

Duality in electron scattering: insights from theory

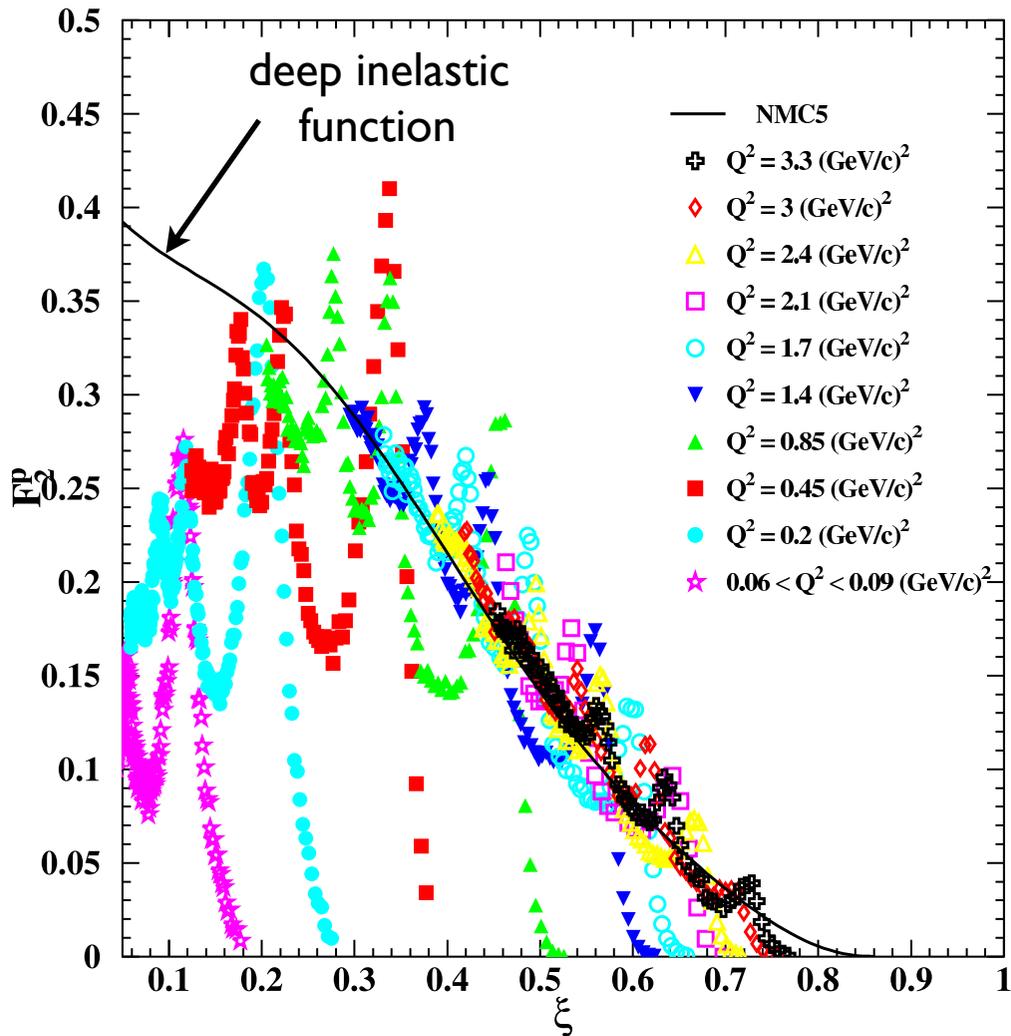
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Outline

- What we (think we) understand from theory
- Duality (global) in QCD
- Local duality
 - insights from models
- What we don't (yet) know ... what do we do next?

Duality in electron-proton scattering



Niculescu et al. (2000)

$$\xi = \frac{2x}{1 + \sqrt{1 + 4M^2x^2/Q^2}}$$

■ average over resonances
(strongly Q^2 dependent)

$\approx Q^2$ independent
scaling function

Duality and QCD

■ Operator product expansion in QCD

→ expand *moments* of structure functions in powers of $1/Q^2$

$$\begin{aligned} M_n(Q^2) &= \int_0^1 dx x^{n-2} F_2(x, Q^2) \\ &= A_n^{(2)} + \frac{A_n^{(4)}}{Q^2} + \frac{A_n^{(6)}}{Q^4} + \dots \end{aligned}$$

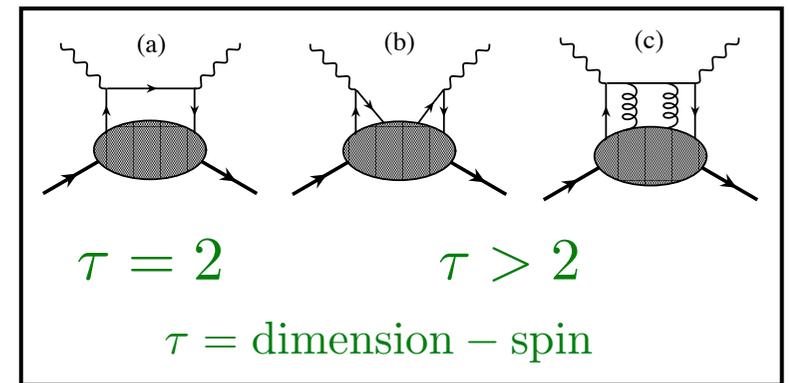


matrix elements of operators
with specific “twist” τ

e.g. $\langle N | \bar{\psi} \gamma^+ \psi | N \rangle$

$\langle N | \bar{\psi} \tilde{G}^{+\nu} \gamma_\nu \psi | N \rangle$

etc.



Duality and QCD

■ Operator product expansion in QCD

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■ If moment \approx independent of Q^2

→ “higher twist” terms $A_n^{(\tau>2)}$ small

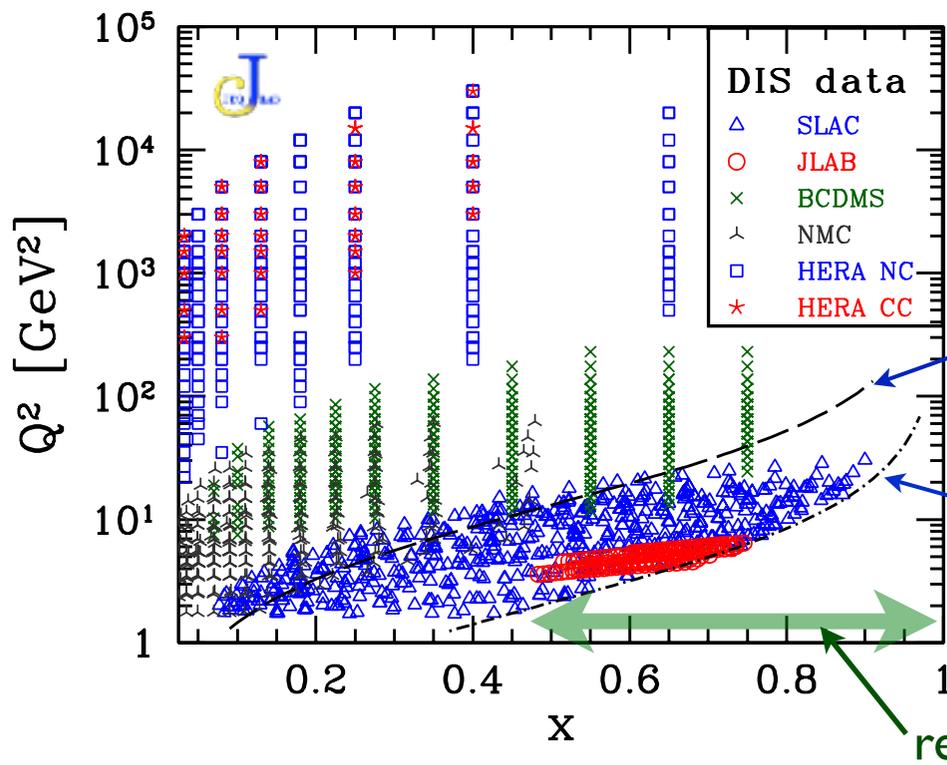
■ Duality \longleftrightarrow suppression of higher twists

Duality and QCD

- Note: at finite Q^2 , from kinematics *any* moment of *any* structure function (of *any* twist) must, by definition, include the resonance region

$$W^2 = M^2 + Q^2 \frac{(1-x)}{x} \quad \longrightarrow \quad x_{\text{res}} = \frac{Q^2}{W_{\text{res}}^2 - M^2 + Q^2}$$

$$W_{\text{res}} = 2 \text{ GeV} \quad \Longrightarrow \quad x_{\text{res}} \approx 0.24 \text{ at } Q^2 = 1 \text{ GeV}^2$$



$$W^2 > 12.25 \text{ GeV}^2$$

$$W^2 > 3 \text{ GeV}^2$$

resonances

Duality and QCD

- Note: at finite Q^2 , from kinematics *any* moment of *any* structure function (of *any* twist) must, by definition, include the resonance region
- Resonance and DIS regions are intimately connected
 - resonances an *integral* part of scaling structure function
 - e.g.* in large- N_c limit, spectrum of zero-width resonances is “maximally dual” to quark-level (smooth) structure function

Resonances & twists

- Total “higher twist” is *small* at scales $Q^2 \sim \mathcal{O}(1 \text{ GeV}^2)$
- On average, nonperturbative interactions between quarks and gluons not dominant (at these scales)
 - nontrivial interference between resonances?
- Can we understand this dynamically, at quark level?
- Can we use resonance region data to learn about *leading twist* structure functions (and *vice versa*)?
 - expanded data set has potentially significant implications for global quark distribution studies

Scaling functions from resonances

■ Earliest attempts predate QCD

→ e.g. harmonic oscillator spectrum $M_n^2 = (n + 1)\Lambda^2$
including states with spin = 1/2, ..., n+1/2

(n even: $I = 1/2$, n odd: $I = 3/2$)

Domokos et al. (1971)

→ at large Q^2 magnetic coupling dominates

$$G_n(Q^2) = \frac{\mu_n}{(1 + Q^2 r^2 / M_n^2)^2} \quad r^2 \approx 1.41$$

→ in Bjorken limit, $\sum_n \rightarrow \int dz$, $z \equiv M_n^2 / Q^2$

$$F_2 \sim (\omega' - 1)^{1/2} (\mu_{1/2}^2 + \mu_{3/2}^2) \int_0^\infty dz \frac{z^{3/2} (1 + r^2/z)^{-4}}{z + 1 - \omega' + \Gamma_0^2 z^2}$$

→ scaling function of $\omega' = \omega + M^2 / Q^2$ ($\omega = 1/x$)

Scaling functions from resonances

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(n even: $I = 1/2$, n odd: $I = 3/2$)

Domokos et al. (1971)

→ in $\Gamma_n \rightarrow 0$ limit

$$F_2 \sim (\mu_{1/2}^2 + \mu_{3/2}^2) \frac{(\omega' - 1)^3}{(\omega' - 1 + r^2)^4}$$

cf. Drell-Yan-West relation

$$G(Q^2) \sim \left(\frac{1}{Q^2}\right)^m \iff F_2(x) \sim (1 - x)^{2m-1}$$

→ similar behavior found in many models

Einhorn (1976) ('t Hooft model)

Greenberg (1993) (NR scalar quarks in HO potential)

Pace, Salme, Lev (1995) (relativistic HO with spin)

Isgur et al. (2001) (transition to scaling)

....

Scaling functions from resonances

■ Phenomenological analyses at finite Q^2

→ additional constraints from threshold behavior at $q \rightarrow 0$ and asymptotic behavior at $Q^2 \rightarrow \infty$

$$\left(1 + \frac{\nu^2}{Q^2}\right) F_2^R = M\nu \left[|G_+^R|^2 + 2|G_0^R|^2 + |G_-^R|^2 \right] \delta(W^2 - M_R^2)$$

Davidovsky, Struminsky (2003)

→ 21 isospin-1/2 & 3/2 resonances (with mass < 2 GeV)

$$|G_{\pm}^R(Q^2)|^2 = |G_{\pm}^R(0)|^2 \left(\frac{|\vec{q}|}{|\vec{q}|_0} \frac{\Lambda'^2}{Q^2 + \Lambda'^2} \right)^{\gamma_1} \left(\frac{\Lambda^2}{Q^2 + \Lambda^2} \right)^{m_{\pm}} \quad m_{+,0,-} = 3, 4, 5$$

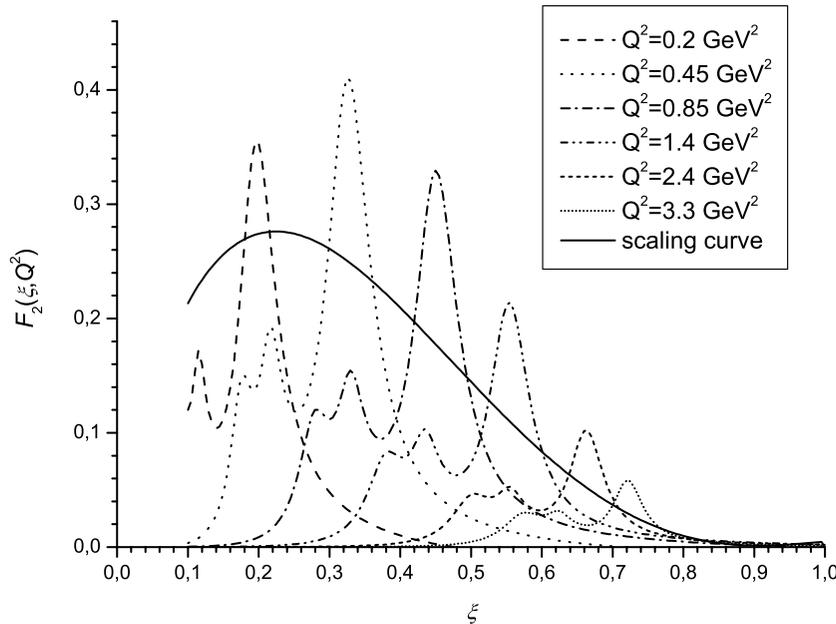
$$|G_0^R(Q^2)|^2 = C^2 \left(\frac{Q^2}{Q^2 + \Lambda''^2} \right)^{2a} \frac{q_0^2}{|\vec{q}|^2} \left(\frac{|\vec{q}|}{|\vec{q}|_0} \frac{\Lambda'^2}{Q^2 + \Lambda'^2} \right)^{\gamma_2} \left(\frac{\Lambda^2}{Q^2 + \Lambda^2} \right)^{m_0}$$

→ in $x \rightarrow 1$ limit,

$$F_2(x) \sim (1 - x)^{m_+}$$

Scaling functions from resonances

■ Phenomenological analyses at finite Q^2



21 isospin-1/2 & 3/2
resonances (mass < 2 GeV)

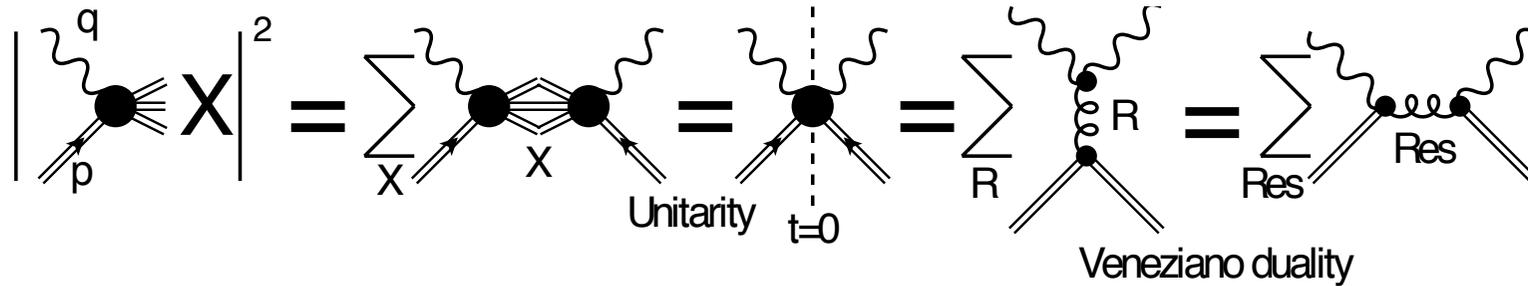
Davidovsky, Struminsky (2003)

→ valence-like structure of dual function suggests
“two-component duality”:

- valence (Reggeon exchange) dual to resonances $F_2^{(\text{val})} \sim x^{0.5}$
- sea (Pomeron exchange) dual to background $F_2^{(\text{sea})} \sim x^{-0.08}$

Scaling functions from resonances

■ Explicit realization of Veneziano & Bloom-Gilman duality



$$V(s, t) = \frac{\Gamma(1 - \alpha(s))\Gamma(1 - \alpha(t))}{\Gamma(2 - \alpha(s) - \alpha(t))}$$

$$\rightarrow s^{\alpha(t)} \quad \text{high } s, \text{ low } |t|$$

→ Veneziano model not unitary,
has no imaginary parts

→ generalization of narrow-resonance approximation,
with nonlinear, complex Regge trajectories

$$D(s, t) = \int_0^1 dz \left(\frac{z}{g}\right)^{-\alpha_s(s(1-z))-1} \left(\frac{1-z}{g}\right)^{-\alpha_t(tz)-1}$$

“dual amplitude with Mandelstam analyticity” (DAMA) model

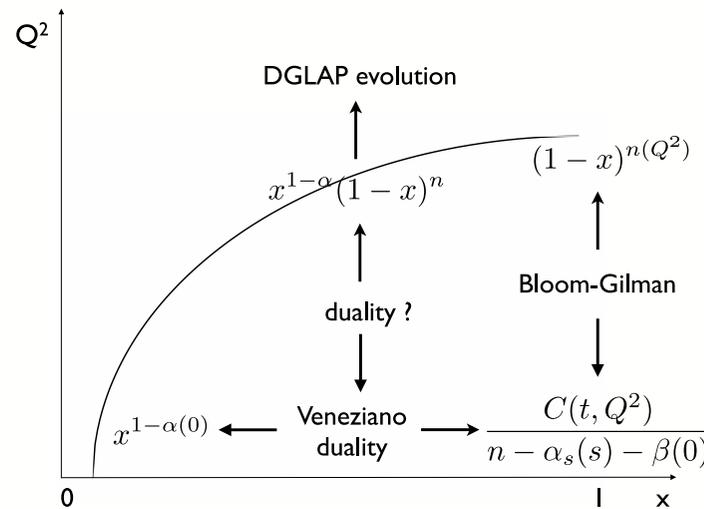
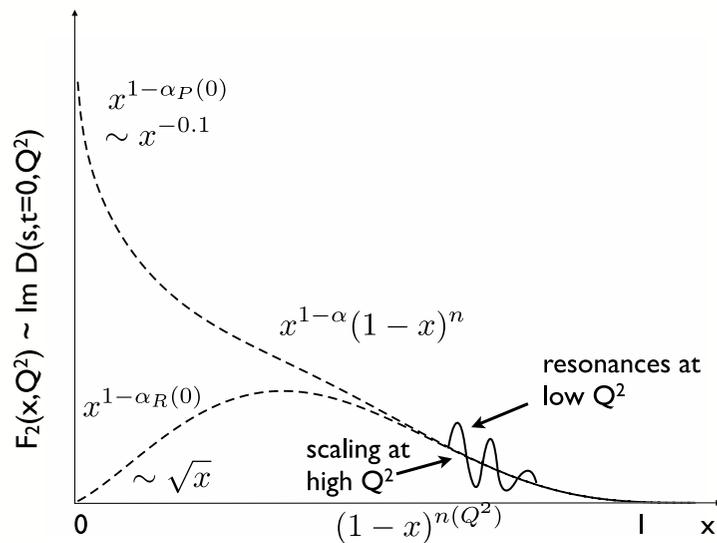
Scaling functions from resonances

■ Explicit realization of Veneziano & Bloom-Gilman duality

→ for large x and Q^2 , have power-law behavior

$$F_2 \sim (1-x)^{2\alpha_t(0) \ln 2g / \ln g}$$

where parameter g can be Q^2 dependent



Jenkowszky, Magas, Londergan, Szczepaniak (2012)

More than one flavor?

- Consider simple quark model with spin-flavor symmetric wave function

low energy

→ *coherent* scattering from quarks $d\sigma \sim \left(\sum_i e_i \right)^2$

high energy

→ *incoherent* scattering from quarks $d\sigma \sim \sum_i e_i^2$

- For duality to work, these must be equal

→ *how can square of a sum become sum of squares?*

■ Dynamical cancellations

→ *e.g.* for toy model of two quarks bound in a harmonic oscillator potential, structure function given by

$$F(\nu, \mathbf{q}^2) \sim \sum_n |G_{0,n}(\mathbf{q}^2)|^2 \delta(E_n - E_0 - \nu)$$

→ charge operator $\sum_i e_i \exp(i\mathbf{q} \cdot \mathbf{r}_i)$ excites
even partial waves with strength $\propto (e_1 + e_2)^2$
odd partial waves with strength $\propto (e_1 - e_2)^2$

→ resulting structure function

$$F(\nu, \mathbf{q}^2) \sim \sum_n \{ (e_1 + e_2)^2 G_{0,2n}^2 + (e_1 - e_2)^2 G_{0,2n+1}^2 \}$$

→ if states degenerate, *cross terms* ($\sim e_1 e_2$) *cancel* when averaged over nearby *even and odd parity* states

■ Dynamical cancellations

- duality is realized by summing over at least one complete set of *even* and *odd* parity resonances
- in NR Quark Model, even & odd parity states generalize to **56** ($L=0$) and **70** ($L=1$) multiplets of spin-flavor SU(6)

representation	${}^2\mathbf{8}[\mathbf{56}^+]$	${}^4\mathbf{10}[\mathbf{56}^+]$	${}^2\mathbf{8}[\mathbf{70}^-]$	${}^4\mathbf{8}[\mathbf{70}^-]$	${}^2\mathbf{10}[\mathbf{70}^-]$	Total
F_1^p	$9\rho^2$	$8\lambda^2$	$9\rho^2$	0	λ^2	$18\rho^2 + 9\lambda^2$
F_1^n	$(3\rho + \lambda)^2/4$	$8\lambda^2$	$(3\rho - \lambda)^2/4$	$4\lambda^2$	λ^2	$(9\rho^2 + 27\lambda^2)/2$

$\lambda(\rho) =$ (anti) symmetric component of ground state wave function

Close, WM (2003, 2009)

■ Dynamical cancellations

→ in SU(6) limit $\lambda = \rho$, with relative strengths of $N \rightarrow N^*$ transitions

$SU(6) :$	$[56, 0^+]^2 8$	$[56, 0^+]^4 10$	$[70, 1^-]^2 8$	$[70, 1^-]^4 8$	$[70, 1^-]^2 10$	<i>total</i>
F_1^p	9	8	9	0	1	27
F_1^n	4	8	1	4	1	18

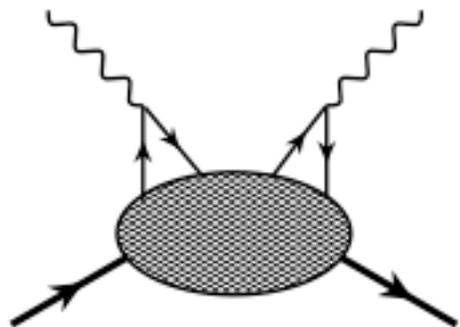
→ summing over all resonances in 56^+ and 70^- multiplets

$$\frac{F_1^n}{F_1^p} = \frac{18}{27} = \frac{2}{3}$$

→ at the quark level, n/p ratio is

$$\frac{F_1^n}{F_1^p} = \frac{4d + u}{d + 4u} = \frac{6}{9} = \frac{2}{3} \quad ! \quad \text{if } u = 2d$$

■ Accidental cancellations of charges?



cat's ears diagram (4-fermion higher twist $\sim 1/Q^2$)

$$\propto \sum_{i \neq j} e_i e_j \sim \left(\sum_i e_i \right)^2 - \sum_i e_i^2$$

↑ *coherent*
↑ *incoherent*

proton HT $\sim 1 - \left(2 \times \frac{4}{9} + \frac{1}{9} \right) = 0 !$

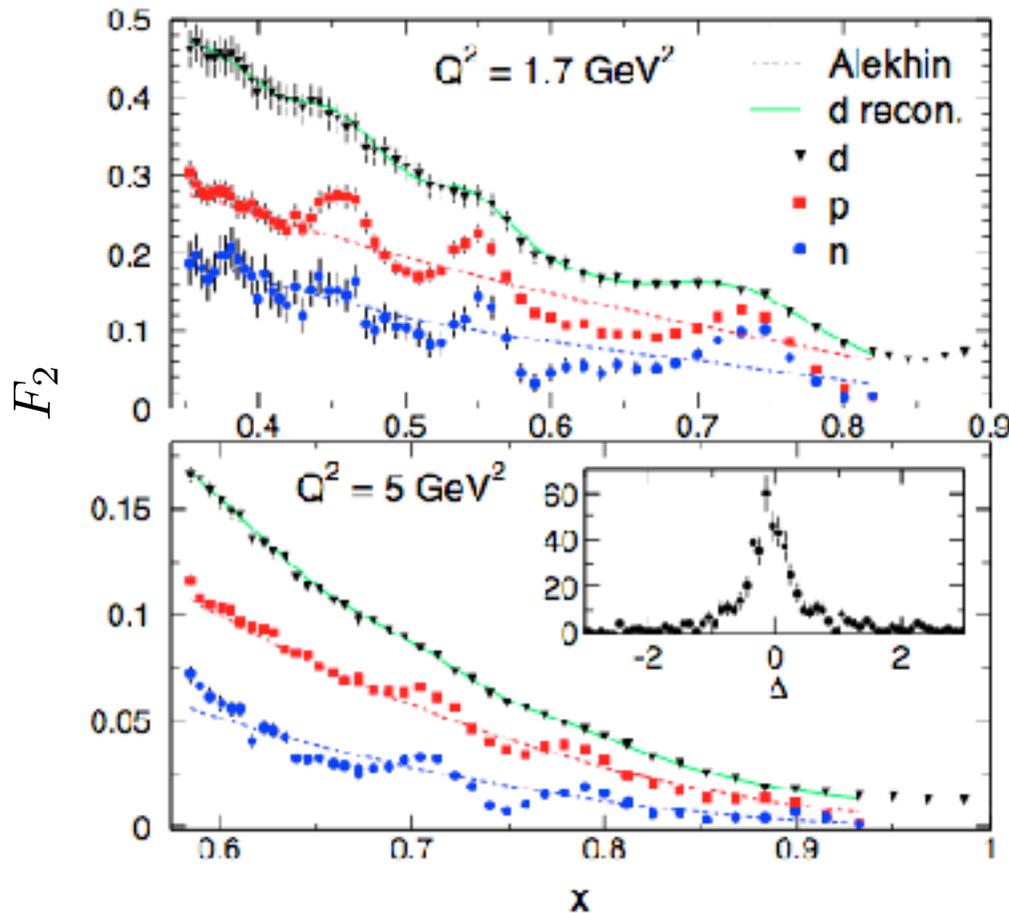
neutron HT $\sim 0 - \left(\frac{4}{9} + 2 \times \frac{1}{9} \right) \neq 0$

→ duality in proton a *coincidence!*

→ should not hold for neutron !!

Duality in electron-neutron scattering

- No free neutron targets, but iterative method allows neutron resonance structure to be extracted from deuteron & proton data



Malace, Kahn, WM, Keppel (2010)

→ *locally*, violations of duality in resonance regions < 15–20% (largest in Δ region)

→ evidence that duality is *not accidental*, but a general feature of resonance–scaling transition!

→ C. Keppel & I. Niculescu talks

Outlook and open questions

- Confirmation of duality (experimentally & theoretically) suggests origin in dynamical cancelations between resonances
 - explore more realistic descriptions based on phenomenological $\gamma^* NN^*$ form factors
- Era of “quantitative duality” — need to define the extent to which duality “works”
- Is duality between (high energy) continuum and resonances, or between total (resonance + background)?
 - “resonance region” vs. “resonances”
 - incorporate nonresonant background in same framework + quantum mechanics
- Where does duality *not* work (and why)?

A photograph of a cave interior, showing a variety of rock formations. On the left, there are several stalactites hanging from the ceiling. The walls and ceiling are covered in textured, layered rock. A bright yellow rectangular box is superimposed over the center of the image, containing the text "Thank you!".

Thank you!