Results from the Hall A GMp Experiment (E12-07-108)

Eric Christy Hampton University / Jefferson Lab





on behalf of the GMp collaboration

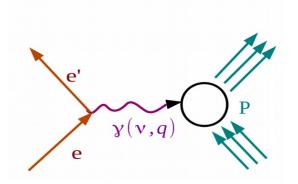
2019 GHP Meeting Denver, CO April 10, 2019

Proton magnetic form factor

- Form factors encode electric and magnetic structure of the nucleon
 - → Form factors characterize the spatial distribution of the electric charge and the magnetization current in the nucleon

|Form Factor|² =
$$\frac{\sigma(\text{Structured object})}{\sigma(\text{Point like object})}$$

In one photon exchange approximation the cross section in *ep* scattering when written in terms of $G_{\scriptscriptstyle M}^{\scriptscriptstyle p}$ and $G_{\scriptscriptstyle E}^{\scriptscriptstyle p}$ takes the following form:



$$\frac{d\sigma}{d\Omega} = \sigma_{Mott} \frac{\varepsilon \left(G_E^p\right)^2 + \tau \left(G_M^p\right)^2}{\varepsilon \left(1 + \tau\right)}, \quad \sigma_{Mott} = \frac{\alpha^2 \cos^2 \frac{\theta}{2}}{4 E^2 \sin^4 \frac{\theta}{2}} \frac{E'}{E}$$

$$\mathcal{J}_{proton} = e\bar{N}(p') \left[\gamma^{\mu} F_1(Q^2) + \frac{i\sigma^{\mu\nu} q_{\nu}}{2M} F_2(Q^2) \right] N(p)$$

$$\mathcal{J}_{\text{proton}} = e\bar{N}(p') \left[\gamma^{\mu} F_1(Q^2) + \frac{i\sigma^{\mu\nu} q_{\nu}}{2M} F_2(Q^2) \right] N(p)$$

Where,

$$G_E = F_1 - \tau F_2 \qquad G_M = F_1 + F_2$$

$$\tau = \frac{Q^2}{4M^2}$$
, $\epsilon = \left[1 + 2(1 + \tau)\tan^2\left(\frac{\theta}{2}\right)\right]^{-1}$

Methods of measurements

Rosenbluth separation method:

→ This method uses different beam energies and angle at fixed Q²

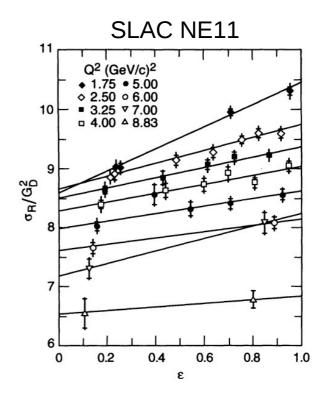
$$\sigma_R = \frac{d\sigma}{d\Omega} \frac{\varepsilon(1+\tau)}{\tau\sigma_{Mott}} = \frac{\varepsilon}{\tau} (G_E^p)^2 + (G_M^p)^2,$$

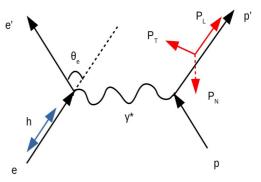
The slope of $\sigma_R(\varepsilon)$ is directly related to G_E^p and the intercept to

Recoil polarization technique:

Polarized electron transfers longitudinal polarization to $G_{\scriptscriptstyle M}^{\scriptscriptstyle p}$, but transverse polarization to $G_{\scriptscriptstyle M}^{\scriptscriptstyle p}$

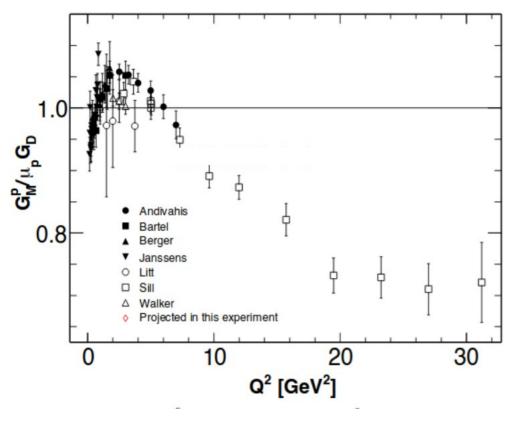
$$\frac{G_E}{G_M} = -\frac{P_t}{P_l} \frac{E_e + E_{e'}}{2M} \tan\left(\frac{\theta_e}{2}\right)$$



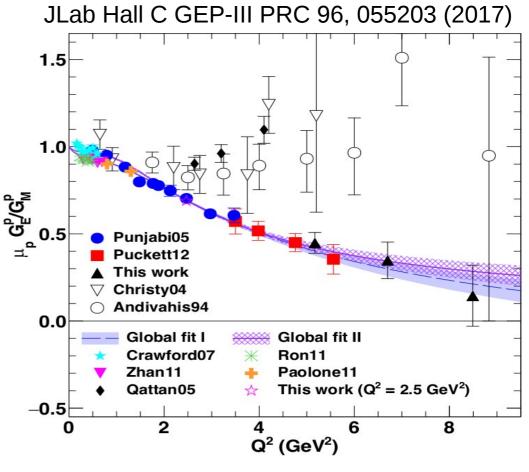


Polarization transfer cannot determine the values of G_E and G_M but can determine the from factor ratio.

Experimental Status of Proton Form Factors



→ Discrepancy in G_E/G_M P-T and Rosenbluth (ϵ) separations

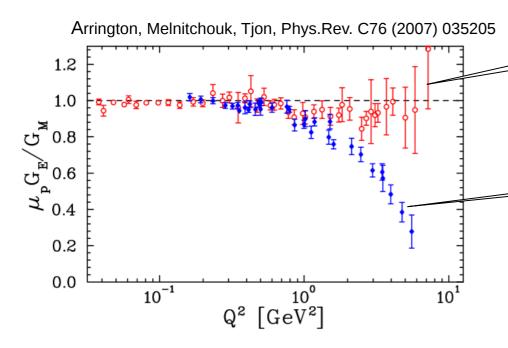


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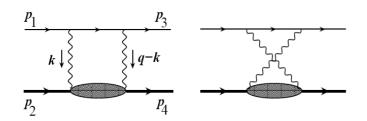
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4

Resolving th Rosenbluth vs P-T discrepancy



Leading explanation is hard 2- γ exchange, *not* included In *standard* radiative corrections of Mo-Tsai, etc.

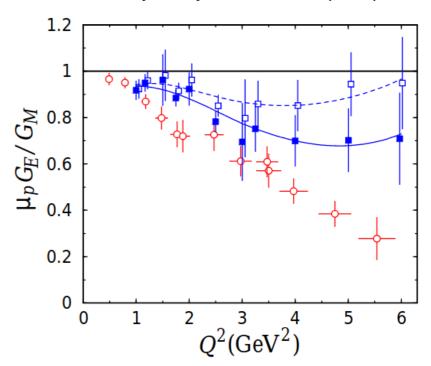


→ Expected to be relatively small for P-T method



Polarization data (PT)

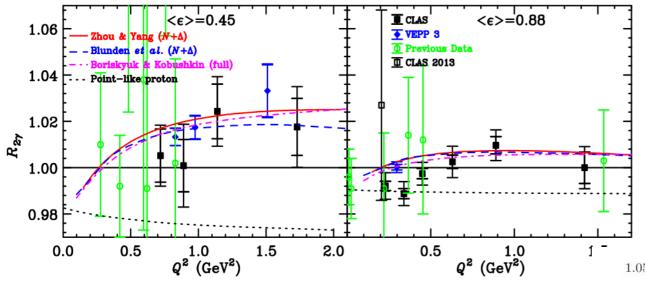
Effect of Hadronic 2-γ corrections of Blunden, Melnitchouk, Tjon Phys.Rev.Lett. 91 (2003) 142304



2-γ contributions from e+p / e-p ratios

Hard 2- γ contribution comes in with different signs for e+p and e-p =>

$$\sigma$$
+/ σ - = R_{2 γ} ~ 1-2 δ _{2 γ}



New data from

- VEPP-3
- CLAS

1.06

1.04

1.02

1.00

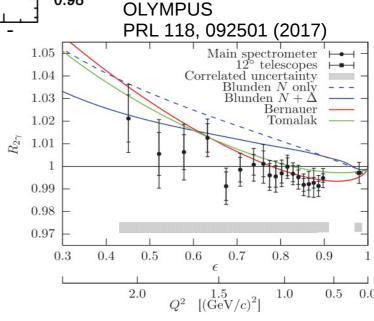
0.98

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OLYMPUS

Conclusions from combined analysis of A. Afanasev, P. G. Blunden, D. Hasell, and B. A. Raue:

- → CLAS and VEPP-3 and OLYMPUS data exclude no TPE hypothesis at >95% confidence level
- Data of insufficient precision to distinguish calculations of 2-γ contributions
- → Renormalization of OLYMPUS results required at twice the estimated uncertainty Eric Christy



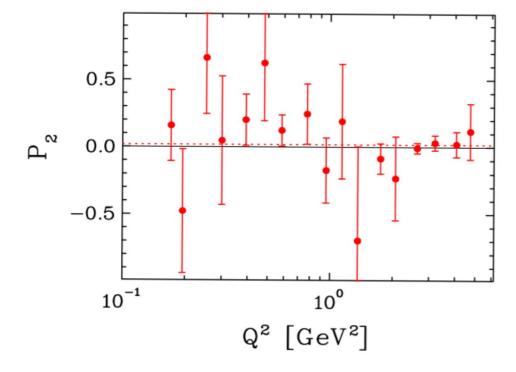
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Non-linearities in existing Rosenbluth data

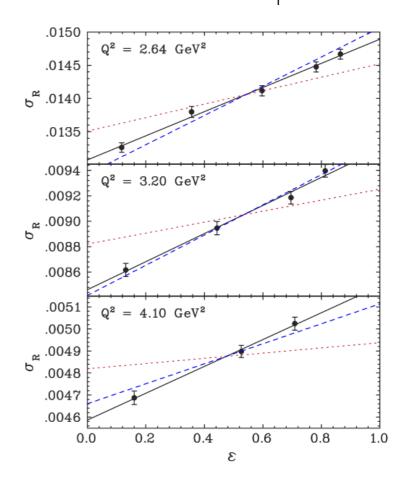
 \rightarrow Existing data indicate *no significant* non-linearities vs ε

Fit of elastic data to quadratic form

$$\sigma_{r} = P_{0} + P_{1}(\epsilon - 0.5) + P_{2}(\epsilon - 0.5)^{2}$$
 $< P_{2} > = 0.019 \pm 0.027$



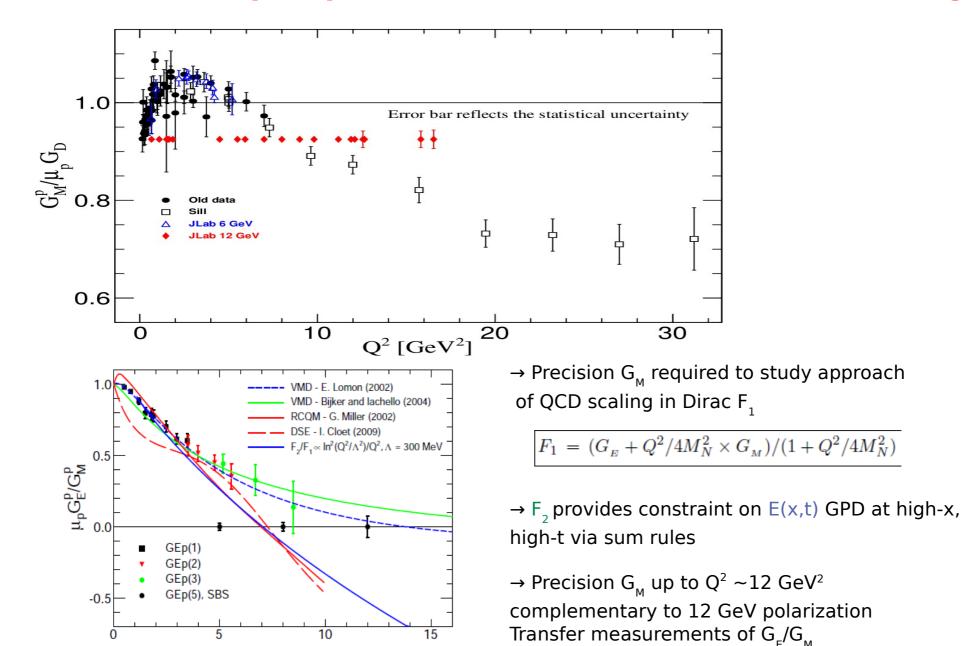
Super-Rosenbluth data also consistent with linear ϵ dependence of $\sigma_{\rm c}$



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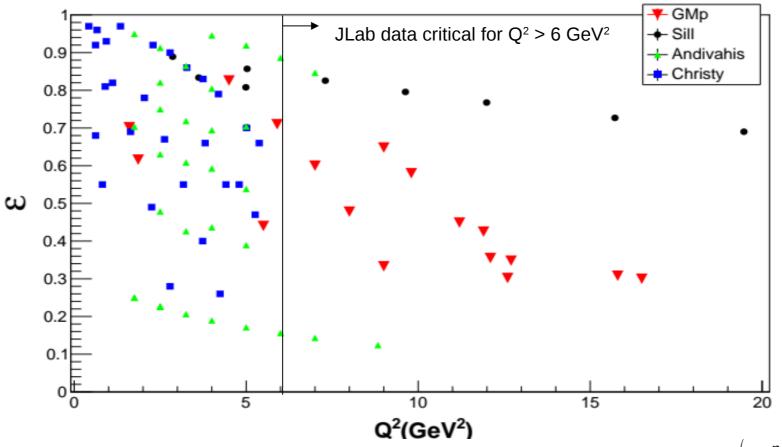
Precision GMp is part of the 12 GeV Form Factor Program



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Q² [GeV²]

GMp and other High Q² data



 \rightarrow GMp12 data at much smaller ϵ than Sill data

$$\frac{d\sigma}{d\Omega} = \sigma_{Mott} \frac{\boldsymbol{\epsilon} (G_E^p)^2 + \tau (G_M^p)^2}{\boldsymbol{\epsilon} (1+\tau)},$$

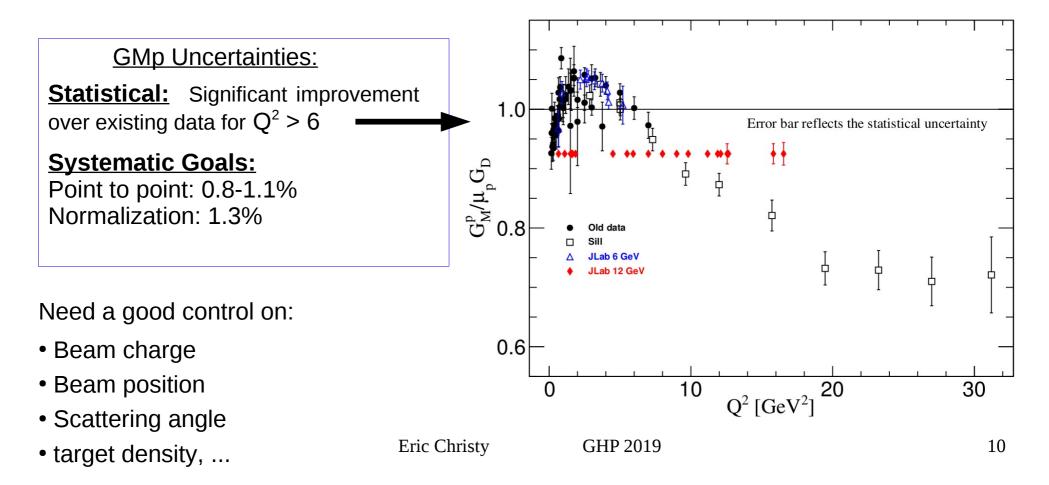
- Less sensitivity to G_E in extracting G_M
- Lever arm in ε provides sensitivity to:
 - 2γ from global fit utilizing $G_{_{\rm F}}$ / $G_{_{\rm M}}$ from polarization transfer

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E12-07-108 Experiment Overview

- Precision measurement of the elastic ep cross-section over the wide range of the Q^2 and extraction of proton magnetic form factor
- \rightarrow To improve the precision of cross section at high Q² by a factor of 3
- To provide insight into scaling behavior of the form factors at high Q²

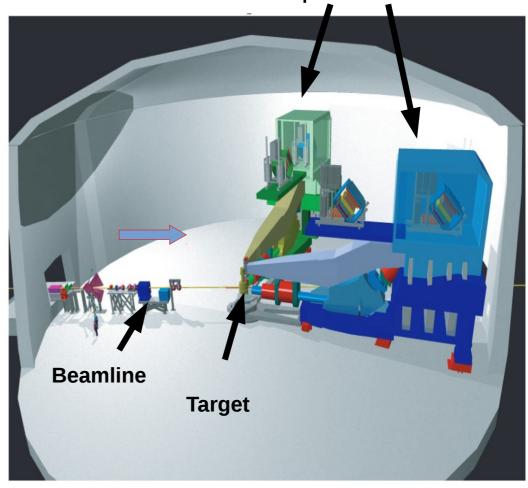


Experimental setup

Jefferson Lab at Newport News Virginia

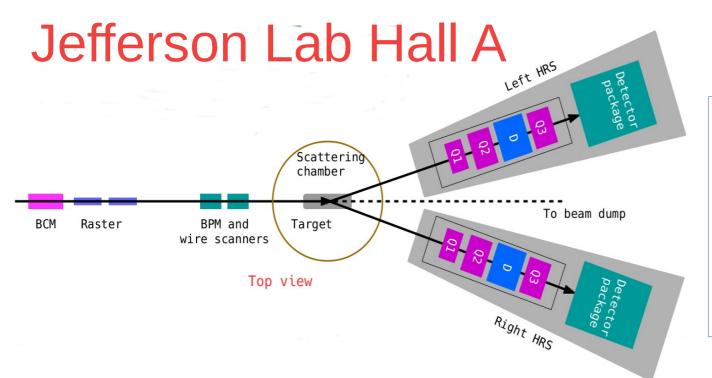
High resolution spectrometers





CEBAF: Continuous Electron Beam Accelerator Facility

Experimental Hall A



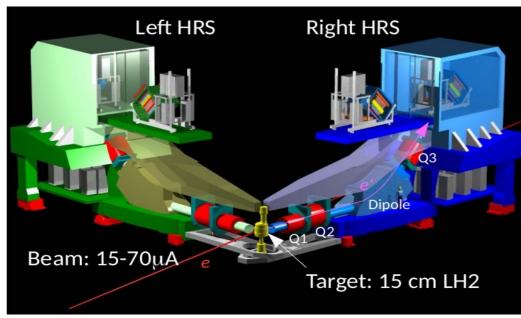
HRS Parameters:

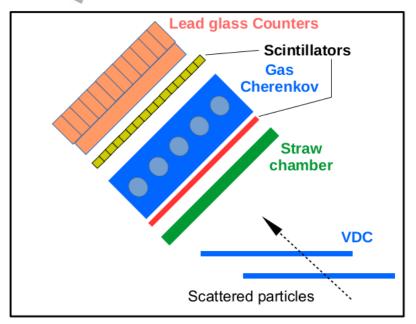
Acceptance: $-4.5\% < \Delta p/p < 4.5\%$, 6 msr

Resolution: δp/p≤2x10⁻⁴

 $\Delta x'_{tar} = 0.5 \text{ mrad (Horizontarl)}$

 $\Delta y'_{tar} = 1.0 \text{ mrad (Vertical)}$





Data collected during GMp

Spring 2015:

E _{beam} (GeV)	HRS	P ₀ (GeV/c)	Θ _{HRS} (deg)	Q ² (GeV/c) ²	Events(k)
2.06	R	1.15	48.7	1.65	157
2.06	L	1.22	45.0	1.51	386
2.06	L	1.44	35.0	1.1	396
2.06	L	1.67	25.0 *	0.66	405

Spring 2016:

* Surveyed angles

E _{beam} (GeV)	HRS	P _o (GeV/c)	Θ _{HRS} (deg)	Q² (GeV/c)²	Events(k)
4.48	R	1.55	52.9	5.5	108
8.84	R	2.10	48.8*	12.7	8
8.84	L	2.50	43.0*	11.9	11
11.02	R	2.20	48.8*	16.5	0.7

Fall 2016: *Most complete systematic studies during this period

E _{beam} (GeV)	HRS	P _o (GeV/c)	Θ _{HRS} (deg)	Q² (GeV/c)²	Events(k)
2.22	R	1.23	48.8*	1.86	356
2.22	L	1.37	42.0*	1.57	2025
8.52	L	2.53	42.0*	11.2	18.9
8.52	L	3.26	34.4	9.8	57.6
8.52	L	3.69	30.9*	9.0	11.6
6.42	L	3.22	30.9*	5.9	48.6
6.42	L	2.16	44.5*	8.0	27.2
6.42	L	3.96	24.3	4.5	30.5
6.42	L	2.67	37.0	7.0	41.4
6.42	R	1.59	55.9*	9.0	11.6
8.52	R	2.06	48.6*	12.1	11
8.52	R	1.80	53.5*	12.6	3.4
10.62	R	2.17	48.8*	15.8	3.6

Measurement of Elastic Cross Section

Cross section:

$$\frac{d\sigma}{d\Omega}(\theta) = \int dE' \frac{N_{\text{det}}(E', \theta) - N_{\text{BG}}(E', \theta)}{\mathcal{L} \cdot \epsilon_{\text{eff}} \cdot \text{LT}} \cdot A(E', \theta) \cdot \text{RC}$$

Reduced cross section:

$$\sigma_{\rm red} = \frac{d\sigma}{d\Omega} \frac{\epsilon(1+\tau)}{\sigma_{\rm Mott}} = \frac{4E^2 \sin^4 \frac{\theta}{2}}{\alpha^2 \cos^2 \frac{\theta}{2}} \frac{E}{E'} \epsilon (1+\tau) \frac{d\sigma}{d\Omega}$$

- Parameters:
 - ullet N_{det} : number of scattered elastic electrons detected
 - N_{BG} : events from background processes
 - \mathcal{L} : Integrated luminosity
 - : Corrections for efficiencies

- LT: live time correction
- $A(E',\theta)$: spectrometer acceptance
- RC: radiative correction factor
- E: beam energy
- θ : Scattering angle

A thorough understanding of all these parameters is crucial for a precision cross section measurement

Extraction of Elastic ep Cross Section

$$\frac{d\sigma}{d\Omega}^{data}(\theta) = \int dE' \frac{N^{data}(E',\theta) - N_{BG}(E',\theta)}{\mathcal{L}^{data}.\epsilon.LT} \cdot \frac{RC^{data}}{A^{data}(E',\theta)} \tag{1}$$

$$\frac{d\sigma}{d\Omega}^{mod}(\theta) = \int dE' \frac{N^{MC}(E',\theta)}{\mathcal{L}^{MC}} \cdot \frac{RC^{MC}}{A^{MC}(E',\theta)}$$

Radiative effects in Monte-Carlo based on improved Mo-Tsai from

R. Ent et. al Phys.Rev. C64 (2001) 054610

$$\frac{d\sigma}{d\Omega}^{data}(\theta)/\frac{d\sigma}{d\Omega}^{mod}(\theta) = \frac{\int^{E_{max}} (N^{data}(E',\theta) - N_{BG}(E',\theta))dE'}{\int^{E_{max}} N^{MC}dE'} \cdot \frac{A^{MC}(E',\theta)}{A^{data}(E',\theta)} \cdot \frac{RC^{data}}{RC^{MC}}$$

Assuming acceptance and ratiative contributions are correctly modeled:

$$\frac{d\sigma}{d\Omega}^{data}(\theta) = \frac{d\sigma}{d\Omega}^{mod}(\theta).\frac{Y^{data}}{Y^{MC}}$$

→ Results were cross checked with acceptance correction method (eq 1) using Rad Cor based on code utilized for later SLAC experiments. Plata
— Simulation

No.85

0.90

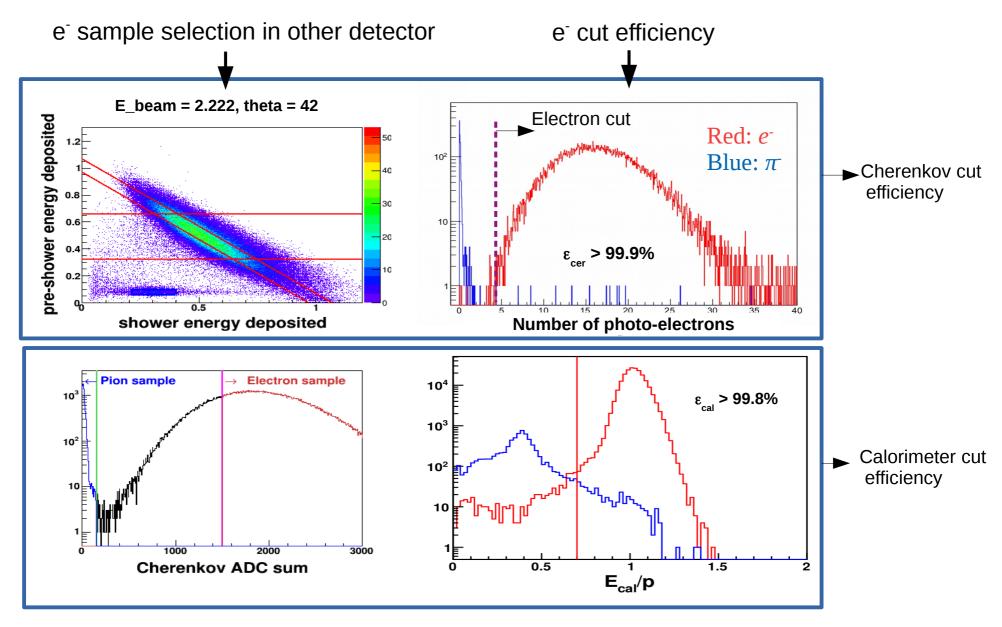
0.95

Invariant mass (GeV/c²)

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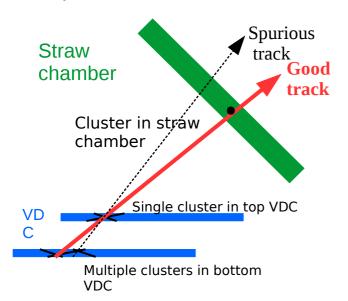
Detector efficiencies



$$\frac{\delta \epsilon}{\epsilon} < 0.1\%$$

VDC Track Reconstruction Efficiency

- Standard Tracking for HRS VDCs utilizes single cluster only in each chamber
- GMp utilized additional Straw Chamber to perform precise checks on efficiency determination

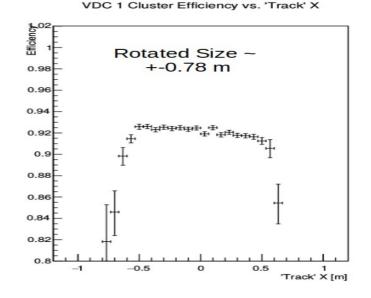


Elastic events were reconstructed with:							
 single cluster in both VDCs single cluster in 1 VDC + SC Longwu Ou (MIT) 							
Kinematic	K3-4	K3-6	K3-7	K3-8	K4-9	K4-10	K4-11
Corrected Yield ratio	1.0016	0.9994	0.999	0.9985	1.0007	1.0021	0.9997
Corrected yields agree to better than 0.2%							

A "coarse" track was formed using scintillator hit and straw chamber. This method enables us to estimate the track intercept at the focal plane without using VDC hits

Barak Schmookler (MIT)

Bashar Aljawrneh (NC A&T)

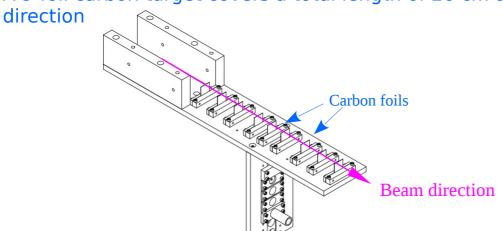


Significant Effort to Improve Optics Calibration

Longwu Ou (MIT)

• **Angle and vertex calibration:** used deep inelastic electrons from multi-foil carbon target

A 9-foil carbon target covers a total length of 20 cm along the beam



Spectrometer entrance

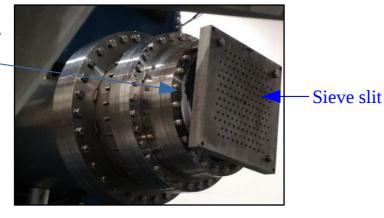
Sieve slit

Multi-foil

target

A 1-inch-thick tungsten sieve slit with high density holes at the spectrometer entrance selects scattered electrons in specific directions

Spectrometer entrance



• Algorithm: Minimization of χ^2 by varying the optics coefficients

$$\chi^2(y_{tg}) = \sum_{\text{events}} (Y_{ijkl} x_{fp}^i \theta_{fp}^j y_{fp}^k \phi_{fp}^l - y_{tg}^{\text{survey}})^2$$

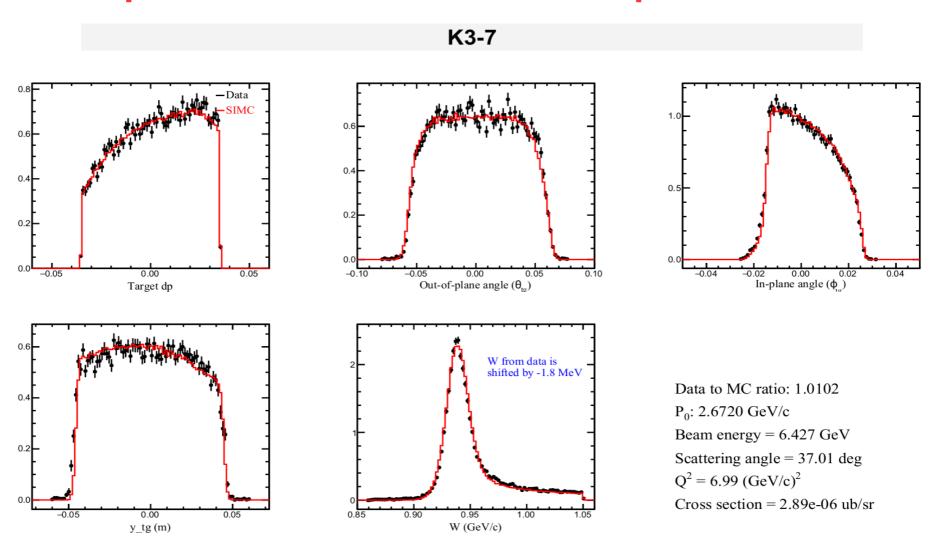
18

Momentum calibration: used elastic electrons from liquid hydrogen target

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Example Data to Monte Carlo Comparison: LHRS

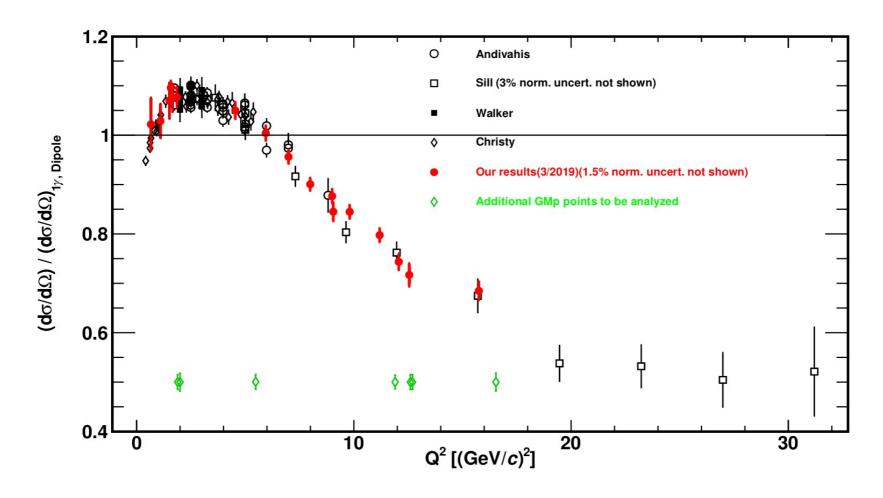


- Excellent comparison after subtraction of target cell endcaps via dummy (~3%)
- Small offsets in W consistent with estimated kinematic uncertainties

Error Budget (LHRS Fall 2016)

Source	$d\sigma/\sigma$ (%) (pt-pt)	$d\sigma/\sigma$ (%) (Norm.)
Beam charge ($\Delta I = 0.06 \ \mu A$)	$0.6(at\ 10\ \mu A)$ - $0.1(at\ 65\ \mu A)$	0.1
Scattering angle ($\Delta\theta = 0.2 \text{ mrad}$)	0.1 - 0.4	0.1 - 0.4
Beam energy ($\Delta E = 5 \times 10^{-4}$)	0.3	0.3
Boiling	$< 0.35 \text{ (at } 10 \ \mu\text{A)} - 0 \text{(at } 60 \ \mu\text{A)}$	$0.35 \; (at \; 60 \; \mu \; A)$
Optics	0.3	0.3
Track Reco	0.2	0.2
PID	0.1	0.1
Trigger	0.2	0.1
Target Length		0.1
Spectrometer acceptance	0.7	0.8
Radiative correction	0.8	1.0
Background subtraction	0.2	0.2
Cross section model		0.1
Total	1.2 - 1.3%	1.4- 1.6%

GMp - E012-07-108 final cross sections

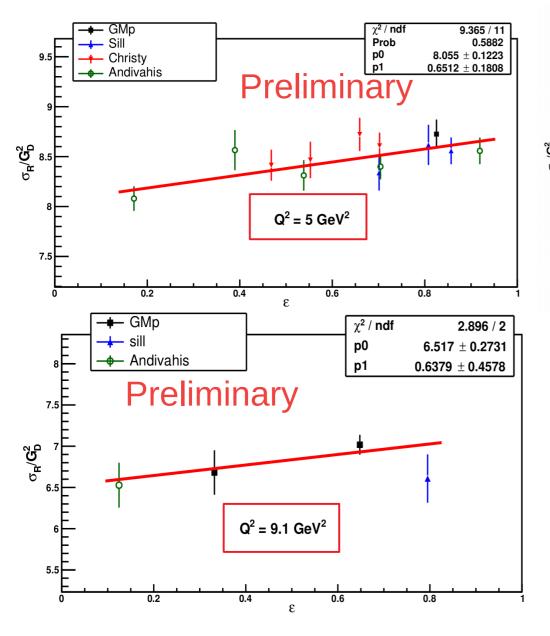


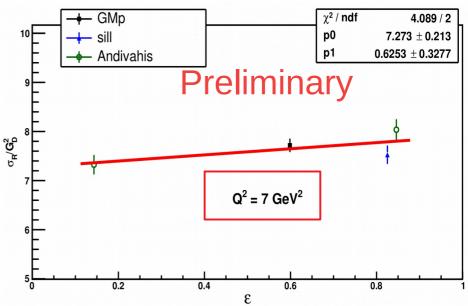
- Cross section relative to 1- γ cross section calculated with $G_E = G_M/\mu = G_{dip}$
- Significant improvement in precision for $Q^2 > 6$.
- Systematic uncertainties on Fall 2016 LHRS data ~1.3% (pt-pt), 1.5% (norm) RHRS (additional 2% from optics)

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Sample GMp Global Rosenbluth separations

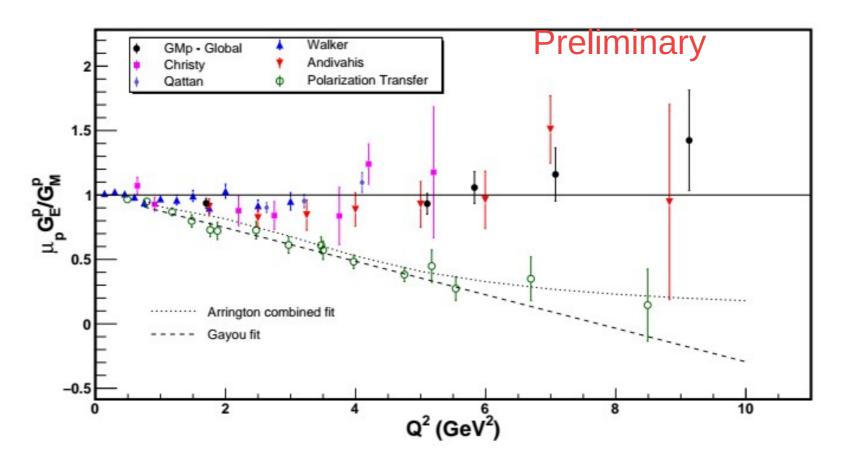




- → utilize all available data
- \rightarrow σ_r centered to common Q² utilizing Cross section fit.

Relative normalization not yet applied

Impact of E12-07-108 data on G_E/G_M at large Q^2



- Lab Hall A GMp12 data significantly reduce uncertainties on G_E/G_M at largest Q^2 => further highlights discrepancy with P-T data up to $Q^2 > 9$
- Full data set provides significantly more sensitivity than shown in select L/T separations

2-γ form factors

P. A. M. Guichon and M. Vanderhaeghen, PRL 91, 142303 (2003).

$$\sigma_{r} = \underline{G_{M}^{2} + 2G_{M} \Re(\delta \widetilde{G}_{M})} + \frac{\epsilon}{\tau} \left[G_{E}^{2} + \frac{4\tau^{2}}{M^{2}} \Re(\widetilde{F}_{3}) (G_{M} + \frac{1}{\tau} G_{E}) + 2G_{E} \Re(\widetilde{G}_{E}) \right]$$

Rosenbluth intercept

Rosenbluth Slope

$$\sigma_{r} = G_{M}^{2} + \frac{\epsilon}{\tau} G_{E}^{2} + 2G_{M} \Re \left(\delta \widetilde{G}_{M}\right) + \epsilon \left[\frac{2}{\tau} G_{E} \Re \left(\delta \widetilde{G}_{E}\right) + \frac{4\tau}{M^{2}} \Re \left(\widetilde{F}_{3}\right) \left(G_{M} + \frac{1}{\tau} G_{E}\right)\right]$$

$$r = \mu G_E / G_M$$

Assuming $2G_{E}\Re\left(\widetilde{G}_{E}
ight)$ is neglible

$$\sigma_r \approx G_M^2 + 2G_M \Re\left(\delta \widetilde{G}_M\right) + \frac{\epsilon}{\tau} \left[\frac{r^2}{\mu^2} G_M^2 + \frac{4\tau^2}{M^2} \Re\left(\widetilde{F}_3\right) G_M \left(1 + \frac{r}{\tau \mu}\right) \right]$$

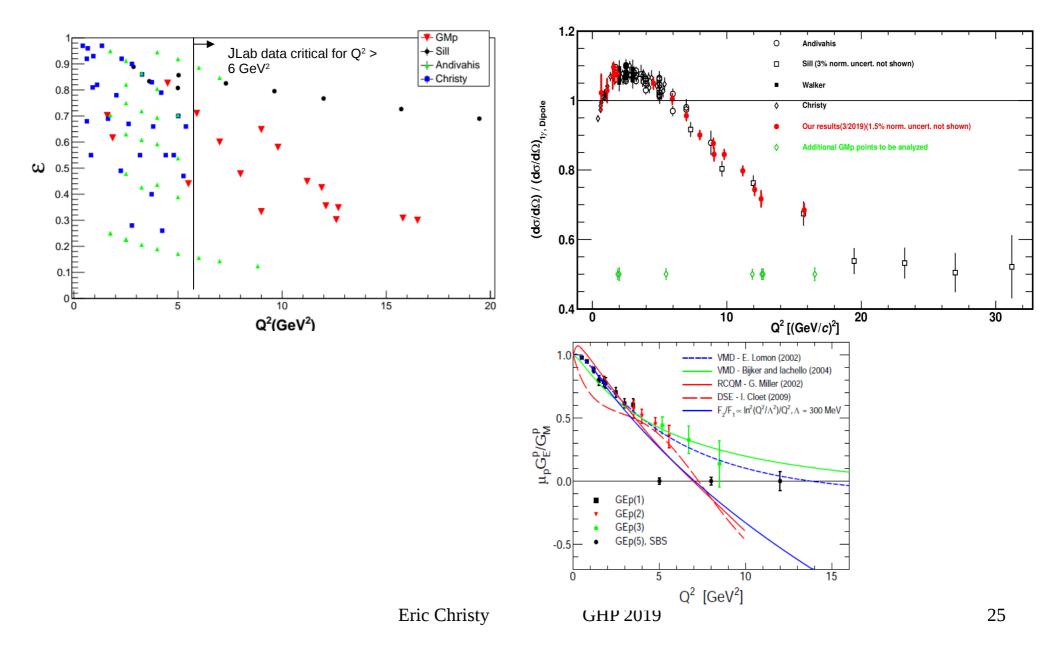
- \rightarrow r constrained by fit to P-T data
- → global fit to cross section data provides access to

$$\boxed{G_M^2(Q^2)} \boxed{ \boxed{\Re(\delta \widetilde{G}_M)(Q^2)}} \text{ And } \boxed{\Re(\widetilde{F}_3)(Q^2)} \boxed{ \blacksquare}$$

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GMp data provides enhanced access to Gmp 2-y Form Factors



Summary

• 12 GeV era GMp experiment in Jefferson Lab Hall A measured e-p elastic cross sections for 21 kinematics with $1 < Q^2 < 16.5 \text{ GeV}^2$

- Final Cross sections for Fall2016 data to be published soon with uncertainties of
 - 1.2 2% pt-pt
 - 1.5% normalization
- Data:
 - → important for JLab 12 GeV Form Factor and GPD program
 - → provides precision normalization for upcoming 12 GeV experiments at JLab
- ϵ coverage complementary to existing data and provides enhanced sensitivity to proton

 G_M and 2- γ Form Factors

→ full power of data through global fits.

GMp (E12-07-108) Analysis Team

- Spokesperson:
 - John Arrington
 - Eric Christy
 - Shalev Gilad
 - Vincent Sulkosky
 - Bogdan Wojtsekhowski
- Postdoc:
 - Kalyan Allada

- Ph.D students (all have defended):
 - Bashar Aljawrneh (NCA&T)
 - Thir Gautam (Hampton U.)
 - Longwu Ou (MIT)
 - Barak Schmookler (MIT)
 - Yang Wang (William & Mary)

Thanks to JLab accelerator team, Hall A target group, and all shift takers for their tremendous effort to make the GMp run successful

Thanks!

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