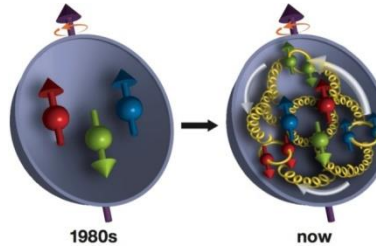


QCD Evolution 2019

13–17 May 2019, Physics Division, Argonne National Laboratory, IL USA



Mechanical properties of spin-1 hadrons

Based on [\[Cosyn, Cotogno, Freese, C.L., arXiv:1903.00408\]](#)

Cédric Lorcé



May 16, ANL, Lemont, USA

Outline



- 1. Energy-momentum tensor and gravitational FFs**
- 2. Mass decomposition and balance equations**
- 3. Angular momentum sum rule**
- 4. Mass radius and inertia tensor**
- 5. Link with Generalized Parton Distributions**
- 6. Summary**

Origin of mass and spin?

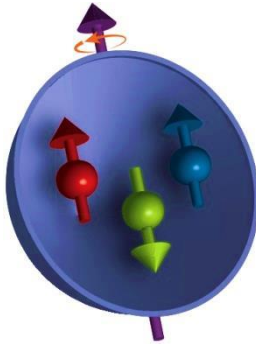
Non-relativistic picture
dominated by **constituents**

Mass

Spin

Until
~ 1980

Spectroscopy



$$M_N \sim \sum_Q M_Q + E_{\text{binding}}$$

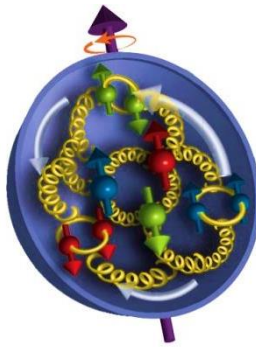
~ 102 % ~ - 2 %

$$J_z^N \sim \sum_Q S_z^Q$$

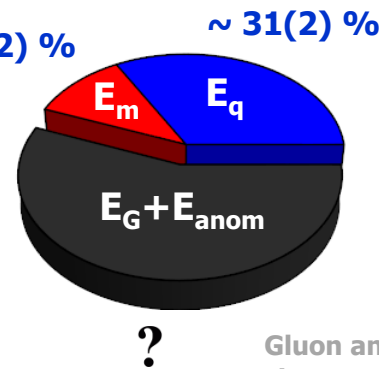
~ 100 %

Relativistic picture
dominated by **dynamics**

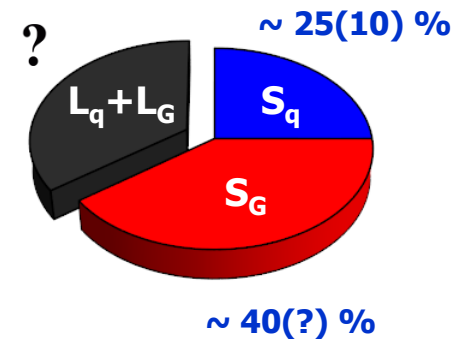
Now
High-energy
scattering



Higgs
mechanism
~ 13(2) %



Orbital angular
momentum



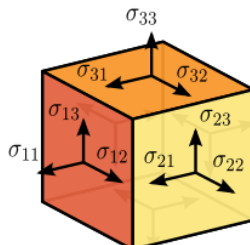
Energy-momentum tensor (EMT)

Mass, spin and pressure all encoded in

$$T^{\mu\nu} = \begin{bmatrix} \text{Energy density} & \text{Momentum density} & & \\ T^{00} & T^{01} & T^{02} & T^{03} \\ T^{10} & T^{11} & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} \\ \text{Energy flux} & \text{Momentum flux} & & \end{bmatrix}$$

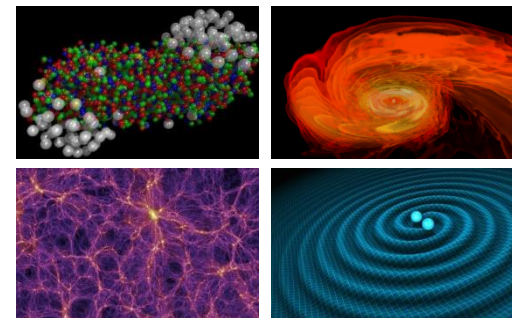
Shear stress

Normal stress (pressure)



Key concept for

- Nucleon mechanical properties
- Quark-gluon plasma
- Relativistic hydrodynamics
- Stellar structure and dynamics
- Cosmology
- Gravitational waves
- Modified theories of gravitation
- ...



Gravitational form factors (GFFs)

See also talks by P. Schweitzer, Y. Hatta, and A. Freese

Symmetrized variables $P = \frac{p' + p}{2}, \quad \Delta = p' - p, \quad t = \Delta^2$

Spin-0

[Pagels (1966)]

[Donoghue, Leutwyler (1991)]

[Ji (1996)]

$$\langle p' | T_a^{\mu\nu}(0) | p \rangle = 2M \left[\frac{P^\mu P^\nu}{M} A_a(t) + \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{M} C_a(t) + M g^{\mu\nu} \bar{C}_a(t) \right]$$

Non-conserved

→ $\sum_a \bar{C}_a(t) = 0$

Spin-1/2

[Kobzarev, Okun (1962)]

[Pagels (1966)]

[Ji (1996)]

[Bakker, Leader, Trueman (2004)]

[Leader, C.L. (2014)]

$$\langle p', s' | T_a^{\mu\nu}(0) | p, s \rangle = \bar{u}(p', s') \Gamma_a^{\mu\nu}(P, \Delta) u(p, s)$$

$$\Gamma_a^{\mu\nu}(P, \Delta) = \frac{P^\mu P^\nu}{M} A_a(t) + \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{M} C_a(t) + M g^{\mu\nu} \bar{C}_a(t) \\ + \frac{P^{\{\mu} i \sigma^{\nu\} \lambda} \Delta_\lambda}{4M} J_a(t) + \frac{P^{[\mu} i \sigma^{\nu] \lambda} \Delta_\lambda}{4M} D_a(t)$$

Gravitational form factors (GFFs)

[Holstein (2006)]

[Abidin, Carlson (2008)]

[Taneja, Kathuria, Liuti, Goldstein (2012)]

[Cosyn, Cotogno, Freese, C.L., arXiv:1903.00408]

[Polyakov, Sun, arXiv:1903.02738]

Spin-1

$$\langle p', s' | T_a^{\mu\nu}(0) | p, s \rangle = 2M \epsilon_\alpha^*(p', s') \mathcal{M}_a^{\mu\nu;\alpha\beta}(P, \Delta) \epsilon_\beta(p, s)$$

$$\begin{aligned} \mathcal{M}_a^{\mu\nu;\alpha\beta}(P, \Delta) = & -g^{\alpha\beta} \left[\frac{P^\mu P^\nu}{M} \mathcal{G}_1^a(t) + \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{M} \left(\frac{1}{4} \mathcal{G}_3^a(t) \right) + M g^{\mu\nu} \left(-\frac{1}{2} \mathcal{G}_8^a(t) \right) \right] \\ & + \frac{P^{\{\mu} g^{\nu\}[\alpha} \Delta^{\beta]}}{4M} \mathcal{G}_5^a(t) + \frac{P^{[\mu} g^{\nu][\alpha} \Delta^{\beta]}}{4M} \mathcal{G}_{10}^a(t) \\ & + M g^{\alpha\{\mu} g^{\nu\}\beta} \left(\frac{1}{4} \mathcal{G}_7^a(t) \right) \\ & + \frac{\Delta^\alpha \Delta^\beta}{2M^2} \left[\frac{P^\mu P^\nu}{M} \mathcal{G}_2^a(t) + \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{M} \left(\frac{1}{4} \mathcal{G}_4^a(t) \right) + M g^{\mu\nu} \left(\frac{1}{2} \mathcal{G}_9^a(t) \right) \right] \\ & + \left[\frac{\Delta^{\{\mu} g^{\nu\}\{\alpha} \Delta^{\beta\}}}{4M} - g^{\alpha\{\mu} g^{\nu\}\beta} \frac{\Delta^2}{4M} - g^{\mu\nu} \frac{\Delta^\alpha \Delta^\beta}{2M} \right] \left(\frac{1}{2} \mathcal{G}_6^a(t) \right) + \frac{\Delta^{[\mu} g^{\nu]\{\alpha} \Delta^{\beta\}}}{4M} \mathcal{G}_{11}^a(t) \end{aligned}$$

Spin-J



Target polarization

Spin-0

$$1 = 1$$

Spin-1/2

Dipole (vector)

$$u(p, s)\bar{u}(p, s') = (\not{p} + M) \frac{\delta_{s's} \mathbb{1} - \mathcal{S}_{s's}^\mu(p) \gamma_\mu \gamma_5}{2}$$

[Bouchiat, Michel (1958)]

Spin-1

[Joos, Kramer (1964)]

[Zwanziger (1965)]

$$\epsilon_\beta(p, s)\epsilon_\alpha^*(p, s') = -\delta_{s's} \frac{1}{3} \left(g_{\beta\alpha} - \frac{p_\beta p_\alpha}{M^2} \right) + \mathcal{S}_{s's}^\mu(p) \frac{i}{2M} \varepsilon_{\mu\beta\alpha\nu} p^\nu - \mathcal{T}_{s's}^{\mu\nu}(p) g_{\mu\beta} g_{\nu\alpha}$$

Dipole (vector)

Quadrupole (tensor)

Onshell constraints $p_\mu \mathcal{S}_{s's}^\mu(p) = 0,$ $p_\mu \mathcal{T}_{s's}^{\mu\nu}(p) = 0,$ $\mathcal{T}_{s's}^{[\mu\nu]}(p) = 0,$ $g_{\mu\nu} \mathcal{T}_{s's}^{\mu\nu}(p) = 0$

Generalization to off-forward case



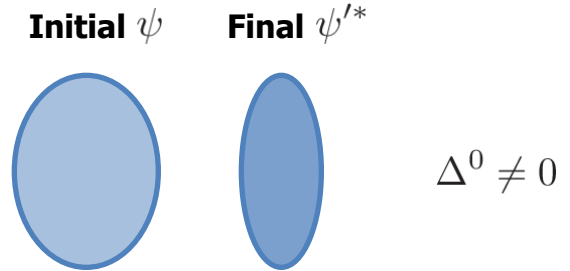
[C.L. (2018)]

[Cosyn, Cotogno, Freese, C.L., arXiv:1903.00408]

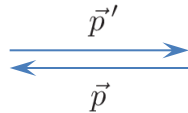
3D distribution in Breit frame

Lorentz factors

$$p^0 = \gamma M, \quad p'^0 = \gamma' M$$



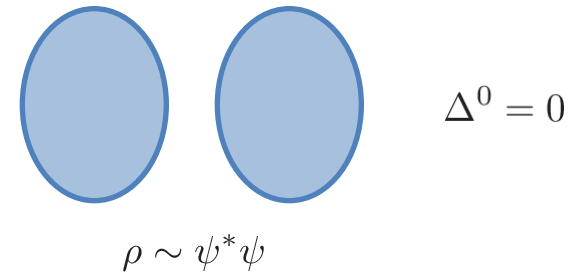
Breit frame



$$\vec{P} = \vec{0}$$



$$\Delta^0 = \frac{\vec{P} \cdot \vec{\Delta}}{P^0} = 0$$



3D distribution

$$\langle T_a^{\mu\nu} \rangle(\vec{r}) = \int \frac{d^3 \Delta}{(2\pi)^3} e^{-i\vec{\Delta} \cdot \vec{r}} \frac{\langle \frac{\vec{\Delta}}{2} | T_a^{\mu\nu}(0) | -\frac{\vec{\Delta}}{2} \rangle}{2P^0}$$

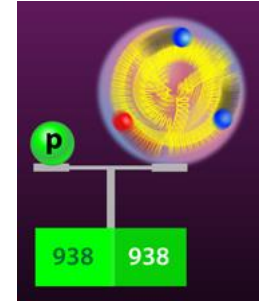
[Polyakov (2003)]
 [C.L., Mantovani, Pasquini (2018)]
 [Polyakov, Schweitzer (2018)]
 [C.L., Moutarde, Trawinski (2019)]

$$P^0 = \sqrt{\frac{\vec{\Delta}^2}{4} + M^2}$$

Mass decomposition and balance equations

1st moment

$$\int d^3r \langle T_a^{\mu\nu} \rangle(\vec{r}) = \frac{\langle \vec{0} | T_a^{\mu\nu}(0) | \vec{0} \rangle}{2M}$$



Forward amplitude

$$\begin{aligned} \langle p, s' | T_a^{\mu\nu}(0) | p, s \rangle &= \delta_{s's} \left\{ 2p^\mu p^\nu \left[\mathcal{G}_1^a(0) + \frac{1}{6} \mathcal{G}_7^a(0) \right] - 2M^2 g^{\mu\nu} \left[\frac{1}{2} \mathcal{G}_8^a(0) + \frac{1}{6} \mathcal{G}_7^a(0) \right] \right\} \\ &\quad - \mathcal{T}_{s's}^{\mu\nu}(p) 2M^2 \frac{1}{2} \mathcal{G}_7^a(0) \end{aligned}$$

Poincaré symmetry

$$\langle p, s' | \sum_a T_a^{\mu\nu}(0) | p, s \rangle = 2p^\mu p^\nu \quad \Rightarrow \quad \sum_a \mathcal{G}_1^a(0) = 1$$

Mass decomposition

$$\sum_a U_a = M \quad U_a = \frac{\langle p | T_a^{00}(0) | p \rangle}{2M} = \left[\mathcal{G}_1^a(0) - \frac{1}{2} \mathcal{G}_8^a(0) \right] M$$

Balance equations

$$\sum_a W_a = 0$$

$$W_a = \frac{\delta^{ij}}{3} \frac{\langle p | T_a^{ij}(0) | p \rangle}{2M} = \left[\frac{1}{2} \mathcal{G}_8^a(0) + \frac{1}{6} \mathcal{G}_7^a(0) \right] M$$

$$\sum_a W_a^{ij} = 0$$

$$W_a^{ij} = \frac{\langle p | T_a^{ij}(0) | p \rangle}{2M} - \delta^{ij} W_a = \mathcal{T}^{ij}(p) \left[-\frac{1}{2} \mathcal{G}_7^a(0) \right] M$$

[C.L. (2018)]

[Cosyn, Cotogno, Freese, C.L., arXiv:1903.00408]

Angular momentum sum rule

2nd moment
$$\int d^3r r^j \langle T_a^{\mu\nu} \rangle(\vec{r}) = \left[-i \nabla_{\Delta}^j \frac{\langle \frac{\vec{\Delta}}{2} | T_a^{\mu\nu}(0) | - \frac{\vec{\Delta}}{2} \rangle}{2P^0} \right]_{\vec{\Delta}=\vec{0}}$$

Orbital and total angular momentum

[Cosyn, Cotogno, Freese, C.L., arXiv:1903.00408]

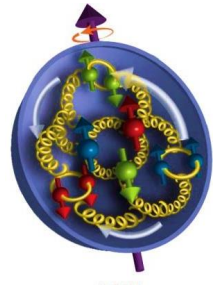
[Polyakov, Sun, arXiv:1903.02738]

$$L_a^i = \varepsilon^{ijk} \int d^3r r^j \langle T_a^{0k} \rangle(\vec{r}) = \frac{s^i}{2} [\mathcal{G}_5^a(0) + \frac{1}{2} \mathcal{G}_7^a(0) + \mathcal{G}_{10}^a(0)]$$

$$J_a^i = \varepsilon^{ijk} \int d^3r r^j \langle T_a^{\{0k\}} \rangle(\vec{r}) = \frac{s^i}{2} [\mathcal{G}_5^a(0) + \underbrace{\frac{1}{2} \mathcal{G}_7^a(0)}_{\text{Missing in}}]$$

Missing in

[Taneja, Kathuria, Liuti, Goldstein (2012)]



Angular momentum sum rule

$$\sum_a J_a^i = s^i \quad \rightarrow \quad \sum_a \mathcal{G}_5^a(0) = 2$$

[Abidin, Carlson (2008)]



Generalization to **arbitrary spin**

[Cotogno, C.L., Lowdon, arXiv:1905.xxxxx]

Mass radius and inertia tensor

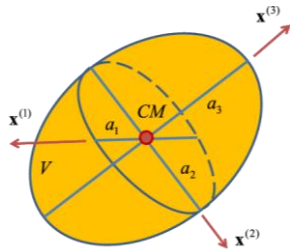
3rd moment

$$\int d^3r r^i r^j \langle T_a^{00} \rangle(\vec{r}) = \left[-\nabla_{\Delta}^i \nabla_{\Delta}^j \frac{\langle \frac{\Delta}{2} |T_a^{00}(0)| - \frac{\Delta}{2} \rangle}{2P^0} \right]_{\vec{\Delta}=\vec{0}}$$

Inertia tensor

$$I_a^{ij} = \int d^3r (\vec{r}^2 \delta^{ij} - r^i r^j) \langle T_a^{00} \rangle(\vec{r})$$

$$= \frac{1}{M} [2\delta^{ij} (\mathcal{A}_a(0) + \frac{1}{3} \mathcal{B}_a(0)) + \mathcal{T}^{ij} \mathcal{B}_a(0)]$$



$$\mathcal{A}_a(t) = -\frac{1}{4} [\mathcal{G}_1^a(t) + 2\mathcal{G}_3^a(t) + \frac{1}{2}\mathcal{G}_8^a(t)] + 2M^2 \frac{d}{dt} [\mathcal{G}_1^a(t) - \frac{1}{2}\mathcal{G}_8^a(t)]$$

$$\mathcal{B}_a(t) = -\mathcal{G}_1^a(t) - \mathcal{G}_3^a(t) + \mathcal{G}_5^a(t) + \frac{1}{2} [\mathcal{G}_6^a(t) + \frac{1}{2}\mathcal{G}_7^a(t) + \mathcal{G}_8^a(t) - \mathcal{G}_9^a(t)]$$

Mass radius

$$\langle \vec{r}^2 \rangle = \frac{\delta^{ij}}{2M} \sum_a I_a^{ij} = \frac{1}{M^2} \sum_a [3\mathcal{A}_a(0) + \mathcal{B}_a(0)]$$

Mass quadrupole moment

$$Q_a^{ij} = \frac{\delta^{ij}}{3} (\delta^{kl} I_a^{kl}) - I_a^{ij} = -\frac{1}{M} \mathcal{T}^{ij} \mathcal{B}_a(0)$$

Link with GPDs

2nd Mellin moment
$$\int_{-1}^1 dx x V_a^\mu(x) = \frac{T_a^{\mu+}(0)}{2(P^+)^2}$$

$$V_q^\mu(x) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \bar{\psi}(-\frac{z^-}{2}) \gamma^\mu \mathcal{W} \psi(\frac{z^-}{2})$$

$$V_G^\mu(x) = \frac{\delta_{\{\alpha\beta\}}^\mu}{2xP^+} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \text{Tr} \left[\mathcal{W} F^{\alpha\lambda}(-\frac{z^-}{2}) \mathcal{W} F_\lambda^\beta(\frac{z^-}{2}) \right]$$

Twist-2
$$\int_{-1}^1 dx x H_2^a(x, \xi, t) = \mathcal{G}_5^a(t)$$

$$\int_{-1}^1 dx x H_5^a(x, \xi, t) = -\frac{t}{4M^2} \mathcal{G}_6^a(t) + \frac{1}{2} \mathcal{G}_7^a(t)$$

Twist-3
$$\int_{-1}^1 dx x G_6^q(x, \xi, t) = -\frac{1}{2} [\mathcal{G}_5^q(t) + \mathcal{G}_{10}^q(t)]$$

[Taneja, Kathuria, Liuti, Goldstein (2012)]
 [Cosyn, Cotogno, Freese, C.L., arXiv:1903.00408]
 [Polyakov, Sun, arXiv:1903.02738]

Angular momentum relations

Spin-1

$$J_z^a = \int_{-1}^1 dx \frac{x}{2} [H_2^a(x, 0, 0) + H_5^a(x, 0, 0)]$$

$$L_z^q = \int_{-1}^1 dx x [-G_6^q(x, 0, 0) + \frac{1}{2} H_5^q(x, 0, 0)]$$

Spin-1/2

$$J_z^a = \int_{-1}^1 dx \frac{x}{2} [H^a(x, 0, 0) + E^a(x, 0, 0)]$$

$$L_z^q = - \int_{-1}^1 dx x G_2^q(x, 0, 0)$$

[Ji (1996)]
 [Penttinen, Polyakov, Shuvaev, Strikman (2000)]
 [Kiptily, Polyakov (2004)]
 [Hatta, Yoshida (2012)]

Summary

$$T^{\mu\nu} = \begin{bmatrix} T^{00} & T^{01} & T^{02} & T^{03} \\ T^{10} & T^{11} & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} \end{bmatrix}$$

Energy density: T^{00}
Momentum density: T^{01}, T^{02}, T^{03}
Energy flux: T^{10}, T^{20}, T^{30}
Momentum flux: $T^{11}, T^{12}, T^{13}, T^{21}, T^{22}, T^{23}, T^{31}, T^{32}, T^{33}$
Shear stress: $T^{12}, T^{21}, T^{13}, T^{31}, T^{23}, T^{32}$
Normal stress (pressure): T^{11}, T^{22}, T^{33}

- **Key information about the structure encoded in EMT**
- **Higher-spin targets involve higher multipole moments**
 ➔ **New quadrupolar effects for spin-1**
- **Breit frame defines the spatial distribution at rest**
- **Mechanical properties can be expressed in terms of spatial moments**
 ➔ **Mass, pressure, moment of inertia, ...**
- **GFFs can be related to GPDs**
 ➔ **Sum rules and relations similar to spin-1/2 case**

You're all welcome!

<https://indico.cern.ch/e/LC2019>



LIGHT CONE 2019

Campus de l'École polytechnique, Palaiseau, France
September 16-20, 2019

Physics topics

- Hadronic structure
- Small-x physics
- QCD at finite temperature
- Few and many-body physics
- Chiral symmetry

Approaches

- Field theories in the front form
- Lattice field theory
- Effective field theories
- Phenomenological models
- Present and future facilities

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Logos: École Polytechnique Université Paris-Saclay, CEA, ANR, Jefferson Lab

Backup slides

Quantum chromodynamics (QCD)

Classical QCD energy-momentum tensor

$$T^{\mu\nu} = \bar{\psi} \gamma^\mu \frac{i \overleftrightarrow{D}^\nu}{2} \psi - G^{a\mu\alpha} G^{a\nu}{}_\alpha + \frac{1}{4} \eta^{\mu\nu} G^2$$

Renormalized trace of the QCD EMT

$$T^\mu{}_\mu = \underbrace{\frac{\beta(g)}{2g} G^2}_{\text{Trace anomaly}} + (1 + \gamma_m) \bar{\psi} \underbrace{m}_{\text{Quark mass matrix}} \psi$$

[Crewther (1972)]
 [Chanowitz, Ellis (1972)]
 [Nielsen (1975)]
 [Adler, Collins, Duncan (1977)]
 [Collins, Duncan, Joglekar (1977)]
 [Nielsen (1977)]

Poincaré invariance

$$\partial_\mu T^{\mu\nu} = 0 \quad \Rightarrow \quad \sum_{i=q,g} A_i(0) = 1, \quad \sum_{i=q,g} \bar{C}_i(t) = 0$$

$$\underbrace{\partial_\mu J^{\mu\alpha\beta}}_{x^\alpha T^{\mu\beta} - x^\beta T^{\mu\alpha} + S^{\mu\alpha\beta}} = 0 \quad \Rightarrow \quad \sum_{i=q,g} B_i(0) = 0, \quad D_q(t) = -G_A^q(t)$$

[Kobzarev, Okun (1962)]
 [Teryaev (1999)]
 [Brodsky, Hwang, Ma, Schmidt (2001)]
 [Leader, C.L. (2014)]
 [Teryaev (2016)]
 [Lowdon, Chiu, Brodsky (2017)]

Textbook decomposition

Forward matrix element

$$\langle P|T^{\mu\nu}(0)|P\rangle = 2P^\mu P^\nu$$

$$\langle P'|P\rangle = 2P^0 (2\pi)^3 \delta^{(3)}(\vec{P}' - \vec{P})$$

Trace decomposition

$$\begin{aligned} 2M^2 &= \langle P|T^\mu{}_\mu(0)|P\rangle \\ &= \underbrace{\langle P|\frac{\beta(g)}{2g} G^2|P\rangle}_{\sim 89\%} + \underbrace{\langle P|(1 + \gamma_m)\bar{\psi}m\psi|P\rangle}_{\sim 11\%} \end{aligned}$$

[Shifman, Vainshtein, Zakharov (1978)]
[Luke, Manohar, Savage (1992)]
[Donoghue, Golowich, Holstein (1992)]
[Kharzeev (1996)]
[Bressani, Wiedner, Filippi (2005)]
[Roberts (2017)]
[Krein, Thomas, Tsushima (2017)]



Manifestly covariant



Reminiscent of Gell-Mann–Oakes–Renner formula for pion



Depends on state normalization



No spatial extension



No clear relation to energy

Ji's decomposition

[Ji (1995)]

Separation of quark and gluon contributions

$$T^{\mu\nu} = \bar{T}^{\mu\nu} + \hat{T}^{\mu\nu} \quad \bar{T}^{\mu\nu} = \bar{T}_q^{\mu\nu} + \bar{T}_g^{\mu\nu} \quad \hat{T}^{\mu\nu} = \hat{T}_m^{\mu\nu} + \hat{T}_a^{\mu\nu}$$

Traceless Pure trace

Forward matrix elements

$$\begin{aligned} \langle P | \bar{T}_i^{\mu\nu}(0) | P \rangle &= 2 \left(P^\mu P^\nu - \frac{1}{4} \eta^{\mu\nu} M^2 \right) A_i(0) \\ \langle P | \hat{T}_i^{\mu\nu}(0) | P \rangle &= \frac{1}{2} \eta^{\mu\nu} M^2 [A_i(0) + 4\bar{C}_i(0)] \end{aligned} \quad \langle O \rangle = \frac{\langle P | \int d^3r O(r) | P \rangle}{\langle P | P \rangle}$$

Ji's decomposition

[Gao *et al.* (2015)]

$$M = \underbrace{M_q}_{\sim 31\%} + \underbrace{M_g}_{\sim 34\%} + \underbrace{M_m}_{\sim 13\%} + \underbrace{M_a}_{\sim 22\%}$$

$\mu = 2 \text{ GeV}$

$$\begin{aligned} M_q &= \langle \bar{T}_q^{00} \rangle |_{\vec{P}=\vec{0}} - \frac{3}{1+\gamma_m} \langle \hat{T}_m^{00} \rangle |_{\vec{P}=\vec{0}} \\ M_g &= \langle \bar{T}_g^{00} \rangle |_{\vec{P}=\vec{0}} \\ M_m &= \frac{4+\gamma_m}{1+\gamma_m} \langle \hat{T}_m^{00} \rangle |_{\vec{P}=\vec{0}} \\ M_a &= \langle \hat{T}_a^{00} \rangle |_{\vec{P}=\vec{0}} \end{aligned}$$



Proper normalization



Clear relation to energy distribution



Scale-dependent interpretation in the rest frame



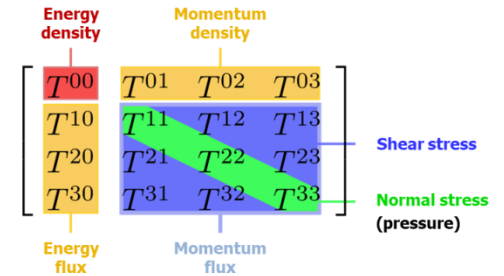
Pressure effects not taken into account

New decomposition

[C.L. (2018)]

Forward matrix element

$$\langle P | T_i^{\mu\nu}(0) | P \rangle = 2P^\mu P^\nu A_i(0) + 2M^2 \eta^{\mu\nu} \bar{C}_i(0)$$



Analogy with relativistic hydrodynamics

Perfect fluid element

$$\Theta_i^{\mu\nu} = (\varepsilon_i + p_i) u^\mu u^\nu - p_i \eta^{\mu\nu}$$



Four-velocity

$$u^\mu = P^\mu / M$$

Energy density

$$\varepsilon_i = [A_i(0) + \bar{C}_i(0)] \frac{M}{V}$$

Isotropic pressure

$$p_i = -\bar{C}_i(0) \frac{M}{V}$$

Nucleon mass decomposition $U_i = \varepsilon_i V$

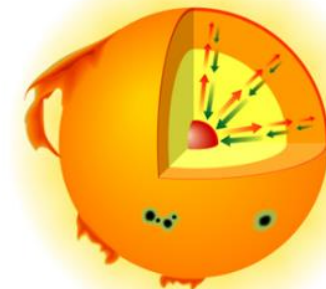
$$M = \underbrace{U_q}_{\sim 44\%} + \underbrace{U_g}_{\sim 56\%}$$

$\mu = 2 \text{ GeV}$

$$p_q = -p_g$$

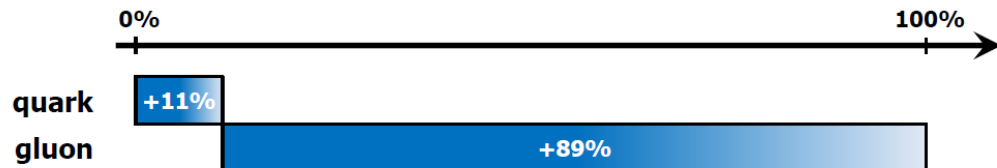
$\sim 11\%$

pressure → quark pressure
gravity → gluon pressure

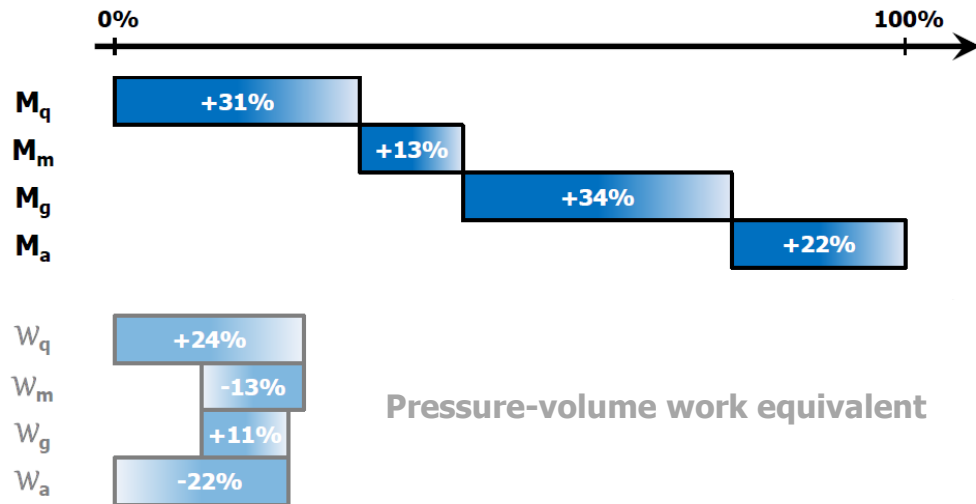


In short

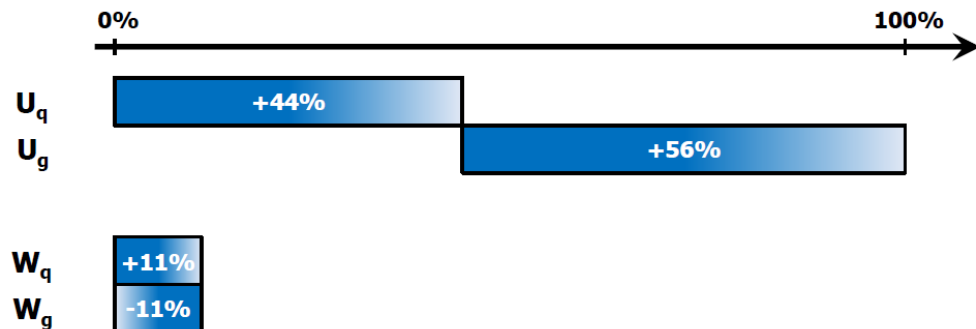
Trace decomposition



Ji's decomposition



New decomposition



Anisotropic medium

[Polyakov (2003)]
 [Goetze *et al.* (2007)]
 [Polyakov, Schweitzer (2018)]
 [C.L., Moutarde, Trawinski, arXiv:1810.09837]

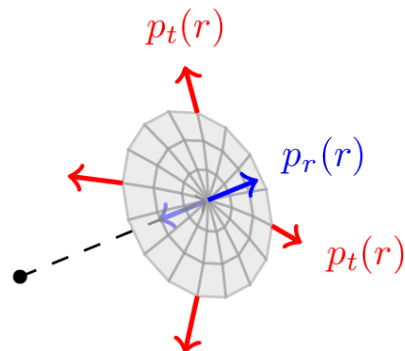
Breit frame amplitude $t = -\vec{\Delta}^2$

$$\frac{\langle \frac{\vec{\Delta}}{2} | T_i^{\mu\nu}(0) | -\frac{\vec{\Delta}}{2} \rangle}{2P^0} = M \left\{ \eta^{\mu 0} \eta^{\nu 0} \left[A_i(t) + \frac{t}{4M^2} B_i(t) \right] + \eta^{\mu\nu} \left[\bar{C}_i(t) - \frac{t}{M^2} C_i(t) \right] + \frac{\Delta^\mu \Delta^\nu}{M^2} C_i(t) \right\}$$

Analogy with relativistic hydrodynamics $r = |\vec{r}|$

Anisotropic fluid

$$\Theta_i^{\mu\nu}(\vec{r}) = [\varepsilon_i(r) + p_{t,i}(r)] u^\mu u^\nu - p_{t,i}(r) \eta^{\mu\nu} + [p_{r,i}(r) - p_{t,i}(r)] \frac{r^\mu r^\nu}{r^2}$$



Isotropic pressure

$$p_i(r) = \frac{p_{r,i}(r) + 2p_{t,i}(r)}{3}$$

Pressure anisotropy

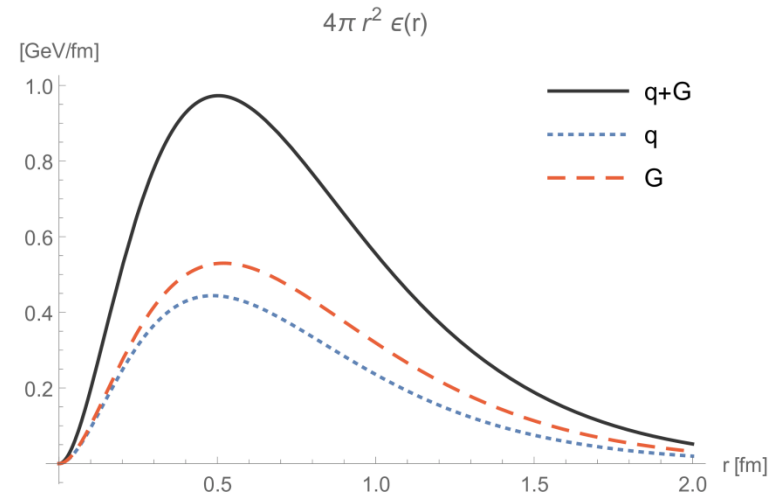
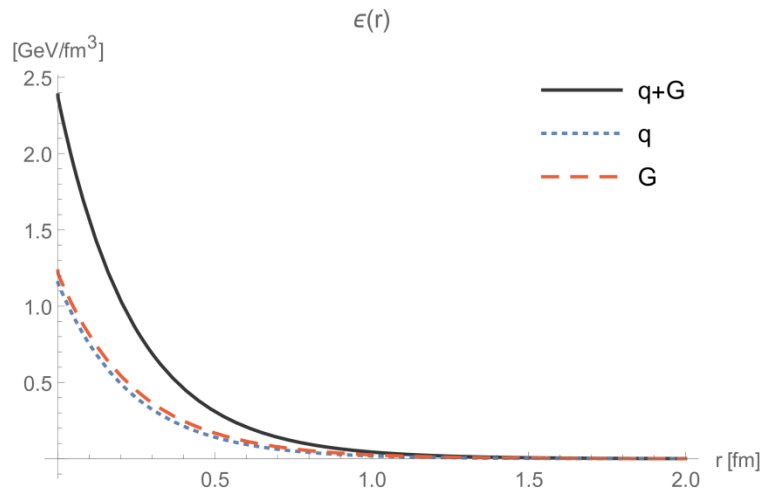
$$s_i(r) = p_{r,i}(r) - p_{t,i}(r)$$

Energy distribution

[C.L., Moutarde, Trawinski, arXiv:1810.09837]

Multipole model for the GFFs

$$F(t) = \frac{F(0)}{(1 + t/\Lambda^2)^n}$$

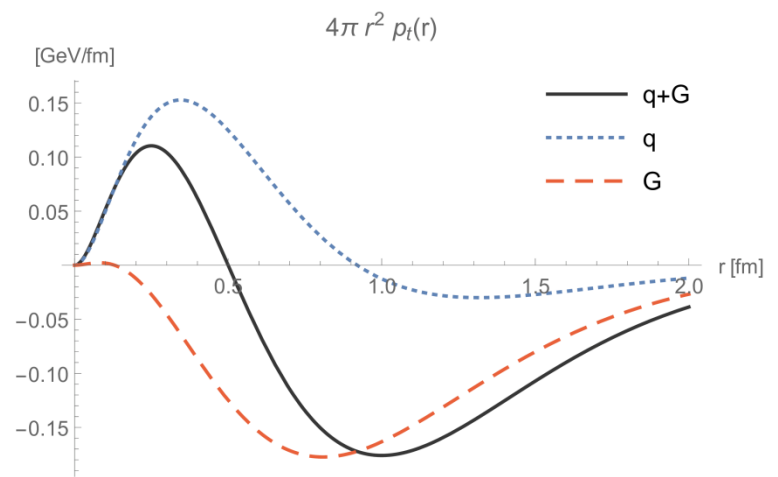
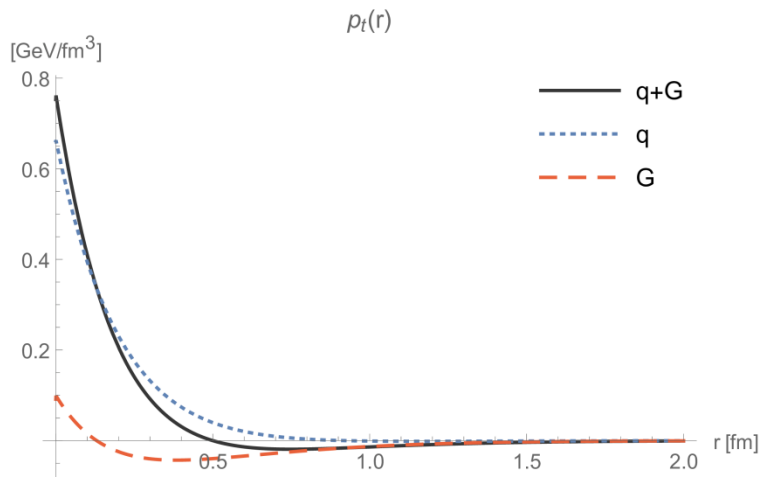
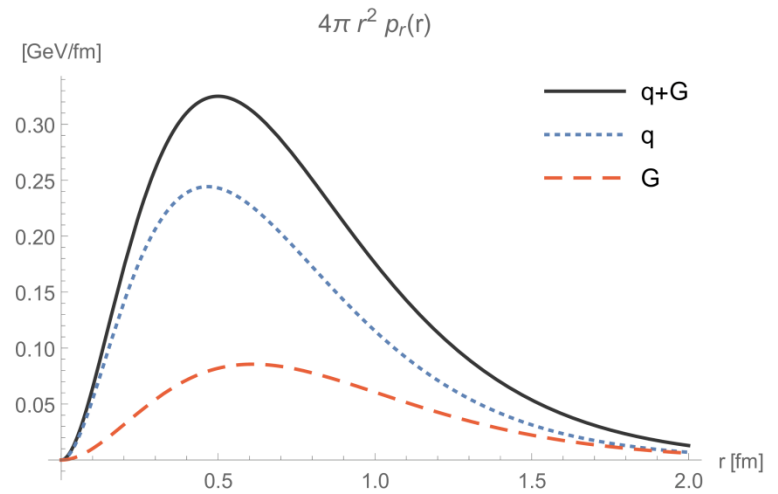
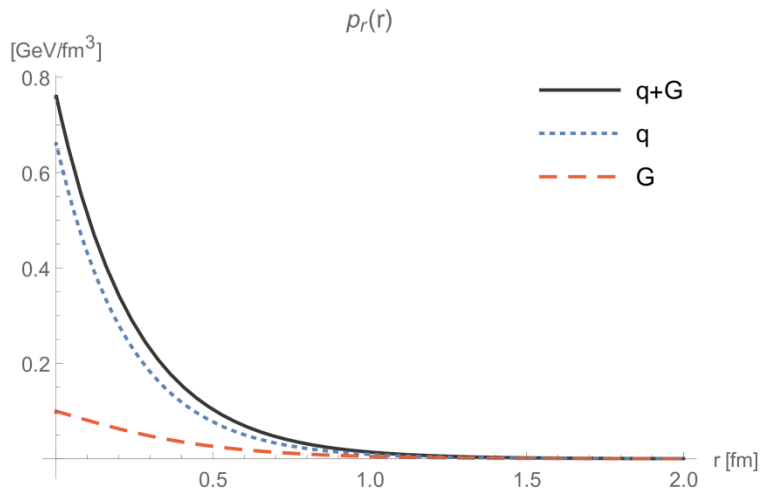


$$\sqrt{\langle r^2 \rangle_M} = 0.91 \text{ fm}$$

$$\sqrt{\langle r^2 \rangle_Q} = 0.84 - 0.88 \text{ fm}$$

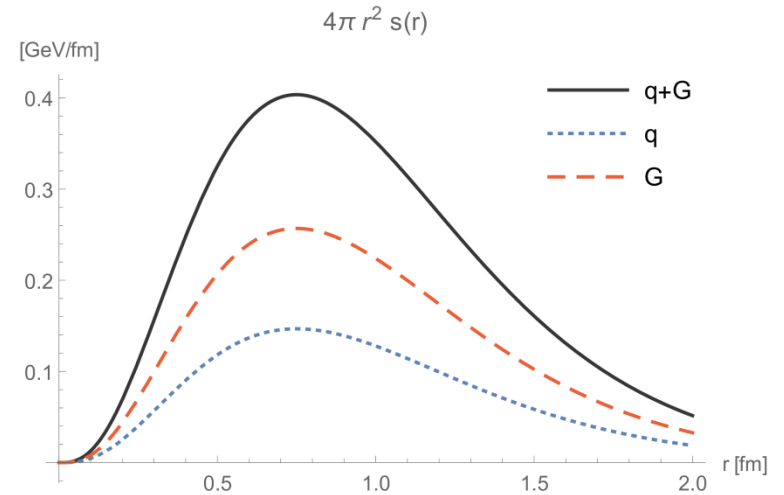
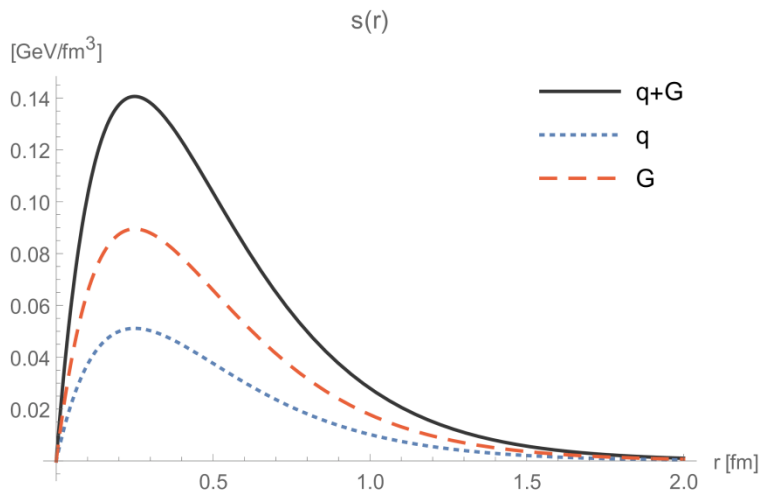
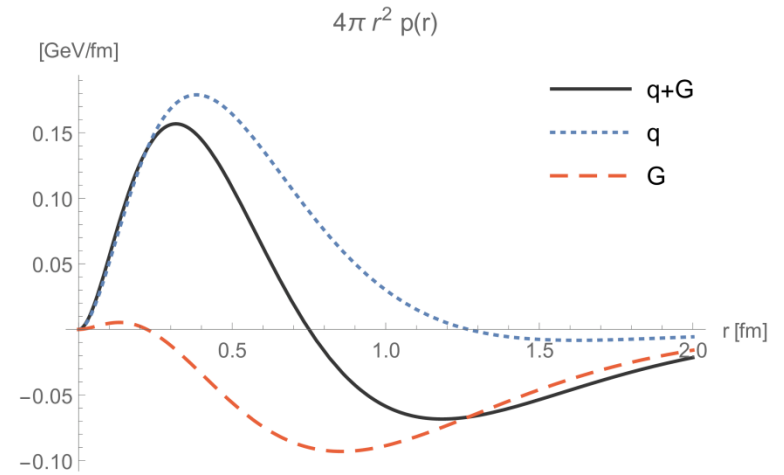
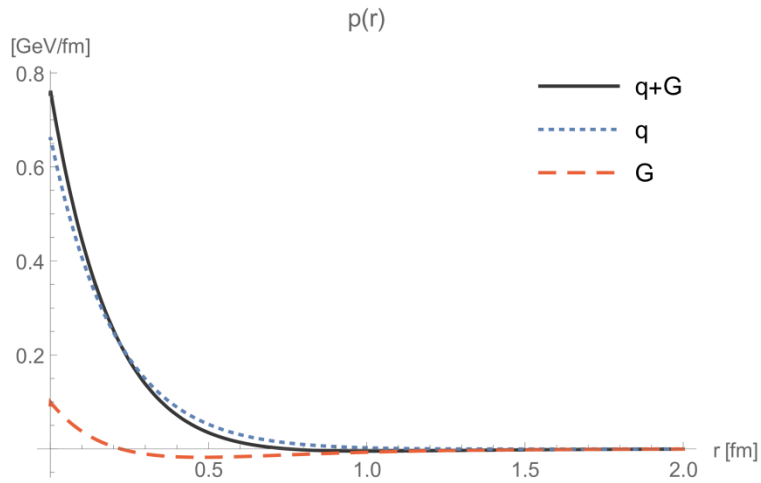
Pressure distribution

[C.L., Moutarde, Trawinski, arXiv:1810.09837]



Pressure distribution

[C.L., Moutarde, Trawinski, arXiv:1810.09837]



Hydrostatic equilibrium

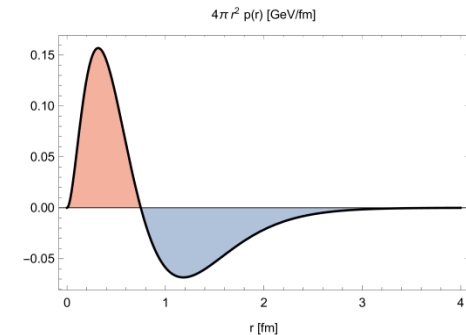
[Polyakov (2003)]
 [Goetze *et al.* (2007)]
 [Polyakov, Schweitzer (2018)]
 [C.L., Moutarde, Trawinski, arXiv:1810.09837]

$$\nabla^i \mathcal{T}^{ij}(\vec{r}) = 0 \quad \Rightarrow \quad \frac{dp_r(r)}{dr} = -\frac{2s(r)}{r}$$

von Laue relation

[von Laue (1911)]

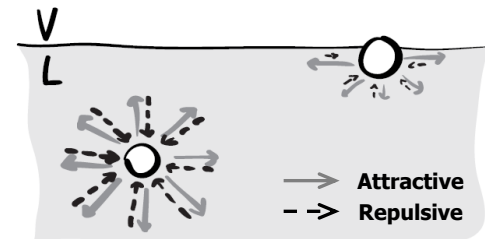
$$\int_0^\infty dr r^2 p(r) = 0$$



Surface tension

[Bakker (1928)]
 [Kirkwood, Buff (1949)]
 [Marchand *et al.* (2011)]

$$\gamma = \int dr s(r)$$



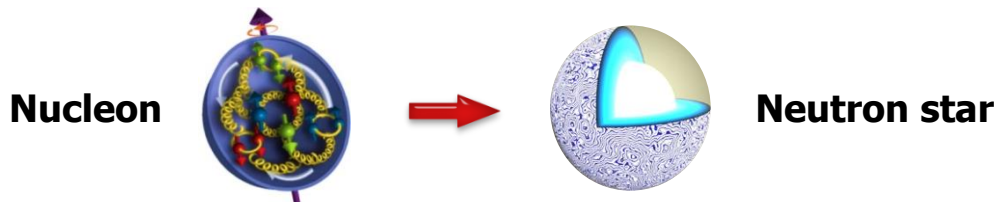
Generalized Young-Laplace relation

$$p(0) = 2 \int_0^\infty dr \frac{s(r)}{r}$$

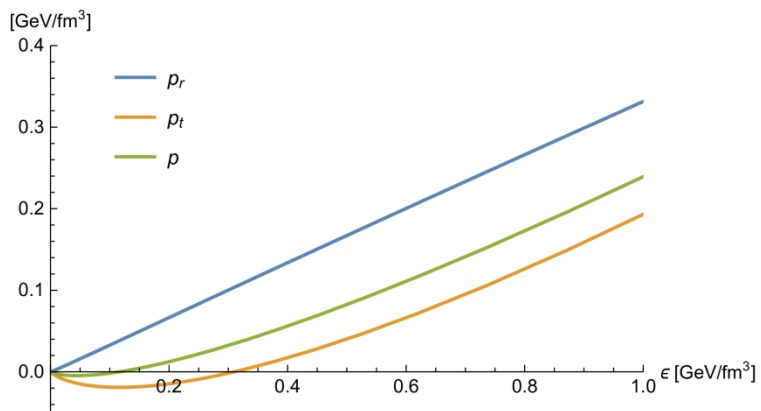
[Thomson (1858)]

$$s(r) \approx \gamma \delta(r - R) \quad \Rightarrow \quad p(0) = \frac{2\gamma}{R}$$

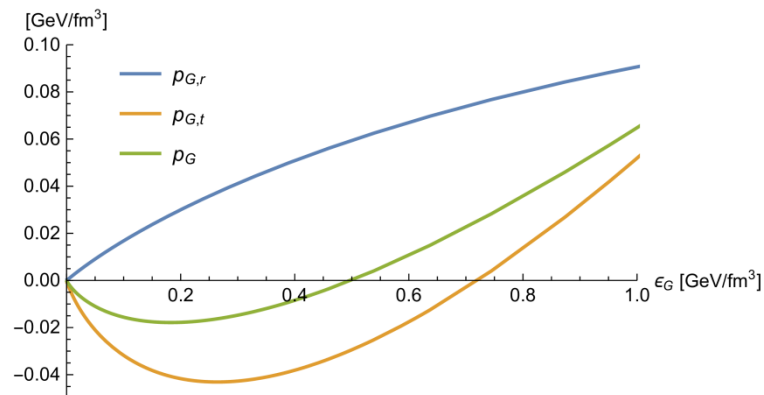
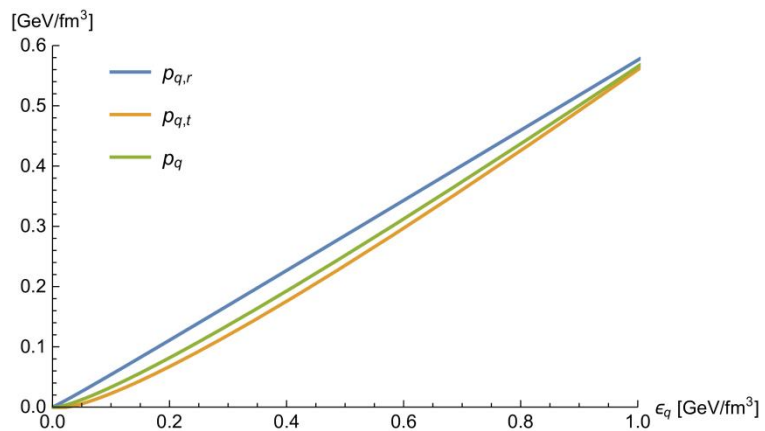
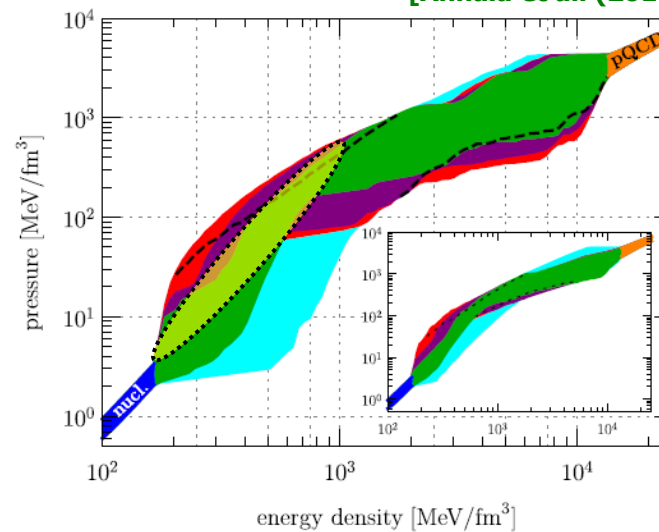
What can we learn?



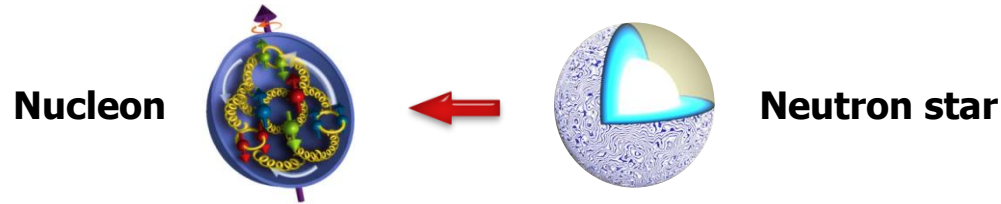
Equation of state



[Annala *et al.* (2018)]



What can we learn?



Stability constraints

Mechanical regularity

- (i) $\varepsilon(0) < \infty$, $p(0) < \infty$ and $s(0) = 0$;
- (ii) $\varepsilon(r) > 0$ and $p_r(r) > 0$;
- (iii) $\frac{d\varepsilon(r)}{dr} < 0$ and $\frac{dp_r(r)}{dr} < 0$.

Speed of sound

- (iv) $0 \leq v_{sr}^2(r) \leq 1$ and $0 \leq v_{st}^2(r) \leq 1$;
- (v) $|v_{st}^2(r) - v_{sr}^2(r)| \leq 1$;
- (vi) $\Gamma(r) = \frac{\varepsilon(r) + p_r(r)}{p_r(r)} v_{sr}^2 > \frac{4}{3}$.

Energy conditions

$$\begin{aligned} \varepsilon(r) + p_i(r) &\geq 0, \\ \varepsilon(r) + p_i(r) &\geq 0 \quad \text{and} \quad \varepsilon(r) \geq 0, \\ \varepsilon(r) + p_i(r) &\geq 0 \quad \text{and} \quad \varepsilon(r) + 3p(r) \geq 0, \\ \varepsilon(r) &\geq |p_i(r)|, \end{aligned}$$

[Wald (1984)]
 [Herrera, Santos (1997)]
 [Poisson (2004)]
 [Abreu, Hernandez, Nunez (2007)]
 [Hawking, Ellis (2011)]