GHP meeting, Denver, April 11, 2019

Twist-3 PDFs/FFs

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based on L. Gamberg, Z. Kang, D. Pitonyak, M.S., S. Yoshida, JHEP 1901, 111 (2019)

Transverse Spin Effects in Single-Inclusive Hard Processes







 $\sqrt{s} = 20 \text{ GeV} [E704 \text{ coll.} (1991)]$

 $\sqrt{s} = 200 \text{ GeV} [STAR coll. (2008)]$

large effects cannot be explained in the standard parton model (using transversity) → collinear Twist-3 Formalism (Efremov, Teryaev, Qiu, Sterman) Collinear twist-3 formalism: several types of (chiral-even&odd) matrix elements compete

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intrinsic twist-3 PDF

$$g_T^q(x) = -\frac{1}{M} \int \frac{d\lambda}{4\pi} e^{i\lambda x} \langle P, \mathbf{S_T} | \bar{q}(0) \, \mathbf{s_T} \gamma_5 \, q(\lambda n) \, | P, \mathbf{S_T} \rangle$$

- sensitive to 'bad quark field components',
- twist-3 characteristics hidden in Dirac structure
- generates the g₂ structure function in DIS
- No probabilistic density interpretation



Collinear twist-3 formalism: several types of (chiral-even&odd) matrix elements compete

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$$G_T^q(z) = -\frac{1}{M_h} \frac{z}{4N_c} \sum_X \int \frac{d\lambda}{2\pi} e^{-i\frac{\lambda}{z}} \langle 0| q(0) | P_h S_{hT}, X \rangle \langle P_h S_{hT}, X | \bar{q}(\lambda n) \, \mathcal{S}_{hT} \gamma_5 | 0 \rangle$$

- Final state hadron spin may be reconstructed, e.g., for a Λ (later)

$$D_T^q(z) = \frac{1}{M_h} \frac{z}{4N_c} \sum_X \int \frac{d\lambda}{2\pi} e^{-i\frac{\lambda}{z}} \langle 0| q(0) | P_h S_{hT}, X \rangle \langle P_h S_{hT}, X | \bar{q}(\lambda n) \epsilon^{P_h n\alpha S_h} \gamma_{\alpha} | 0 \rangle$$



- yet another function, time-reversal violated in fragmentation process (!)

kinematical twist-3 PDFs:
Small transverse quark/gluon momenta k_T:

$$(k_T \times S_T) f_{1T}^{\perp,q}(x,k_T^2) \propto \int \frac{d\lambda d^2 z_T}{(2\pi)^3} e^{i\lambda x + ik_T \cdot z_T} \langle P, S_T | \bar{q}(0) \not \ll \mathcal{W} q(\lambda n + z_T) | P, S_T \rangle$$

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Sivers function
'transhelicity'

Collinear twist-3 formalism: TMD moments are needed

$$\int f_{1T}^{\perp,(1)}(x) = \int d^2k_T \, \frac{k_T^2}{2M^2} f_{1T}^{\perp}(x, k_T^2)$$

$$g_{1T}^{(1)}(x) = \int d^2k_T \ \frac{k_T^2}{2M^2} g_{1T}(x, k_T^2)$$

→ twist-3 characteristics through small transverse parton momentum k_T



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Same for fragmentation: kinematical twist-3 FF with transverse spin:

$$\Delta_{\partial}^{\alpha}(z) = \int d^2 \mathbf{p_T} \, \mathbf{p_T}^{\alpha} \, \Delta(z, z \mathbf{p_T}) \qquad \longrightarrow \qquad G_{1T}^{\perp(1), \Lambda/q}(z) \qquad D_{1T}^{\perp(1), \Lambda/q}(z)$$

Dynamical twist-3: Quark - Gluon - Quark Correlations (ETQS-matrix elements)



$$2M \, i\epsilon^{Pn\rho S} F_{FT}^{q}(x,x') = \int \frac{d\lambda}{2\pi} \int \frac{d\mu}{2\pi} \mathrm{e}^{i\lambda x'} \mathrm{e}^{i\mu(x-x')} \langle P, S_{T} | \bar{q}(0) \not n \, igF^{n\rho}(\mu n) \, q(\lambda n) | P, S_{T} \rangle$$

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<u>'dynamical twist - 3'</u>

 \rightarrow 3 - parton correlator: suppression by additional propagator

 \rightarrow <u>dependence on two parton momenta x, x'</u>: 2-dimensional support, richer parton dynamics

→ so far: only "diagonal support" π F_{FT}(x,x) = f_{1T}^{⊥(1)}(x) constraint by data

→ <u>'integrated' F_{FT}(x,x')</u>: average transverse color Lorentz force on struck quark [Burkardt, PRD88, 114502]

$$F^{n\rho} = [\vec{E} + \vec{n} \times \vec{B}]^{\rho} \propto \int dx \int dx' F_{FT}(x, x') \propto \int dx \, x^2 \, g_T(x)$$

'dynamical' twist-3 FF with transverse spin:

 $\Delta_F^{\alpha}(z,z') \sim \langle 0 | q(\lambda m) g F^{m\alpha}(\mu m) | P_{\Lambda}, S_{\Lambda}; X \rangle \langle P_{\Lambda}, S_{\Lambda}; X | \bar{q}(0) | 0 \rangle$ $\implies \hat{D}_{FT}^{\Lambda/q}(z,z'), \, \hat{G}_{FT}^{\Lambda/q}(z,z')$



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$$\frac{1}{z} - \frac{1}{z}$$

Same operator as for dynamical twist-3 PDFs, but: No gluonic or fermonic poles

$$FF(z,z) = 0$$
 FF

$$FF(z,0) = 0$$

$$\frac{\partial}{\partial z'}FF(z,z')\Big|_{z'=z} = 0$$

Support properties different

$$z \le z' < \infty$$

- No time reversal: dynamical twist-3 FFs are complex

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More 'dynamical' twist-3 functions (in pp at LO, in ep/e+e- at NLO) triple-gluon correlations qq-gluon correlation





QCD EoM relation & Lorentz-Invariance Relations [Kanazawa, Koike, Metz, Pitonyak, MS, PRD 2016]

QCD EoM for Twist-3 PDFs

$$g_{1T}^{(1)}(x) = x g_T(x) - \frac{m_q}{M} h_1(x)$$
$$+ \int_{-1}^1 dx' \frac{F_{FT}(x, x') - G_{FT}(x, x')}{x - x'}$$

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$$D_{1T}^{\perp(1)}(z) + \frac{D_T(z)}{z} = \int_0^1 d\beta \frac{\Im[\hat{D}_{FT}(z, z/\beta)] - \Im[\hat{G}_{FT}(z, z/\beta)]}{1 - \beta}$$
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LIR for Twist-3 PDFs based on translation invariance

 $g_T(x) = g_1(x) + \frac{d}{dx}g_{1T}^{(1)}(x) - 2\int_{-1}^1 dx' \,\frac{G_{FT}(x,x')}{(x-x')^2}$

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Two equations, three functions \rightarrow eliminate 'intrinsic & kinematical twist-3'

$$g_T(x) = \int_x^1 \frac{dy}{y} g_1(y) + \frac{m_q}{M} \left(\frac{1}{x} h_1(x) - \int_x^1 \frac{dy}{y^2} h_1(y) \right)$$
$$+ \int_x^1 \frac{dy}{y^2} \int_{-1}^1 dz \left[\frac{(1 - y\delta(y - x))F_{FT}(y, z)}{y - z} - \frac{(3y - z - y(y - z)\delta(y - x))G_{FT}(y, z)}{(y - z)^2} \right]$$

EoM & LIR relation crucial for gauge invariance, invariance of LC vector n

Λ^{\uparrow} - production in electron-positron annihilation at NLO

"Parton Model like" at LO











Typical NLO features:

* infrared safe (cancellation of $1/\epsilon^2$ - poles in dim. reg.)

$$\hat{\sigma}_{\text{virt}} + \hat{\sigma}_{\text{real}} = \mathcal{O}(1/\varepsilon) \qquad \hat{\sigma}^{q/g} \propto -\frac{1}{\varepsilon} P_{q/g\,q}(w) + \mathcal{O}(\varepsilon^0)$$





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$$D_{1,\text{bare}}^{\Lambda/q}(z) = D_{1,\text{ren}}^{\Lambda/q}(z) + \frac{\alpha_s}{2\pi} \frac{S_{\varepsilon}^{\text{MS}}}{\varepsilon} \sum_{i=q,g} \int_z^1 \frac{dw}{w} P_{iq}(w) D_{1,\text{ren}}^{\Lambda/i}(\frac{z}{w}) + \mathcal{O}(\alpha_s^2)$$

<u>Transverse Λ polarization at LO</u>

<u>'intrinsic' & 'kinematical' twist-3 FF:</u>

'dynamical' twist-3 FF:



<u>Single-Transverse Λ spin asymmetry</u>

- ♦ Unique effect driven by a single fragmentation function D_T → absent in DIS (1γ)
- * EoM needed at LO to preserve e.m. current conservation of hadronic tensor ($q_{\mu} W^{\mu\nu} = 0$) (EoM not optional!)

<u>Transverse Λ polarization at NLO</u>

[Gamberg, Kang, Pitonyak, M.S., Yoshida, JHEP 2019]

- Study the NLO dynamics for twist-3 fragmentation in the simplest process
- Different compared to twist-3 distributions (no pole contributions)

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Virtual & Real diagrams (qg/q - channel here, gg/g, qqb/g not shown)





E.o.M. - relations are crucial: Eliminate 'intrinsic' twist-3 contributions: only then color gauge invariance at NLO! ✓ Imaginary parts: In the dynamical fragmentation process & loop diagrams Infrared 1/ε² - poles cancel ✓ 1/ε - poles of imaginary parts of loops cancel through E.o.M. ✓ 1/ε - collinear poles of real parts of loops through MSbar - renormalization (?)





$$E_{h} \frac{d\sigma_{U}^{\text{EoM},2}}{d^{d-1}\vec{P}_{h}}(S_{h}) = (4\pi^{2}z_{h}^{2})^{\varepsilon} \frac{2\alpha_{\text{em}}^{2}N_{c}}{z_{h}s^{2}} \frac{2M_{h}\epsilon^{P_{h}mlS_{h}}}{s} (2v-1) \times \\ \sum_{q=u,\overline{u},...} e_{q}^{2} \left[-2D_{1T}^{\perp(1),q}(z_{h}) + 2\int_{0}^{1} d\beta \frac{\Im[\hat{D}_{FT}^{*} - \hat{G}_{FT}^{*}](z_{h}, \frac{z_{h}}{\beta})}{1-\beta} \right] \\ + \frac{\alpha_{s}}{2\pi}S_{\varepsilon} \int_{z_{h}}^{1} \frac{dw}{w^{2}} \int_{0}^{1} d\beta \left\{ \hat{\sigma}_{D_{1T}^{\perp(1)}}^{q;\text{EoM}}(w) D_{1T}^{\perp(1),q}(\frac{z_{h}}{w}) \right\}$$

























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read off evolution equations from collinear divergences for quark twist-3 FF $D_T(z)$

$$\begin{split} \frac{\partial}{\partial \ln \mu^2} \Big(D_T^f(z;\mu) \Big) &= \frac{z}{2} \int_z^1 \frac{\mathrm{d}w}{w^2} \int_0^1 \mathrm{d}\beta \left[P_{1,f \to f}^{[1]}(w) \, D_{1T}^{\perp(1),f}(\frac{z}{w};\mu) + P_{1,f \to g}^{[1]}(w) \, D_{1T}^{\perp(1),g}(\frac{z}{w};\mu) \right. \\ &+ P_{2,f \to fg}^{[1]}(w,\beta) \, \frac{\Im[\hat{D}_{FT}^{fg} - \hat{G}_{FT}^{fg}](\frac{z}{w},\beta;\mu)}{1 - \beta} + P_{3,f \to fg}^{[1]}(w,\beta) \, \frac{2\,\Im[\hat{D}_{FT}^{fg}](\frac{z}{w},\beta;\mu)}{(1 - \beta)^2} \\ &+ \sum_{f'=q',\bar{q}'} P_{4,f \to f'\bar{f}'}^{[1]}(w,\beta)\,\Im[\hat{D}_{FT}^{f'\bar{f}'}(\frac{z}{w},\beta;\mu)] + \sum_{f'=q',\bar{q}'} P_{5,f \to f'\bar{f}'}^{[1]}(w,\beta)\,\Im[\hat{G}_{FT}^{f'\bar{f}'}(\frac{z}{w},\beta;\mu)] \\ &+ P_{6,f \to gg}^{[1]}(w,\beta) \, \frac{\Im[\hat{N}_2^s(\frac{z}{w},\beta;\mu)]}{\beta^2(1 - \beta)^2} + P_{7,f \to gg}^{[1]}(w,\beta) \, \frac{\Im[\hat{N}_2^s(\frac{z}{w},\beta;\mu)]}{\beta^2(1 - \beta)^2} + P_{8,f \to gg}^{[1]}(w,\beta) \, \frac{\Im[\hat{N}_1(\frac{z}{w},\beta;\mu)]}{\beta^2(1 - \beta)^2} \Big] \end{split}$$

ordinary DGLAP splitting functions

$$P_{1,f\to f}^{[1]}(w) = -2\frac{C_F\alpha_s}{2\pi} \left(\frac{1+w^2}{(1-w)_+} + \frac{3}{2}\delta(1-w)\right) \qquad \qquad P_{1,f\to g}^{[1]}(w) = 4\frac{C_F\alpha_s}{2\pi} \left(\frac{1+(1-w)^2}{w}\right) = \frac{1}{2\pi} \left(\frac{1$$

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Final proof of one-loop factorization: Need to derive evolution equation directly from correlator! Previous work on *unpolarized* chiral-odd twist-3 fragmentation: [Belitsky, Kuraev, NPB 1996; Ma, Zhang, PLB 2017]

"The Gribov-Lipatov reciprocity fulfilled for two-particle cut-vertices only!"

Summary

- Transverse Spin Polarization: Long history, measured in ep/ppcollisions, theoretical treatment more complicated
- * We can learn about the parton dynamics in the nucleon and fragmentation process, non-perturbative QCD EoM and LIR are crucial
- * Λ production in e⁺e⁻: NLO completed,
 - \rightarrow calculate 'splitting functions' for polarized Λ fragmentation function
 - → Double Spin Asymmetry A_{LT} equally important...
- * A lot more work to do...