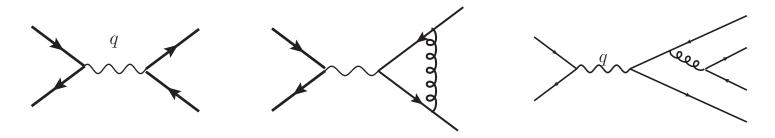
From the space-time structure of partonic scattering to parton-hadron duality

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- Aim: relation between coordinate-space happenings at partonic level, quantum mechanical interference, and resonance phenomena.
- A view of dynamical parton-hadron interface.
- For simplicity: $e^+e^- \rightarrow \text{partons} \rightarrow \text{hadrons}$.
- Rough quantitative analysis with naive string model.

Inspired by: Gupta-Quinn, Poggio-Quinn-Weinberg, JCC's QCD book.

Are pQCD predictions absolute first principles predictions?

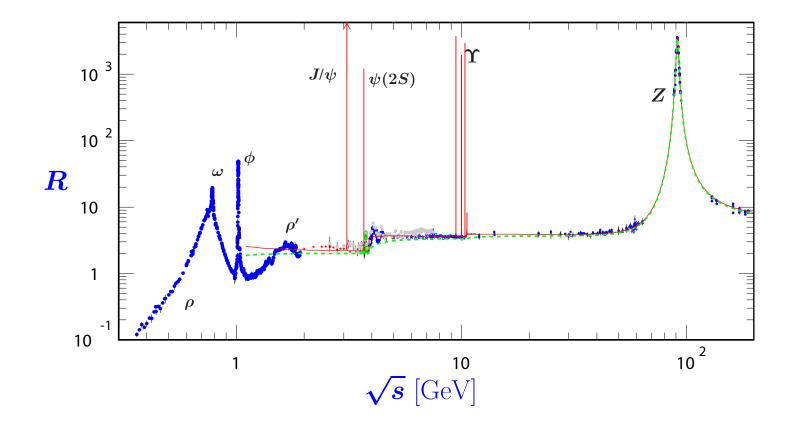


- Standard successful absolute pQCD predictions for total cross section, jet distributions, etc. (No fragmentation functions here.)
- (One parameter, e.g., $\alpha_s(m_Z)$.)
- Textbook scattering: On-shell quarks and gluons: at $t \to \infty$ in final state.
- Only hadrons in real final state, so we are really referring to *finite* t for partonic quasi-particles.
- There's something missing in justification of pQCD predictions.

Gupta and Quinn argument [PRD 25, 838 (1982)]

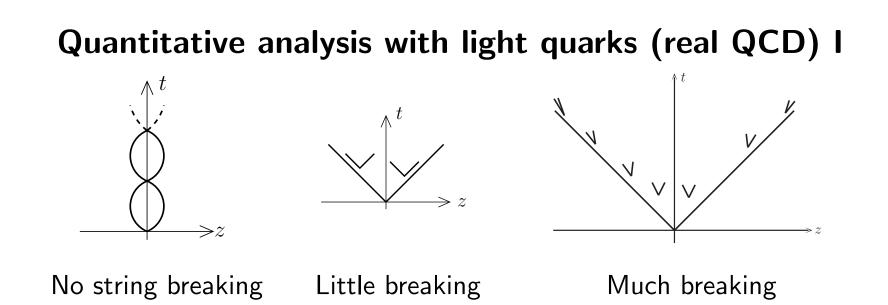
- Consider alternate world where all quarks are heavy ($\alpha_s(M_q) \ll 1$).
- Still have confining linear potential κr from gluons.
- $q\bar{q}$ pair production perturbatively small. Until Q is extremely large, string doesn't break.
- $\sigma_{had. tot.}(Q)$ is series of long-lived resonances, probably with lots of glueballs.
- Jet calculations surely fail.
- Actual cross section are *not* equal to the perturbative results, even though we can take $M/Q \rightarrow 0$ in perturbative calculations.

 $R(e^+e^- \rightarrow \text{hadrons})$



$$[R(e^+e^- \to \text{hadrons}) = \sigma(e^+e^- \to \text{hadrons}) / \sigma(e^+e^- \to \mu^+\mu^-)]$$

Duality workshop, Sep. 24, 2018



- Use very elementary string model of Artru-Mennessier [NPB 70,93 (1974)] with parameters from data.
- String tension κ (energy/unit length); pair production rate \mathcal{P} per unit t, z:

$$\kappa = 0.8 \,\mathrm{Gev/fm}, \qquad \mathcal{P} = 1.9 \,\mathrm{fm}^{-2}$$

- Analyze ranges of Q
 - Where is there an identifiable short-distance scattering?
 - Where does string breaking dominate over bounce back?

• String tension κ (energy/unit length); pair production rate \mathcal{P} per unit t, z:

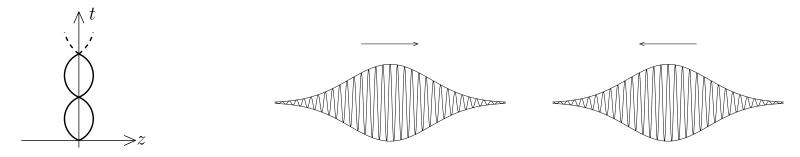
$$\kappa = 0.8 \,\mathrm{Gev}/\mathrm{fm}, \qquad \mathcal{P} = 1.9 \,\mathrm{fm}^{-2}$$

• Hence:

Q (GeV)	$z_{\sf max}$ (fm)	$t_{\rm bounce}~({\rm fm}/c)$	$\langle pairs angle$ in triangle
1	0.6	1.2	0.7
2	1.2	2.4	2.7
10	6	12	68 (?!?)
10	60	120	

• Once string has broken somewhere, later nearby breaking doesn't occur. (Stay inside space-like region $z^2 - t^2 \lesssim 1 \, {\rm fm}^2$ w.r.t. hard scattering.)

Derivation of validity of pQCD predication for smoothed total cross section I



• To restrict interaction to region much less than "bounce-back" time but reasonably much larger than 1/Q, use wave packet in-state:

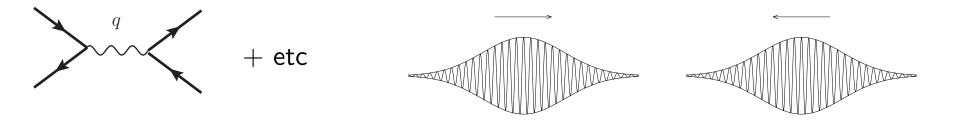
$$|\psi\rangle = \sum_{\mathbf{l}_1, \mathbf{l}_2, \lambda_1, \lambda_2} \left| e^+(\mathbf{l}_1, \lambda_1), \, e^-(\mathbf{l}_2, \lambda_2); \, \operatorname{in} \right\rangle \, \psi_1(\mathbf{l}_1, \lambda_1) \, \psi_2(\mathbf{l}_2, \lambda_2)$$

with

$$1/Q \ll \Delta z \ll t_{\rm bounce}, \qquad 1/t_{\rm bounce} \ll \Delta Q \ll Q.$$

• Aim: relate to experimental measurements with precise beam energy.

Derivation of validity of pQCD predication for smoothed total cross section II



• Hadronic part of final state:

$$|\phi; \text{hadron, out}\rangle = \int d^4x \, j^{\mu}(x) \, |0\rangle \, e^{-iq \cdot x} \phi_{\mu}(x)$$

 $\phi_{\mu}(x)$ is from: lepton wave functions, vertex to photon, photon propagator.

- Oscillation factor $e^{-iq \cdot x}$ extracted $\implies \phi_{\mu}(x)$ is envelope of a kind of wave packet.
- ... (see my QCD book)
- Then $R = \sigma(had.) / \sigma(\mu^+ \mu^-)$ with wave-packet beams is

$$\bar{R}(Q^2, \Delta^2) = \int \mathrm{d}s \ F(s - Q^2, \Delta^2) \, R_{\mathsf{std.}}(s)$$

Derivation of validity of pQCD predication for smoothed total cross section III q + etc

• Then $R = \sigma(had.)/\sigma(\mu^+\mu^-)$ with wave-packet beams is

$$\bar{R}(Q^2, \Delta^2) = \int \mathrm{d}s \ F(s - Q^2, \Delta^2) \ R_{\mathsf{std.}}(s)$$

- Δ is momentum width <u>of wave function</u>;
- Q is central value of $\sqrt{(l_1 + l_2)^2}$ for wave function;
- $F(s Q^2, \Delta^2)$ obtained from wave function, positive, peaked at $s = Q^2$.
- $\bar{R}(Q^2, \Delta^2)$ is local average of usual R, which is normalization times Fourier transform of

$$\langle 0|j^{\mu}(x)j^{\nu}(y)|0\rangle = \sum_{X} \langle X|j^{\mu}(x)|0\rangle^{*} \langle X|j^{\nu}(y)|0\rangle$$

QM interference between production of QCD stuff at different x and y.

• Wave function gives restriction to times of x - y less than bounce back.

Ranges of Q

- $Q < Q_0$ pQCD not applicable.
- $Q_0 < Q < Q_1$ Resonances, but parton-hadron duality.
- $Q_1 < Q$ String breaks, normal pQCD methods including jets, fragmentation functions, etc

Where should parton-hadron duality apply or not apply?

- Applies where we aren't sensitive to detailed final state:
 - $e^+e^- \rightarrow$ hadrons total cross section
 - DIS
- Where it may well not apply, or will need modification:
 - Jet distributions.
 - Hadron-inclusive cross sections. (Anything with a fragmentation function.)
- Not clear:
 - Drell-Yan,

Note: Carefully examined proofs of factorization: They often apply to locally averaged cross sections.

Overview

- Parton-hadron duality is at some level and in some cases a theorem.
- It applies when "bounce-back" from string tension doesn't lose to pair production rate in string, if there's such a range.
- String model gives quantitative link between microscopic parton propagation etc, and measured cross sections.
- Future: better mesh coordinate-space and semi-classical view with momentum space and fully quantum view.