# First look at KaonLT experiment data

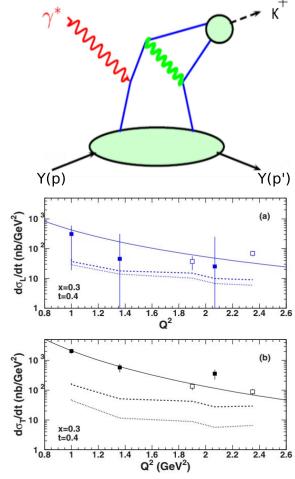
Richard Trotta, Tanja Horn, Garth Huber, Pete Markowitz,
Stephen Kay, Vijay Kumar, Vladimir Berdnikov, Mireille Muhoza,
Nathan Heinrich,
and the KaonLT collaboration





## L/T separated data for verifying reaction mechanism

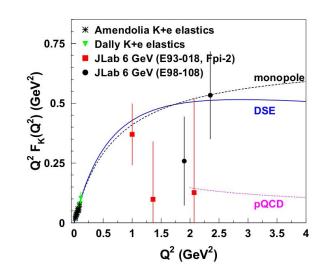
- Jlab 6 GeV data demonstrated the technique of measuring the Q<sup>2</sup> dependence of L/T separated cross sections at fixed x/t to test QCD Factorization
  - $\circ$  Consistent with expected scaling of  $\sigma_{\rm L}$  to leading order Q<sup>-6</sup> but with relatively large uncertainties
- Separated cross sections over a large range in Q<sup>2</sup> are essential for:
  - Testing factorization and understanding dynamical effects in both Q<sup>2</sup> and –t kinematics
  - Interpreting non-perturbative contributions in experimentally accessible kinematics



M. Carmignotto et al., PhysRevC 97(2018)025204

#### **Meson Form Factors**

- Pion and kaon form factors are of special interest in hadron structure studies
  - Pion lightest QCD quark system and crucial in understanding dynamic generation of mass
  - Kaon next simplest system containing strangeness
- Clearest case for studying transition from non-perturbative to perturbative regions
- Jlab 6 GeV data showed FF differs from hard QCD calculation
  - Evaluated with asymptotic valence-quark Distribution Amplitude (DA), but large uncertainties
- 12 GeV FF extraction data require:
  - measurements over a range of t, which allow for interpretation of kaon pole contribution



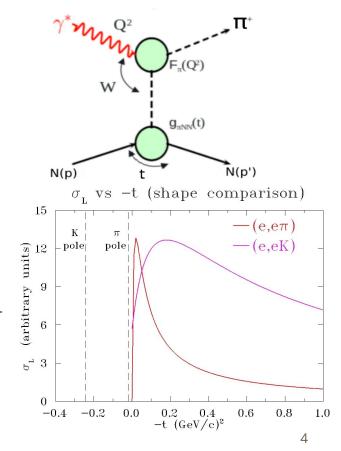
M. Carmignotto et al., PhysRevC **97**(2018)025204 F. Gao et al., Phys. Rev. D 96 (2017) no. 3, 034024

#### Experimental Determination of the $\pi/K+$ Form Factor

- At larger Q<sup>2</sup>,  $F_{\pi^+}^{2}$  must be measured indirectly using the "pion cloud" of the proton via the p(e,e' $\pi^+$ )n process
  - At small –t, the pion pole process dominates  $\sigma_{\rm p}$
  - o In the Born term model,  $F_{n+}^{2}$  appears as

$$\frac{d\sigma_L}{dt} \propto \frac{-t}{(t-m_\pi^2)} g_{\pi NN}^2(t) Q^2 F_\pi^2(Q^2, t)$$

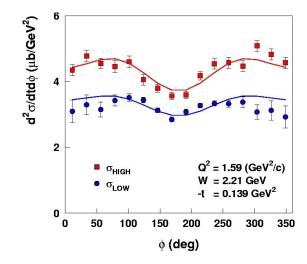
- Requirements:
  - $\circ$  Full L/T separation of the cross section isolation of  $\sigma_{i}$
  - Selection of the pion pole process
  - Extraction of the form factor using a model
  - Validation of the technique model dependent checks



#### L/T Separation Example

$$2\pi \frac{d^2\sigma}{dtd\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

 $\sigma_L$  will give us  $F^2_{K_1}$ 



T. Horn et al., PhysRevC 97(2006)192001

- σ<sub>L</sub> is isolated using the Rosenbluth separation technique
- Measure the cross section at two beam energies and fixed W, Q<sup>2</sup>, -t
- Simultaneous fit using measured azimuthal angle (φ) allows for extracting L, T, LT, and TT
  - Careful evaluation of the systematic uncertainties is important due to the 1/ε amplification in the σ<sub>L</sub> extraction
- Must have magnetic spectrometers for such precision cross section measurements
  - This is only possible in Hall C at JLab

#### L/T Separation Example

$$2\pi \frac{d^2\sigma}{dtd\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos2\phi$$

$$\frac{\cos\phi}{\cos\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos2\phi$$

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$$\frac{\cos\phi}{\sin\phi} = \varepsilon \frac{d\sigma_{L}}{dt} + \frac{d\sigma_{L}}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{LT}}{dt} \cos\phi$$

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$$\frac{\cos\phi}{\sin\phi} = \varepsilon \frac{d\sigma_{L}}{dt} + \frac{d\sigma_{L}}{dt} + \frac{d\sigma_{L}}{dt} \cos\phi$$

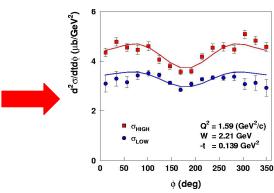
$$\frac{\cos\phi}{\sin\phi} = \varepsilon \frac{d\sigma_{L}}{dt} + \frac{d\sigma_{L}}{dt} + \frac{d\sigma_{L}}{dt} \cos\phi$$

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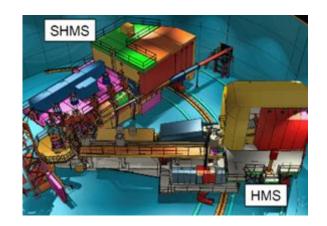
- Three SHMS angles for azimuthal (φ) coverage to determine the interference terms (LT, TT)
- Using the two beam energies (ε) to separate longitudinal (L) from transverse (T) cross section



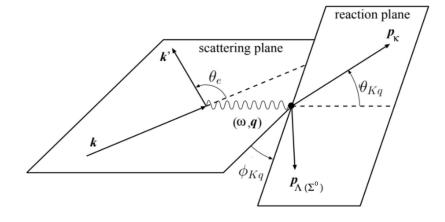
Fit using measured  $\epsilon$  and  $\varphi$  dependence

## Review E12-09-011 (KaonLT) Goals

- Q<sup>2</sup> dependence will allow studying the scaling behavior of the separated cross sections
  - First cross section data for Q<sup>2</sup> scaling tests with kaons
  - Highest Q<sup>2</sup> for L/T separated kaon electroproduction cross section
  - First separated kaon cross section measurement above W=2.2 GeV

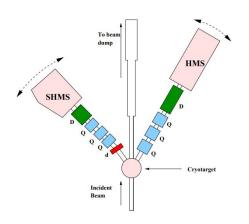


- t-dependence allows for detailed studies of the reaction mechanism
  - Contributes to understanding of the non-pole contributions, which should reduce the model dependence
  - Bonus: if warranted by data, extract the kaon form factor



## **Kaon LT - Data Collected**

 The p(e, e'K<sup>+</sup>)Λ,Σ<sup>0</sup> experiment ran in Hall C at Jefferson Lab over the fall and spring.



E	Q <sup>2</sup>	W	х	$\varepsilon_{\rm high}^{}/\varepsilon_{\rm low}^{}$
(GeV)	(GeV <sup>2</sup> )	(GeV)		
10.6/8.2	5.5	3.02	0.40	0.53/0.18
10.6/8.2	4.4	2.74	0.40	0.72/0.48
10.6/8.2	3.0	3.14	0.25	0.67/0.39
10.6/6.2	3.0	2.32	0.40	0.88/0.57
10.6/6.2	2.115	2.95	0.21	0.79/0.25
4.9/3.8	0.5	2.40	0.09	0.70/0.45

#### **Experimental Details**

Hall C: k<sub>e</sub>=3.8, 4.9, 6.4, 8.5, 10.6 GeV

#### SHMS for kaon detection :

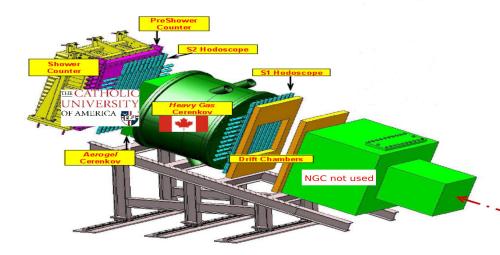
- o angles, 6 30 deg
- momenta, 2.7 6.8 GeV/c

#### HMS for electron detection :

- angles,10.7 31.7 deg
- o momenta, 0.86 5.1 GeV/c

#### Particle identification:

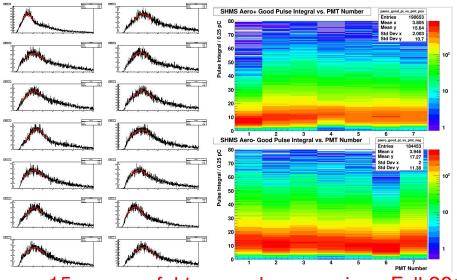
- Dedicated Aerogel Cherenkov detector for kaon/proton separation
  - Four refractive indices to cover the dynamic range required by experiments
- Heavy gas Cherenkov detector for kaon/pion separation

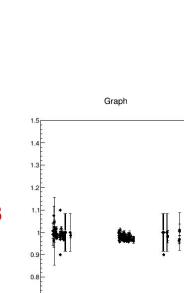


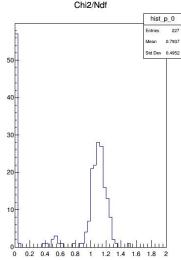
n	π <sub>thr</sub> (GeV/c)	K <sub>thr</sub> (GeV/c)	P <sub>thr</sub> (GeV/c)
1.030	0.57	2.00	3.80
1.020	0.67	2.46	4.67
1.015	0.81	2.84	5.40
1.011	0.94	3.32	6.31

#### **Aerogel Cherenkov detector in SHMS**









~15 successful tray exchanges since Fall 2018

- Aerogel performance as expected
- Trays require some optimization before next use - prevent damage from crane operation

Analysis by V. Berdnikov

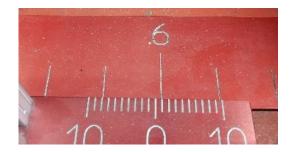
4500 5000 5500 6000 6500 7000 7500 8000 8500

1(

#### SHMS small angle operation

- Some issues with opening and small angle settings at beginning of run
  - o SHMS at 6.01°
  - HMS at 12.7°

[12/17/18]









#### **KaonLT Event Selection**

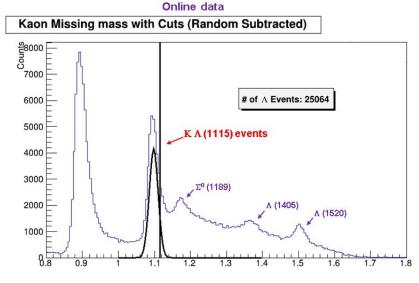
Isolate Exclusive Final States through missing mass

$$M_x = \sqrt{(E_{det} - E_{init})^2 - (p_{det} - p_{init})^2}$$

- Coincidence measurement between kaons in SHMS and electrons in HMS
  - simultaneous studies of KΛ and KΣ<sup>0</sup> channels...and a few others...
- Kaon pole dominance tests through

$$\frac{\sigma_L(\gamma^* p \to K^+ \Sigma^0)}{\sigma_L(\gamma^* p \to K^+ \Lambda)}$$

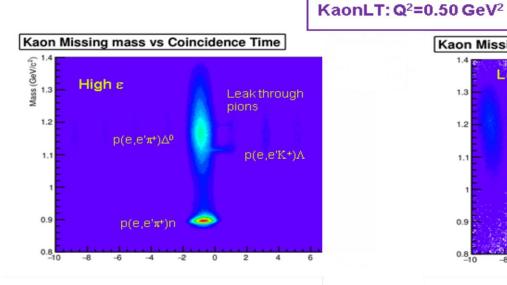
• Should be similar to ratio of coupling constants  $g^2_{pK\Sigma}/g^2_{pK\Lambda}$  in t-channel

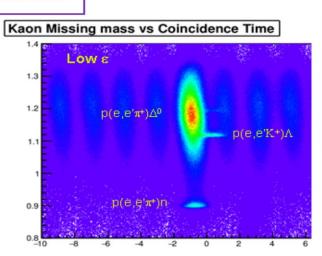


Plot by R. Ambrose, S. Kay, R. Trotta

#### Interesting Physics in the other channels

• Large difference in L/T ratio between p(e.e' $\pi$ )n and p(e,e' $\pi$ ) $\Delta$ 0 final states – G. Huber hclog #3640187



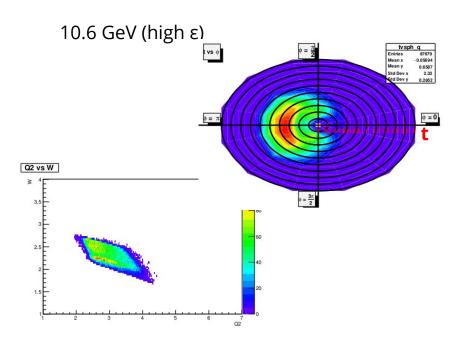


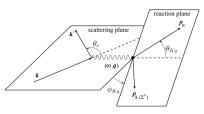
Plots by R. Ambrose, S. Kay, R. Trotta

- Large increase in neutron missing mass peak at high epsilon is evidence of the pion-pole process at low  $Q^2$  and small –t, which suggests  $\sigma_L >> \sigma_T$
- $\Delta^0$  exclusive longitudinal cross section expected to be at best  $\sigma_L \sim \sigma_T$

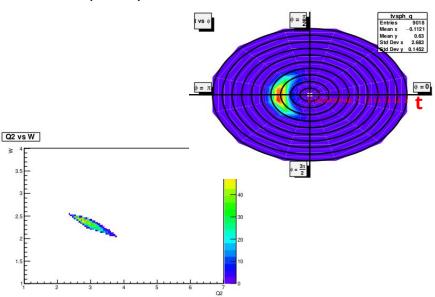
## Comparison of high and low $\varepsilon$ [Q<sup>2</sup>=3.0, W=2.32, x=0.40]

- [10.6 Gev (high  $\varepsilon$ ), 6.2 Gev (low  $\varepsilon$ )]
- Left  $(\theta_{high} = 21.18, \theta_{low} = 16.28)$





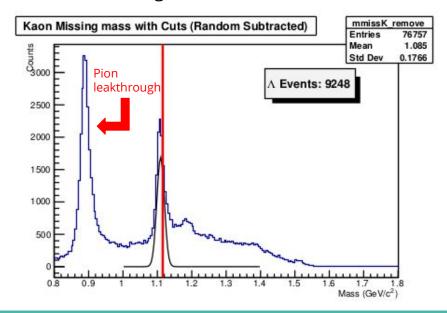
6.2 GeV (low ε)



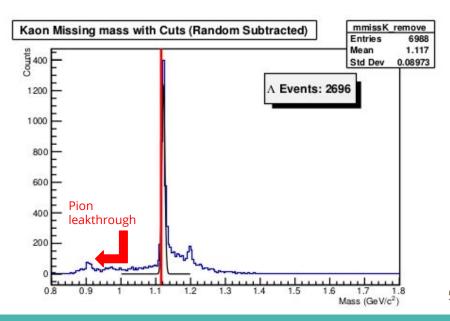
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10.6 GeV (high ε)



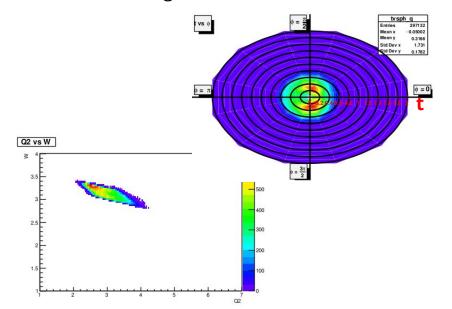
6.2 GeV (low ε)



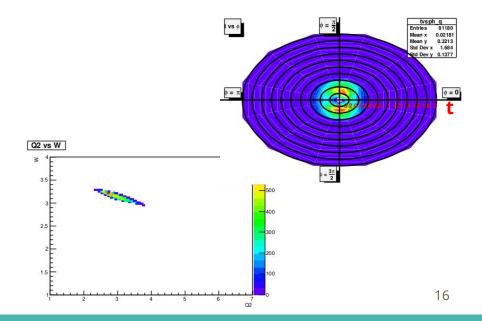
## Comparison of high and low $\varepsilon$ [Q<sup>2</sup>=3.0, W=3.14, x=0.25]

- [10.6 Gev (high ε), 8.2 Gev (low ε)]
- Center ( $\theta_{high}$ =9.42, $\theta_{low}$ =6.89)

10.6 GeV (high ε)



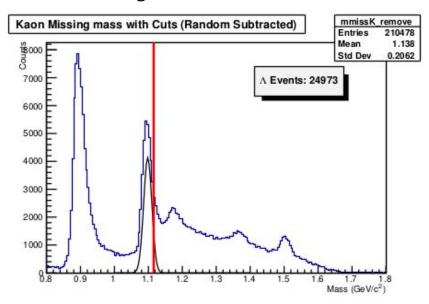
8.2 GeV (low ε)



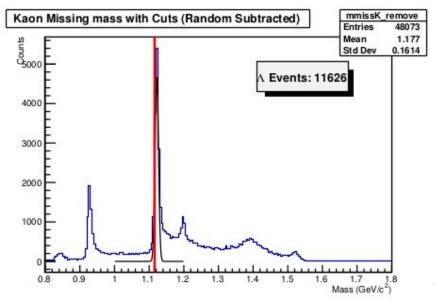
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10.6 GeV (high ε)



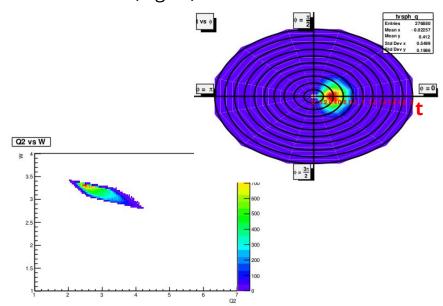
8.2 GeV (low  $\varepsilon$ )



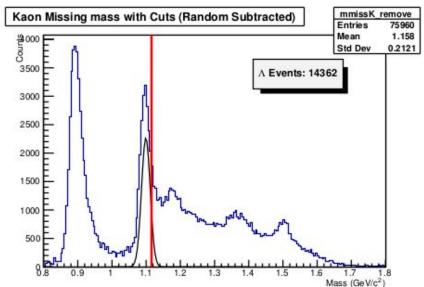
## Comparison of high and low $\varepsilon$ [Q<sup>2</sup>=3.0, W=3.14, x=0.25]

- [10.6 Gev (high ε)]
- Right ( $\theta_{high}$ =6.65)

10.6 GeV (high ε)



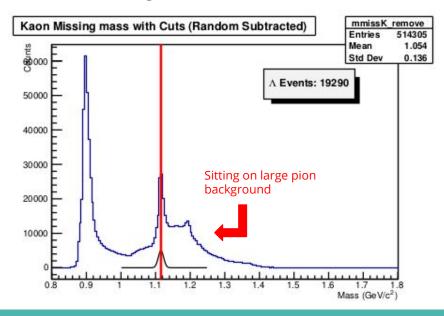




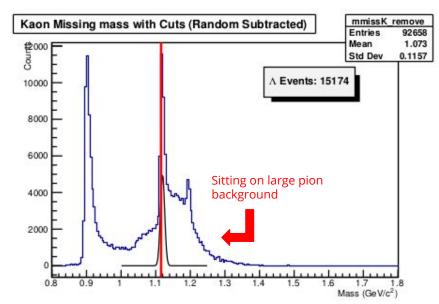
## Comparison of high and low $\varepsilon$ [Q<sup>2</sup>=0.5, W=2.40, x=0.09]

- [4.9 Gev (high ε), 3.8 Gev (low ε)]
- Center  $(\theta_{high} = 8.86, \theta_{low} = 6.79)$

#### 4.9 GeV (high ε)



#### 3.8 GeV (low $\varepsilon$ )



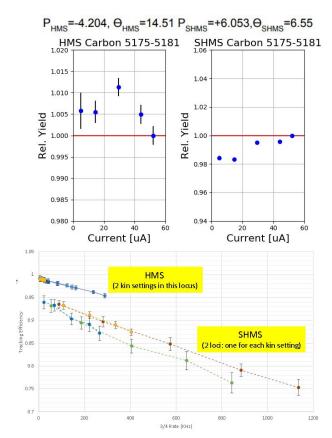
#### **Analysis Phases**

#### **Current Phase**

- 1. Calibrations
- Calorimeter, aerogel, HG cer, HMS cer, DC, Quartz plan of hodo
- Assure we are replaying to optimize our physics settings
- Efficiencies and offsets
  - Luminosity and elastics
- First iteration of cross section
  - Bring everything together
- 4. Fine tune
  - Fine tune values to minimize systematics
- 5. Repeat previous step
  - Repeat until acceptable cross sections are reached
- 6. Possible attempt at form factor extraction
  - Fit the data to a model and iterate

### Phase 1: Early luminosity analysis (pre-calibrations)

- Understanding efficiencies from luminosity scans has been ongoing with only one run period having been looked at
- Relative yield has been reduced to ~2% spread for carbon target
- Tracking efficiencies are a big contributor
  - At a given ¾ rate, HMS tracking efficiency is
     ~4% higher than that of the SHMS
  - HMS tracking efficiency is mostly independent of kinematic setting – not the case for the SHMS
  - SHMS tracking efficiency extrapolates to ~95% at 0 KHz – hadron tracking efficiency low by 4-6%



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Phase 1: Calibration of SHMS HGC detector and hodoscope time walking

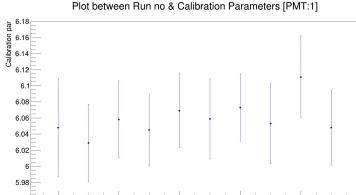
Plot between Run no & Calibration

#### Heavy Gas Cherenkov

 Run dependence of calibration parameters for PMT1 to check the consistency of calibration.

#### Hodoscope timewalking

Plot from PMT 2+ on S1X plane, similar for others



8092

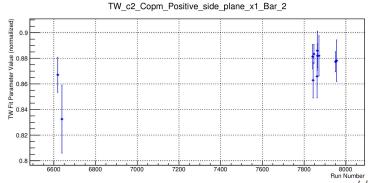
8094

8096

Run no

8088

8090



Analysis by V. Kumar, N. Heinrich and M. Muhoza

#### Conclusion

- Kaon can provide an interesting way to expand previous data of charged pion form factor with access to the mechanism involving strangeness
- E12-09-011 has completed its 2018-19 run
- Potential to extract the Kaon form factor from the L/T separated cross sections to the highest Q<sup>2</sup> achievable at Jlab
  - Full azimuthal coverage, good phase space matching and favorable rates to allow Kaon cross section separation
- Provide much needed data for Q<sup>2</sup> scaling at fixed x and -t in Kaon electroproduction to validate QCD factorization for hadron imaging studies
- Currently in the first phase of analysis

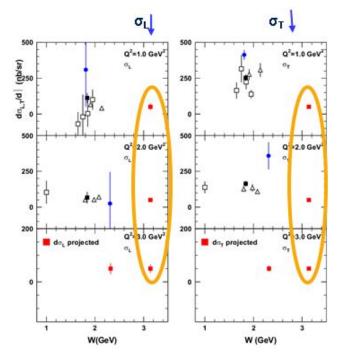
## **Extra Slides**

#### **KaonLT Sample Projections**

 E12-09-011: Separated L/T/LT/TT cross section over a wide range of Q<sup>2</sup> and t

E12-09-011 spokespersons: T. Horn, G. Huber, P. Markowitz

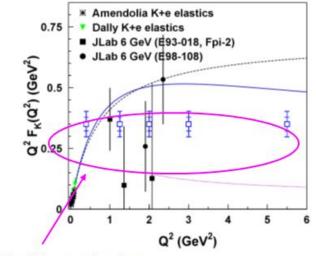
- JLab 12 GeV Kaon Program features:
  - First cross section data for Q<sup>2</sup> scaling tests with kaons
  - Highest Q<sup>2</sup> for L/T separated kaon electroproduction cross section
  - First separated kaon cross section measurement above W=2.2 GeV



blue points from M. Carmignotto, PhD thesis (2017)

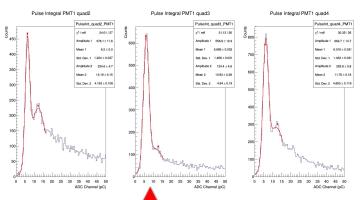
## KaonLT: Projections for $F_{K+}(Q^2)$ Measurements

- E12-09-011: primary goal L/T separated kaon cross sections to investigate hard-soft factorization and non-pole contributions
- Possible K<sup>+</sup> form factor extraction to highest possible Q<sup>2</sup> achievable at JLab
  - Extraction like in the pion case by studying the model dependence at small t
  - Comparative extractions of  $F_{\pi}^2$  at small and larger t show only modest model dependence
    - larger t data lie at a similar distance from pole as kaon data

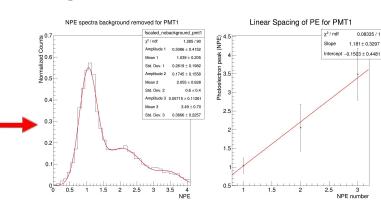


Possible extractions from 2018/19 run

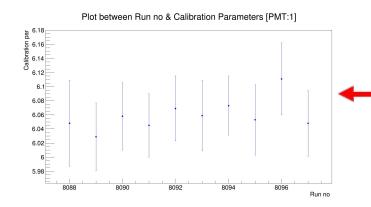
#### P1: Calibration of HGC Detector (SHMS)



To see the second & third photo-electron, we fitted the scaled histogram with Poisson function and subtracted the higher photoelectron.



Showing the SPE in HGC for PMT1 FADC and fit it with a Gaussian function to get the mean of peaks.



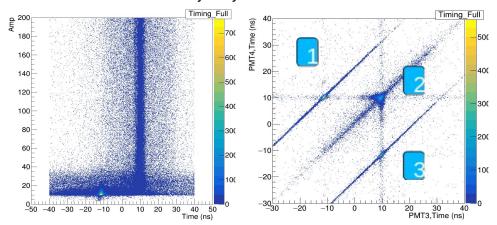
Run dependence of calibration parameters for the PMT1 to check the consistency of calibration.

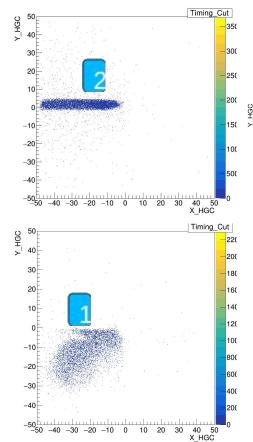
1.181 ± 0.3297

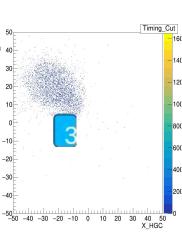
NPE number

## **HGC Timing Study**

- In addition to main timing peak at +10ns, there is an unexpected second peak at -10ns.
- To better understand the origin of the unexpected peak, plot b/w Timing vs Amplitude.
  - o 2nd peak corresponds to small pulses only.
- We also checked the tracking position in focal plane coordinates.
  - Interesting correlation between hit position and timing remains a mystery.







#### P1: SHMS Hodoscope Time Walk Calibration

#### In order to correct for time walk we:

- Plot ADC amplitude against TDC ADC time
- Fit This Function:  $f_{TW} = c_1 + \frac{1}{\left(\frac{a}{TDC_{Thrs.}}\right)^{c_2}}$
- Subtract second term
- Check parameter stability over run periods 6600 - 8000, stable within error
- Plots from PMT 2+ on 1x plane, similar for others

