Luminosity reduction due to phase modulations at the HL-LHC crab cavities

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Introduction

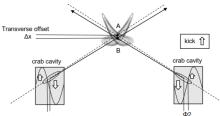
- Crab cavities are installed around IP1 & IP5 to compensate geometric reduction in luminosity due to the beam crossing angle at the IP.
- There will be two crab cavities per beam on either side of IP1 and IP5. Sixteen cavities in total.

Table 1: List of parameters for the baseline HL-LHC (HL-LHC V1.2 [1])

Proton energy at collision	7 TeV
Number of protons in a bunch N	$2.2 \times 10^{11} \text{ ppb}$
Number of bunch n_b	2748
r.m.s. bunch length $(4\sigma_z)$	1.2 ns
Longitudinal emittance	2.5 eVs
Transverse normalized emittance $\epsilon_{n(x,y)}$ (r.m.s)	2.5 µm
Full crossing angle	510 µrad
Nominal crab cavity voltage	3.4 MV/cavity
Crab cavity RF frequency	400.79 MHz

Introduction

- Existing klystron power is insufficient to run the HL-LHC with the current LLRF control algorithm.
- New algorithm (Full-Detuning) avoids this limitation by accepting a phase modulation of accelerating RF cavity, which is operational since 2017 [2].
- Crab cavity (C.C) cannot follow the phase modulations (up to 100 ps pk-pk), resulting in a phase error w.r.t the individual bunch centroid.
- Phase errors at C.C creates significant sinusoidal bunch distortions at the IP.



Investigate the impact of phase modulations on the peak luminosity

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Analytical solution of bunch distributions

- **1** Define Gaussian bunch distribution at C.C: $\rho_{c.c}(x_{c.c}, y_{c.c}, z_{c.c})$
 - A thin crab kick: $\delta x'_{c.c} = \frac{eV_1}{E_s} \sin(kz_{c.c} + \phi)$ $(\delta x_{c.c} = 0)$
- 2 Transport the bunch from C.C to IP: $\rho_{IP}(x_{IP}, y_{IP}, z_{IP})$
 - $\pi/2$ phase advance between C.C to IP
 - $\alpha_{IP} = 0$
- **3** Rotate coordinate system at IP by $\theta/2$: $\rho_{IP}(\tilde{x}_{IP}, \tilde{y}_{IP}, \tilde{z}_{IP})$

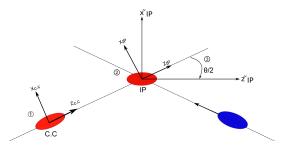


Figure 1: Coordinate system around the IP.

Details are presented in [3].

Bunch distributions without phase modulations at the IP

Table 2: Twiss parameters of baseline optics in the HL-LHC (HL-LHC V1.2 [1]).

$\beta_{x,y}^*$ at IP5	0.20, 0.20 m	
$\beta_{x,y}$ at C.C before IP5		
$\alpha_{x,y}^*$ at IP5	0,0	
$\alpha_{x,y}$ at C.C before IP5	-14.0, -36.7	

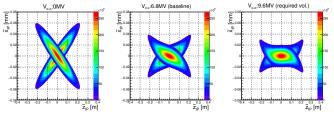


Figure 2: Bunch distributions at IP with crab voltages at 0 MV (without crab cavity), 6.8 MV (baseline) and 9.6 MV (required for full correction) without phase modulations.

- Voltages quoted are for two CCs (total voltage on one IP side).
- The baseline gives partial compensation of the crossing angle. See [4].

Bunch distributions at IP with coherent phase modulation

The Full-Detuning phase modulation depends on the filling pattern.
 With same filling scheme in both rings, the phase modulation of colliding bunches are identical in IP1 and IP5

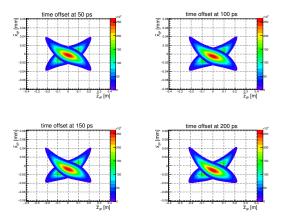


Figure 3: Bunch distributions at IP with crab voltages at 6.8 MV with phase modulations.

Luminosity calculation with phase modulations

Translating bunches from the IP along the beam line (z_{IP}) : $z_{IP} \pm ct$

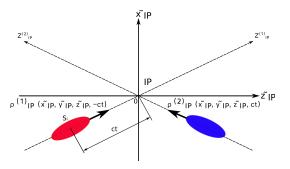


Figure 4: Schematic view of colliding pairs around the IP.

Integrated peak luminosity:

$$L=2\cdot\cos^{2}\frac{\theta}{2}N_{1}N_{2}f_{rev}n_{b}\int\int\int\int_{-\infty}^{\infty}\rho_{IP}^{(1)}\left(\tilde{x}_{IP},\tilde{y}_{IP},\tilde{z}_{IP},-ct\right)\cdot\rho_{IP}^{(2)}\left(\tilde{x}_{IP},\tilde{y}_{IP},\tilde{z}_{IP},ct\right)d\tilde{x}_{IP}d\tilde{y}_{IP}d\tilde{z}_{IP}d(ct).$$

$$(1)$$

 N_1 , N_2 : number of protons in bunches, n_b : number of bunches

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Linear transfer map from MADX

- Twiss parameters and first-order transfer map are created by MADX
- Transfer-map: IP1&IP5, 2 main accelerating RF cavities around IP4, one pair of C.C around IP1 and IP5
- HL-LHC V1.2 [1] (round optics $\beta_x^* = \beta_y^*$)

PYTRACK simulations

- Inject protons (10⁵) at IP1 and observe bunch distributions at IP5
- Transverse initial bunch distribution: Gaussian
- Longitudinal initial bunch profiles: Gaussian and q-Gaussian (n=2.5) are generated by BLonD code [5].
- q-Gaussian profile is a close match to the measured LHC longitudinal profile. The definition of q-Gaussian is in [4].
- Ramping up crab voltage over 1000 turns (two synchrotron periods, F_s =23Hz, F_{rev} =11kHz) linearly to keep quasi-static synchrotron motion

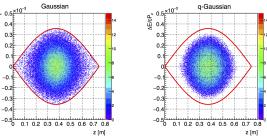


Figure 5: Longitudinal initial bunch distributions generated by BLonD.

Bunch distributions with coherent and incoherent phase modulations

- Coherent phase modulations: Identical filling pattern B1-B2
- Incoherent phase modulation can be caused by imperfect phase alignment of the crabbing and uncrabbing cavity pairs or by different fillings for the two beams (not planned).
- The incoherent phase modulation results in a larger degradation in luminosity because the cores of the two colliding bunches do not see the same kick, resulting in a transverse offset at the IP.

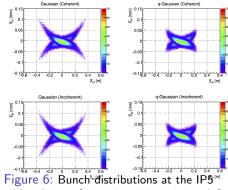


Figure 6: Bunch-distributions at the IP5 with coherent ($\phi_1=\phi_2=100$ ps=0.25rad @ 400.79 MHz) and incoherent ($\phi_2=100$ ps and $\phi_1=0$) phase modulations.

Comparisons of peak luminosity

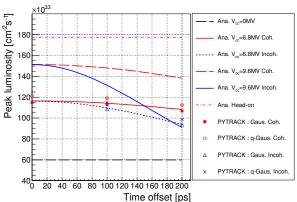


Figure 7: Peak luminosity with phase modulations in time. The peak luminosity is computed by numerical integration of Eq. 1 using Python.

- Coherent phase modulations at 100 ps: 2 % reduction (Analytical)
- Incoherent phase modulations at 100 ps: 6 % reduction (Analytical)

Summary

- Analytical model of peak luminosity with phase modulations has been derived. This study was presented at HL-LHC collaboration meeting in 2017 [6] and published at NIMA [3].
- Simulations have been applied with PYTRACK to compare to the analytical results, which is in good agreements.
- Reduction of peak luminosity is less than 2 % with the expected coherent phase modulation at 100 ps for both beams. In the limiting scenario of an incoherent phase modulation at 100 ps, it sums to 6 % of reductions.
- Phase modulations for the nominal HL-LHC operations are acceptable for physics runs.

Back Up

Luminosity calculation with phase modulations

$$\begin{split} L = & \frac{\cos^2\frac{\theta}{2}N^2 f_{\text{rev}} n_b}{\pi^{5/2} \sigma_{x1}^* \sigma_{x2}^* \sigma_{z1} \sigma_{z2} \sqrt{8(\sigma_{y1}^{*2} + \sigma_{y2}^{*2})}} \int \int \int_{-\infty}^{\infty} e^{-\left(\frac{\left(\tilde{x}_{\text{IP}} \cos\frac{\theta}{2} - \tilde{z}_{\text{IP}} \sin\frac{\theta}{2}\right)^2 + \tilde{C}_1}{2\sigma_{x1}^{*2}} + \frac{\left(\tilde{x}_{\text{IP}} \cos\frac{\theta}{2} + \tilde{z}_{\text{IP}} \sin\frac{\theta}{2}\right)^2 + \tilde{C}_2}{2\sigma_{x2}^{*2}}\right)}{e^{-\left(\frac{\left(\tilde{x}_{\text{IP}} \sin\frac{\theta}{2} + \tilde{z}_{\text{IP}} \cos\frac{\theta}{2} - ct\right)^2}{2\sigma_{z1}^2} + \frac{\left(-\tilde{x}_{\text{IP}} \sin\frac{\theta}{2} + \tilde{z}_{\text{IP}} \cos\frac{\theta}{2} + ct\right)^2}{2\sigma_{z2}^2}\right)} d\tilde{x}_{\text{IP}} d\tilde{z}_{\text{IP}} d(ct). \end{split}$$

where \tilde{C}_1 and \tilde{C}_2 are

$$\begin{split} \tilde{C}_1 = & \left(\beta^* \beta_{c.c} \frac{eV_1}{E_s} \sin \left(k \left(\tilde{x}_{IP} \sin \frac{\theta}{2} + \tilde{z}_{IP} \cos \frac{\theta}{2} - ct \right) + \phi_1 \right) \\ + & 2 \sqrt{\beta^* \beta_{c.c}} \left(\tilde{x}_{IP} \cos \frac{\theta}{2} - \tilde{z}_{IP} \sin \frac{\theta}{2} \right) \right) \frac{eV_1}{E_s} \sin \left(k \left(\tilde{x}_{IP} \sin \frac{\theta}{2} + \tilde{z}_{IP} \cos \frac{\theta}{2} - ct \right) + \phi_1 \right) \end{split}$$

$$\begin{split} \tilde{C}_2 = & \left(\beta^* \beta_{c.c} \frac{eV_1}{E_s} \sin \left(k \left(-\tilde{x}_{\textit{IP}} \sin \frac{\theta}{2} + \tilde{z}_{\textit{IP}} \cos \frac{\theta}{2} + ct \right) - \phi_2 \right) \right. \\ & \left. - 2 \sqrt{\beta^* \beta_{c.c}} \left(\tilde{x}_{\textit{IP}} \cos \frac{\theta}{2} + \tilde{z}_{\textit{IP}} \sin \frac{\theta}{2} \right) \right) \frac{eV_1}{E_s} \sin \left(k \left(-\tilde{x}_{\textit{IP}} \sin \frac{\theta}{2} + \tilde{z}_{\textit{IP}} \cos \frac{\theta}{2} + ct \right) - \phi_2 \right) \end{split}$$

q-Gaussian distribution

q-Gaussian distribution [4]:

$$F(J) = F_0 \left(1 - \frac{J}{J_0} \right)^n, \tag{2}$$

where the normalized coefficient F_0 , the action $J=\epsilon/(2\pi)$ (longitudinal emittance ϵ) and the J_0 corresponds to the initial full longitudinal emittance (ϵ_0)

We define the bunch length as 4- σ which is equivalent to the full width at half maximum (FWHM)

For Further Reading



[1] R. Thomas and L. Medina,

"Parameter update for the nominal HL-LHC: Standard, BCMS, and 8b+4e''.

HL-LHC parameters V6.1.0, 2017



[2] T. Mastoridis, P. Baudrenghien and J. Molendijk,

"Cavity voltage phase modulation to reduce the high-luminosity Large Hadron Collider rf power requirements"

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[3] E. Yamakawa, R. Apsimon, P. Baudrenghien, R. Calaga, A.C. Dexter

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NIM-A, vol. 908, 2018.

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- [4] E. Metral, et at.,
 - "Update of the HL-LHC operational scenarios for proton operation", CERN-ACC-NOTE-2018-0002, 2018.
 - [5] H. Timko, H. Muller, J.A. Lasheen, D. Quartullo, "Benchmarking the beam longitudinal dynamics code BLonD", IPAC'16, 2016.
 - [6] E.Yamakawa, P. Baudrenghien, R. Calaga, R. Apsimon and A. C. Dexter,
 - "Crab cavity studies including full detuning & offsets due to long range beam-beam interactions",
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