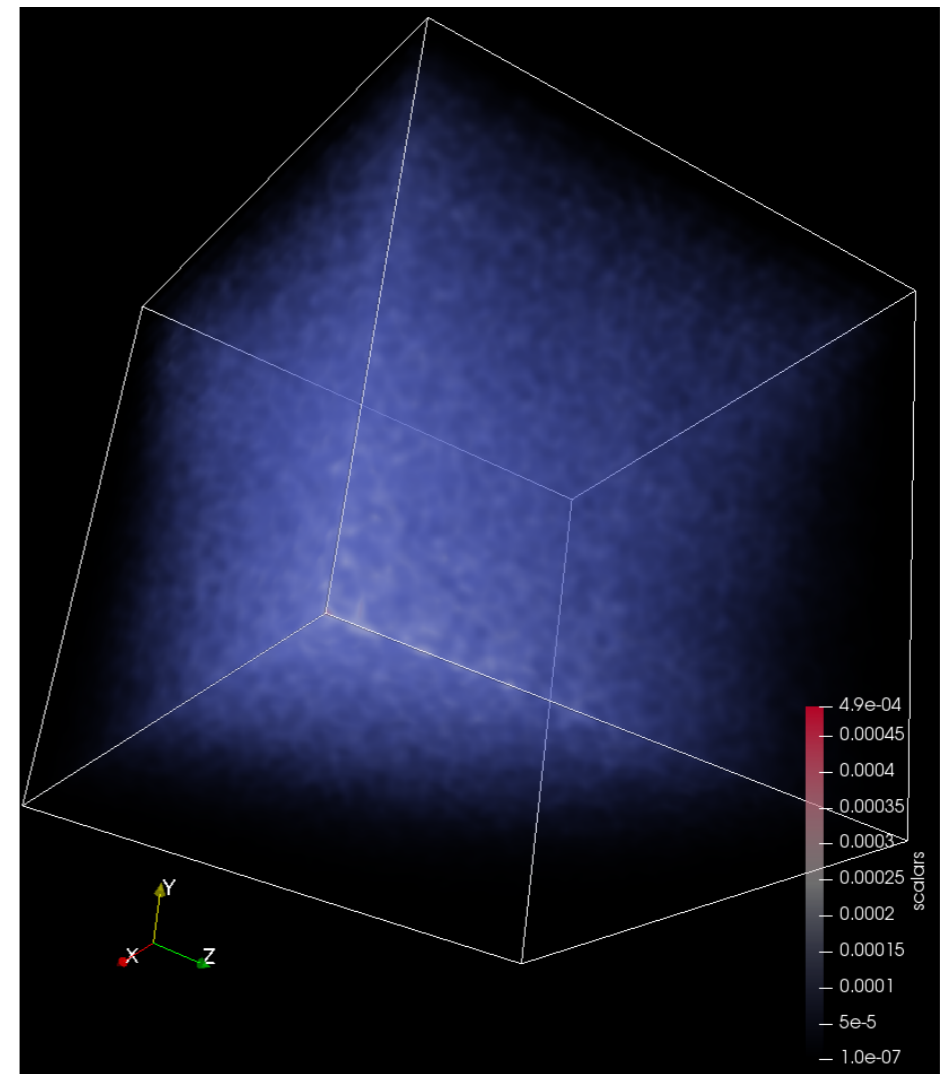
$$\tau = (U_1, U_2, \dots, U_d) \cdot s$$

where  $U_i$  are unitary matrices and  $s$  is the core tensor.

Example from one of the generated tensors:



$\mathcal{T}$



$\mathcal{S}$

# Higher-order Singular Value Decomposition

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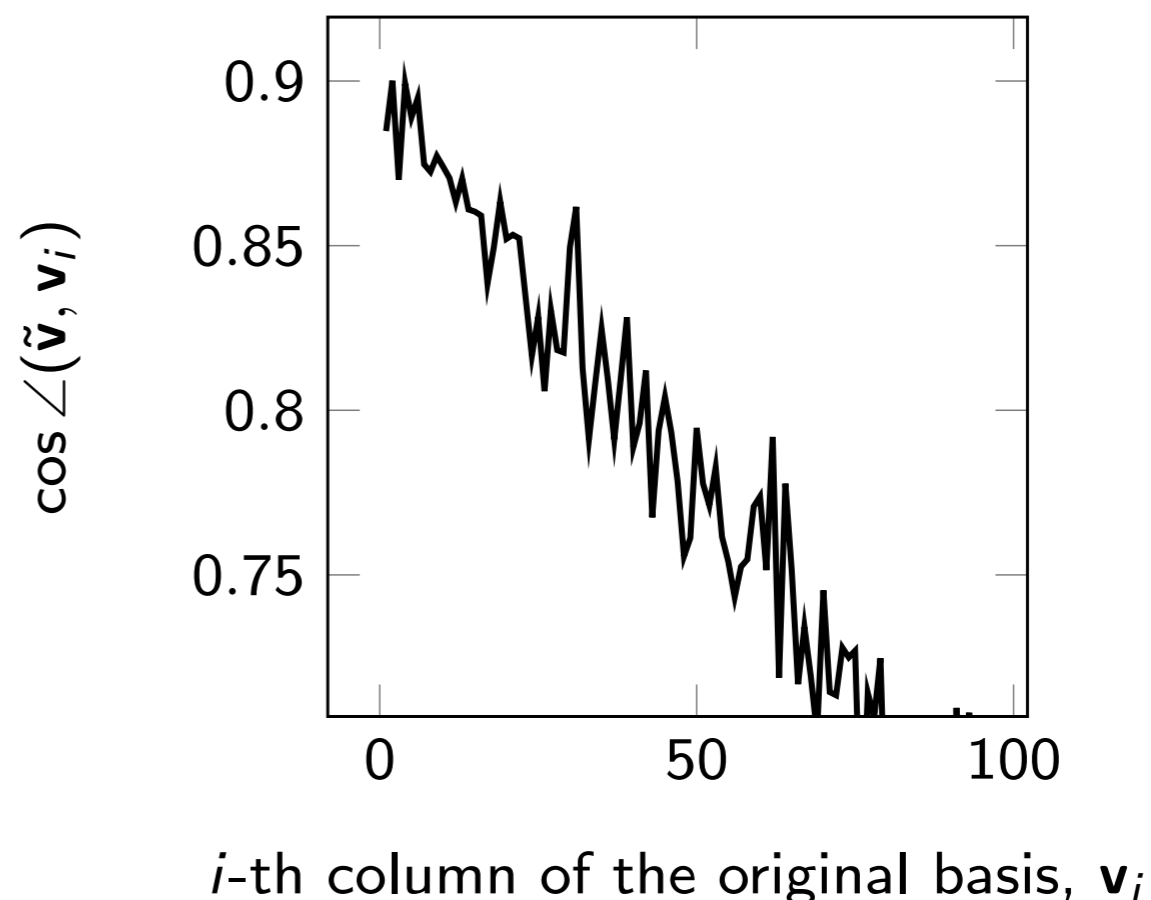
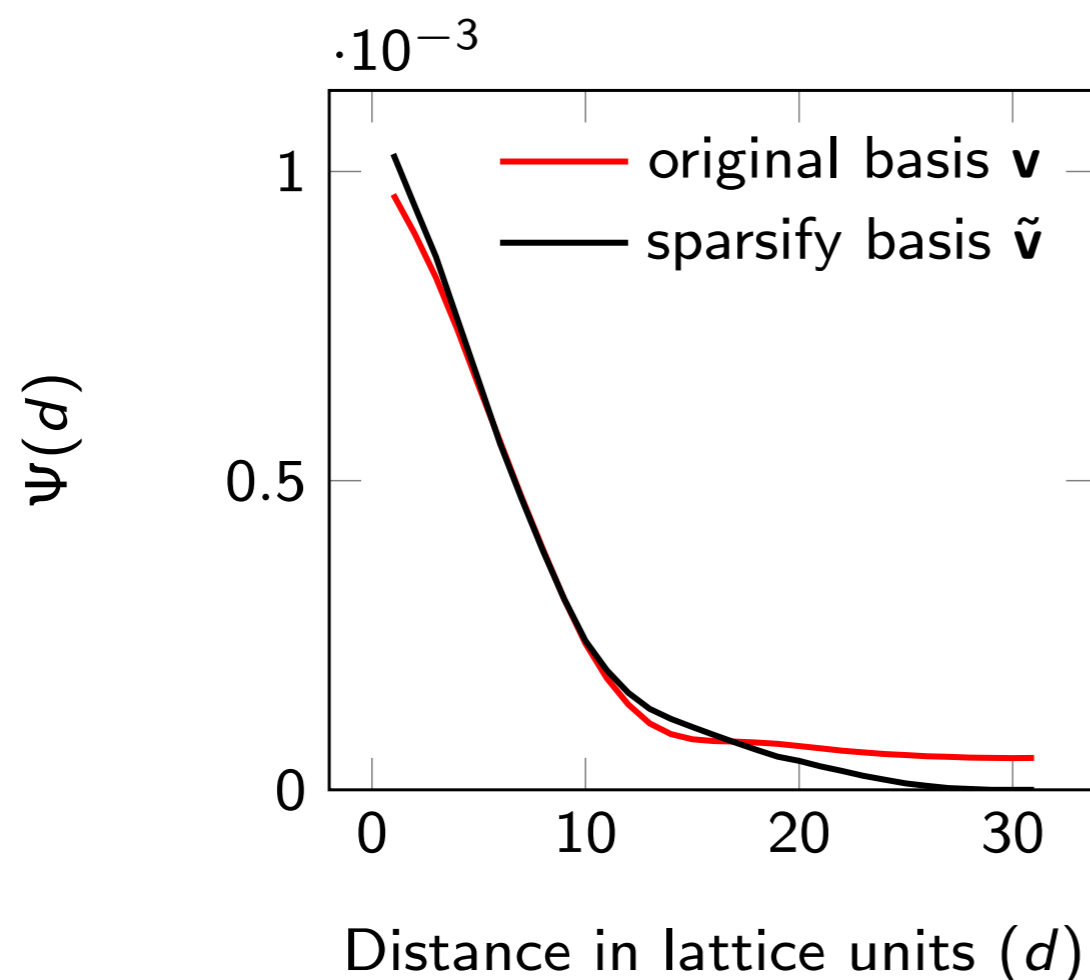
- HOSVD accelerates the contractions of tensors if near null rows, columns and fibers can be dropped from the core tensors
- In the studied cases, the core tensors can be reduced 20% in dimension
- But to keep the symmetries of the original tensors requires contracting the core tensors six times per original contraction
- At the end, we expect modest reduction of time, about 2-3 times

# Smearing basis sparsification

- Instead of sparsifying the tensors, we propose to sparsify the smearing basis  $V(t)$  that creates the tensors
- The smearing basis consist of the eigenvectors of the lattice Laplacian with the smallest energies
- On the local scale those eigenvectors look similar (local coherence)
- Given a regular partition of a lattice with fully connected components, the eigenvectors can be approximated as the linear combinations of vectors with support on one of the components

# Smearing basis sparsification

- We are able to generate a new basis  $\tilde{V}(t)$  that has similar spatial distribution to the original one, and
- lower energy eigenvectors are better represented



$$\psi(d) = \sqrt{\sum_{|\mathbf{x}-\mathbf{y}|=d, \alpha \in C, i \leq K} |\mathbf{v}_{\mathbf{x}\alpha i} \mathbf{v}_{\mathbf{y}\alpha i}^\dagger|^2}$$

- This basis will accelerate the contractions 16-64 times!



# Inverse Problem

with

*Joe Karpie*  
*Alexander Rothkopf*  
*Savas Zafeiropoulos*

## Factorization:

$$\mathfrak{M}(\nu, z^2) = \int_0^1 dx q(x, \mu) \mathcal{K}(x\nu, z^2 \mu^2) + \mathcal{O}(z^2 \Lambda_{qcd}^2)$$



**perturbative kernel (known to NLO)**

## Model for lattice matrix elements

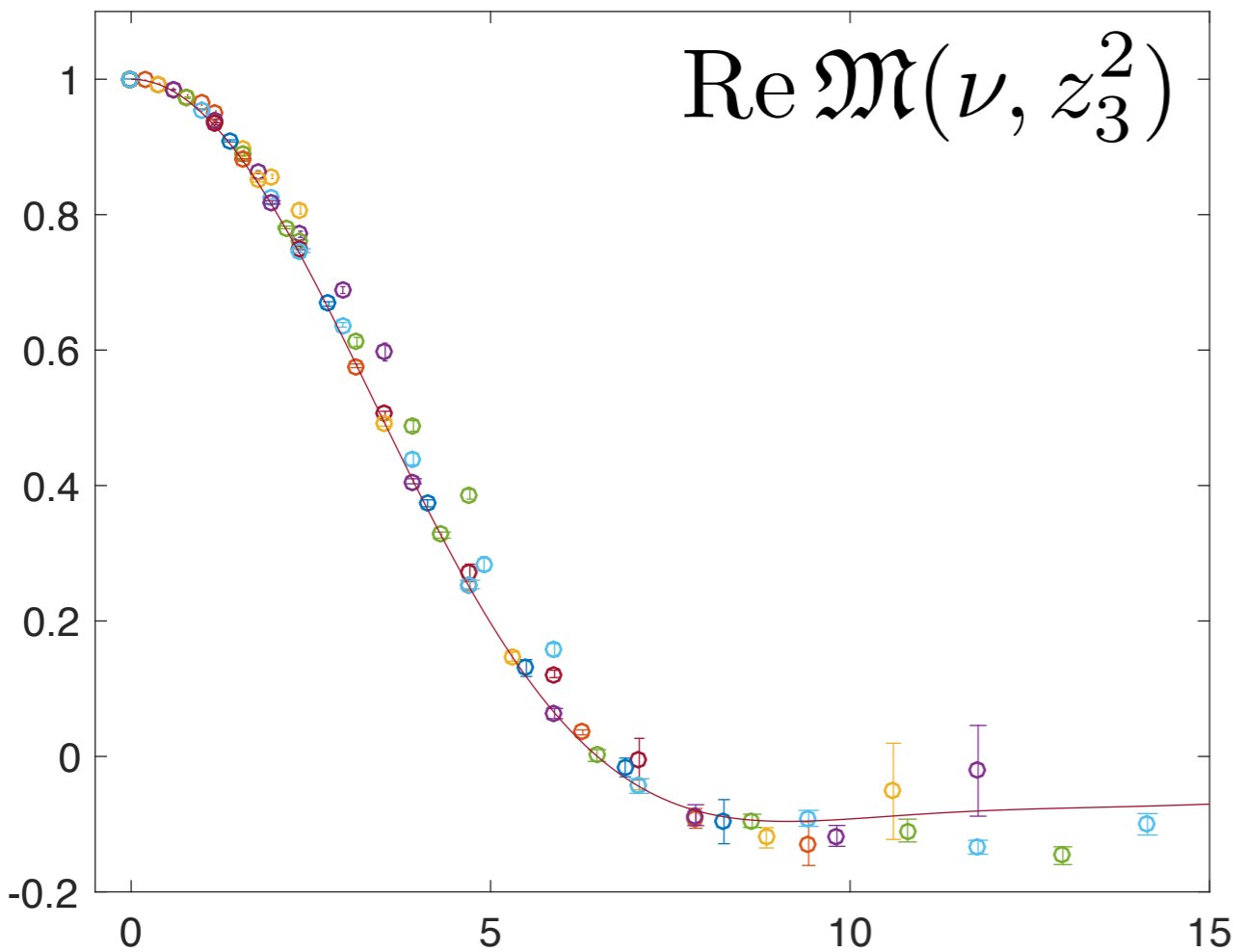
$$\mathfrak{M}(\nu, z^2) = N \int_0^1 dx q(x, \mu) \mathcal{K}(x\nu, z^2 \mu^2) e^{-z^2 \Lambda^2 x(1-x)}$$



**Primordial TMD**

$$\mathfrak{M}(\nu, z^2) = N \int_0^1 dx q(x, \mu) \mathcal{K}(x\nu, z^2 \mu^2) e^{-z^2 \Lambda^2 x(1-x)}$$

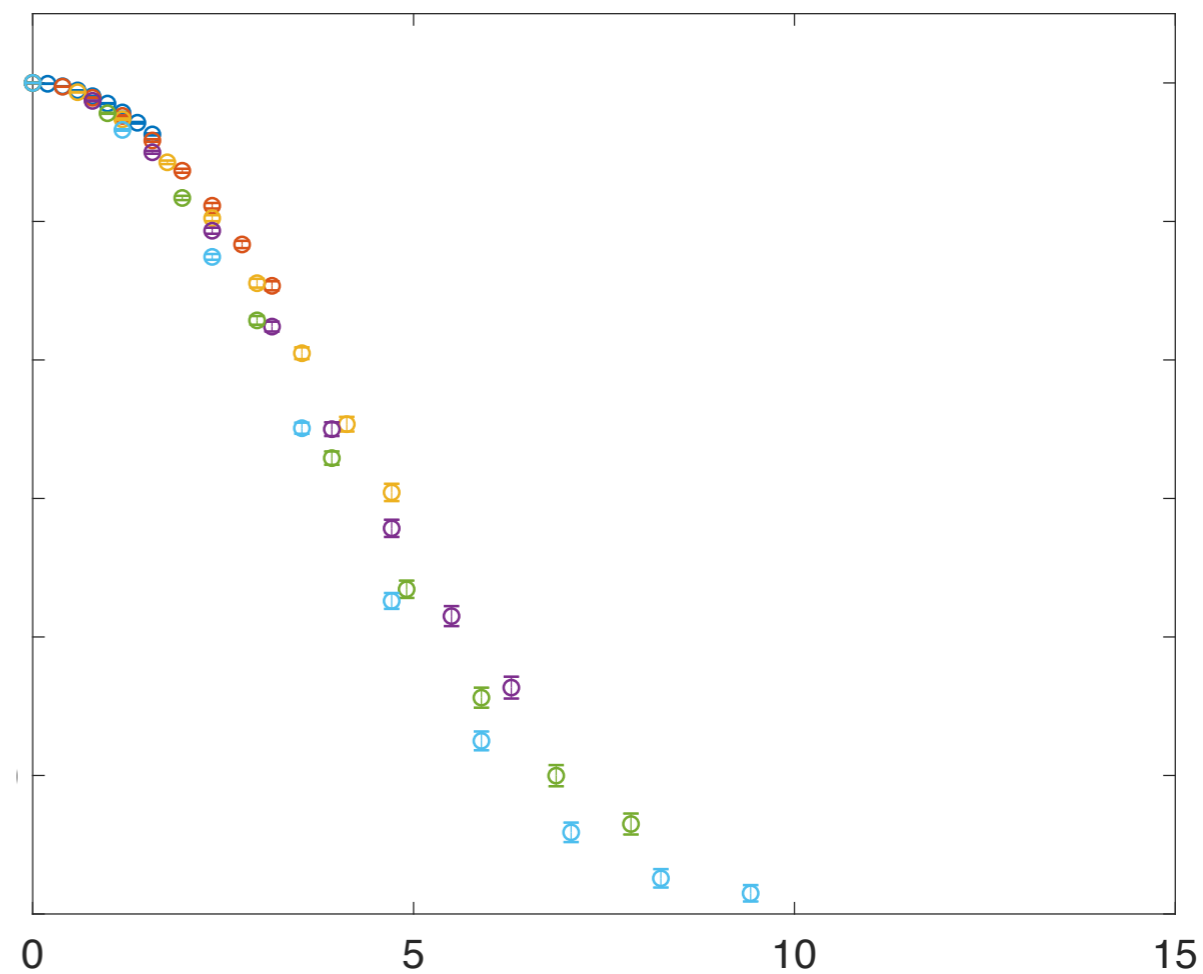
- Use the phenomenological PDFs from CJ15, NNPDF, etc. for  $q(x, \mu)$
- Explore models for the “TMD” that produce data that look similar to the lattice QCD results
- Reconstruct the PDF from mock matrix element data and assess the fidelity of the reconstruction by comparing with the known input PDF



**Lattice QCD result**

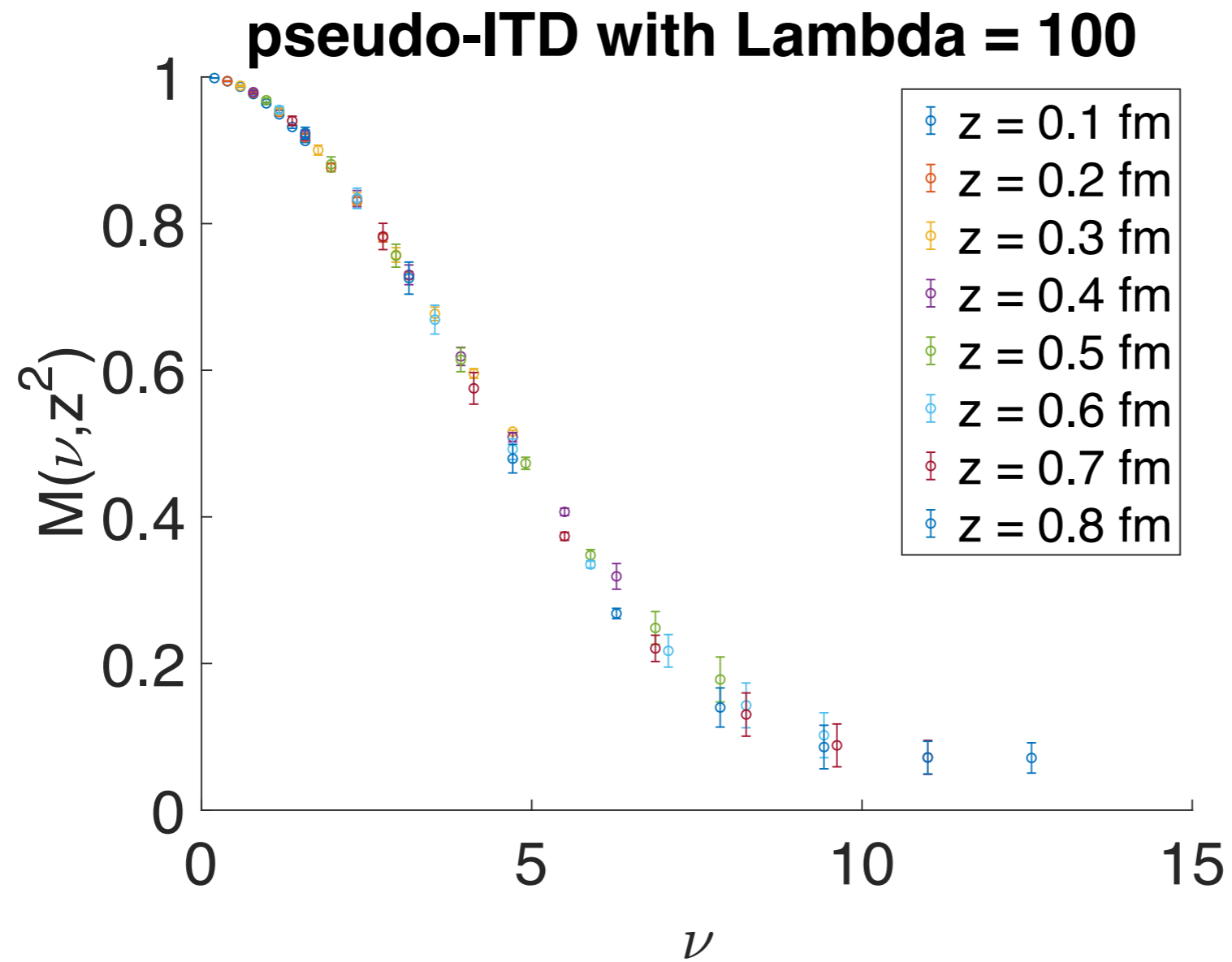
Orginos et al. (Phys.Rev. D96 (2017) no.9, 094503 )

$\nu$



**Mock data produced by model  
using CJ15 pdfs (  $\mu=2$  GeV )**

**LO kernel**



**NLO kernel with CJ15 PDFs**

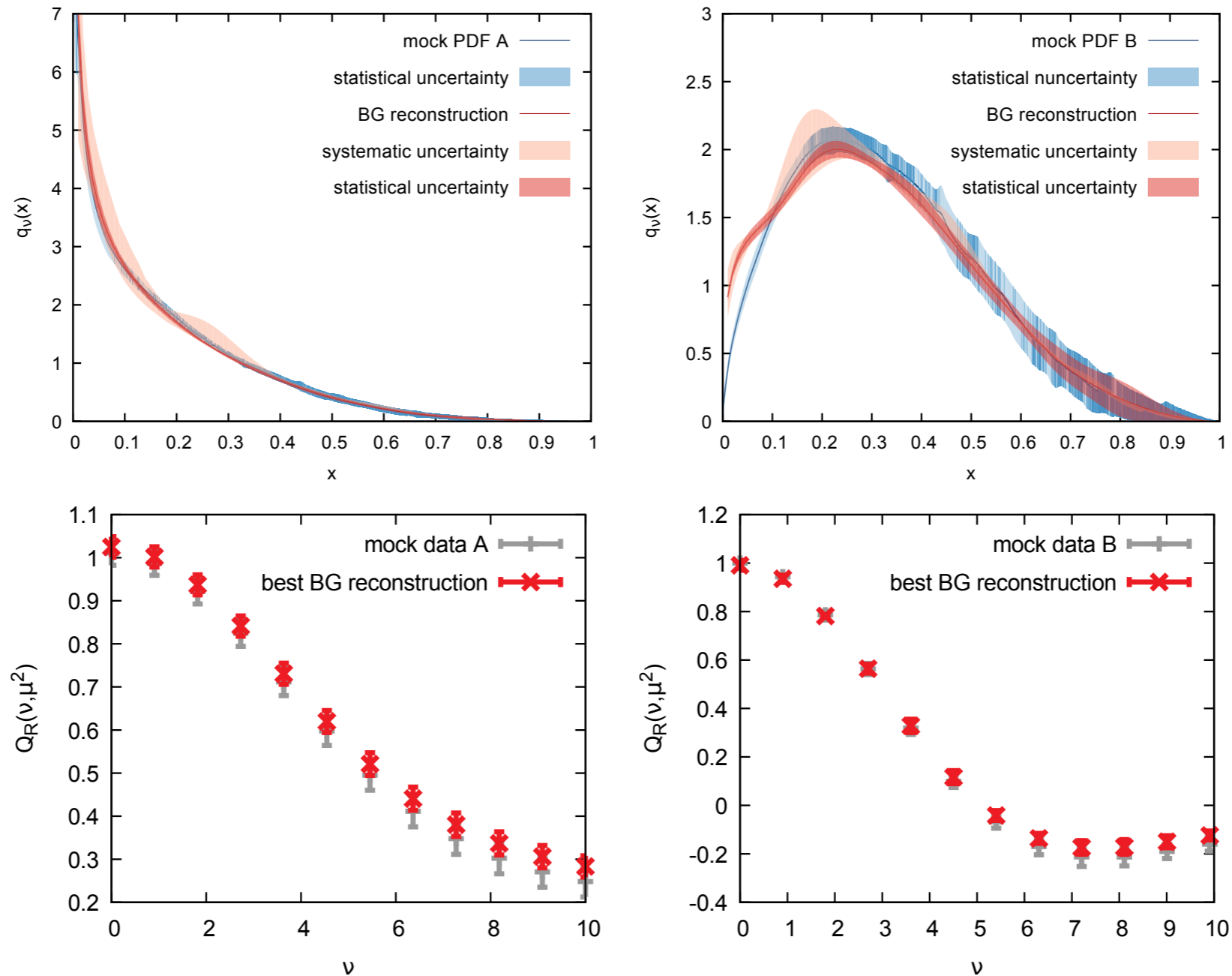
# Methods to be studied

- Backus-Gilbert
- Bayesian reconstruction
- Neural network parametrization
- Simple functional form parameterizations
- Understand how to properly handle and remove the polynomial corrections.

# Some tests

Karpie et. al JHEP 1904 (2019) 057

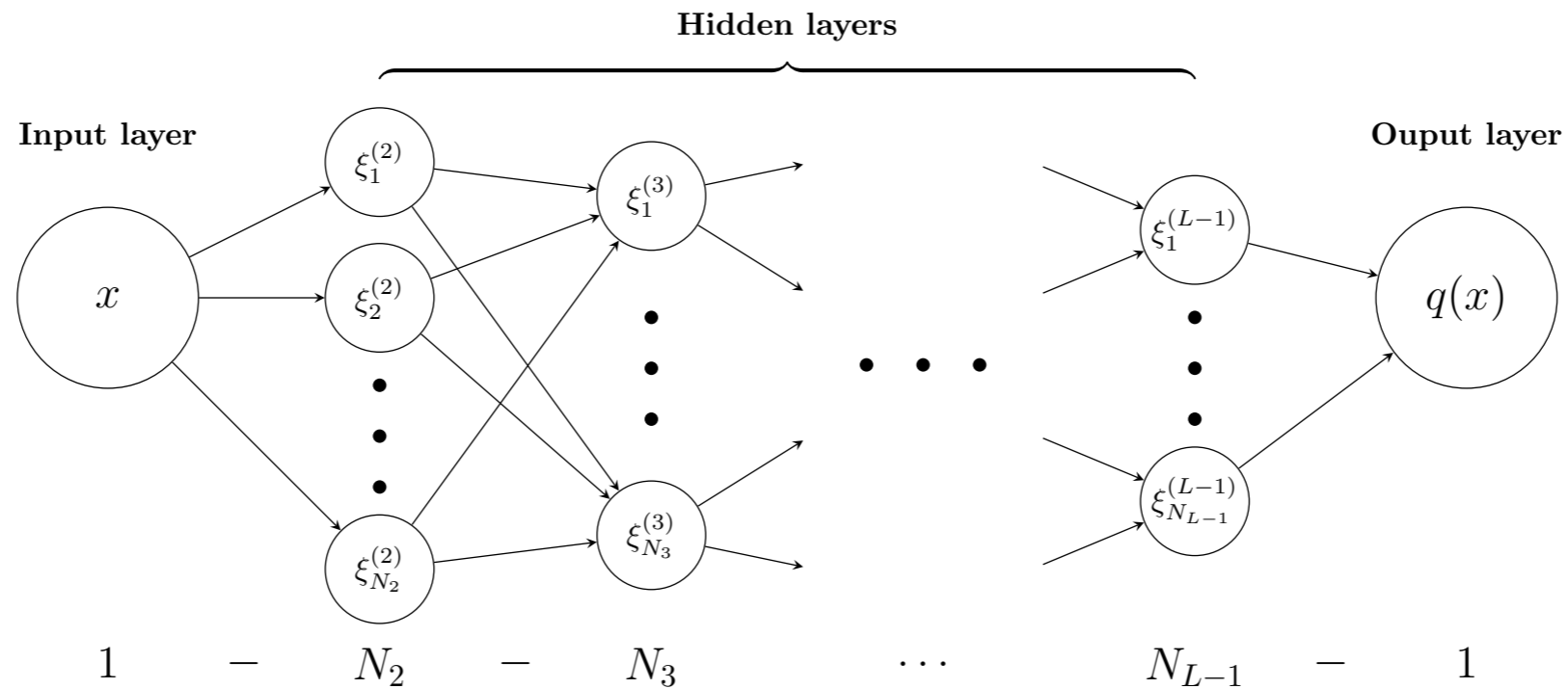
**Ignoring ( $z^2$ ) corrections**



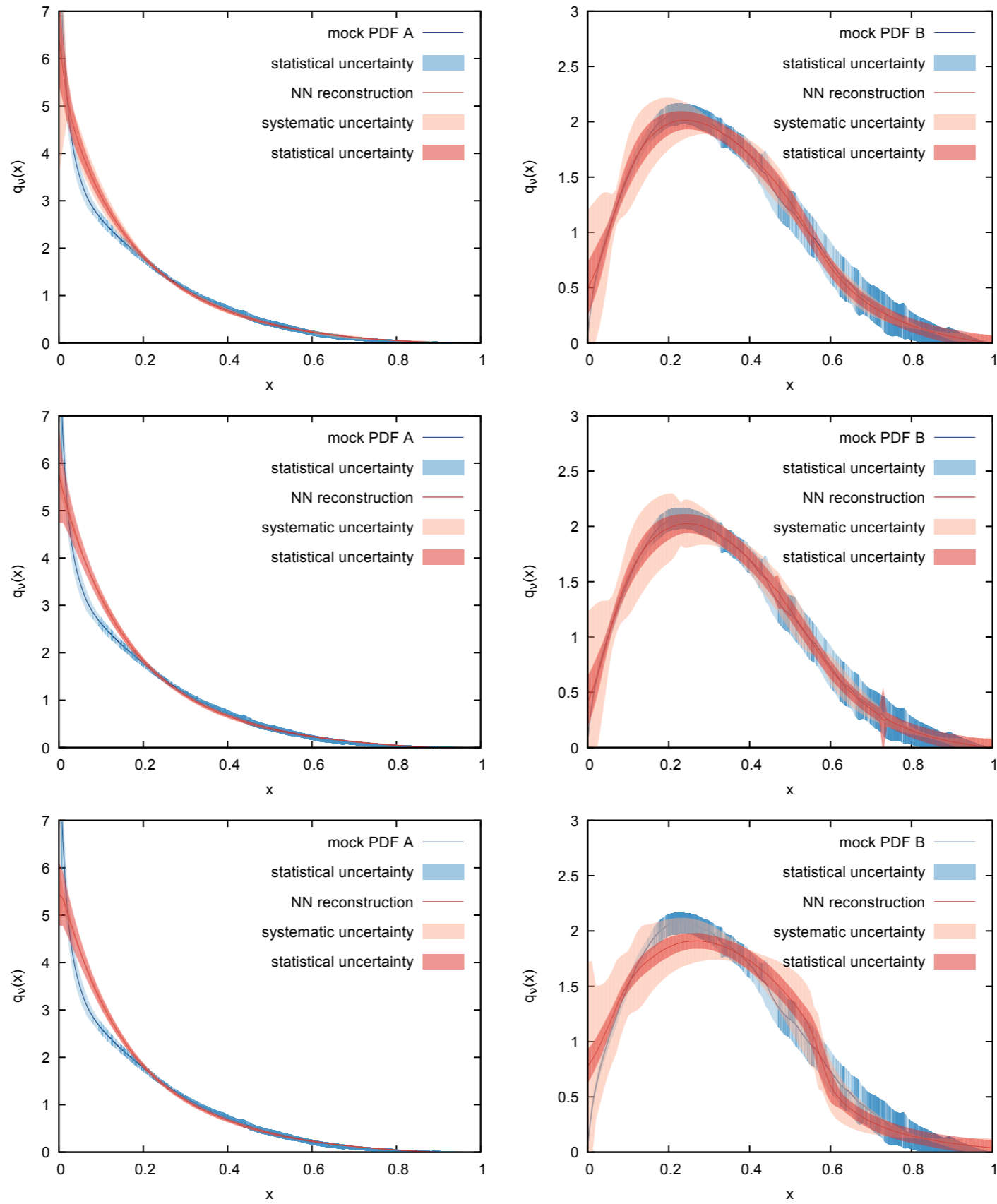
**Figure 6.**  $x$ -space PDF's reconstructed using the Backus-Gilbert (BG) method from  $N_\nu = 12$  Ioffe-time data points on the interval  $\nu = [0, 10]$  (top) as well as the input data (gray crosses) compared to the data arising from the reconstructed PDF (red crosses) in the bottom panels. The plots in the left column show the results for mock data based on a phenomenological PDF, while the right column from the modified scenario where the PDF vanishes at the origin. The reconstruction was performed with preconditioning exponents  $a = -0.35$  and  $b = 2$  for scenario A and  $a = 0.3$  and  $b = 2$  for scenario B.



# Neural-Network reconstruction



**Flexible parametrization of  $q(x)$**



Karpie et. al JHEP 1904 (2019) 057

**Figure 9.** The genetically trained neural nets. The blue band is the original data. The red band is the reconstructed PDF with statistical and systematic errors. The left column is with NNPDF data. The right column is with modified data. The first row has a network geometry of 1-3-1. The second row has a geometry of 1-4-1. The third row has a geometry of 1-2-2-1.

# Summary

- Ideas for improving the efficiency of the matrix element computation have been developed
  - Study of their performance is under way
- We have constructed models for producing mock data to study the effectiveness of various methods for solving the inverse problem at hand
  - Study the fidelity of the reconstruction is under way