Measurement of Polarization Observables for the Σ^+ Hyperon

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06/20/2019

Outline

1. Introduction
   - The Reaction, $\gamma p \to K^0\Sigma^+$
   - The Spectroscopy of Baryon Resonances
   - The Formalism of Hyperon Polarization
2. Experimental Approach and Data Analysis
   - The Experiment
   - The Reaction
   - The Extraction of Polarization Observables
3. Summary and Outlook
   - Summary

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Why $\gamma p \to K^0\Sigma^+$?

- Photoproduction of neutral kaons offers advantage over charged ones since photons cannot couple directly to (vanishing) charge of the meson.
- Data on isospin related channels $K^0\Sigma^+$ and $K^+\Sigma^0$ allow for disentanglement of contributions from N^* and Δ^* resonances.

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- Hyperon decay allows measurement of asymmetries, which allow for the extraction of hyperon recoil polarization P.
- Trade-off however is low cross-sections, leading to less statistics.
- The determination of the polarization observables allows for an understanding of the intermediate steps involved in the reaction.

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- These observables are sensitive to interference from different states.

		Status as seen in										
Particle	J^P	overall	$N\gamma$	$N\pi$	$\Delta\pi$	$N\sigma$	$N\eta$	ΛK	ΣK	$N\rho$	$N\omega$	$N\eta'$
N	$1/2^+$	****										
$N(1440)$	$1/2^+$	****	****	****	****	**	****					
$N(1520)$	$3/2^-$	****	****	****	****	***						
$N(1535)$	$1/2^-$	****	****	****	***	*	****					
$N(1650)$	$1/2^-$	****	****	****	***	*	****	*				
$N(1675)$	$5/2^-$	****	****	****	****	***						
$N(1680)$	$5/2^+$	****	****	****	****	***						
$N(1700)$	$3/2^-$	***	**	***	***	*	*			*		
$N(1710)$	$1/2^+$	****	****	****	*		***	**	*	*	*	
$N(1720)$	$3/2^+$	****	****	****	***	*	*	****	*	*	*	
$N(1860)$	$5/2^+$	**	*	**		*	*					
$N(1875)$	$3/2^-$	***	**	**	*	**	*	*	*	*	*	
$N(1880)$	$1/2^+$	***	**	*	**	*	*	**	**		**	
$N(1895)$	$1/2^-$	****	****	*	*	*	****	**	**	*	*	****
$N(1900)$	$3/2^+$	****	****	**	**	*	*	**	**		*	**
$N(1990)$	$7/2^+$	**	**	**			*	*	*			
$N(2000)$	$5/2^+$	**	**	*	**	*	*				*	
$N(2040)$	$3/2^+$	*		*								
$N(2060)$	$5/2^-$	***	***	**	*	*	*	*	*	*	*	
$N(2100)$	$1/2^+$	***	**	***	**	**	*	*		*	*	**
$N(2120)$	$3/2^-$	***	***	**	**	**		**	*		*	*
$N(2190)$	$7/2^-$	****	****	****	****	**	*	**	*	*	*	
$N(2220)$	$9/2^+$	****	**	****			*	*	*			
$N(2250)$	$9/2^-$	****	**	****			*	*	*			
$N(2300)$	$1/2^+$	**		**								
$N(2570)$	$5/2^-$	**		**								
$N(2600)$	$11/2^-$	***		***								
$N(2700)$	$13/2^+$	**		**								

**** Existence is certain.
*** Existence is very likely.
** Evidence of existence is fair.
* Evidence of existence is poor.

Figure 1: Particle Data Group.

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- Circularly polarized photons may have their polarization transferred fully or partially to spin orientation of hyperons within reaction plane.
- The C_x and C_z double polarization observables allow for a characterization of the transferred polarization from incident beam to recoiling hyperon along the orthonormal axes in the scattering plane.

Polarization Observables

Spin-Dependent Cross-Section for $K^0\Sigma^+$ Photoproduction

$$\rho_Y \frac{d\sigma}{d\Omega_{K^+}} = \left.\frac{d\sigma}{d\Omega_{K^+}}\right|_{\text{unpol}} \{1 + \sigma_y P + P_\odot(C_x\sigma_x + C_z\sigma_z)\}$$

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Polarization Components

$$P_{\Sigma^+_x} = P_\odot C_x$$

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- Transverse (induced) polarization $P_{\Sigma_y^+}$ is equivalent to P observable.
- The $\hat{x}$ and $\hat{z}$ components of hyperon polarization are proportional to C_x, C_z via degree of beam polarization $P_\odot$.

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Scattering Frame

$$\gamma(k) + p(q_1) \to K^0(q_2) + \Sigma^+(q_3),$$

$$\hat{y} = \frac{\vec{k} \times \vec{q_2}}{|\vec{k} \times \vec{q_2}|}, \qquad \hat{z} = \frac{\vec{k}}{|\vec{k}|}, \qquad \hat{x} = \vec{y} \times \vec{z}.$$

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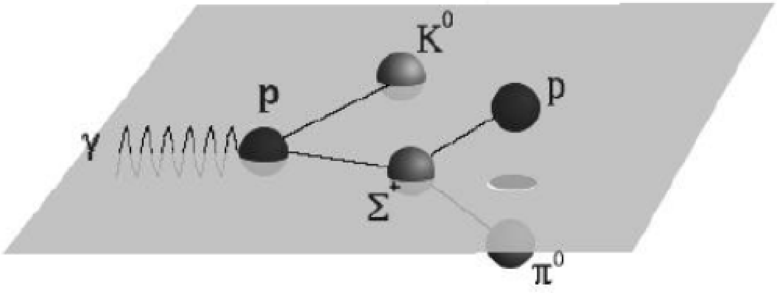

Figure 2: Scattering plane of $\gamma p \to K^0\Sigma^+$.

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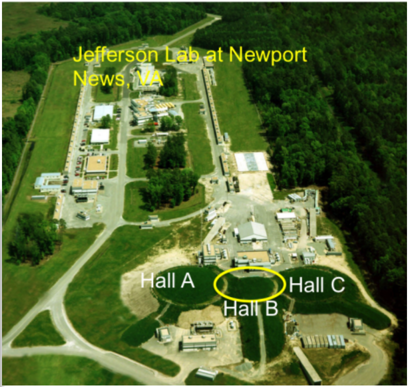

Figure 3: Hall B at JLab, home to the g12 experiment.

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g12 Data Cuts

- **Vertex Cut:** -110.0 cm $<$ z-vertex $<$ -70.0 cm.
- **Timing Cut:** $|\Delta_{\mathrm{TBID}}| < 1$ ns.
- **Particle ID Cut:** $\Delta\beta = |\beta_{\mathrm{c}} - \beta_{\mathrm{m}}| \leq 3\sigma$.
- **Fiducial Cut:** nominal scenario.

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Missing Momentum

$$x^\mu = k^\mu + P^\mu - \sum_{i=1}^{2,3} p_i^\mu.$$

Missing Mass

$$m_X^2 = x^\mu x_\mu.$$

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- Q-factor used as event weight in order to determine signal contribution to physical distributions.
- Due to $K^0 \rightarrow \pi^+\pi^-$ and $\Sigma^+ \rightarrow p\pi^0$ correlation, reference quantity can be either invariant mass.

Q-factor background subtraction

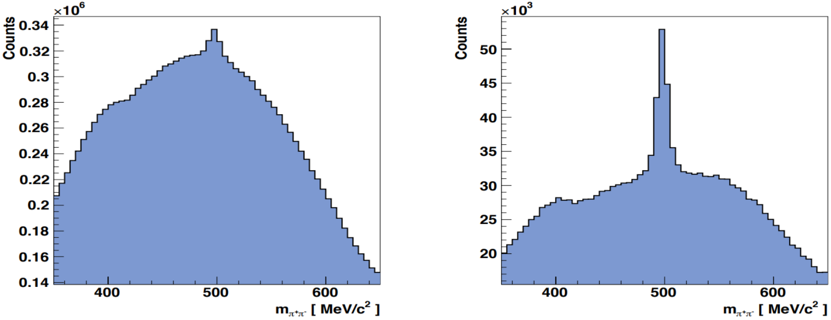

Figure 4: **Left:** Raw distribution of all g12 $\pi^+\pi^-\pi^0$ events. **Right:** Same invariant mass $\pi^+\pi^-$ after Σ^+ mass-cut has been employed.

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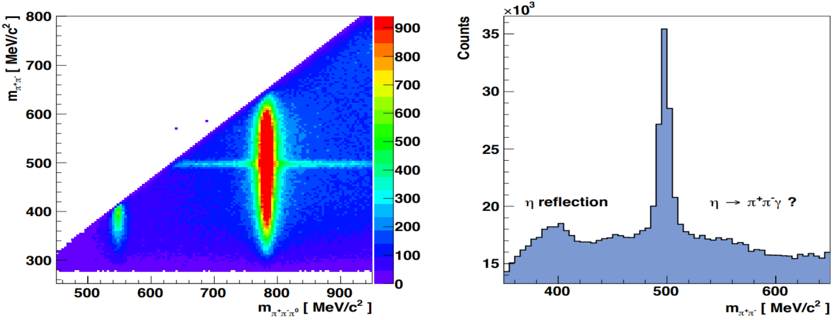

Figure 5: **Left:** Invariant $\pi^+\pi^-\pi^0$ mass vs. $\pi^+\pi^-$ mass for all g12 $\pi^+\pi^-\pi^0$ events. **Right:** Same invariant $\pi^+\pi^-$ mass distribution after the ω and Σ^+ mass-cuts.

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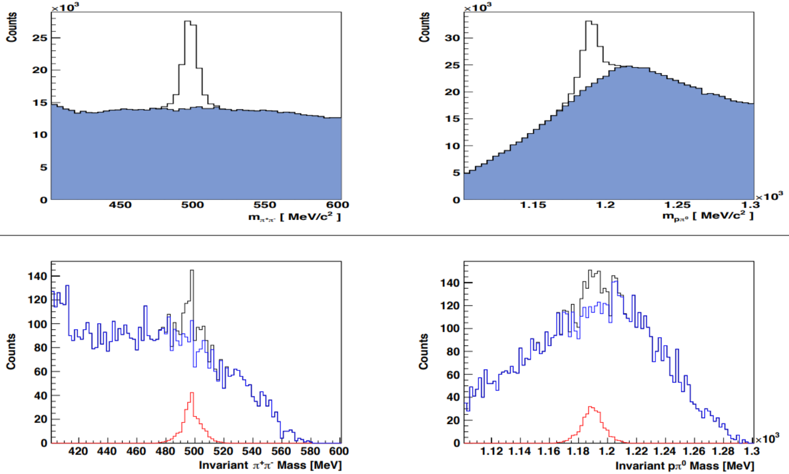

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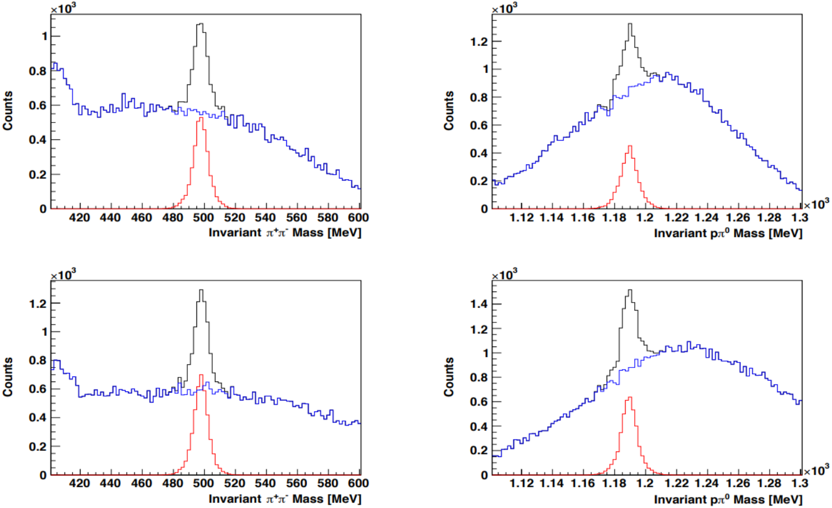

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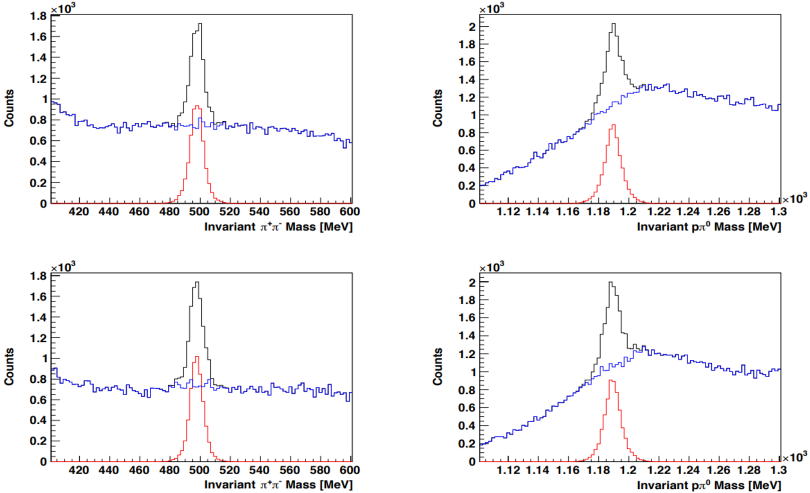

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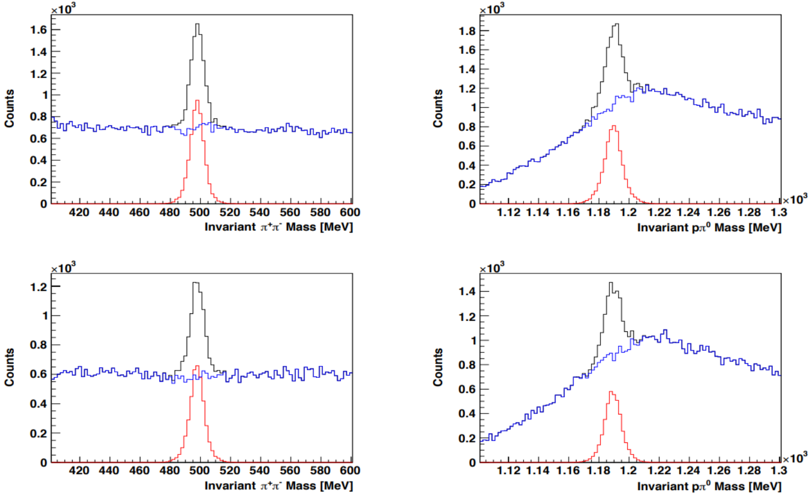

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- P observable does not require beam's polarization data; P can be seen as the "C_y" observable.

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- For the induced polarization, we integrate over all events above and below the reaction plane so that no direct angular dependence on the proton is needed for either observable C_x and C_z.

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- Results already approved and included within analysis notes.

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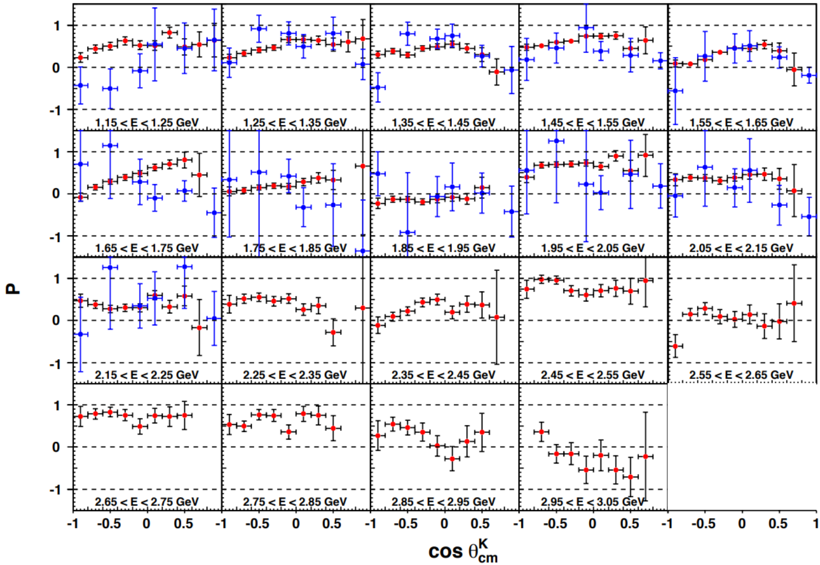

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  - **Maximum-likelihood fit:** Simultaneous extraction of all observables P, C_x and C_z.

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- Likewise one can also extract double-polarization observables by forming the following asymmetry

$$A(\cos\theta^p_x, \cos\theta^p_z) = \frac{N_+ - N_-}{N_+ + N_-} = \alpha P_{\odot} C_x \cos\theta^p_x + \alpha P_{\odot} C_z \cos\theta^p_z.$$

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Single Event Probability Distribution Function

$$\mathcal{P}(\cos\theta_x^p, \cos\theta_z^p, \cos\theta_y^p | C_x, C_z, P) =$$
$$1 \pm P_{\odot}\alpha(C_x \cos\theta_x^p + C_z \cos\theta_z^p) + \alpha P \cos\theta_y^p$$

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Probability Distribution Function

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Outline

Summary and Future Analysis

- The photoproduction of $K^0\Sigma^+$ is an unexplored and important channel to study due to its parity violation, allowing for the construction of an asymmetry, thus allowing for measurement of polarization observables.

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- Polarization observables allow for the disentanglement of N^* and Δ^* resonances.
- Determination of observables leads an understanding of the intermediate steps involved in the reaction of interest.
- Double polarization observables C_x and C_z have never measured for the $K^0\Sigma^+$ channel. Future analysis will carry out measurement of these observables using methods highlighted in this presentation.

Thank you so much for your time!