

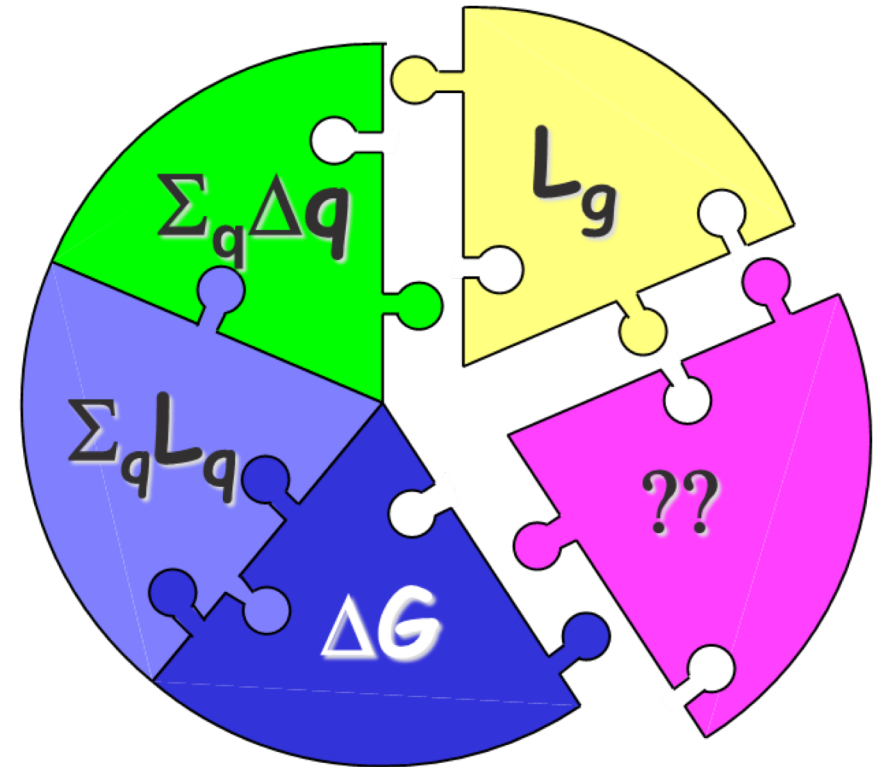
Before-Dinner Discussion

Yuri Kovchegov & Andrea Signori

Proton Spin

Proton Spin Puzzle

- Current experimental status of helicity PDFs
- OAM
- Lattice results
- Small-x
- What is the roadmap to solving the spin puzzle



$$S = \frac{1}{2} = \frac{1}{2} \Delta \Sigma + \Delta G + L$$

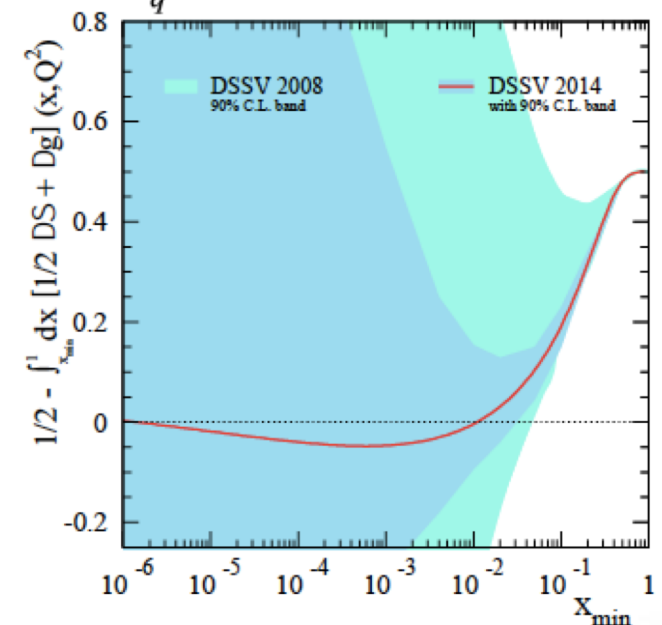
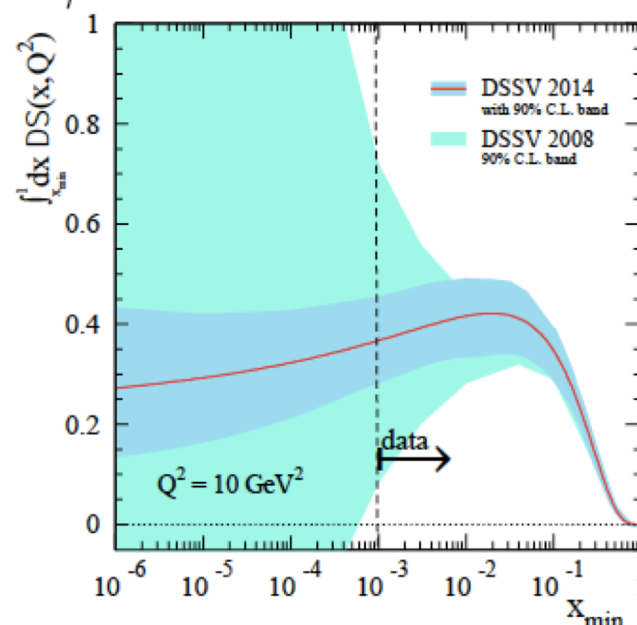
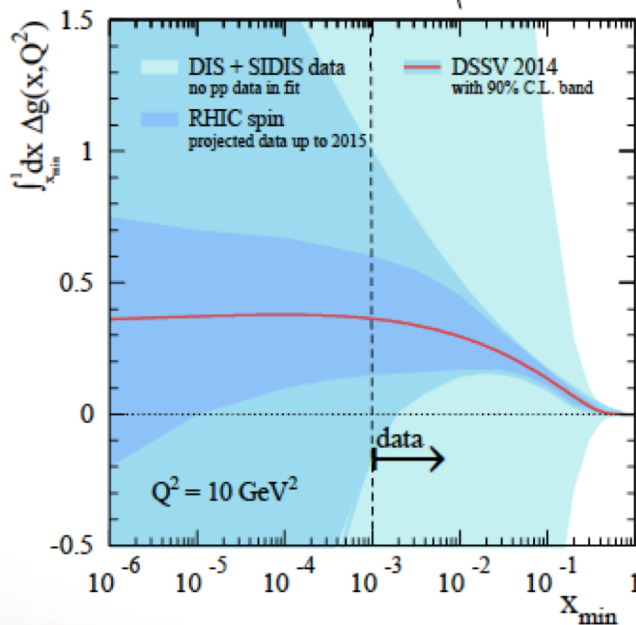
What forms the Spin of the Proton



Spin is more than the number $\frac{1}{2}$! It is the interplay between the intrinsic properties and interactions of quarks and gluons

What do we know:

$$\frac{1}{2}\hbar = \left\langle P, \frac{1}{2} \left| J_{QCD}^z \right| P, \frac{1}{2} \right\rangle = \underbrace{\frac{1}{2} \int_0^1 dx \Delta \Sigma(x, Q^2)}_{\text{total quark spin}} + \underbrace{\int_0^1 dx \Delta G(x, Q^2)}_{\text{gluon spin}} + \underbrace{\int_0^1 dx \left(\sum_q L_q^z + L_g^z \right)}_{\text{angular momentum}}$$



$\frac{1}{2}$ - Gluon
40%

Quarks
30%

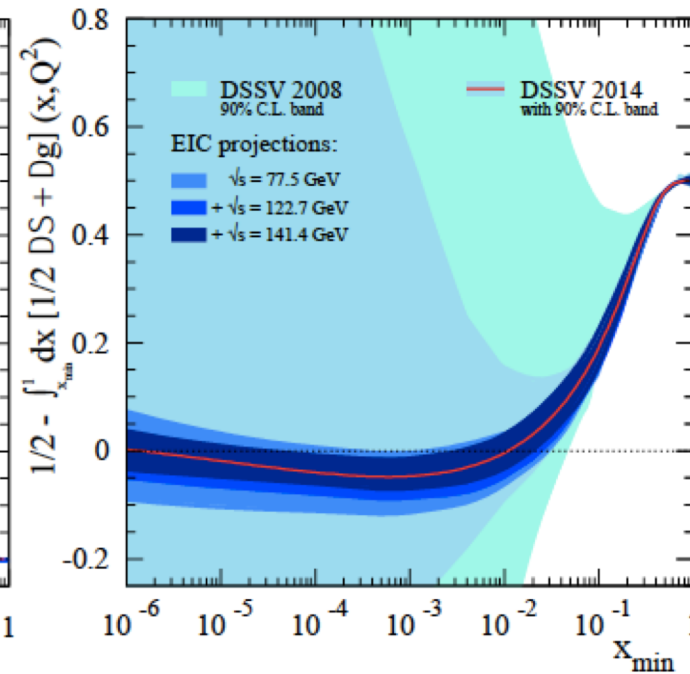
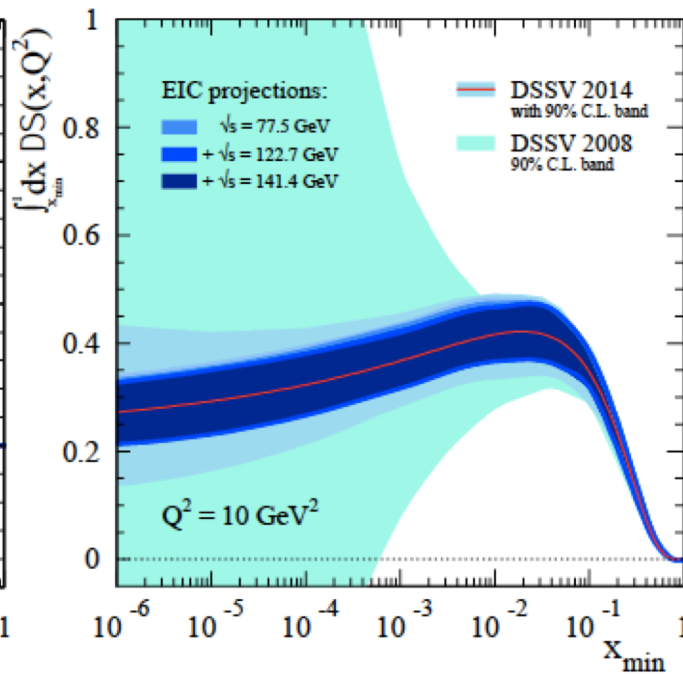
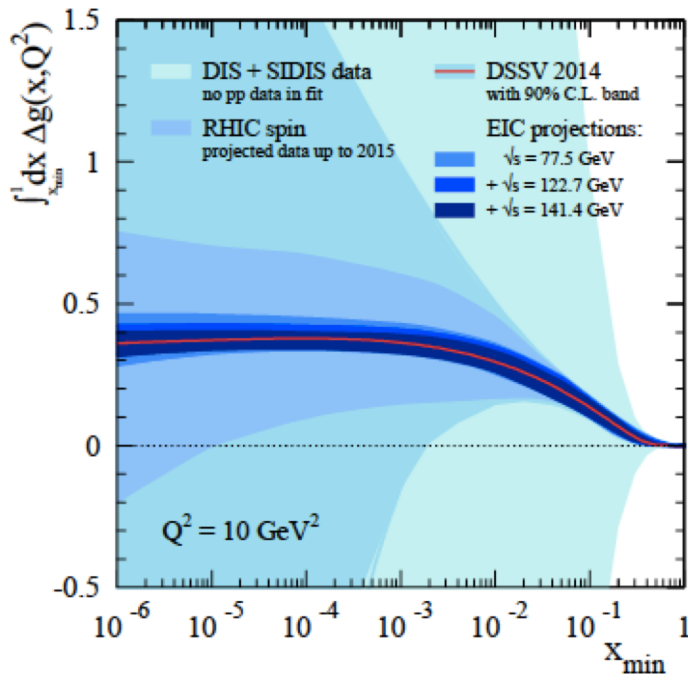
= orbital angular momentum

INT EIC program

Where does the Spin of the proton hide

“Helicity sum rule:”

$$\frac{1}{2}\hbar = \left\langle P, \frac{1}{2} \left| J_{QCD}^z \right| P, \frac{1}{2} \right\rangle = \underbrace{\sum_q \frac{1}{2} S_q^z}_{\text{total quark spin}} + \underbrace{S_g^z}_{\text{gluon spin}} + \underbrace{\sum_q L_q^z + L_g^z}_{\text{angular momentum}}$$



1/2 - Gluon - Quarks = orbital angular momentum



INT EIC program

'PDF' for OAM?

Take the staple-shaped Wilson line and define

$$L^{q,g} = \int dx \int d^2b_{\perp} d^2k_{\perp} (\vec{b}_{\perp} \times \vec{k}_{\perp})_z W^{q,g}(x, \vec{b}_{\perp}, \vec{k}_{\perp})$$



$$L^{q,g}(\boldsymbol{x}) = \int d^2b_{\perp} d^2k_{\perp} (\vec{b}_{\perp} \times \vec{k}_{\perp})_z W^{q,g}(\boldsymbol{x}, \vec{b}_{\perp}, \vec{k}_{\perp})$$

Agrees with [Harindranath, Kundu \(1998\)](#); [Hagler, Schafer \(1998\)](#) in the light-cone gauge.

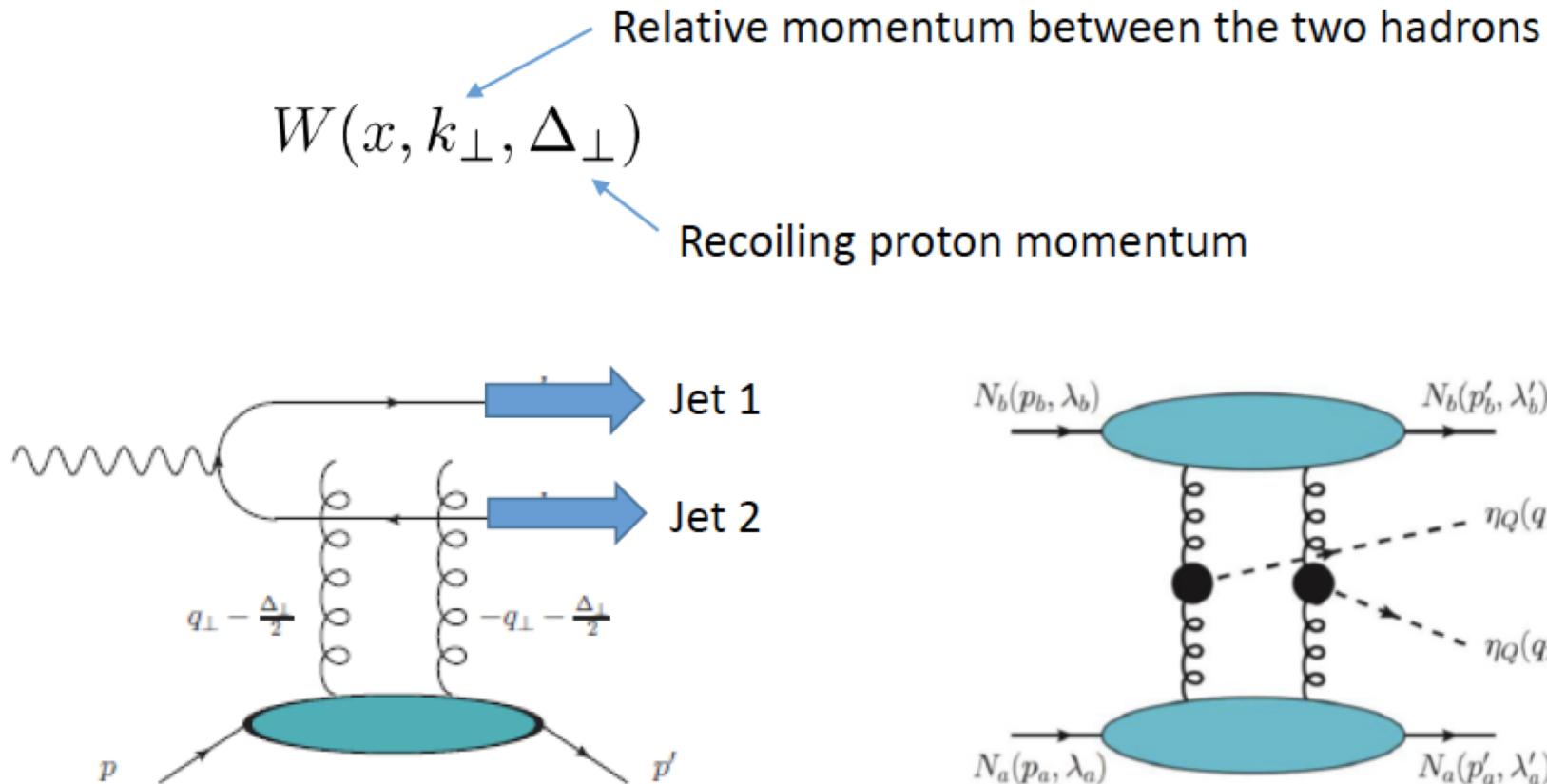
Warning: This is NOT the usual (twist-two) PDF.

Observables for OAM

Y. Hatta, INT EIC program, 2018

Essentially the measurement of Wigner/GTMD

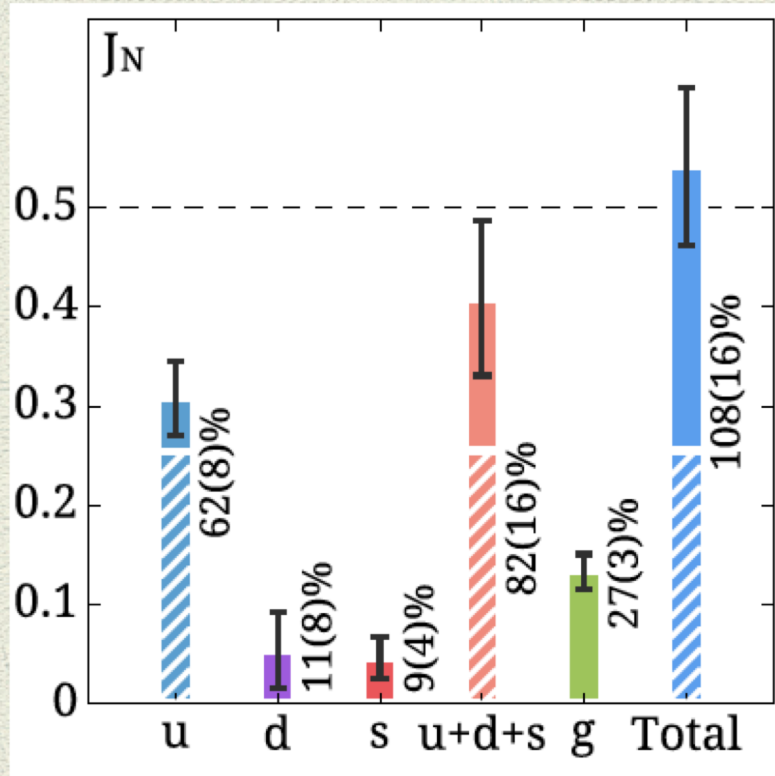
Tag two hadrons (jets) in the final state, together with the recoiling proton



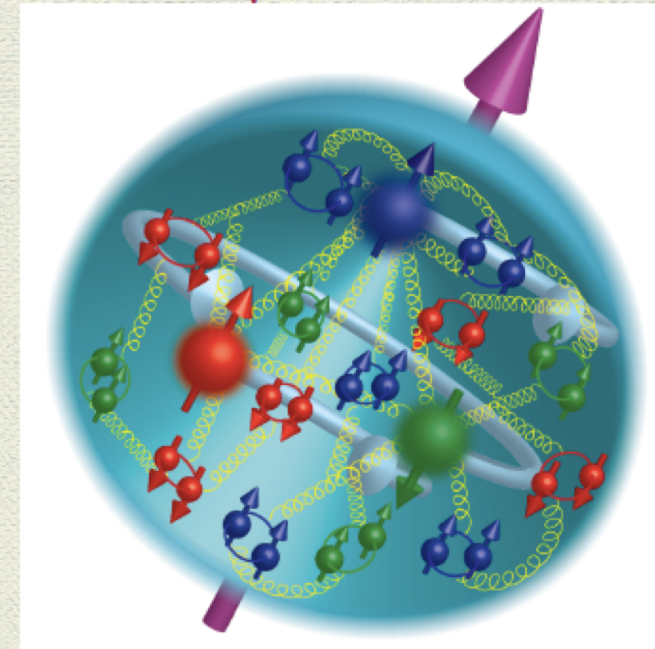
Ji, Yuan, Zhao (2017); YH, Nakagawa, Xiao, Yuan, Zhao (2017),
Bhattacharya, Metz, Zhou (2017); Bhattacharya, Metz, Ojha, Tsai, Zhou (2018)

The proton spin from LQCD

[C. Alexandrou et al., Phys. Rev. Lett. 119, 142002 (2017), [arXiv:1706.02973]]



Better understanding of the spin distribution



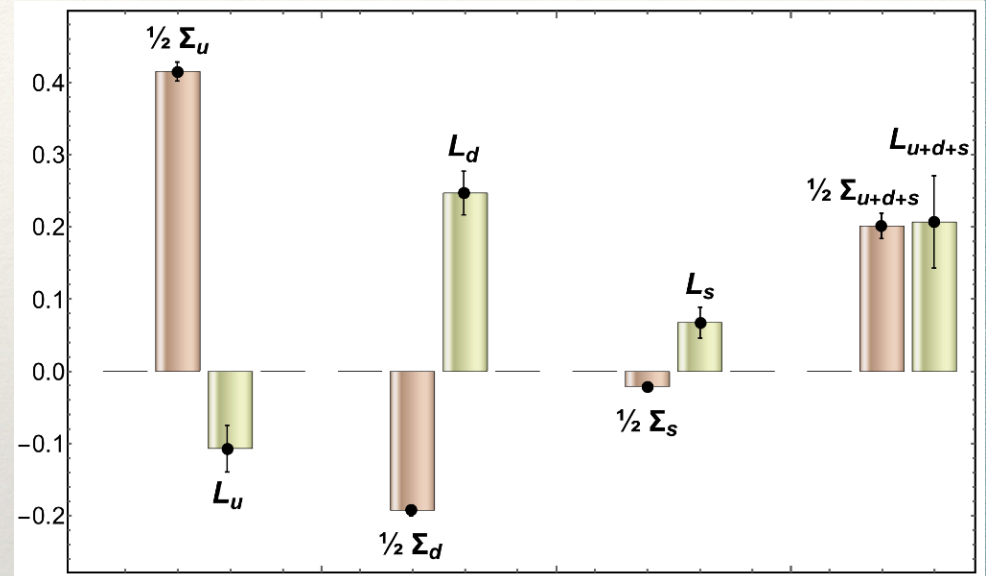
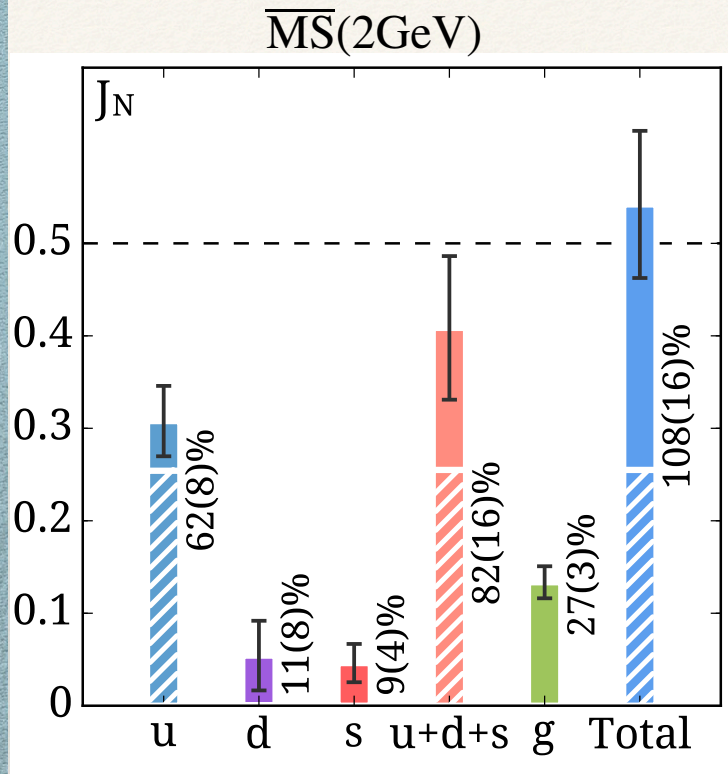
Designed by Z.-E. Meiziani

Striped segments: valence quark contributions (connected)
Solid segments: sea quark & gluon contributions (disconnected)

* Satisfaction of spin and momentum sum rule is not forced

Spin decomposition

[C. Alexandrou et al., Phys. Rev. Lett. 119, 142002 (2017), [arXiv:1706.02973]]



Quark orbital angular momentum extracted indirectly ($L_q = J_q - \Sigma_q$)

- ★ Striped segments: connected contributions (connected)
- ★ Solid segments: disconnected contributions (quark & gluon)

Satisfaction of spin and momentum sum rule is not forced

Helicity PDFs at Small-x

- Theoretical calculations by Bartels, Ermolaev and Ryskin (BER, 1996) and by YK, Pitonyak and Sievert (2015-17, KPS).
- KPS results (large- N_c):

$$\Delta q(x, Q^2) \sim \left(\frac{1}{x}\right)^{\alpha_h^q} \quad \text{with} \quad \alpha_h^q = \frac{4}{\sqrt{3}} \sqrt{\frac{\alpha_s N_c}{2\pi}} \approx 2.31 \sqrt{\frac{\alpha_s N_c}{2\pi}}$$
$$\Delta G(x, Q^2) \sim \left(\frac{1}{x}\right)^{\alpha_h^G} \quad \text{with} \quad \alpha_h^G = \frac{13}{4\sqrt{3}} \sqrt{\frac{\alpha_s N_c}{2\pi}} \approx 1.88 \sqrt{\frac{\alpha_s N_c}{2\pi}}$$

- BER results at large- N_c (for both quark and gluon helicity distributions):

$$\alpha_h^{BER} = \sqrt{\frac{17+\sqrt{97}}{2}} \sqrt{\frac{\alpha_s N_c}{2\pi}} \approx 3.66 \sqrt{\frac{\alpha_s N_c}{2\pi}}$$

For finite N_c and $N_f=4$ BER have $3.66 \rightarrow 3.45$.

Helicity PDFs at Small-x

$$PDF(x) \sim \left(\frac{1}{x}\right)^\alpha$$

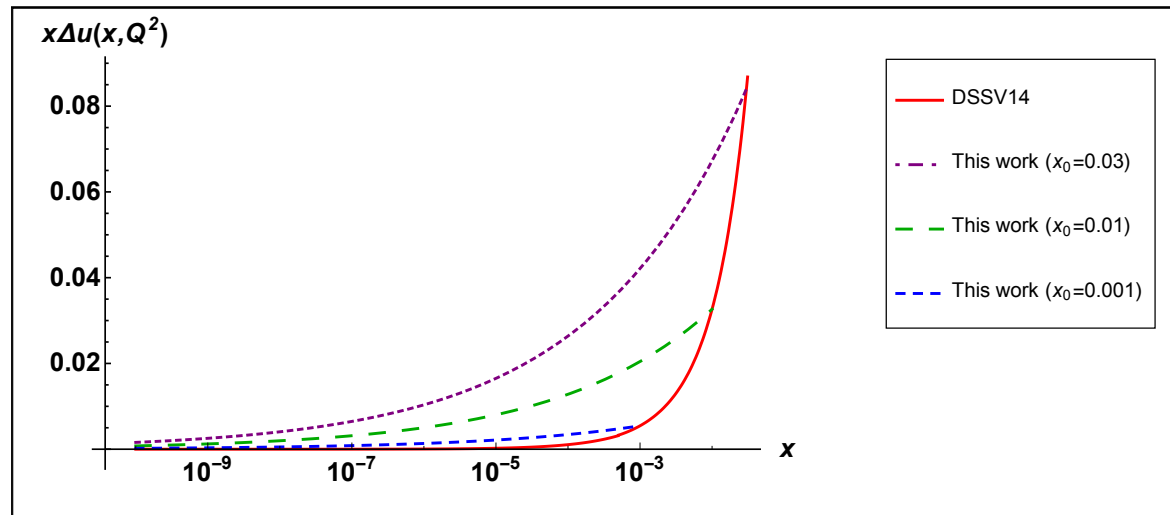
- The summary of the existing powers of x is

Observable	Evolution	Intercept	$Q^2 = 3 \text{ GeV}^2$ $\alpha_s = 0.343$	$Q^2 = 10 \text{ GeV}^2$ $\alpha_s = 0.249$	$Q^2 = 87 \text{ GeV}^2$ $\alpha_s = 0.18$
Unpolarized flavor singlet structure function F_2	LO BFKL Pomeron	$1 + \frac{\alpha_s N_c}{\pi} 4 \ln 2$	1.908	1.659	1.477
Unpolarized flavor non-singlet structure function F_2	Reggeon	$\sqrt{\frac{2\alpha_s C_F}{\pi}}$	0.540	0.460	0.391
Flavor singlet structure function g_1^S	us (Pure Glue)	$2.31 \sqrt{\frac{\alpha_s N_c}{2\pi}}$	0.936	0.797	0.678
	BER (Pure Glue)	$3.66 \sqrt{\frac{\alpha_s N_c}{2\pi}}$	1.481	1.262	1.073
	BER ($N_f = 4$)	$3.45 \sqrt{\frac{\alpha_s N_c}{2\pi}}$	1.400	1.190	1.011
Flavor non-singlet structure function g_1^{NS}	BER and us (large- N_c)	$\sqrt{\frac{\alpha_s N_c}{\pi}}$	0.572	0.488	0.415

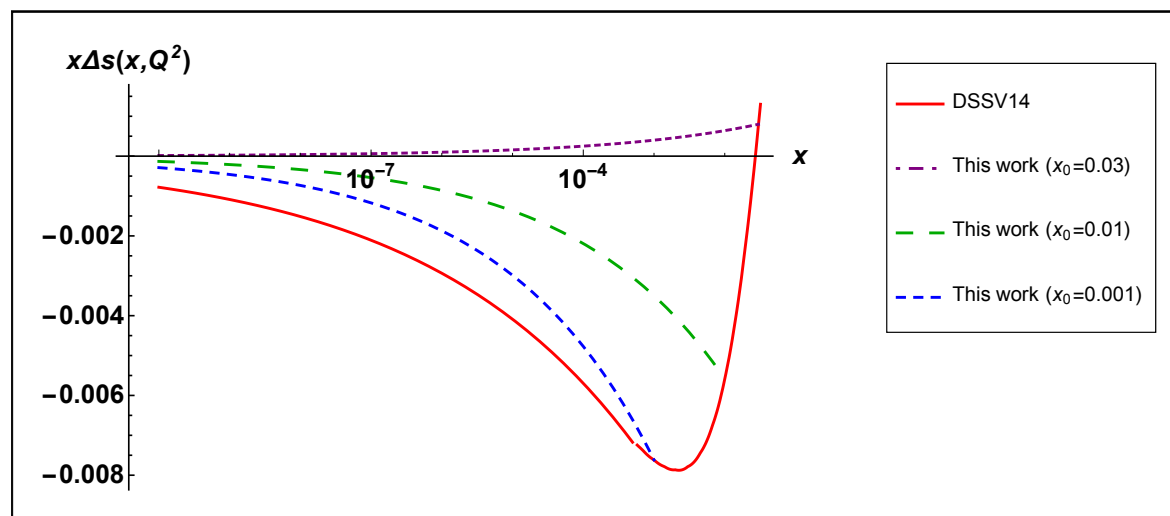
- Similar disagreement exists for OAM distributions at small x.
- However, even the smaller KPS powers lead to potentially important contributions to the proton spin coming from small x. (next)

Impact on proton spin

- We have attached a $\Delta\tilde{\Sigma}(x, Q^2) = N x^{-\alpha_h}$ curve to the existing hPDF's fits at some ad hoc small value of x labeled x_0 :

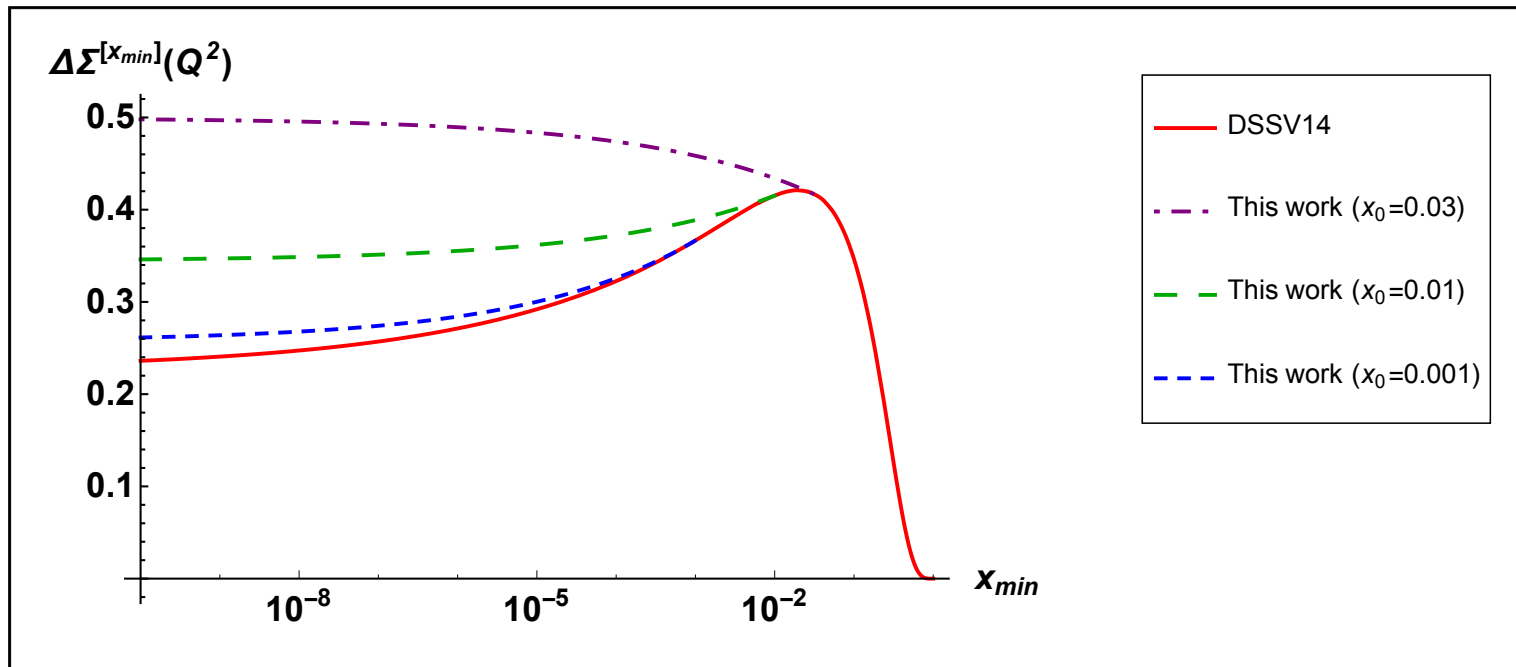


“ballpark”
phenomenology



Impact on proton spin

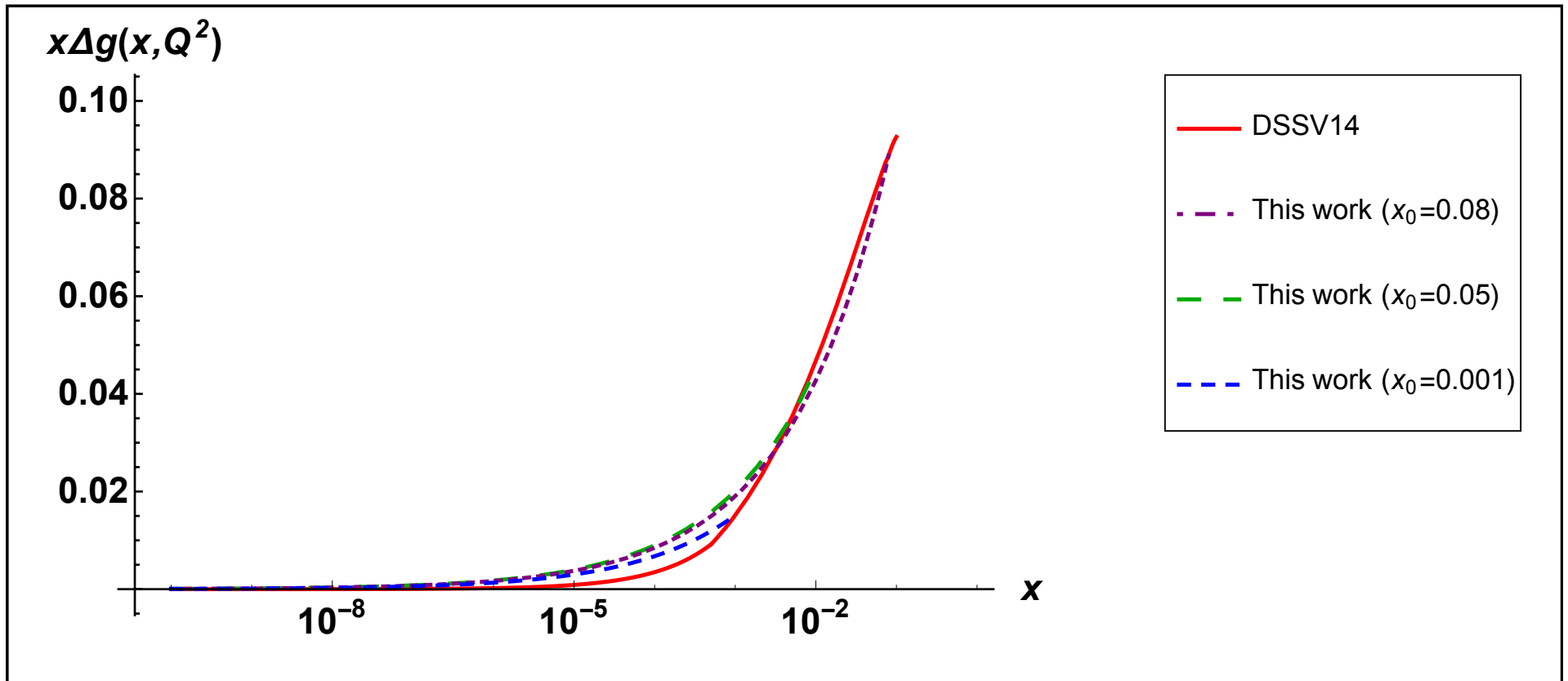
- Defining $\Delta\Sigma^{[x_{min}]}(Q^2) \equiv \int_{x_{min}}^1 dx \Delta\Sigma(x, Q^2)$ we plot it for $x_0=0.03, 0.01, 0.001$:



- We observe a moderate to significant enhancement of quark spin.
- More detailed phenomenology is needed in the future.

Impact of our ΔG on the proton spin

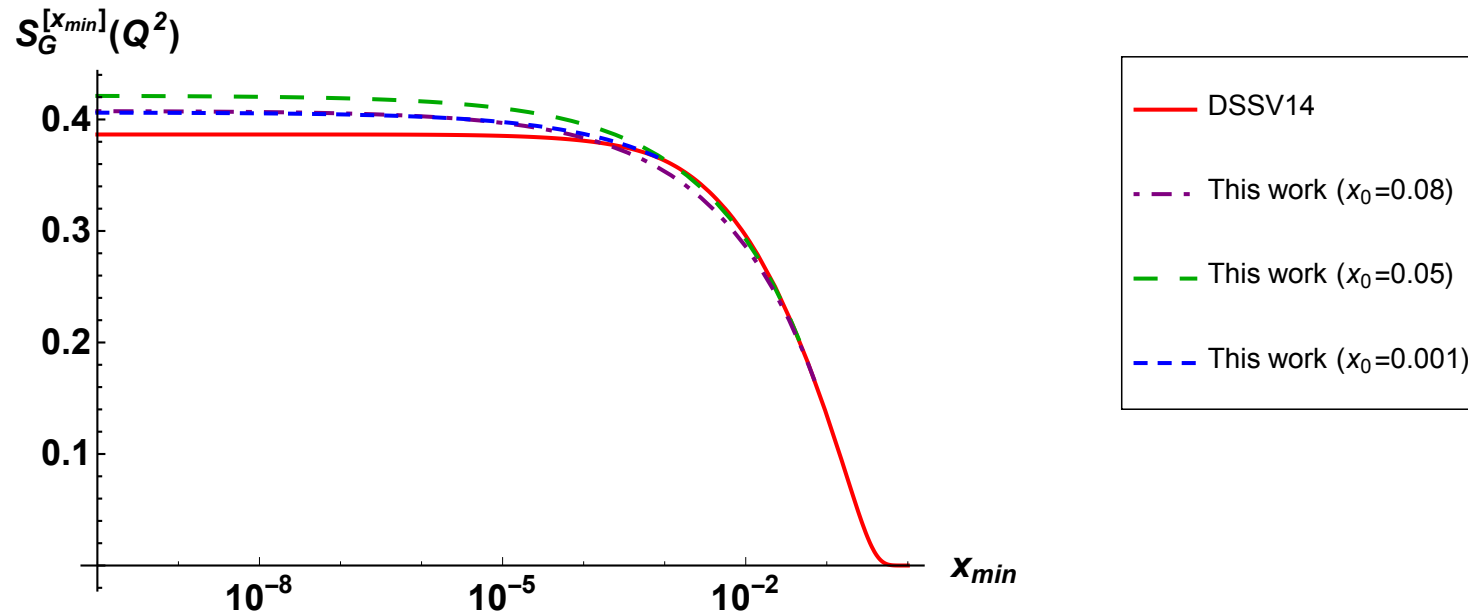
- We have attached a $\Delta\tilde{G}(x, Q^2) = N x^{-\alpha_h^G}$ curve to the existing hPDF's fits at some ad hoc small value of x labeled x_0 :



“ballpark”
phenomenology

Impact of our ΔG on the proton spin

- Defining $S_G^{[x_{min}]}(Q^2) \equiv \int_{x_{min}}^1 dx \Delta G(x, Q^2)$ we plot it for $x_0=0.08, 0.05, 0.001$:



- We observe a moderate enhancement of gluon spin.
- More detailed phenomenology is needed in the future.