

QCD Evolution 2019, Argonne National Lab, May 14, 2019

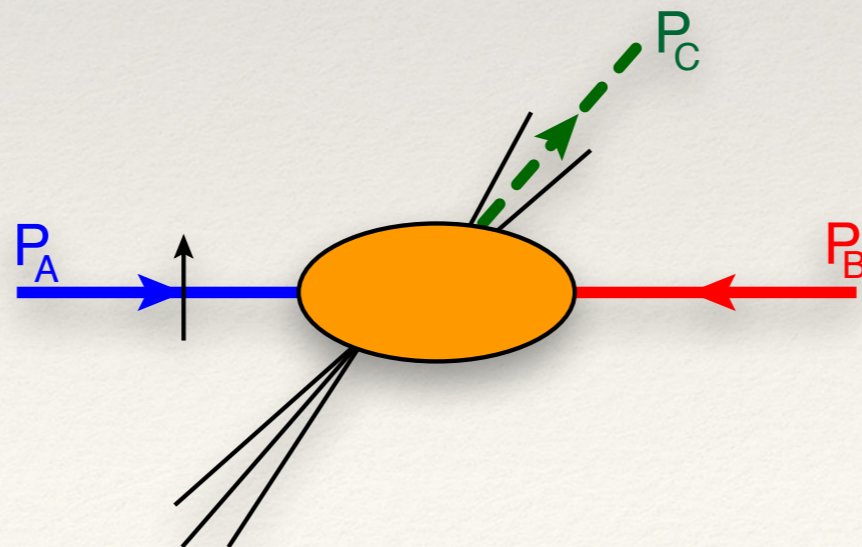
The transverse nucleon spin
asymmetry in photon SIDIS in
the collinear twist-3 formalism

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with W. Albaltan, A. Prokudin

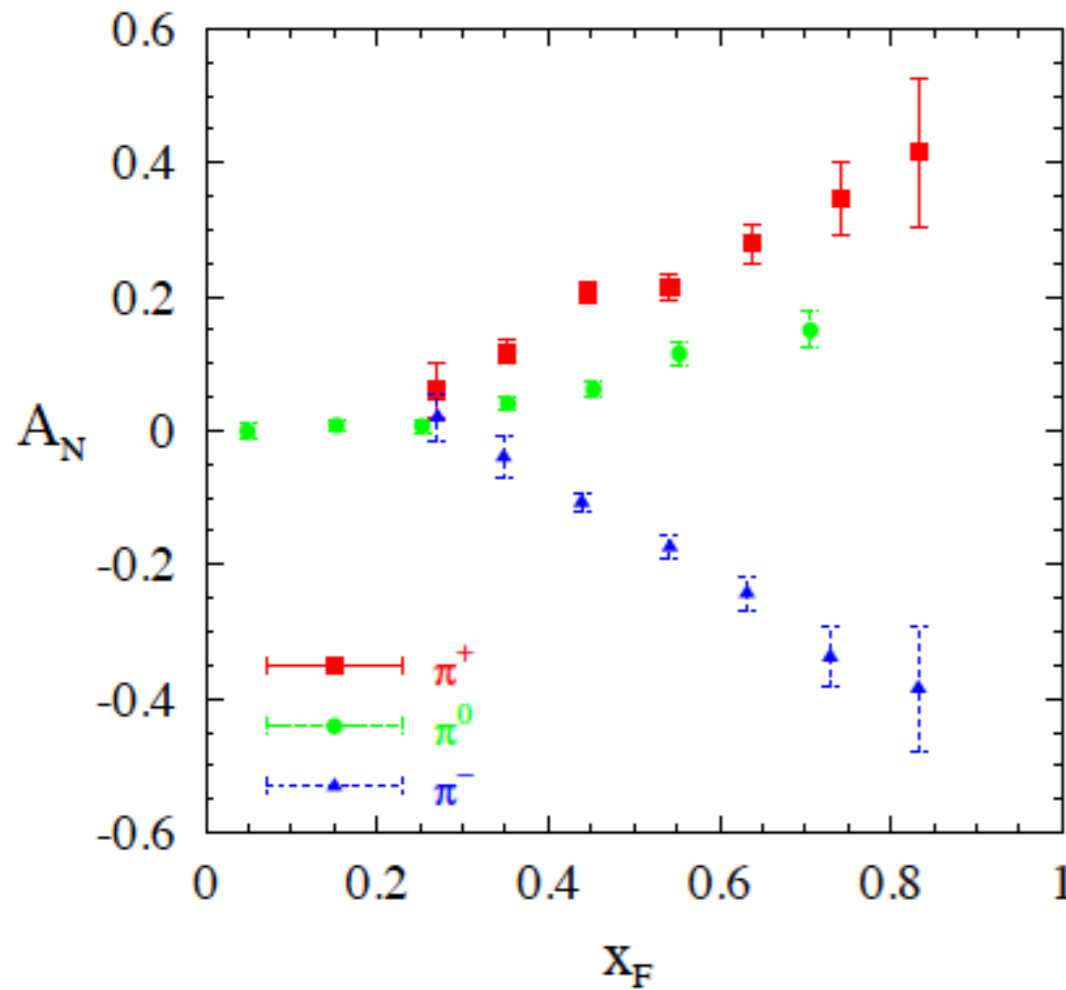
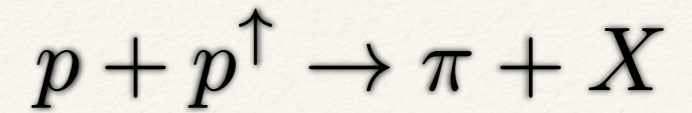
Transverse Spin Effects in Single-Inclusive Hard Processes

$$P_A^\uparrow + P_B \rightarrow P_C + X$$

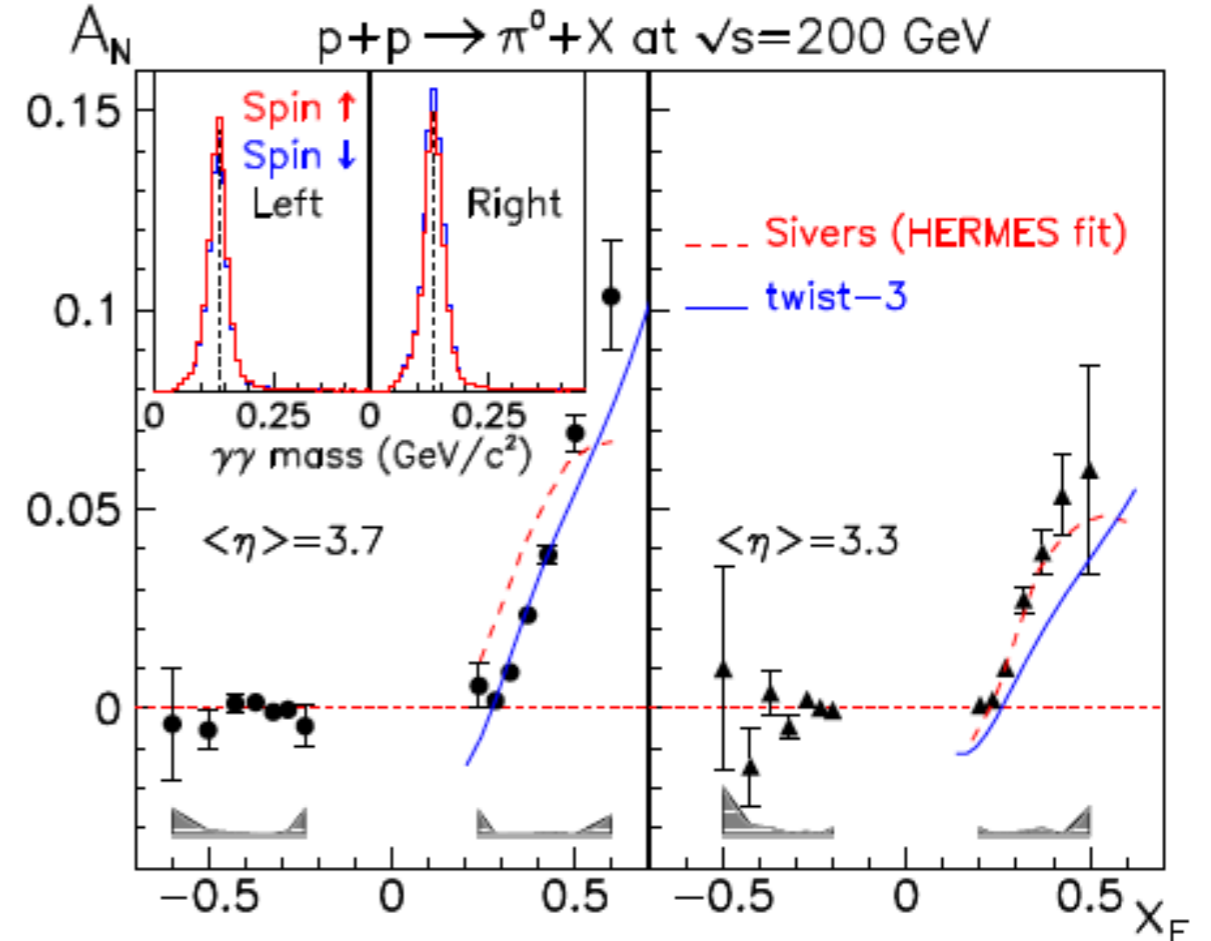


“Show-off” Transverse SSA

$$A_N = \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow}$$



$\sqrt{s} = 20 \text{ GeV}$ [E704 coll. (1991)]



$\sqrt{s} = 200 \text{ GeV}$ [STAR coll. (2008)]

large effects

cannot be explained in the standard parton model
(using transversity)

→ collinear Twist-3 Formalism
(Efremov, Teryaev, Qiu, Sterman)

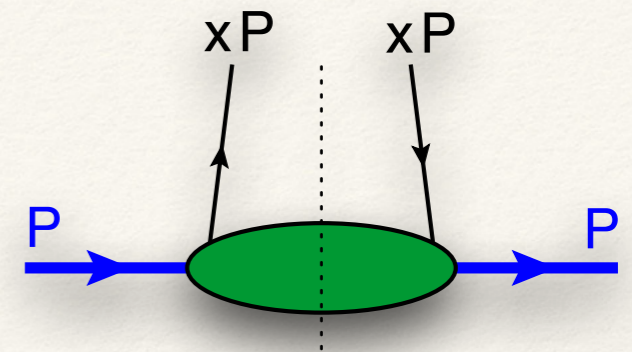
Collinear twist-3 formalism: several types of matrix elements compete

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intrinsic twist-3 PDF

$$g_T^q(x) = -\frac{1}{M} \int \frac{d\lambda}{4\pi} e^{i\lambda x} \langle P, S_T | \bar{q}(0) \not{\epsilon}_T \gamma_5 q(\lambda n) | P, S_T \rangle$$

- sensitive to 'bad quark field components',
- twist-3 characteristics hidden in Dirac structure
- generates the g_2 structure function in DIS
- No probabilistic interpretation

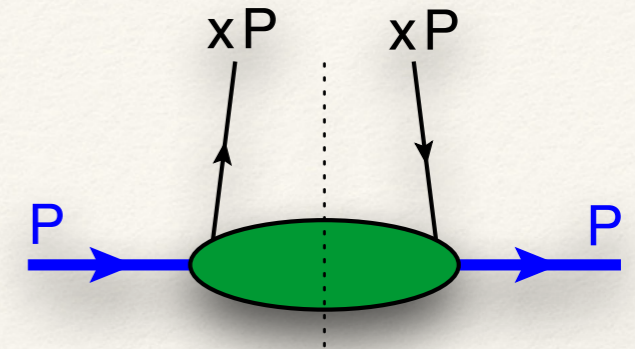


Collinear twist-3 formalism: several types of matrix elements compete

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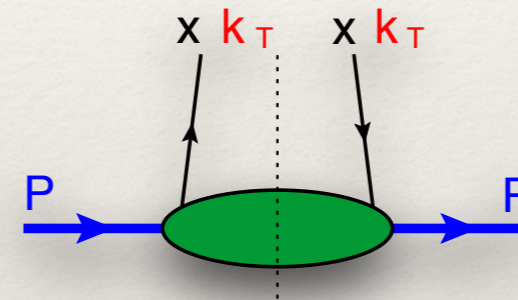
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kinematical twist-3 PDFs:

Small transverse quark/gluon momenta k_T :



$$(\mathbf{k}_T \times S_T) f_{1T}^{\perp,q}(x, k_T^2) \propto \int \frac{d\lambda d^2 z_T}{(2\pi)^3} e^{i\lambda x + i\mathbf{k}_T \cdot \mathbf{z}_T} \langle P, S_T | \bar{q}(0) \not{n} \mathcal{W} q(\lambda n + \mathbf{z}_T) | P, S_T \rangle$$

Sivers function

$$(\mathbf{k}_T \cdot S_T) g_{1T}^q(x, k_T^2) \propto \int \frac{d\lambda d^2 z_T}{(2\pi)^3} e^{i\lambda x + i\mathbf{k}_T \cdot \mathbf{z}_T} \langle P, S_T | \bar{q}(0) \not{n} \gamma_5 \mathcal{W} q(\lambda n + \mathbf{z}_T) | P, S_T \rangle$$

'transhelicity'

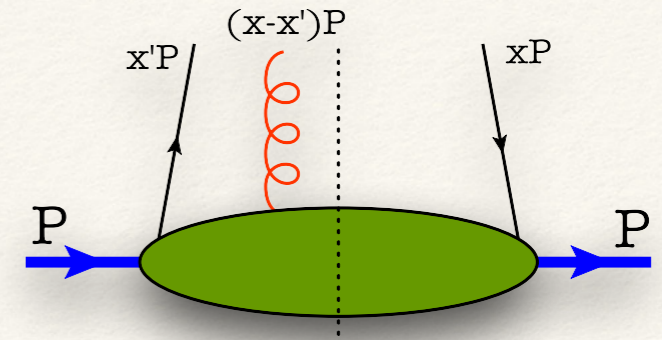
Collinear twist-3 formalism: TMD moments are needed

$$f_{1T}^{\perp,(1)}(x) = \int d^2 k_T \frac{k_T^2}{2M^2} f_{1T}^{\perp}(x, k_T^2)$$

$$g_{1T}^{(1)}(x) = \int d^2 k_T \frac{k_T^2}{2M^2} g_{1T}(x, k_T^2)$$

→ twist-3 characteristics through small transverse parton momentum k_T

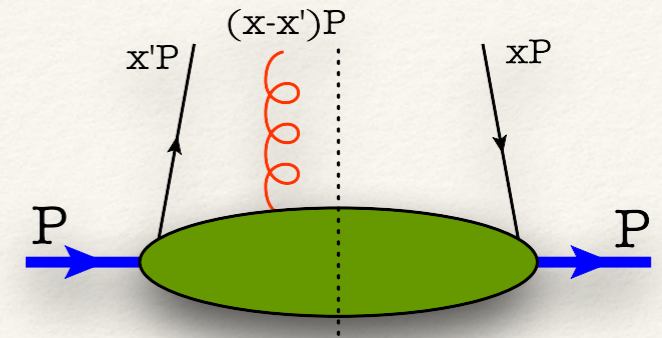
Dynamical twist-3: Quark - Gluon - Quark Correlations
(ETQS-matrix elements)



$$2M i\epsilon^{Pn\rho S} F_{FT}^q(x, x') = \int \frac{d\lambda}{2\pi} \int \frac{d\mu}{2\pi} e^{i\lambda x'} e^{i\mu(x-x')} \langle P, S_T | \bar{q}(0) \not{n} igF^{n\rho}(\mu n) q(\lambda n) | P, S_T \rangle$$

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‘dynamical twist - 3’

→ 3 - parton correlator: suppression by additional propagator

→ Quark-Gluon-Quark correlation functions
drive x-dependence of TMDs like Sivers function, transhelicity, etc.

→ so far: only “diagonal support” $\pi F_{FT}(x, x) = f_{1T}^{\perp(1)}(x)$ constraint by SIDIS data

→ ‘integrated’ $F_{FT}(x, x')$: average transverse color Lorentz force on struck quark
[Burkardt, PRD88, 114502], see talk by M. Burkardt

$$F^{n\rho} = [\vec{E} + \vec{n} \times \vec{B}]^\rho \propto \int dx \int dx' F_{FT}(x, x') \propto \int dx x^2 g_T(x)$$

QCD EoM relation & Lorentz-Invariance Relations

[Kanazawa, Koike, Metz, Pitonyak, MS, PRD 2016]

QCD EoM for Twist-3 PDFs

$$g_{1T}^{(1)}(x) = x g_T(x) - \frac{m_q}{M} h_1(x) + \int_{-1}^1 dx' \frac{F_{FT}(x, x') - G_{FT}(x, x')}{x - x'}$$

$$\pi F_{FT}^q(x, x) = f_{1T}^{\perp(1),q}(x)$$

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LIR for Twist-3 PDFs

based on translation invariance

$$g_T(x) = g_1(x) + \frac{d}{dx} g_{1T}^{(1)}(x) - 2 \int_{-1}^1 dx' \frac{G_{FT}(x, x')}{(x - x')^2}$$

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Two equations, three functions → eliminate ‘intrinsic & kinematical twist-3’

$$g_T(x) = \int_x^1 \frac{dy}{y} g_1(y) + \frac{m_q}{M} \left(\frac{1}{x} h_1(x) - \int_x^1 \frac{dy}{y^2} h_1(y) \right) + \int_x^1 \frac{dy}{y^2} \int_{-1}^1 dz \left[\frac{(1-y\delta(y-x)) F_{FT}(y,z)}{y-z} - \frac{(3y-z-y(y-z)\delta(y-x)) G_{FT}(y,z)}{(y-z)^2} \right]$$

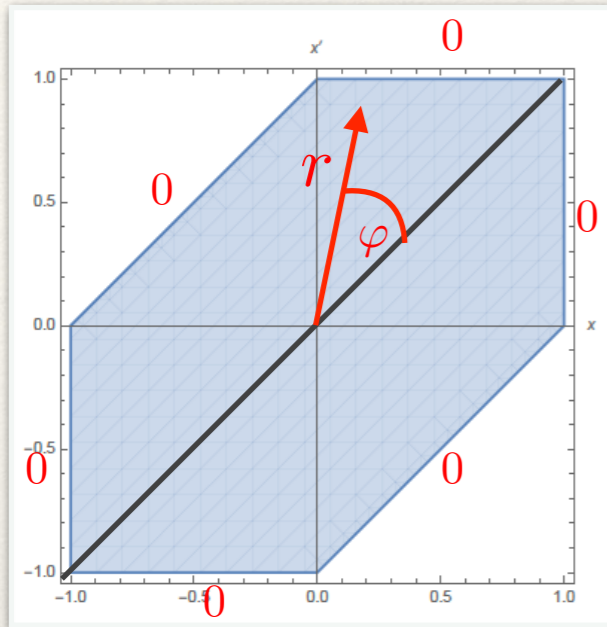
EoM & LIR relation crucial for gauge invariance, invariance of LC vector n

Support properties

$$-1 \leq x, x' \leq 1$$

$$|x - x'| \leq 1$$

and continuous

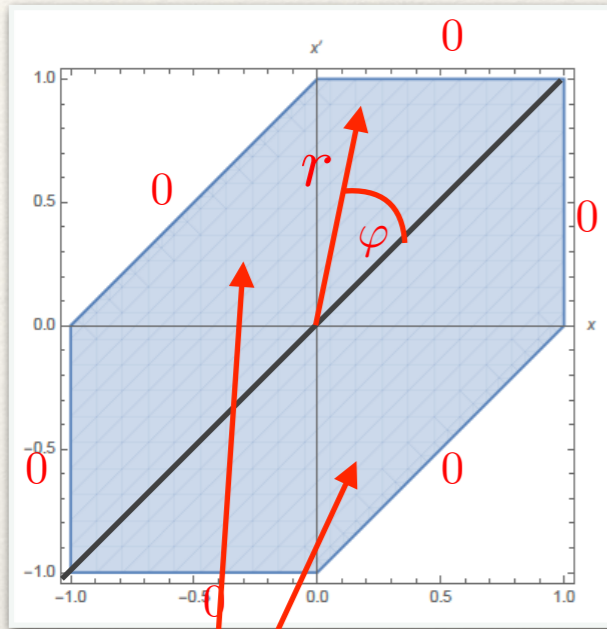


$$F_{FT}(x, x') = +F_{FT}(x', x) \implies \sum_n a_n(r) \cos(n\varphi)$$

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$$\pi F_{FT}^q(-x, -x) = f_{1T}^{\perp(1),\bar{q}}(x)$$

Fixes a_0, a_1

Gluons poles 'known' from SIDIS experiments

$$F_{FT}(x, x') = \left(\frac{(1-x^2)(1-x'^2)(1-(x-x'))}{(1-xx')^2} \right)^\delta (a_0(r) + a_1(r) \cos(\varphi))$$

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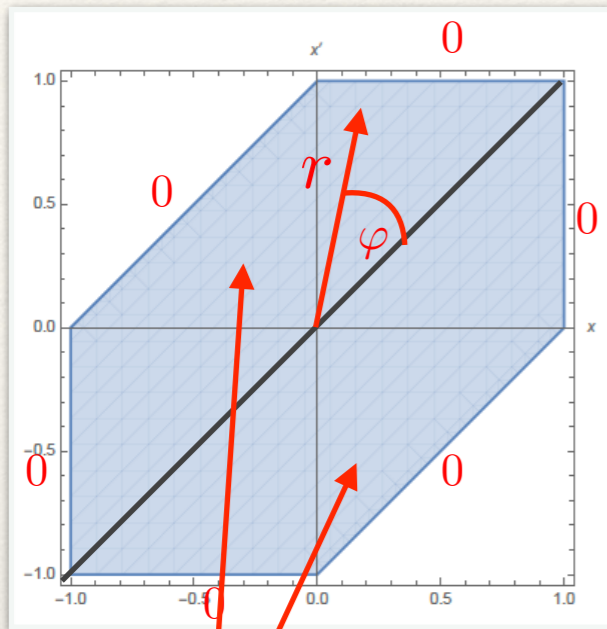
Model ansätze

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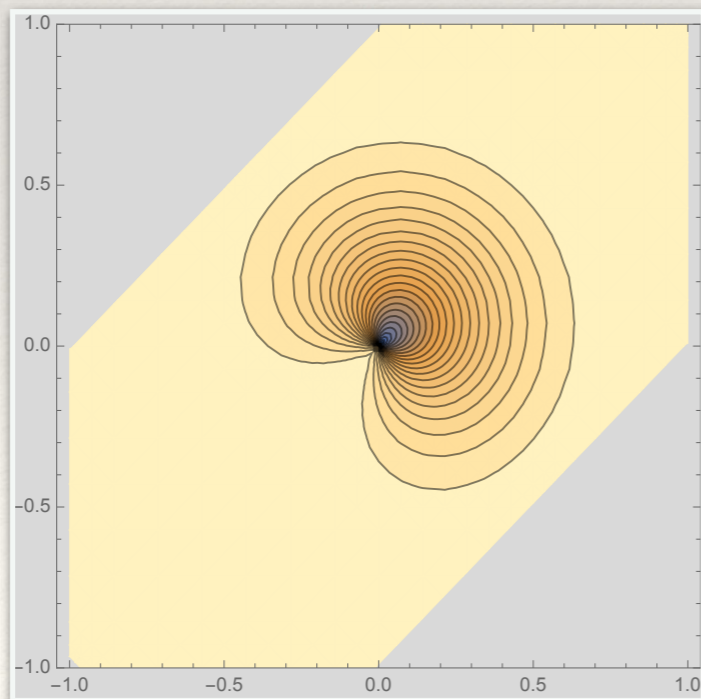
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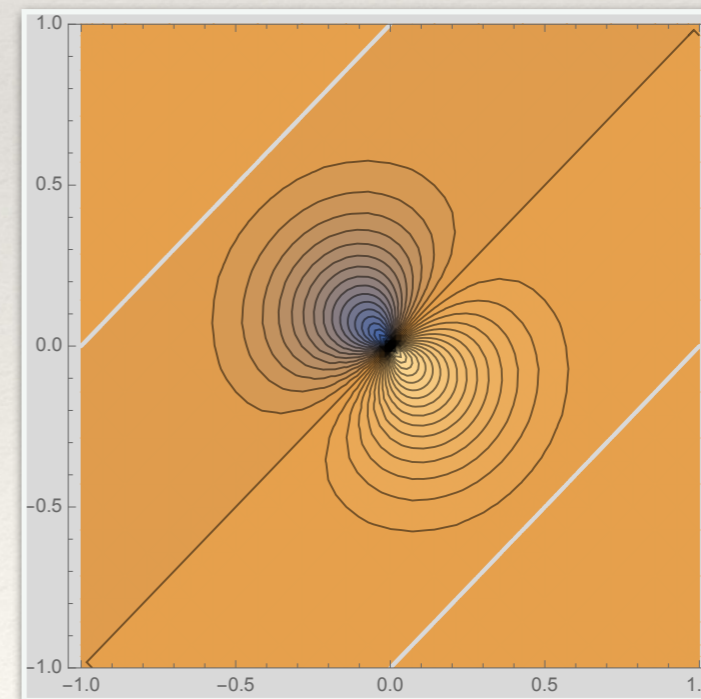
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$F_{FT}(x, x')$



$G_{FT}(x, x')$



Transverse Spin Asymmetries in Photon SIDIS

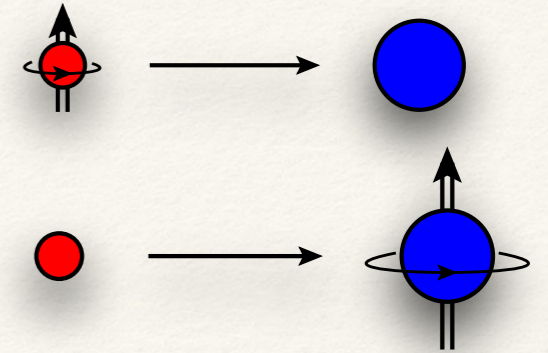
Transverse SSA in DIS

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$A_N = 0$ in 1- γ exchange [Christ, Lee, 1966]!

\Rightarrow **New effects** for a Two-Photon Exchange!

$$A_N = \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow}$$

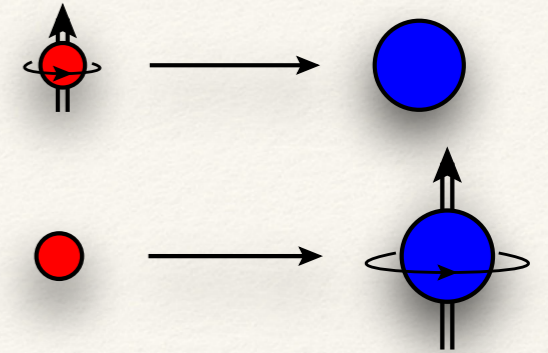


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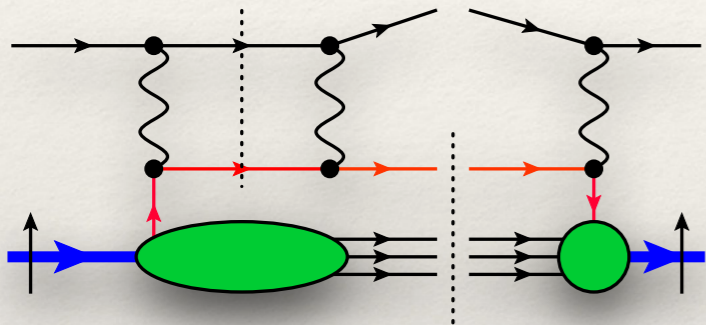
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Transverse Target Single Spin Asymmetry in DIS

[Metz, M.S., Goeke, PLB 2006]



$$A_{UT} \propto \alpha_{\text{em}} \frac{M}{Q} \left(\frac{a}{\varepsilon} + b \right) \sum_q e_q^3 x_B g_T^q(x_B)$$

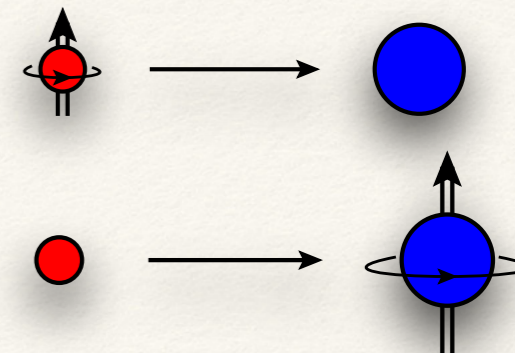
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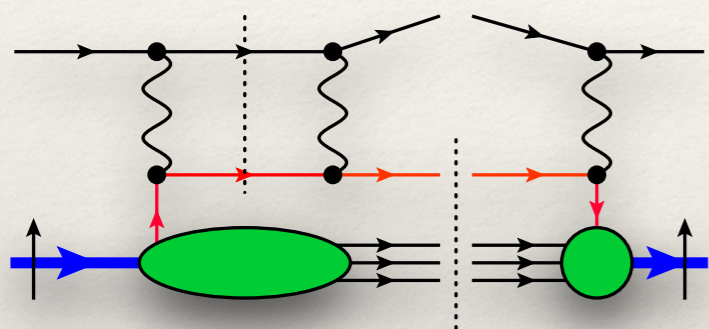
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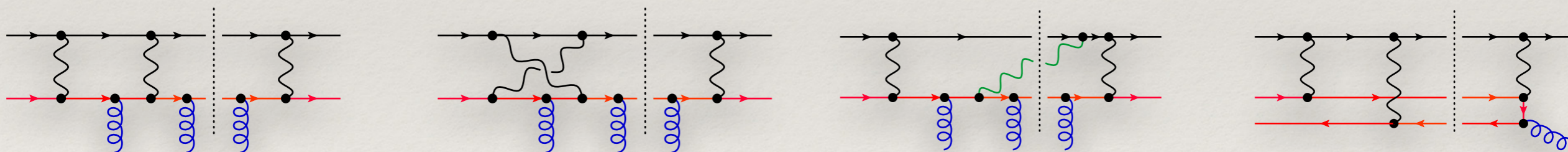
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Inclusion of 'Quark - Gluon Correlations'



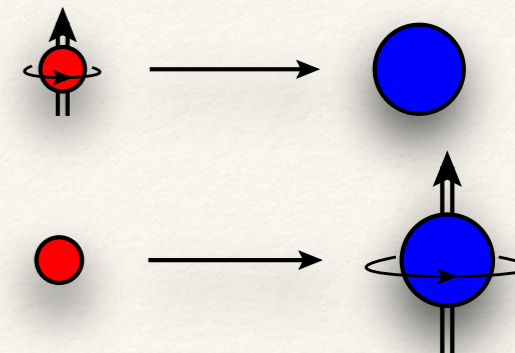
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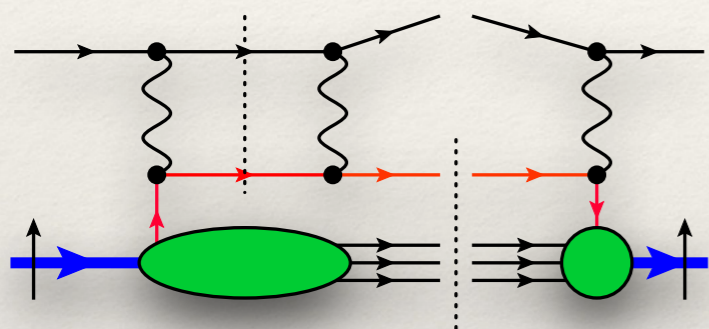
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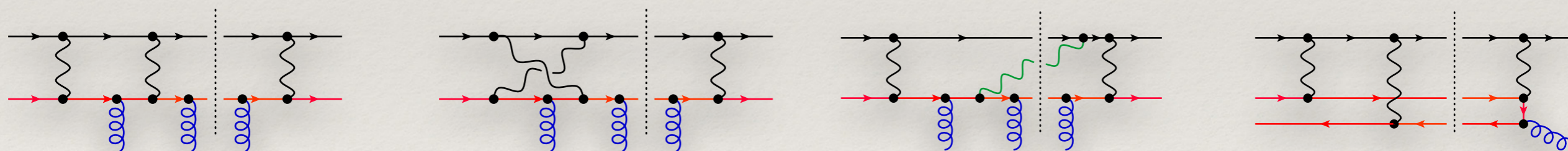
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Result: [MS; Metz, Pitonyak, Schäfer, Schlegel, Vogelsang, PRD 2012, 2013]

$$A_{UT} \propto \alpha_{\text{em}} \frac{M}{Q} \left[\int dx' (\hat{\sigma}_+ F_{FT}(x_B, x') + \hat{\sigma}_- G_{FT}(x_B, x')) + \hat{\sigma}^\gamma F_{FT}^\gamma(x_B, x_B) \right]$$

Photon SIDIS: $e + p \longrightarrow e + \gamma + X$

[Albaltan, Prokudin, M.S., in preparation]

Idea: Circumvent Christ - Lee theorem: hard, isolated real photon emission

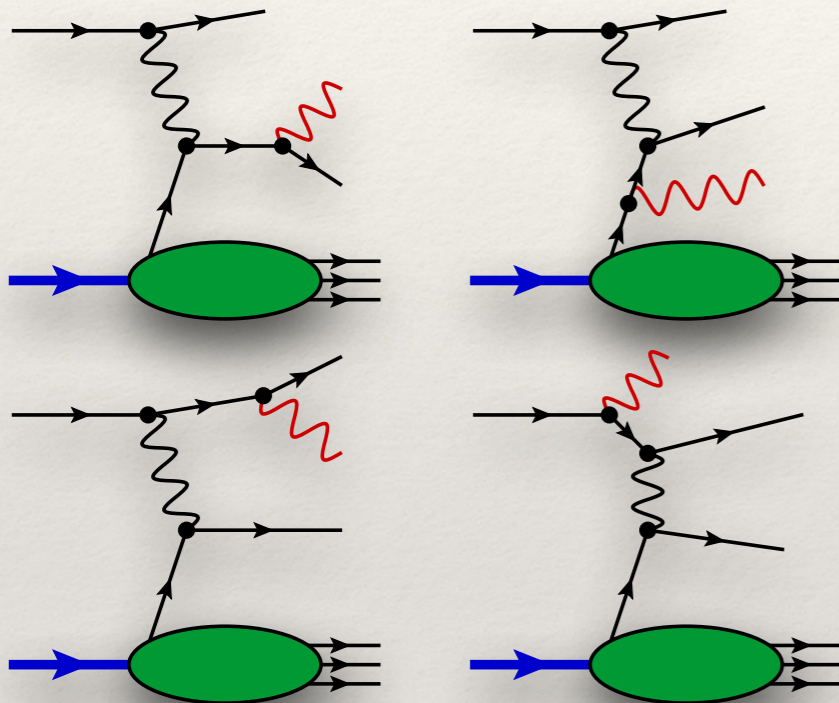
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unpolarized cross section in the parton model

[Brodsky, Gunion, Jaffe, PRD 1972]



- avoid photon fragmentation: isolated photons
- collinear factorization:
information on final quark is integrated out
- LO result:

$$E_\gamma E_e \frac{d\sigma_{UU}}{d^3\vec{P}_\gamma d^3\vec{P}_e} = \sum_q \left[e_q^2 \hat{\sigma}_2 + e_q^3 \hat{\sigma}_3 + e_q^4 \hat{\sigma}_4 \right] f_1^q(\bar{x})$$

- two scales:

$$Q^2 = -(l - l' - P_\gamma)^2$$

$$\tilde{Q}^2 = -(l - l')^2$$

- two 'Bjorken-x':

$$x_B = \frac{Q^2}{2P \cdot (l - l' - P_\gamma)}$$

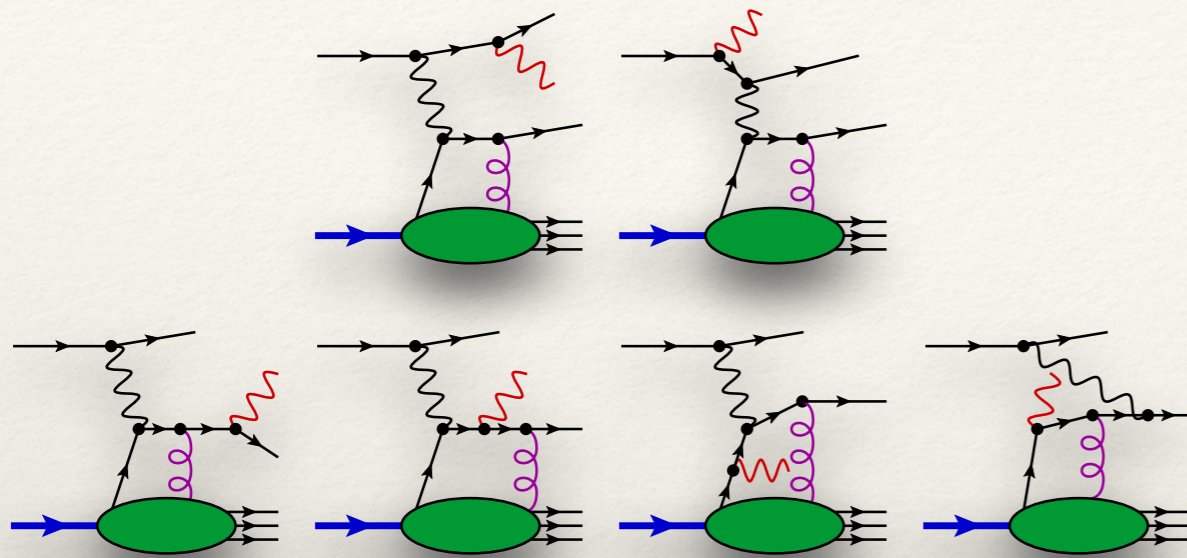
$$\tilde{x}_B = \frac{\tilde{Q}^2}{2P \cdot (l - l')}$$

BGJ - criterion for parton model dominance:

$$Q^2, \tilde{Q}^2, |Q^2 - \tilde{Q}^2| \gg M^2$$

Transverse SSA in photon SIDIS

Include *intrinsic, kinematical & dynamical* twist - 3 contributions



At tree-level (LO):

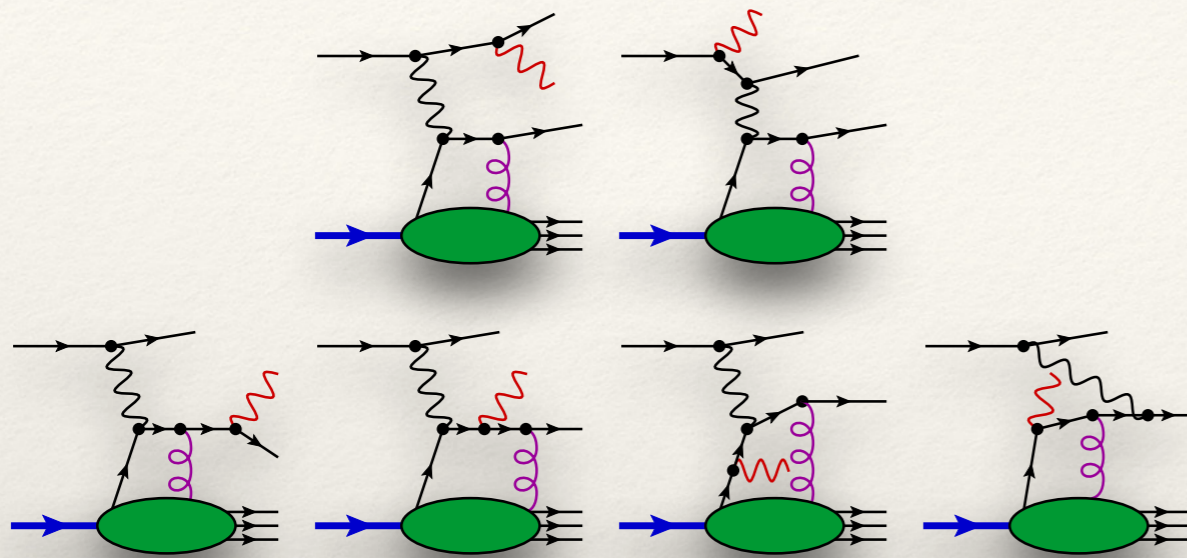
No contribution from g_T and $g_{1T}^{(1)}$
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Quark - Gluon correlations:

- 1) **Soft Gluon Poles:** $F_{FT}(x_B, x_B)$
- 2) **Soft Fermion Poles:** $F_{FT}(x_B, 0)$
- 3) **Hard Poles:** $F_{FT}(x_B, \tilde{x}_B)$

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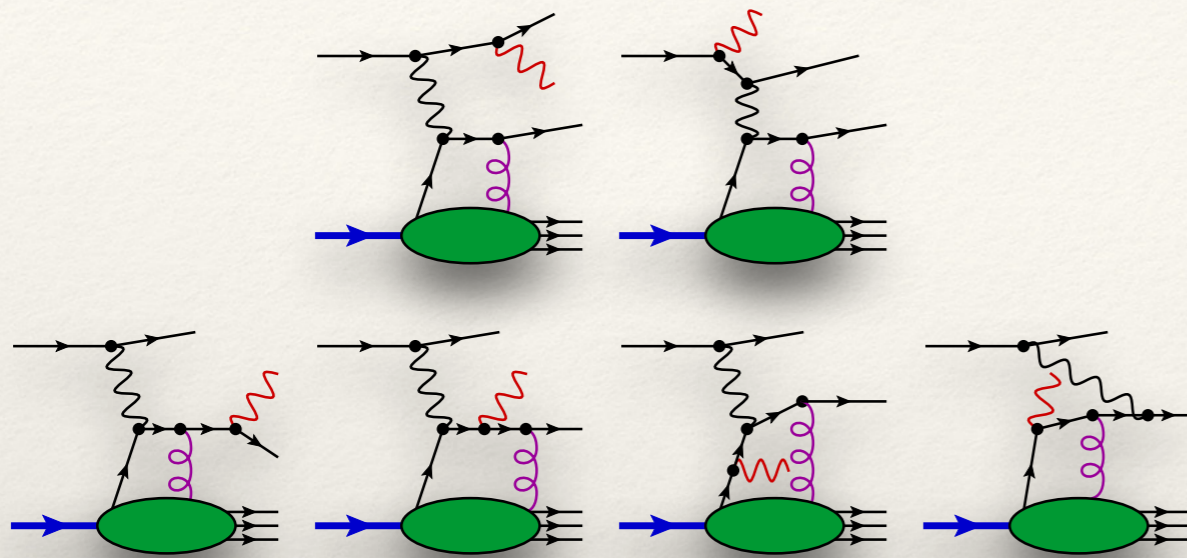
$$E_\gamma E' \frac{d\sigma_{UT}}{d^3\vec{P}_\gamma d^3\vec{P}_e} = \sum_q \left[\hat{\sigma}_{+,HP} F_{FT}^q(x_B, \tilde{x}_B) + \hat{\sigma}_{+,SFP} F_{FT}^q(x_B, 0) + \hat{\sigma}_{-,HP} G_{FT}^q(x_B, \tilde{x}_B) + \hat{\sigma}_{-,SFP} G_{FT}^q(x_B, 0) \right]$$

Soft Gluon Poles vanish !

Bethe-Heitler contribution vanishes
(Christ - Lee theorem)!

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\Rightarrow unique process to directly study “off-diagonal” support
of twist - 3 Quark - Gluon Correlation functions!

Is the transverse nucleon SSA feasible in an experiment
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Comparison to DIS cross section at EIC: suppression by $\alpha = 1/137$

Event rate smaller

→ $(1/200 - 1/1000) \times$ DIS 'total' (binned) cross section at EIC

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Size of SSA

→ unknown, any estimate would be pure speculation

→ Probably small, no BH contribution, (partial) cancellation of Compton & interference contribution (charge signs)

→ but too small?

Any experimental information would help...

Summary

- ❖ Transverse Spin Polarization: Long history, measured in ep/pp-collisions, theoretical treatment more complicated
- ❖ We can learn about the parton dynamics in the nucleon, e.g., transverse forces, non-perturbative QCD EoM and LIR are crucial
- ❖ Photon SIDIS: May be able to scan the support of dynamical twist-3 functions point-by-point at LO.
- ❖ Experimental opportunity at EIC, COMPASS, JLab
 - input would help our understanding of quark-gluon correlation
 - valuable for evolution of qgq functions and TMDs.