## The impact of $\eta$ photoproduction on the resonance spectrum

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## Degrees of freedom: Quarks or hadrons?

## The Missing Resonance Problem

- above 1.8 GeV much more states are predicted than observed,
"Missing resonance problem"
Lattice calculation (single hadron approximation):

[Edwards et al., Phys.Rev. D84 (2011)]
- only 15 established $N^{*}$ states (PDG 2015)
- $\sim 48 \%$ of the states have ${ }^{* * * *}$ or ${ }^{* * *}$ status (PDG 1982: $58 \%$ with ${ }^{* * * *}$ or ${ }^{* * *}$ )
$N^{*}$ spectrum in a relativistic quark model:


Löring et al. EPJ A 10, 395 (2001), experimental spectrum: PDG 2000
Overviews: Crede, Roberts, Rep. Prog. Phys. 76 (2013) Aznauryan et al., Int. J. Mod. Phys. E 22 (2013)

## Hybrid Baryons



- QCD at low energies
- Non-perturbative dynamics

Q1: how many are there?
Q2: what are they?
$\rightarrow$ mass generation \& confinement
$\rightarrow$ rich spectrum of excited states (missing resonance problem)
(2-quark/3-quark, hadron molecules, exotics,...)


## New results for the baryon spectrum

Quark-diquark with reduced pseudoscalar + vector diquarks: GE, Fischer, Sanchis-Alepuz, PRD 94 (2016)



- Scale $\Lambda$ set by $f_{\pi}$
- Current-quark mass $m_{q}$ set by $m_{\pi}$
- c adjusted to $\rho-a_{1}$ splitting
- $\eta$ doesn't change much

Using ONLY meson-baryon degrees of freedom (no explicit quark dynamics):

## Manifestly gauge invariant approach based on full BSE solution <br> [Ruic, M. Mai, U.-G. Meissner PLB 704 (2011)]



## Gauge invariance

- Exact unitary meson-baryon scattering amplitude T with parameters, fixed to reproduce:
- $\pi N$-partial wave $S_{11}$ and $S_{31}$ for $\sqrt{s}<1560 \mathrm{MeV}$

Arndt et al. (2012)

- $\pi^{-} p \rightarrow \eta n$ differential cross sections

Prakhov et al. (2005)



$\rightarrow$ Making the "Missing resonance problem" worse ?!

Phenomenology

## Resonances or not?


$\pi N \rightarrow \pi N$
EPECUR/SAID PRC 93 (2016)

[CBELSA/TAPS EPJA 53 (2017)]


## Current state in $\eta$ photoproduction: Multipoles from different groups



From: EtaMAID2018
[Tiator et al., EPJA54 (2018)] Analyzes:

$$
\begin{gathered}
\gamma p \rightarrow \eta p \\
\gamma p \rightarrow \eta^{\prime} p \\
\gamma n \rightarrow \eta n \\
\gamma n \rightarrow \eta^{\prime} n
\end{gathered}
$$

EtaMAID2018
BnGa [PLB 772 (2017)]
JuBo (dotted) [EPJA 54 (2018)] KSU [1804.06031]

Review: Krusche, Wilkins, [Prog.Part.Nucl.Phys. 80 (2014)]


## The Julich-Bonn Dynamical Coupled-Channel Approach

e.g. EPJ A 49, 44 (2013)

Dynamical coupled-channels (DCC): simultaneous analysis of different reactions
The scattering equation in partial-wave basis

$$
\begin{aligned}
&\left\langle L^{\prime} S^{\prime} p^{\prime}\right| T_{\mu \nu}^{\prime}|L S p\rangle=\left\langle L^{\prime} S^{\prime} p^{\prime}\right| V_{\mu \nu}^{\prime}|L S p\rangle+ \\
& \sum_{\gamma, L^{\prime \prime} S^{\prime \prime}} \int_{0}^{\infty} d q q^{2}\left\langle L^{\prime} S^{\prime} p^{\prime}\right| V_{\mu \gamma}^{\prime \prime}\left|L^{\prime \prime} S^{\prime \prime} q\right\rangle \frac{1}{E-E_{\gamma}(q)+i \epsilon}\left\langle L^{\prime \prime} S^{\prime \prime} q\right| T_{\gamma \nu}^{\prime \prime}|L S p\rangle
\end{aligned}
$$



## JuBo: Channels and Analytic Structure

Channels included:


- (2-body) unitarity and analyticity respected
- 3-body $\pi \pi N$ channel:
- parameterized effectively as $\pi \Delta, \sigma N, \rho N$
- $\pi N / \pi \pi$ subsystems fit the respective phase shifts
$\square$ branch points move into complex plane



## JuBo: Data base

[D. Roenchen, M. D., U.-G. Meißner, EPJ A 54, 110 (2018)

| Reaction | Observables (\# data points) | p./channel |
| :---: | :---: | :---: |
| $\begin{aligned} & \pi N \rightarrow \pi N \\ & \pi^{-} p \rightarrow \eta n \\ & \pi^{-} p \rightarrow K^{0} \Lambda \\ & \pi^{-} p \rightarrow K^{0} \Sigma^{0} \\ & \pi^{-} p \rightarrow K^{+} \Sigma^{-} \\ & \pi^{+} p \rightarrow K^{+} \Sigma^{+} \end{aligned}$ | PWA GW-SAID WI08 (ED solution) $\begin{aligned} & d \sigma / d \Omega \text { (676), } P \text { (79) } \\ & d \sigma / d \Omega \text { (814), } P(472), \beta \text { (72) } \\ & d \sigma / d \Omega \text { (470), } P(120) \\ & d \sigma / d \Omega \text { (150) } \\ & d \sigma / d \Omega \text { (1124), } P(551), \beta \text { (7) } \end{aligned}$ | $\begin{array}{r} 3,760 \\ 755 \\ 1,358 \\ 590 \\ 150 \\ 1,682 \end{array}$ |
| $\begin{aligned} & \gamma p \rightarrow \pi^{0} p \\ & \gamma p \rightarrow \pi^{+} n \\ & \gamma p \rightarrow \eta p \\ & \gamma p \rightarrow K^{+} \Lambda \end{aligned}$ | $\begin{aligned} & d \sigma / d \Omega(10743), \Sigma(2927), P(768), T(1404), \Delta \sigma_{31}(140), \\ & G(393), H(225), E(467), F(397), C_{x_{\mathrm{L}}^{\prime}}(74), C_{z_{\mathrm{L}}^{\prime}}(26) \\ & d \sigma / d \Omega(5961), \Sigma(1456), P(265), T(718), \Delta \sigma_{31}(231), \\ & G(86), H(128), E(903) \\ & d \sigma / d \Omega(5680), \Sigma(403), P(7), T(144), F(144), E(129) \\ & d \sigma / d \Omega(2478), P(1612), \Sigma(459), T(383), \\ & C_{x^{\prime}}(121), C_{z^{\prime}}(123), O_{x^{\prime}}(66), O_{z^{\prime}}(66), O_{x}(314), O_{z}(314), \end{aligned}$ | $\begin{aligned} & 17,564 \\ & 9,748 \\ & 6,507 \\ & 5,936 \end{aligned}$ |
|  | in total 48,050 |  |

## Resonance Couplings

Resonance states: Poles in the $T$-matrix on the $2^{\text {nd }}$ Riemann sheet


- $\operatorname{Re}\left(E_{0}\right)=$ "mass", $-2 \operatorname{lm}\left(E_{0}\right)=$ "width"
- elastic $\pi N$ residue ( $\left|r_{\pi N}\right|, \theta_{\pi N \rightarrow \pi N}$ ), normalized residues for inelastic channels $\left(\sqrt{\Gamma_{\pi N} \Gamma_{\mu}} / \Gamma_{\text {tot }}, \theta_{\pi N \rightarrow \mu}\right)$
- photocouplings at the pole: $\tilde{A}_{\text {pole }}^{h}=A_{\text {pole }}^{h} e^{i \vartheta^{h}}, h=1 / 2,3 / 2$

Inclusion of $\gamma p \rightarrow K^{+} \Lambda$ in JüBo ("JuBo2017-1"): 3 additional states

|  | $z_{0}[\mathrm{MeV}]$ | $\frac{\Gamma_{\pi N}}{\Gamma_{\text {tot }}}$ | $\frac{\Gamma_{\eta N}}{\Gamma_{\text {tot }}}$ | $\frac{\Gamma_{K \Lambda}}{\Gamma_{\text {tot }}}$ | $\frac{\Gamma_{K \Sigma}}{\Gamma_{\text {tot }}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{N}(1900) 3 / 2^{+}$ | $1923-i 108.4$ | $1.5 \%$ | $0.78 \%$ | $2.99 \%$ | $69.5 \%$ |
| $\mathrm{~N}(2060) 5 / 2^{-}$ | $1924-i 100.4$ | $0.35 \%$ | $0.15 \%$ | $13.47 \%$ | $27.02 \%$ |
| $\Delta(2190) \mathbf{1} / 2^{+}$ | $2191-i 103.0$ | $33.12 \%$ |  |  | $3.78 \%$ |

- $N(1900) 3 / 2^{+}$: s-channel resonances, seen in many other analyses of kaon photoproduction (BnGa), 3 stars in PDG
- $N(2060) 5 / 2^{-}$: dynamically generated, 2 stars in PDG, seen e.g. by BnGa
- $\Delta(2190) 3 / 2^{+}$: dyn. gen., no equivalent PDG state


## How to quantify the impact of new measurements?

Consider correlations of helicity couplings extracted from experiment
[D. Sadasivan,M.D., M. Mai, in preparation]


## Results from analysis of world data of $\eta$ photoproduction

[D. Sadasivan,M.D., M. Mai, in preparation]

Here $A=|A| e^{i \phi}$ defined at the resonance pole.



## Bulk properties of uncertainties from different data sets

| Helicity Coupling | All | No E | No F | No T | No $\Sigma$ |
| :--- | :---: | :---: | :---: | :---: | :---: |
| Number of Data Points | 6425 | 6369 | 6281 | 6281 | 6022 |
| Generalized Variance | 0.0494 | 0.0521 | 0.1288 | 0.1239 | 6.664 |
| $\sqrt{\operatorname{Tr} C}$ | 10.4965 | 10.51 | 12.00 | 11.423 | 19.85 |
| Multicollinearity | 8.173 | 8.203 | 9.280 | 9.5323 | 10.371 |
| Condition number | 133.61 | 132.10 | 173.664 | 164.1 | 322.66 |

C=Covariance Matrix

Generalized Variance
= Det[C] ~Volume of the Error Ellipsoid

| Helicity Coupling | No artificial data | Cx | Cz | Cx and Cz |
| :--- | :---: | :---: | :---: | :---: |
| Number of Data Points | 6425 | 6569 | 6569 | 6713 |
| Generalized Variance | 0.0494 | 0.03758 | 0.0362 | $\underline{0.0132}$ |
| $\sqrt{\operatorname{Tr} C}$ | 10.4965 | 10.72 | 10.487 | 10.102 |
| Multicollinearity | 8.173 | 7.599 | 6.770 | 6.157 |
| Condition number | 133.61 | 112.47 | 109.69 | 107.683 |



- Allows to trace quantitatively the impact of data sets and observables
- Helpful in design of new measurements
- Correlations allow to assess quality of theory predictions


## Resonances and other structures

CLAS/JuBo (M. D., D. Rönchen), Phys.Lett. B755 (2016)

- First-ever measurement of observable $E$ in $\eta$ photoproduction, enabled through the CLAS FROST target


Is this a new narrow baryonic resonance?
$\rightarrow$ Conventional explanation in terms of interference effects.


## Summary

- Complicated phenomenology of excited baryons through coupled-channel and three-body effects
$\rightarrow$ Conceptual progress needed to connect to lattice QCD calculations.
- $\eta$ photoproduction ideally suitable to study excited baryons
$\rightarrow$ Isospin filter, good channel for missing resonances
- Global analyses of pion and photon-induced reactions
$\rightarrow$ Jülich-Bonn analysis confirms new states in analysis of photoproduction
- Model selection techniques to extract minimal spectrum of excited baryons

Spare slides

## Spectrum of N* resonances



- Most new resonances by Bonn-Gatchina group; [Slide: V. Crede/Nstar 2017, slight modifications]
- Many from kaon photoproduction [See also: Crede, Roberts, Rep. Prog. Phys. 76 (2013)]

Spectrum of N* resonances


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[See also: Crede, Roberts, Rep. Prog. Phys. 76 (2013)]

$$
S=\mathbb{1}+i T
$$

Unitarity: $S S^{\dagger}=1 \Leftrightarrow-i\left(T-T^{\dagger}\right)=T T^{\dagger}$

- 3-body unitarity:
discontinuities from $t$-channel exchanges
$\rightarrow$ Meson exchange from requirements of the $S$-matrix



## Other cuts

- to approximate left-hand cut $\rightarrow$ Baryon $u$-channel exchange
- $\sigma, \rho$ exchanges from crossing plus analytic continuation.

$\vec{q}=\overrightarrow{p_{1}}-\overrightarrow{p_{3}}$

$\vec{q}=\vec{q}_{1}-\vec{p}_{4}$


$$
\vec{q}=\vec{p}_{1}+\vec{p}_{2}=0
$$

## Visible influence of new states



$N(1900) 3 / 2^{+}, N(2060) 5 / 2^{-}$in $\sigma_{\text {tot }}$ in $\pi^{-} p \rightarrow K^{+} \Sigma^{-}$


## Amplitude parametrization



Disp. rel. (Aznauryan, Burkert,..) KT equations, t-channel analyticity; Restoration of crossing symmetry via dispersion relations (Aitchison, Kubis, Szczepaniak, Tiator)


Integral-equation implementation of amplitude

- Giessen


【
$T=V+V G T$,
Genuine Resonance:



Unitarity loop G:

- Re G $\rightarrow 0$ : K-matrix
- V point-like: SAID Integral equation: Julich-Bonn, ANL-Osaka

Amplitude parametrization

- Giessen



Disp. rel. (Aznauryan, Burkert,..) KT equations, t-channel analyticity; Restoration of crossing symmetry via dispersion relations (Aitchison, Kubis, Szczepaniak, Tiator, ...)

## Impact of data

## CBELSA/TAPS

## Impact of new data



Data: CBELSA/TAPS Collaboration (T: Hartmann et al. PLB 748, 212 (2015) , E: Gottschall et al. PRL 112,
4012003 (2014), G: Thiel et al. PRL 109, 102001 (2012), Thiel et al. arXiv:1604.02922)
Predictions: black solid lines: BnGa, red dash-dotted: SAID, blue dashed: JüBo, green dotted: MAID

## Impact of new data



Data: CBELSA/TAPS Collaboration (T: Hartmann et al. PLB 748, 212 (2015) , E: Gottschall et al. PRL 112, 012003 (2014), G: Thiel et al. PRL 109, 102001 (2012), Thiel et al. arXiv:1604.02922)

Fits: black solid lines: BnGa, red dash-dotted: SAID, blue dashed: JüBo

## Impact of new data



- Multipole solutions approach each other
- Remaining discrepancies

Julich-Bonn, BnGa, SAID
$\operatorname{var}(1,2)=\frac{1}{2} \sum_{i=1}^{16}\left(\mathcal{M}_{1}(i)-\mathcal{M}_{2}(i)\right)\left(\mathcal{M}_{1}^{*}(i)-\mathcal{M}_{2}^{*}(i)\right) .(31)$


Fit to world data on $\pi N \rightarrow \pi N, \eta N, K \Lambda, K \Sigma$ ( $\sim 10^{5} \exp$. points) [Rönchen, M.D. et al., EPJA 49 (2013)]

Selected results for $\pi^{-} p \rightarrow K^{0} \Lambda$ [almost complete experiment]


## Re-measuring hadron-induced reactions

Fits: D. Rönchen, M.D., et al., EPJ A49 (2013)

$\rightarrow$ Physics Opportunities with meson beams, Briscoe, M.D., Haberzettl, Manley, Naruki, Strakovsky, Swanson, EPJ A51 (2015)

## Toward Data-driven Analyses

[M.D., Revier, Rönchen, Workman, arXiv:1603.07265, PRC 2016]

- Multi-channel analyses to detect faint resonance signals
- All groups use GW/SAID partial waves for $\pi N \rightarrow \pi N$
- The chi-square obtained in fits to single-energy solutions is not related to chi-square of a fit to data $\rightarrow$ Statistical interpretation of resonance signals difficult.
- Provide online covariance matrices etc. to allow other groups to perform correlated chi-square fits.

Slight adaptation of their code allows other groups to obtain a $\chi^{2}$ (almost) as if they fitted to $\pi N \rightarrow \pi N$ directly.

$$
\begin{aligned}
& \begin{aligned}
\chi^{2}(\mathbf{A}) & =\chi^{2}(\hat{\mathbf{A}})+(\mathbf{A}-\hat{\mathbf{A}})^{T} \hat{\Sigma}^{-1}(\mathbf{A}-\hat{\mathbf{A}}) \\
& +\mathcal{O}(\mathbf{A}-\hat{\mathbf{A}})^{3}
\end{aligned} \\
& \begin{array}{l}
\text { Covariance matrices etc. can be downloaded } \\
\text { on the SAID and JPAC web pages. }
\end{array} \\
& \hline
\end{aligned}
$$

## Amplitude reconstruction from complete experiments and

 truncated partial-wave expansions[Workman, Tiator, Wunderlich, M.D., H. Haberzettl, PRC (2017)]

How do complete experiment and truncated partial wave complete experiment compare. Depending on which partial-wave content is admitted in the amplitude?

| Set | Included Partial Waves | CEA | TPWA | Complete Sets for TPWA |
| :---: | :---: | :---: | :---: | :---: |
| 1 | $L=0\left(E_{0+}\right)$ | 1(1) | 1(1)1 | $I[1]$ |
| 2 | $J=1 / 2\left(E_{0+}, M_{1-}\right)$ | 4(4) | $\begin{aligned} & 4(4) 1 \\ & 4(3) 2 \end{aligned}$ | $\begin{aligned} & I[1], \check{P}[1], \check{C}_{x}[1], \check{C}_{z}[1] \\ & I[2], \check{P}[1], \check{C}_{x}[1] \end{aligned}$ |
| 3 | $L=0,1\left(E_{0+}, M_{1-}, E_{1+}\right)$ | 6(6) | $\begin{aligned} & 6(6) 1 \\ & 6(4) 2 \\ & 6(3) 3 \end{aligned}$ | $\begin{aligned} & I[1], \check{\Sigma}[1], \check{T}[1], \check{P}[1], \check{F}[1], \check{G}[1] \\ & I[2], \check{\Sigma}[1], \check{T}[2], \check{P}[1] \\ & I[3], \check{\Sigma}[1], \check{T}[2] \end{aligned}$ |
| 4 | $L=0,1\left(E_{0+}, M_{1-}, E_{1+}, M_{1+}\right)$ <br> full set of $4 S, P$ wave multipoles | $\dagger$ | $\begin{aligned} & 8(5) 2 \\ & 8(4) 3 \\ & \hline \end{aligned}$ | TPWA at 1 angle not possible $\begin{aligned} & I[2], \check{\Sigma}[1], \check{T}[2], \check{P}[2], \check{F}[1] \\ & I[3], \check{\Sigma}[1], \check{F}[2], \check{H}[2] \end{aligned}$ |
| 5 | $L=0,1,2\left(E_{0+}, M_{1-}, E_{1+}, E_{2-}\right)$ | 8(8) | $\begin{aligned} & 8(8) 1 \\ & 8(4) 2 \\ & 8(3) 3 \end{aligned}$ | $\begin{aligned} & I[1], \check{\Sigma}[1], \check{T}[1], \check{P}[1], \check{F}[1], \check{G}[1], \check{C}_{x}[1], \check{O}_{x}[1] \\ & I[2], \check{\Sigma}[2], \check{T}[2], \check{P}[2] \\ & I[3], \check{\Sigma}[2], \check{T}[3] \end{aligned}$ |
| 6 | $J \leq 3 / 2\left(E_{0+}, M_{1-}, E_{1+}, M_{1+}, E_{2-}, M_{2-}\right)$ | $\dagger$ | $\begin{aligned} & 12(5) 3 \\ & 12(4) 4 \end{aligned}$ | TPWA at 1 or 2 angles not possible $\begin{aligned} & I[3], \check{\Sigma}[2], \check{T}[3], \check{P}[2], \check{F}[2] \\ & I[4], \check{\Sigma}[2], \check{F}[3], \check{H}[3] \end{aligned}$ |
| 7 | $L=0,1,2\left(E_{0+}, \ldots, M_{2+}\right)$ <br> full set of $8 S, P, D$ wave multipoles | $\dagger$ | $\begin{array}{\|l} \left\lvert\, \begin{array}{l} 16(6) 3 \\ 16(5) 4 \end{array}\right. \\ \hline 16(4) 5 \end{array}$ | $\begin{aligned} & \text { TPWA at } 1 \text { or } 2 \text { angles not possible } \\ & I[3], \check{\Sigma}[3], \check{T}[3], \check{P}[3], \check{F}[3], \check{G}[1] \\ & I[4], \check{\Sigma}[3], \check{T}[3], \check{P}[3], \check{F}[3] \\ & \hline I[5], \check{\Sigma}[3], \check{F}[4], \check{H}[4] \quad \text { Four are } \end{aligned}$ |

Order: \# of different measurements, \# of different observables \# of different angles

Data: Akondi et al. (A2 at MAMI) PRL 113, 102001 (2014)

-=-=- prediction
fit

| Beam | Target | Recoil |
| :---: | :---: | :---: |
| 0 | $+y$ | 0 |
| 0 | $-y$ | 0 |



| Beam | Target | Recoil |
| :---: | :---: | :---: |
| +1 | $+x$ | 0 |
| -1 | $+x$ | 0 |

