

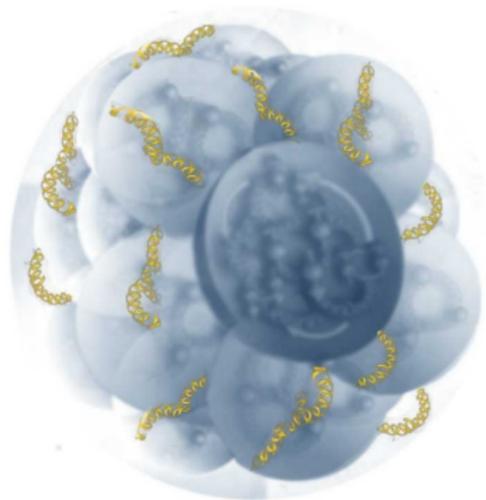
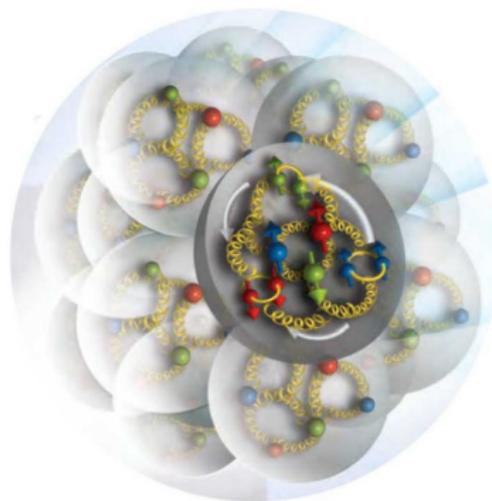
Search for Exotic Glue in Nuclei

Gluonic Transversity in Polarized DIS

J. Maxwell

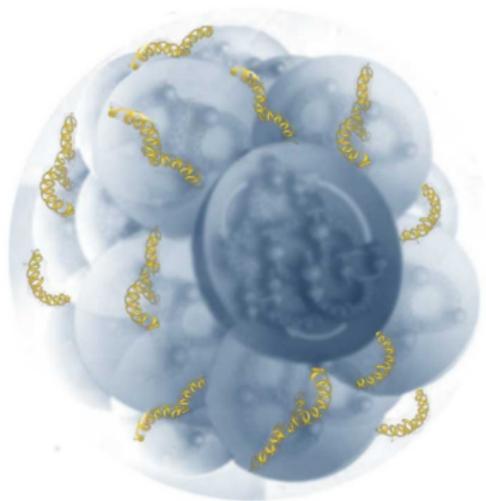
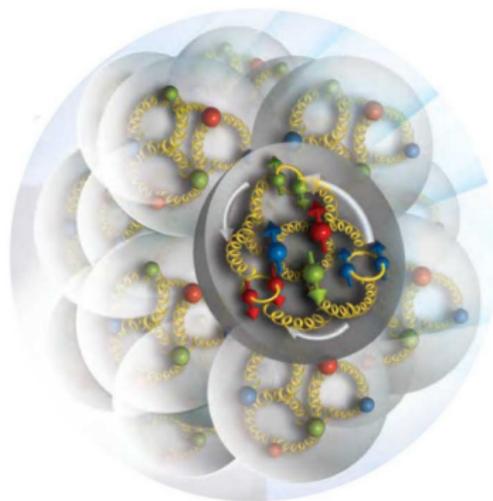
 **Jefferson Lab**

CLAS Collaboration Meeting
June 19th, 2019



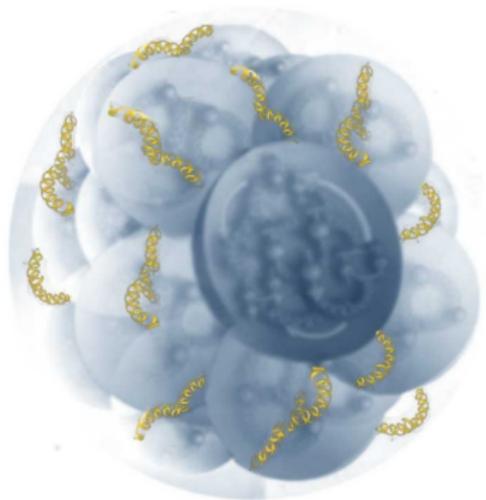
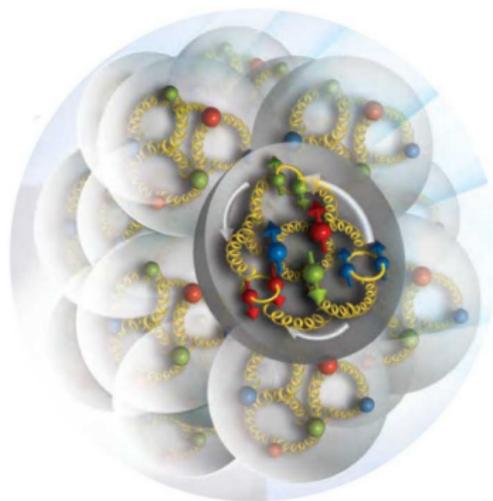
Outline

- 1 Double Helicity-Flip Structure Function
 - Lattice Calculations
 - Measurement Approaches
- 2 Jefferson Lab Measurement
 - JLab Polarized Target
 - Hall B?
- 3 Gluonometry at the EIC



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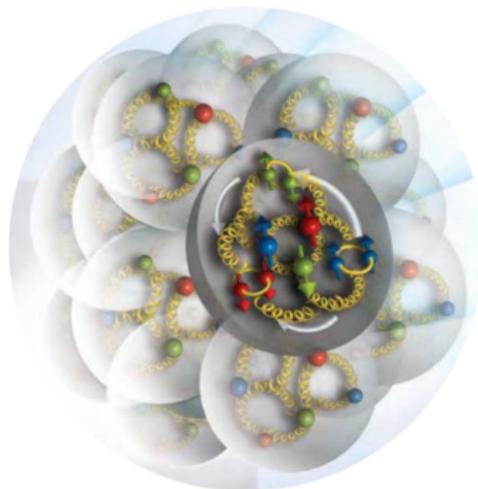


Gluon Structure of Nuclei

- Understanding glue is a key challenge of NP and central goal of EIC

Studying gluons is tricky

- Gluon does not couple to photon
 - Probed indirectly by electron scattering from nuclei
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- A nuclear glue effect, free from contributions of any nucleon, could offer invaluable view of nuclear structure
 - "Nuclear Gluonometry" (Jaffe, Manohar, 1989) offers a probe sensitive **only** to gluonic states in the nucleus: $\Delta(x, Q^2)$

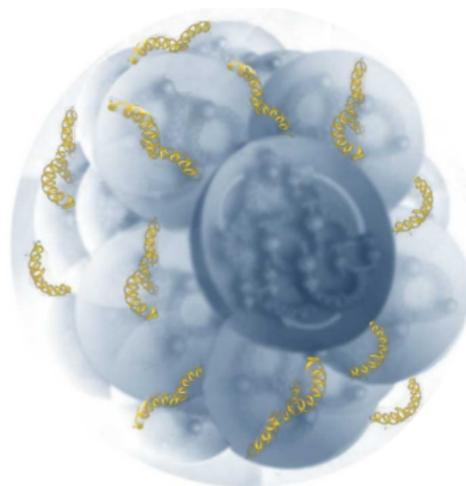


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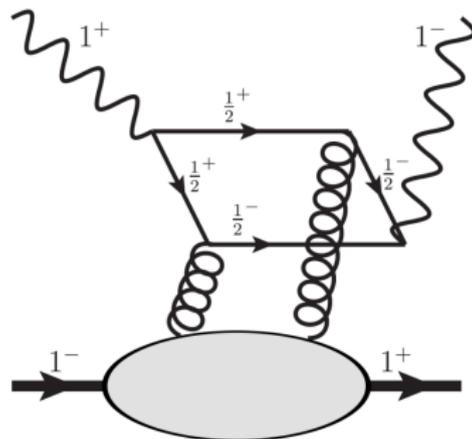
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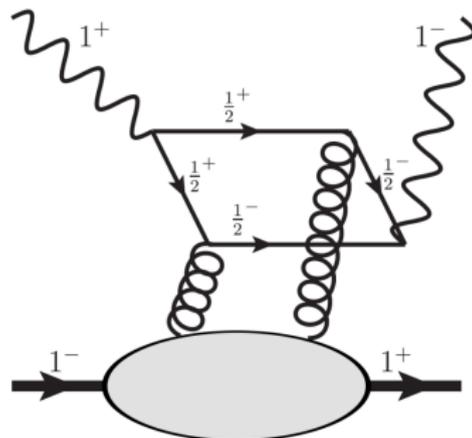
Double Helicity-Flip Structure Function $\Delta(x, Q^2)$

- $\Delta(x, Q^2)$ corresponds to helicity amplitude $A_{+-,-+}$
 - Photon helicity flip of two
 - Unavailable to bound nucleons or pions
 - Purely gluonic observable
- Hadrons: Gluonic Transversity
- Nuclei: Exotic Glue
 - Gluons not associated with an individual nucleon
- Unpolarized e beam on transversely polarized nuclei, spin ≥ 1
 - Primary challenge of measurement is polarized target or source
- Moments calculable in Lattice QCD



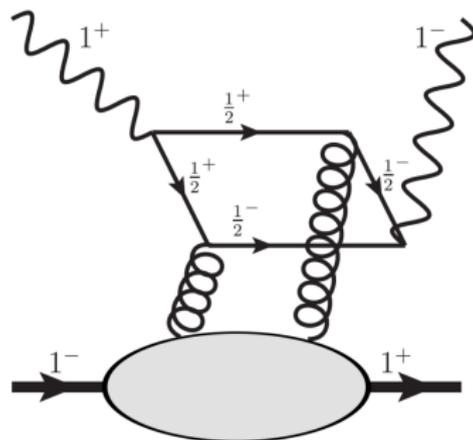
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Lattice QCD Guidance for Δ

- Initial calculations for first moment of Δ on spin-1 $\phi (s\bar{s})$
 - $m_\pi = 405$ MeV
 - Gave definitive signal¹
- Following year, first moment of Δ calculated on non-physical d
 - $m_\pi = 806$ MeV
 - Again definitive signal was seen²
- Results have generated significant interest in an observable mostly ignored since 1989
- Calculation with physical d underway

¹Detmold, Shanahan, P.Rev.D 94, 2016

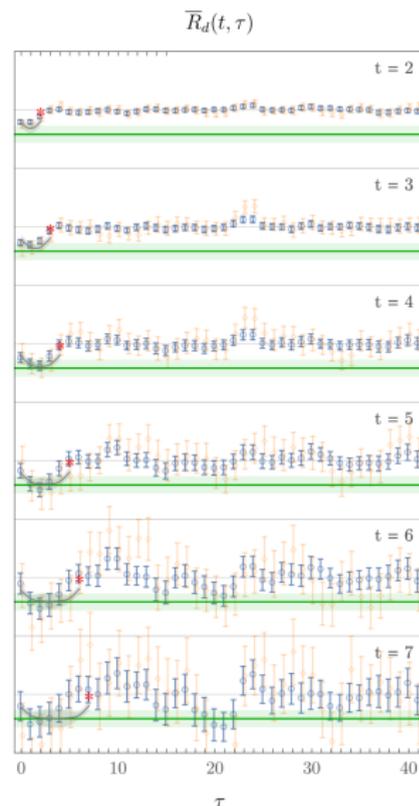
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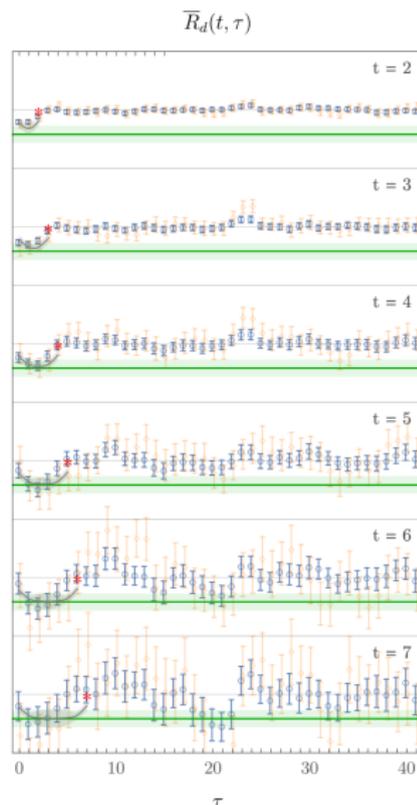


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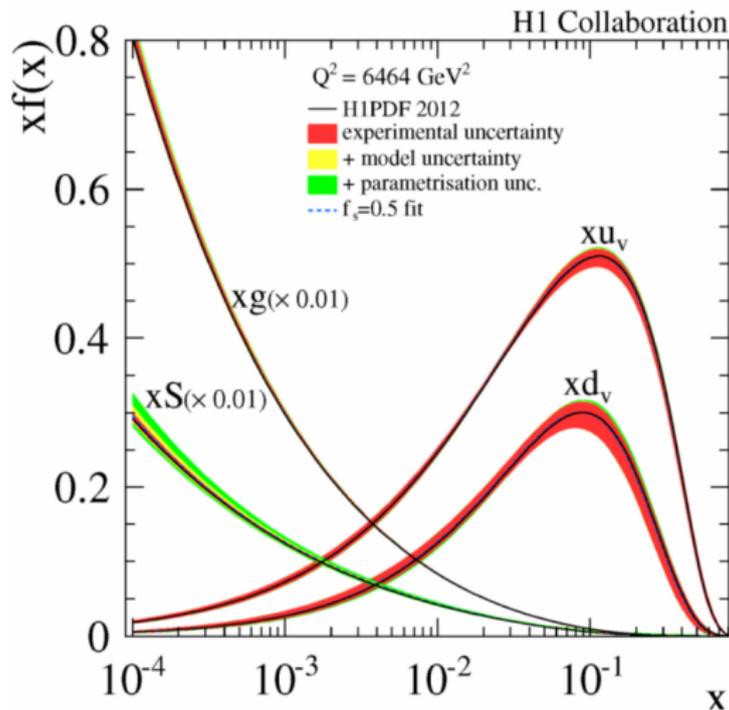
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Where do we start looking?



Measuring $\Delta(x, Q^2)$ via DIS

- Transversely aligned, spin-1 target and unpolarized electron incident from $-z$
- In the Bjorken limit, double helicity component of the hadronic tensor $W_{\mu\nu, \alpha\beta}^{\Delta=2}(E, E')$ becomes (dropping higher twist structure functions)¹:

$$\lim_{Q^2 \rightarrow \infty} \frac{d\sigma}{dx dy d\phi} = \frac{e^4 M E}{4\pi^2 Q^4} \left(xy^2 F_1(x, Q^2) + (1-y)F_2(x, Q^2) - \frac{x(1-y)}{2} \Delta(x, Q^2) \cos 2\phi \right)$$

¹Jaffe, Manohar, Phys Letters B 223 (2) (1989).

For a spin-1 target polarized at angle θ_m from the z -axis and electron incident from $-z$, target spin $\lambda_m = (1, 0, -1)$:

$$\frac{d\sigma}{dx dy d\phi}(\lambda_m) = \frac{2y\alpha^2}{Q^2} \left(F_1 + \frac{2}{3}a_m b_1 + \frac{1-y}{xy^2} \left(F_2 + \frac{2}{3}a_m b_2 \right) - \frac{1-y}{y^2} c_m \sin^2 \theta_m \Delta(x, Q^2) \cos(2\phi) \right)$$

with

$$a_m = \frac{1}{4} c_m (3 \cos^2 \theta_m - 1)$$

$$c_m = 3|\lambda_m| - 2$$

Differences of cross sections: N_+, N_0, N_- for $\lambda_m = (1, 0, -1)$

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Differences of cross sections: N_+, N_0, N_- for $\lambda_m = (1, 0, -1)$

Average over Polarization: $N_+ + N_- + N_0 \Rightarrow \bar{\sigma}$

- $c_+ + c_0 + c_- = 0$

$$\frac{d\bar{\sigma}}{dx dy d\phi} = \frac{2y\alpha^2}{Q^2} \left(F_1 + \frac{1-y}{xy^2} F_2 \right)$$

- Of course, no Δ dependence
- Δ also cancels out of vector polarization difference
 $(N_+ - N_0) + (N_0 - N_-) = N_+ - N_-$
 - $c_+ - c_- = 0$

Tensor Polarization: $(N_+ - N_0) - (N_0 - N_-) \Rightarrow \Delta\sigma$

- $c_+ - 2c_0 + c_- = 6$

$$\frac{d\Delta\sigma}{dx dy d\phi} = \frac{2y\alpha^2}{Q^2} \left((3 \cos^2 \theta_m - 1) \left(b_1 + \frac{1-y}{xy^2} b_2 \right) - \frac{1-y}{y^2} 6 \sin^2 \theta_m \Delta(x, Q^2) \cos(2\phi) \right)$$

- Tensor structure functions b_1, b_2 contribute significantly
- Unless! $(3 \cos^2 \theta_m - 1) = 0 \Rightarrow \theta_m = 54.7^\circ$

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Difference of Polarized and Unpolarized:

$$N_+ - \bar{N} = N_+ - \frac{1}{3}(N_+ + N_- + N_0) = \frac{1}{3}(N_+ - N_0) \Rightarrow \hat{\sigma}$$

- $c_+ - c_0 = 1$

$$\frac{d\hat{\sigma}}{dx dy d\phi} = \frac{2y\alpha^2}{Q^2} \left(\frac{1}{6}(3 \cos^2 \theta_m - 1)(b_1 + \frac{1-y}{xy^2}b_2) - \frac{1-y}{y^2} \sin^2 \theta_m \Delta(x, Q^2) \cos(2\phi) \right)$$

- Again tensor structure functions b_1, b_2 contribute significantly unless $\theta_m = 54.7^\circ$

3 ways to measure $\Delta(x, Q^2)$

$$(3 \cos^2 \theta_m - 1) \left(b_1 + \frac{1-y}{xy^2} b_2 \right) - \frac{1-y}{y^2} \sin^2 \theta_m \Delta(x, Q^2) \cos(2\phi)$$

- ① Leverage $\cos(2\phi)$ to isolate $\Delta(x, Q^2)$ dependence
 - Need azimuthal detector acceptance
- ② Form tensor asymmetry: $\mathcal{A} = \frac{1}{A} \frac{N_+ + N_- - 2N_0}{N_+ + N_- + 2N_0}$
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 - Change polarization to produce N_+, N_- and N_0 yields
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Transverse Polarized Target Nuclei

- Need a spin ≥ 1 nucleus, but this is a **multi-nucleonic** effect
 - Expected larger in compact nuclei (like EMC effect?)
 - Perhaps explains enhanced LQCD signal with larger m_π , more compact d ?
- Deuteron? Should be investigated, but may not offer best chance for discovery.
 - Expect two nucleons to good approximation
- Something heavier: Li? $\alpha + d$
- Practical limitations from available polarized targets
 - Long history of polarized p and d in solid targets
 - Lithium Hydride and Deuteride: ${}^6\text{LiH}$, ${}^6\text{LiD}$, also ${}^7\text{LiH}$
 - Ammonia: ${}^{14}\text{NH}_3$, ${}^{14}\text{ND}_3$, also ${}^{15}\text{NH}_3$
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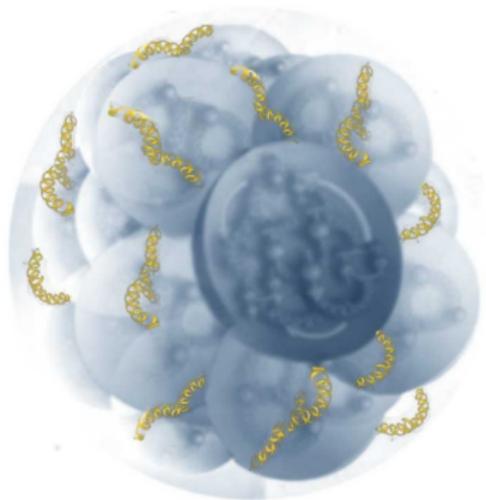
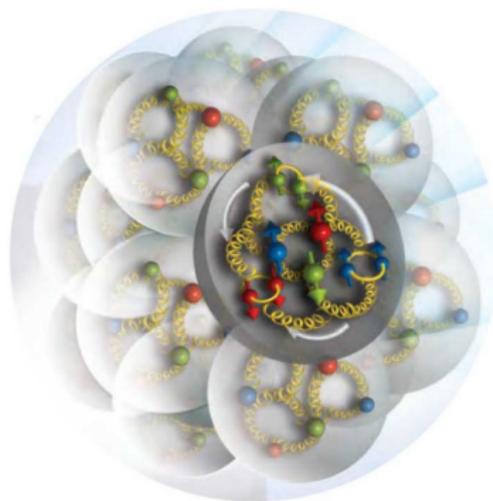
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JLab: Fixed, Polarized Target Approach?

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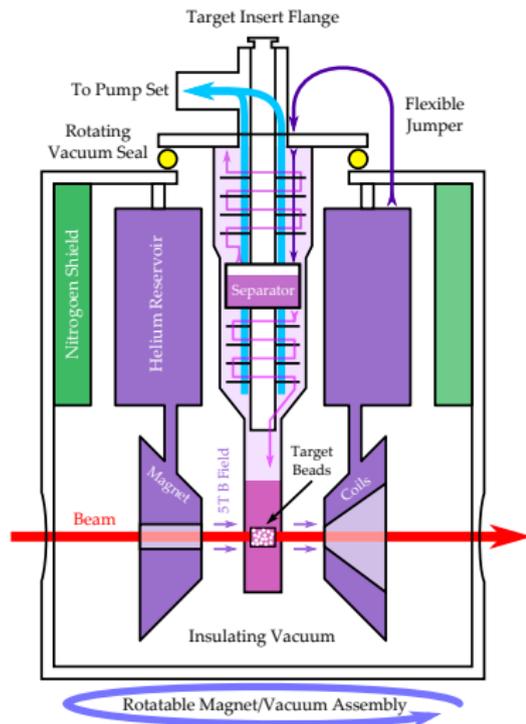
Kinematic Reach with 12 GeV CEBAF in Hall C

- 11 GeV, unpolarized e^- on fixed, polarized $^{14}\text{NH}_3$
- Preliminary SHMS Monte Carlo (Gaskell, Arrington)
 - Transverse (not 54.7° !) UVa magnet (M. Jones)

θ	E (GeV)	E' (GeV)	Q^2 (GeV $^2/c^2$)	x	Rate (Hz)
10.5	11	5	1.842	0.164	170
10.5	11	4	1.474	0.112	152
10.5	11	3	1.105	0.074	138
10.5	11	2	0.737	0.044	100
15	11	5	3.748	0.333	28
15	11	4	2.999	0.228	30
15	11	3	2.249	0.15	32
15	11	2	1.499	0.089	34

JLab Solid Polarized Target

- Dynamic Nuclear Polarization
 - 5 T field, 1 K ^4He evap. fridge
 - Dope material with paramagnetic radicals (NH_3 : NH_2 or H)
 - Leverage $e - p$ spin coupling
 - μ -waves drive polarizing transitions
 - e relaxes to flip-flop with new p
- Irradiated Ammonia: 95% p , 40% d
 - Beam current < 100 nA
 - P decay: anneals and replacement
- Workhorse DIS technique at SLAC, JLab; 2012's g_p^2 most recently², CLAS12 polarized target under testing



²Pierce, Maxwell, NIM A 738 (2014).

Polarization, Tensor *Alignment* and DNP

$$P = (N_+ - N_0) + (N_0 - N_-)$$

$$= N_+ - N_-$$

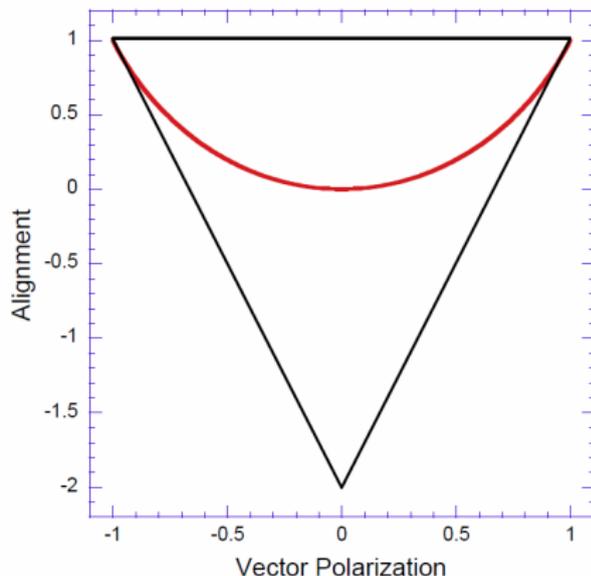
$$A = (N_+ - N_0) - (N_0 - N_-)$$

$$= 1 - 3N_0$$

- Polarization and alignment can be anywhere in the black triangle
- At *equal spin temperature*, can be only on red curve:

$$A = 2 - \sqrt{4 - 3P^2}.$$

- For $P = 40\% \Rightarrow A = 13\%$

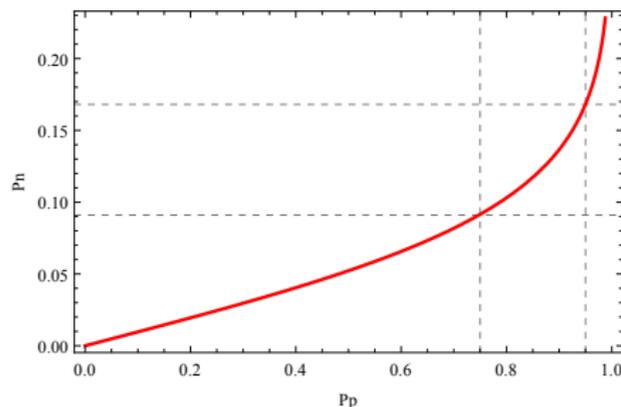


Nitrogen Polarization in Ammonia: Not Easy

- We can also relate polarization of N to p at EST:

$$P_N = \frac{4 \tanh((\omega_N/\omega_p) \operatorname{arctanh}(P_p))}{3 + \tanh^2((\omega_N/\omega_p) \operatorname{arctanh}(P_p))}$$

- At 95% p : 17% N
 - $P_N = 17\% \Rightarrow A_N = 2\%$
- NMR measurement is difficult
 - Peaks too far apart for one NMR scan (2.4 MHz)
 - Overcome at SMC with 2 sweeps, changing B field³



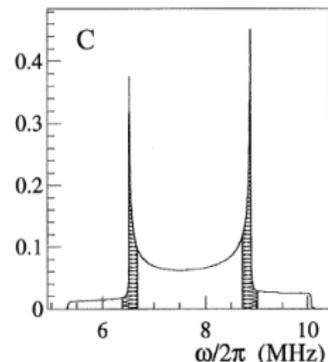
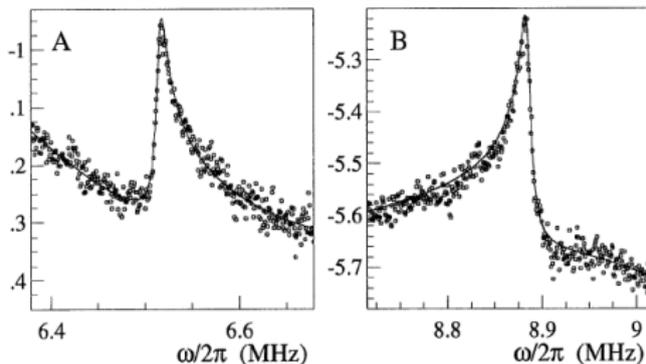
³B. Adeva, NIM A 419 (1998).

Nitrogen Polarization in Ammonia: Not Easy

- We can also relate polarization of N to p at EST:

$$P_N = \frac{4 \tanh((\omega_N/\omega_p) \operatorname{arctanh}(P_p))}{3 + \tanh^2((\omega_N/\omega_p) \operatorname{arctanh}(P_p))}$$

- At 95% p : 17% N
 - $P_N = 17\% \Rightarrow A_N = 2\%$
- NMR measurement is difficult
 - Peaks too far apart for one NMR scan (2.4 MHz)
 - Overcome at SMC with 2 sweeps, changing B field³

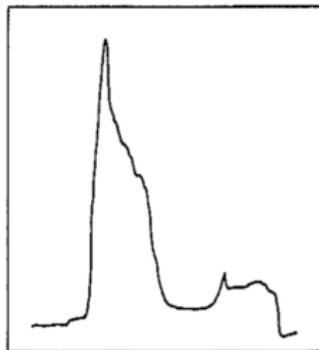


³B. Adeva, NIM A 419 (1998).

Techniques to Improve P_N, A_N

- “RF Hole Burning”⁴
 - Used for ND₃, crucial for b_1 exp.
 - Vast separation of NMR peaks in N will help.
- Cross Spin Transfer
 - Drop magnetic field (56mT) for cross relaxation of p Zeeman, N quadrupolar systems
 - SMC: 40% $P_N \Rightarrow 12\%$ A_N
- RF Spin Transfer
 - Same effect ultimately
 - RF to drive $p - N$ flip-flops
 - Allow dynamic pumping of N while μ -waves pump p

⁴P. Delheij, NIM A 251 (1986).

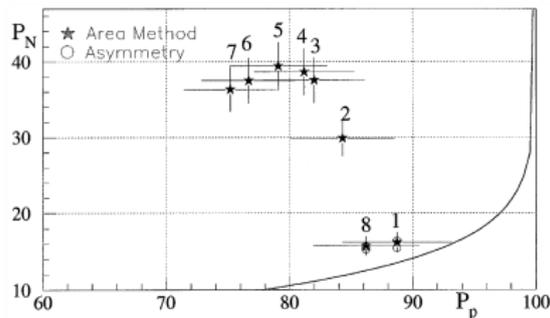


Starting polarization = 37%
RF-induced alignment “up to 20%”

W. Meyer and E. Schiling, Proceedings of 4th Int’l Workshop on Polarized Target Materials and Techniques, Bonn (1984) 165.

Techniques to Improve P_N, A_N

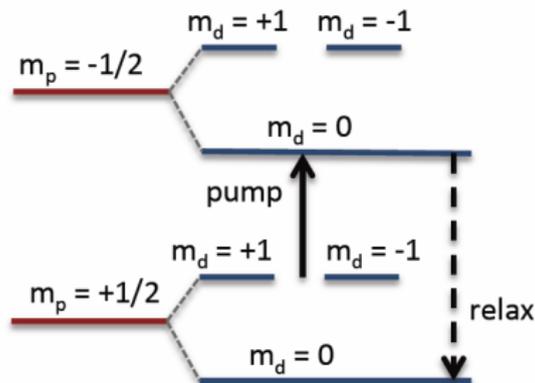
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Techniques to Improve P_N, A_N

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Jefferson Lab Letter of Intent 12-16-006 (Hall C)

- ~ 30 PAC days with solid polarized target⁵
 - Run with approved measurement of b_1 in Hall C
 - Ballpark 1% statistical error
 - Heavily dependent on achieved polarization
 - Systematic uncertainty on target polarization 4-5%
- LOI Reception, PAC 44
 - Encouragement with charges
 - Guidance on size of Δ from Lattice QCD
 - Polarized N target yet to be proven
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Nuclear Gluonometry in Hall B?

$$(3 \cos^2 \theta_m - 1) \left(b_1 + \frac{1-y}{xy^2} b_2 \right) - \frac{1-y}{y^2} \sin^2 \theta_m \Delta(x, Q^2) \cos(2\phi)$$

- ① $\cos(2\phi)$ offers $\Delta(x, Q^2)$ sensitivity
 - Need to produce polarizing field transverse to beam
 - Actually, only need to rotate out to 54.7°
- ② Form tensor asymmetry: $\mathcal{A} = \frac{1}{A} \frac{N_+ + N_- - 2N_0}{N_+ + N_- + 2N_0}$
 - Set target at $\theta_m = 54.7^\circ$
 - Yields at N_+ , N_- and N_0 separated in time: systematic headaches
- ③ Form difference of vector polarized and unpolarized cross sections
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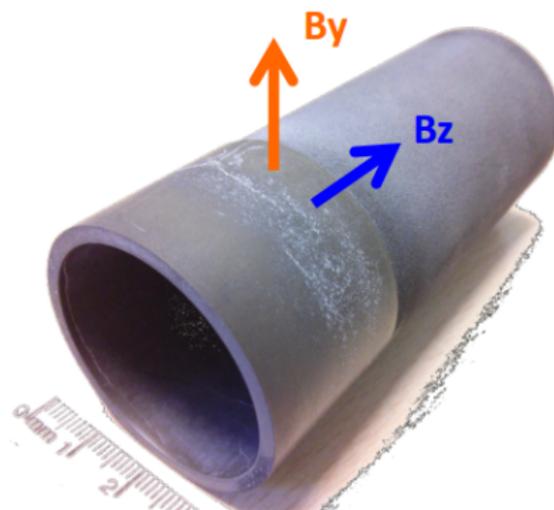
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Transverse Targets in CLAS12

- HD-Ice: transverse polarized target with bulk superconductor⁶
 - Lock in field at desired angle, then introduce into CLAS12
 - Bulk superconductor both shields longitudinal field and holds transverse
- For a DNP target, field uniformity requirement is more restrictive
 - Reduce target cell size?
- Second, transverse magnet inside CLAS12 5 T solenoid?
 - 3.5 T would rotate to 54.7°
 - BUT, need to keep $\int B \cdot dl$ low

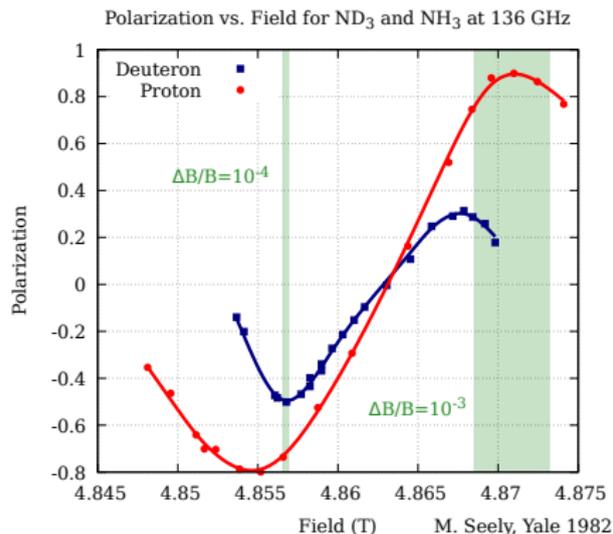


⁶Statera, NIM A 882 (2018)

Transverse Targets in CLAS12

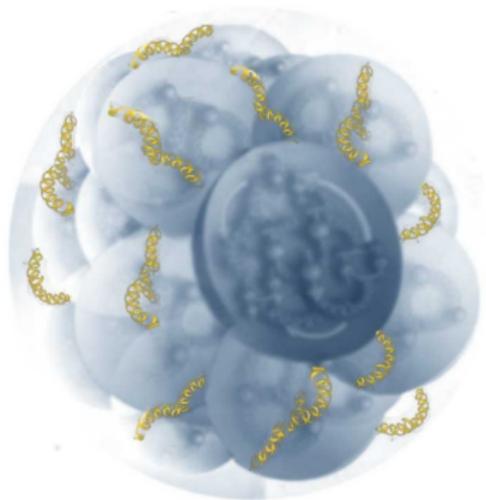
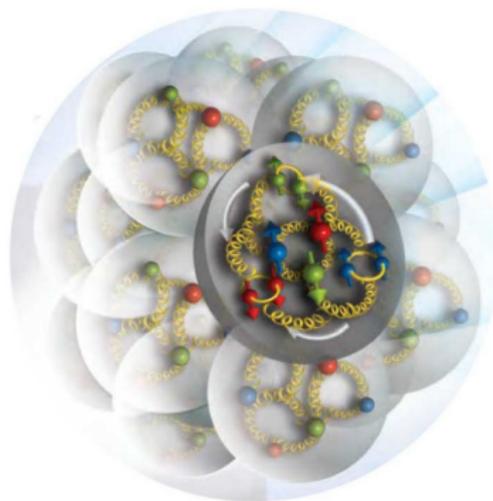
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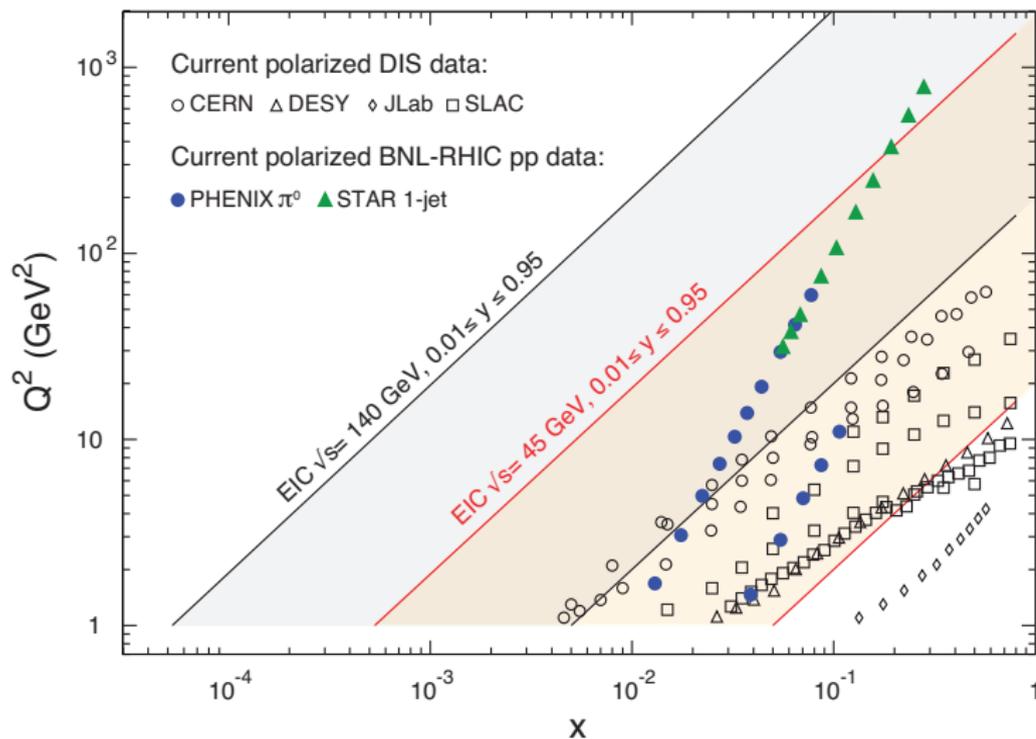


Outline

- 1 Double Helicity-Flip Structure Function
 - Lattice Calculations
 - Measurement Approaches
- 2 Jefferson Lab Measurement
 - JLab Polarized Target
 - Hall B?
- 3 Gluonometry at the EIC



Kinematic Reach at Electron-Ion Collider



EIC white paper

Polarized Ion Beam Possibilities

At EIC, $\Delta(x, Q^2)$ search becomes a problem of available ion sources and their corresponding depolarizing resonances.

Nucleus	Spin	Technique	Pol.	Flux	G
^2H	1	OP, ABS	100%	$1\mu\text{A}$	-0.14
^6Li	1	OP, ABS	88%	$2.4\mu\text{A}$	-0.18
^7Li	$\frac{3}{2}$	OP, ABS			1.53
^8Li	2	TFM	$\sim 1\%$		
^{10}B	3	Not known			
^{23}Na	$\frac{3}{2}$	OP, ABS	77%	$6.5\mu\text{A}$	0.55

Towards Design of an Optimized EIC Experiment

- Exploration of Δ in x , Q^2 , S , & A
 - How does effect change for different nuclear spin ≥ 1 ?
 - Spin-1/2 species important cross-check
 - How does effect change for different atomic masses?
 - Spin-1 ${}^6\text{Li}$ vs. Spin-3/2 ${}^7\text{Li}$
- Simulate measurement for Inclusive DIS on Nuclei
- Estimate running time for given statistical uncertainties
 - Species choice informed by simulation
 - Loss of luminosity compared to JLab made up for by lack of dilution, kinematic coverage
- Fixed target experiment at JLab would be a crucial first exploration, provide high- x data for the production of moments

Summary

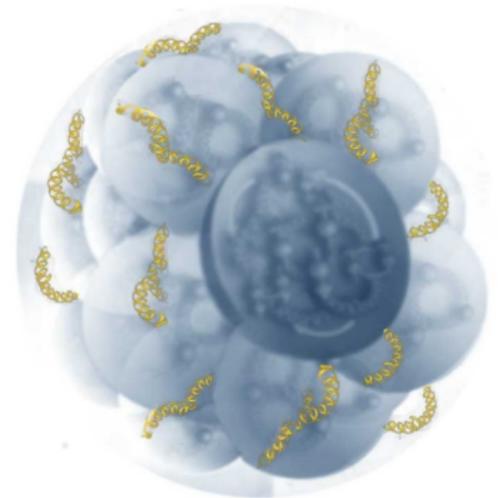
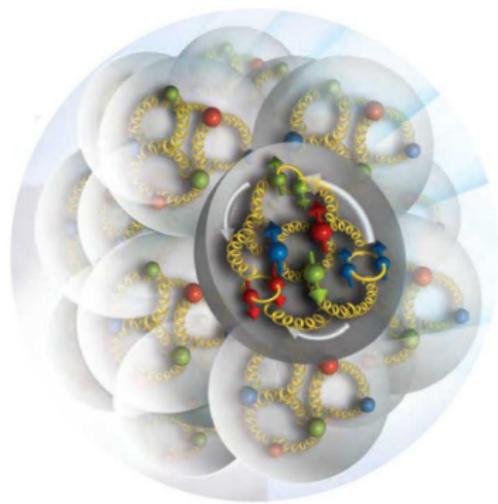
- $\Delta(x, Q^2)$ offers a rare look at gluonic components in the nucleus
 - Significant Lattice QCD result drives interest
 - Need spin ≥ 1 , polarized, nuclear target
 - Low x , where glue dominates, region of interest
- Jefferson Lab experiment still in pre-proposal stage
 - Polarized ^{14}N target fertile for development
 - Transverse target in Hall B worth investigating
 - SoLID also a possibility
- EIC capable of thorough search
 - Vast low x exploration
 - Polarized ion sources needed, Li and Na most attractive
 - Spin manipulation of polarized, “heavy” ions crucial
 - Initial investigations towards measurements at eRHIC (R.Milner) and JLEIC (JDM) begun

JLab Nuclear Gluonometry Collab:

- JLab: M. Jones, C. Keith, J. Maxwell, D. Meekins
- MIT: W. Detmold, R. Jaffe, R. Milner, P. Shanahan
- Univ. of Virginia: D. Crabb, D. Day, D. Keller, O. Rondon
- Oak Ridge: J. Pierce

Thanks to A. Zelenski, V. Morozov,
Y. Furletova

Thank you for your attention!



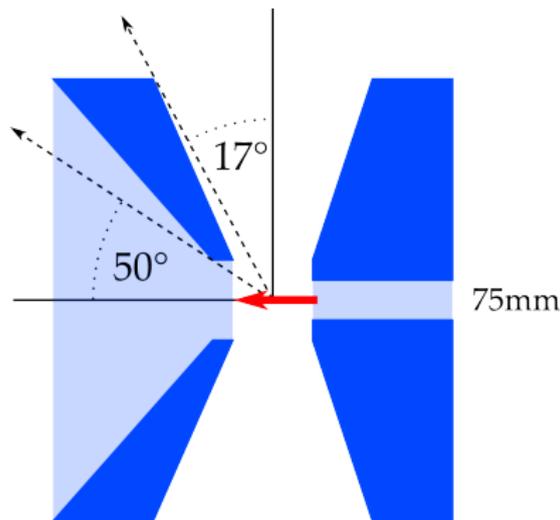
5 T Split-Pair Target Magnet

- Can we get $\theta_m = 54.7^\circ$
- Old Hall C Magnet, with largest opening angles, retired in 2012
 - Better than 10^{-4} uniformity in $3 \times 3 \times 3 \text{ cm}^3$ volume
- g_2^p ran with modified Hall B magnet
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 - $\sim \$500\text{k}$



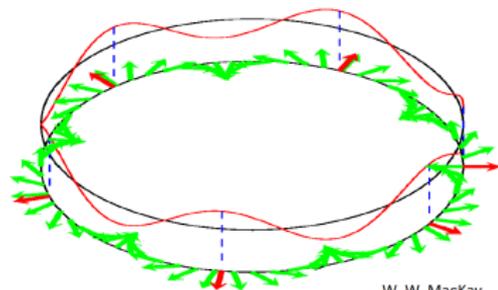
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Spin Manipulation in Ring

- Depolarizing resonances when spin precession frequency = frequency of perturbing B field⁷
- Imperfection: $\nu_s = G\gamma = n$
- Intrinsic: $\nu_s = G\gamma = Pn + \nu_y$
- Anomalous g -factor G
 - ⁷Li: G of 1.53 (like proton's 1.79) \Rightarrow easy
 - ⁶Li: G of -0.18 (like deuteron's -0.14) \Rightarrow hard
 - ²³Na: G of 0.55 could work at RHIC with more snakes
 - Figure-8 makes for easier manipulation at lower G

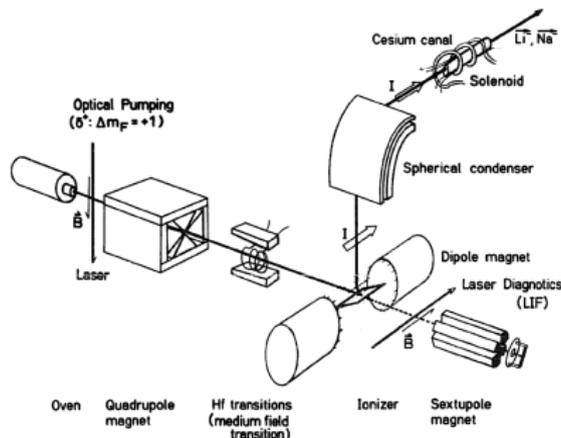


W. W. MacKay

⁷Bai, Courant *et al.*, BNL-96726-2012-CP, 2012.

Spin Polarized Alkali Sources

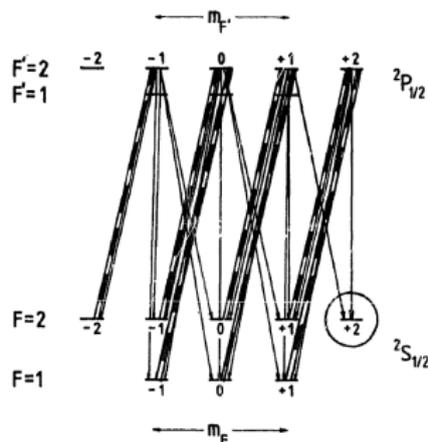
- Improved Heidelberg Source adds OP (1986)⁸
 - Laser pumped, modulated to pump both multiplets
 - ${}^6\text{Li}$: $A = 85\%$, ${}^{23}\text{Na}$: $A = 77\%$
 - Polarization limited due to lack of full ionization



⁸H. Reich, NIM A288 (1989)

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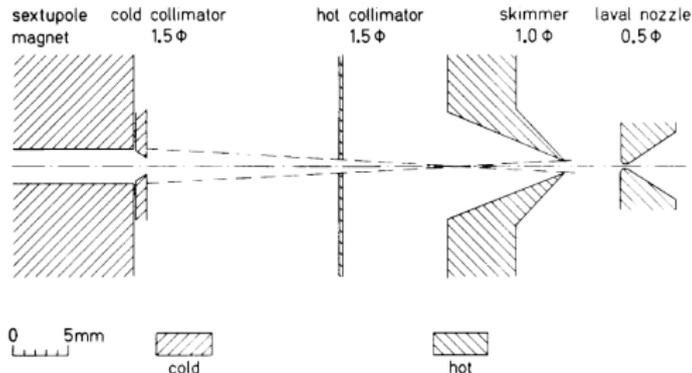
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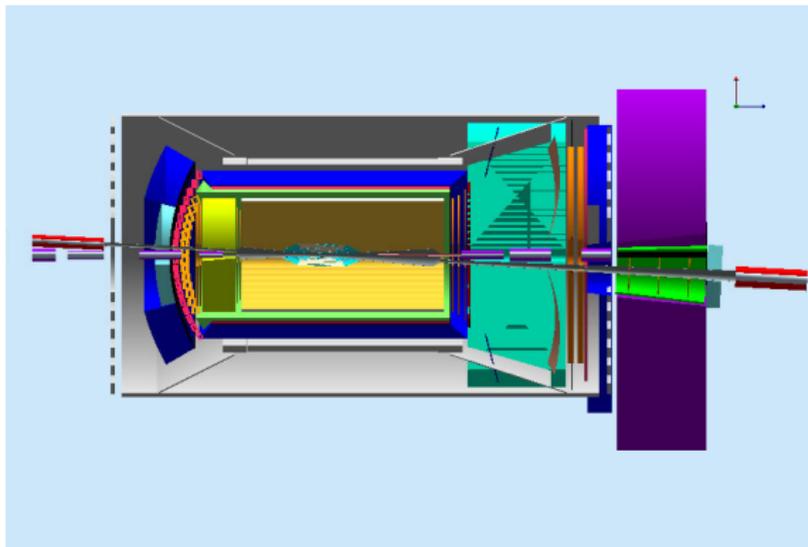
- Heidelberg Atomic Beam Polarized Source (1975)⁹
 - Laval nozzle, Sextupole Stern–Gerlach give $m = +1/2$
 - RF used for adiabatic transitions to fill other states
 - Surface ionization, heated tungsten strip
 - ${}^{6,7}\text{Li}$: $0.57 < |P| < 0.65$, 200 nA
 - ${}^{23}\text{Na}$: 50% losses to P and current in ionization



⁹E. Steffens, NIM 143 (1977)

Starting JLEIC Modeling

- First work with JLEIC detector geometry, eA kinematics
- BeAGLE generator for eD , eLi , eNa (thanks to V. Morozov, M.Baker), GEMC geometry (thanks to Y.Furletova)



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