Color Confinement and Supersymmetric Properties of Hadron Physics from Light-Front Holography



with Guy de Tèramond, Hans Günter Dosch, C. Lorce, K. Chiu, R. S. Sufian, A.Deur, T. Liu Yerevan, Armenia 24-28 September 2018

Correlations in Partonic and Hadronic Interactions 2018

Profound Questions for Hadron Physics

- Origin of the QCD Mass Scale
- Color Confinement
- Spectroscopy: Tetraquarks, Pentaquarks, Gluonium, Exotic States
- Universal Regge Slopes: n, L, both Mesons and Baryons
- Massless Pion: Bound State
- Dynamics and Spectroscopy
- QCD Coupling at all Scales
- QCD Vacuum Do QCD Condensates Exist?



The leading Regge trajectory: Δ resonances with maximal J in a given mass range. Also shown is the Regge trajectory for mesons with J = L+S.

E. Klempt and B. Ch. Metsch

2012



Supersymmetry in QCD

- A hidden symmetry of Color SU(3)c in hadron physics
- QCD: No squarks or gluinos!
- Emerges from Light-Front Holography and Super-Conformal Algebra
- Color Confinement
- Massless Pion in Chiral Limit

Partonic Correlations Yeravan September 25, 2018

Color Confinement, Hadron Dynamics, and Hadron Spectroscopy from Light-Front Holography and Superconformal Algebra





Measurements of hadron LF wavefunction are at fixed LF time

Like a flash photograph

Fixed
$$\tau = t + z/c$$

$$x_{bj} = x = \frac{k^+}{P^+}$$

Invariant under boosts! Independent of P^{μ}

Light-Front Wavefunctions: rigorous representation of composite systems in quantum field theory



Invariant under boosts! Independent of P^{μ}

Causal, Frame-independent. Creation Operators on Simple Vacuum, Current Matrix Elements are Overlaps of LFWFS

Light-Front QCD

Physical gauge: $A^+ = 0$

Exact frame-independent formulation of nonperturbative QCD!

$$\begin{split} L^{QCD} &\to H_{LF}^{QCD} \\ H_{LF}^{QCD} &= \sum_{i} [\frac{m^{2} + k_{\perp}^{2}}{x}]_{i} + H_{LF}^{int} \\ H_{LF}^{int}: \text{ Matrix in Fock Space} \\ H_{LF}^{QCD} |\Psi_{h} \rangle &= \mathcal{M}_{h}^{2} |\Psi_{h} \rangle \\ |p, J_{z} \rangle &= \sum_{n=3} \psi_{n}(x_{i}, \vec{k}_{\perp i}, \lambda_{i}) |n; x_{i}, \vec{k}_{\perp i}, \lambda_{i} \rangle \end{split}$$

Eigenvalues and Eigensolutions give Hadronic Spectrum and Light-Front wavefunctions

LFWFs: Off-shell in P- and invariant mass



Bound States in Relativistic Quantum Field Theory:

Light-Front Wavefunctions Dirac's Front Form: Fixed $\tau = t + z/c$

Fixed
$$\tau = t + z/c$$

 $\psi(x_i, \vec{k}_{\perp i}, \lambda_i)$
 $x = \frac{k^+}{P^+} = \frac{k^0 + k^3}{P^0 + P^3}$

Invariant under boosts. Independent of P^{μ}

$$\mathbf{H}_{LF}^{QCD}|\psi\rangle = M^2|\psi\rangle$$

Direct connection to QCD Lagrangian

LF Wavefunction: off-shell in invariant mass

Remarkable new insights from AdS/CFT, the duality between conformal field theory and Anti-de Sitter Space





Sign reversal in DY!

Final-State Interactions Produce Pseudo T-Odd (Sivers Effect)

Hwang, Schmidt, sjb Collins

- Leading-Twist Bjorken Scaling!
- Requires nonzero orbital angular momentum of quark
- Arises from the interference of Final-State QCD Coulomb phases in S- and Pwaves;
- Wilson line effect -- Ic gauge prescription
- Relate to the quark contribution to the target proton anomalous magnetic moment and final-state QCD phases
- QCD phase at soft scale!
- New window to QCD coupling and running gluon mass in the IR
- **QED S and P Coulomb phases infinite -- difference of phases finite!**
- Alternate: Retarded and Advanced Gauge: Augmented LFWFs

Dae Sung Hwang, Yuri V. Kovchegov, Ivan Schmidt, Matthew D. Sievert, sjb





Mulders, Boer Qiu, Sterman Pasquini, Xiao, Yuan, sjb



Hoyer, Marchal, Peigne, Sannino, sjb

QCD Mechanism for Rapidity Gaps







Schmidt, Yang; sjb

Reggeon Contribution to DDIS Constructive Interference! Phase from

signature factor

Nuclear Antishadowing not universal!

Shadowing and Antishadowing in Lepton-Nucleus Scattering

Shadowing: Destructive Interference
 of Two-Step and One-Step Processes
 Pomeron Exchange

• Antishadowing: Constructive Interference of Two-Step and One-Step Processes! Reggeon and Odderon Exchange

 Antishadowing is Not Universal!
 Electromagnetic and weak currents: different nuclear effects !
 Potentially significant for NuTeV Anomaly} Jian-Jun Yang Ivan Schmidt Hung Jung Lu sjb





Two-Step Process in the q+=0 Parton Model Frame Illustrates the LF time sequence



One-Step / Two-Step Interference

Study Double Virtual Compton Scattering $\gamma^* A \to \gamma^* A$

Cannot reduce to real phase matrix element of local operator! No Sum Rules!

OPE matrix elements & LFWFs are real for stable hadrons, nuclei

$$\begin{array}{c} \text{Light-Front QCD} \\ \mathcal{L}_{QCD} \longrightarrow H_{QCD} \\ \downarrow \\ (H_{LF}^{0} + H_{LF}^{I}) |\Psi \rangle = M^{2} |\Psi \rangle \\ \downarrow \\ [\vec{k}_{\perp}^{2} + m^{2} + V_{\text{eff}}^{LF}] \psi_{LF}(x, \vec{k}_{\perp}) = M^{2} \psi_{LF}(x, \vec{k}_{\perp}) \\ \hline \\ [-\frac{d^{2}}{d\zeta^{2}} + \frac{1 - 4L^{2}}{4\zeta^{2}} + U(\zeta)] \psi(\zeta) = \mathcal{M}^{2} \psi(\zeta) \\ \hline \\ \text{AdS/QCD:} \\ \hline \\ U(\zeta) = \kappa^{4} \zeta^{2} + 2\kappa^{2} (L + S - 1) \end{array}$$

Semiclassical first approximation to QCD

Fixed $\tau = t + z/c$



Coupled Fock states

Elímínate hígher Fock states and retarded interactions

Effective two-particle equation

Azímuthal Basís
$$\zeta, \phi$$

Single variable Equation $m_q = 0$

Confining AdS/QCD potential!

Sums an infinite # diagrams

de Tèramond, Dosch, sjb

Ads/QCD Soft-Wall Model

 $e^{\varphi(z)} = e^{+\kappa^2 z^2}$





$$\left[-\frac{d^2}{d\zeta^2} - \frac{1 - 4L^2}{4\zeta^2} + U(\zeta)\right]\psi(\zeta) = M^2\psi(\zeta)$$



Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

Single variable (

Confinement scale:

Unique **Confinement Potential!**

Conformal Symmetry of the action

$\kappa \simeq 0.5 \ GeV$

de Alfaro, Fubini, Furlan: Fubini, Rabinovici:

Scale can appear in Hamiltonian and EQM without affecting conformal invariance of action!

GeV units external to QCD: Only Ratios of Masses Determined

Maldacena





 \bullet Isomorphism of SO(4,2) of conformal QCD with the group of isometries of AdS space

$$s^2 = \frac{R^2}{z^2} (\eta_{\mu\nu} dx^{\mu} dx^{\nu} - dz^2),$$

 $x^{\mu} \rightarrow \lambda x^{\mu}, \ z \rightarrow \lambda z$, maps scale transformations into the holographic coordinate z.

• AdS mode in z is the extension of the hadron wf into the fifth dimension.

d

• Different values of z correspond to different scales at which the hadron is examined.

$$x^2 \to \lambda^2 x^2, \quad z \to \lambda z.$$

 $x^2 = x_\mu x^\mu$: invariant separation between quarks

• The AdS boundary at $z \to 0$ correspond to the $Q \to \infty$, UV zero separation limit.

AdS/CFT

Dílaton-Modífied Ads

$$ds^{2} = e^{\varphi(z)} \frac{R^{2}}{z^{2}} (\eta_{\mu\nu} x^{\mu} x^{\nu} - dz^{2})$$



- Soft-wall dilaton profile breaks conformal invariance $e^{\varphi(z)} = e^{+\kappa^2 z^2}$
- Color Confinement in z
- •Introduces confinement scale к
- Uses AdS₅ as template for conformal theory

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Color Confinement, Hadron Dynamics, and Hadron Spectroscopy from Light-Front Holography and Superconformal Algebra



$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$

Ads Soft-Wall Schrödinger Equation for bound state of two scalar constituents:

$$\left[-\frac{d^2}{dz^2} - \frac{1 - 4L^2}{4z^2} + U(z)\right]\Phi(z) = \mathcal{M}^2\Phi(z)$$

$$U(z) = \kappa^4 z^2 + 2\kappa^2 (L + S - 1)$$

Derived from variation of Action for Dilaton-Modified AdS₅ **Identical to Single-Variable Light-Front Bound State Equation in** ζ !



Light-Front Holography: Unique mapping derived from equality of LF and AdS formula for EM and gravitational current matrix elements and identical equations of motion

Meson Spectrum in Soft Wall Model

$$m_{\pi} = 0$$
 if $m_q = 0$

Pion: Negative term for J=0 cancels positive terms from LFKE and potential

Massless pion!

- Effective potential: $U(\zeta^2) = \kappa^4 \zeta^2 + 2\kappa^2 (J-1)$
- LF WE

$$\left(-rac{d^2}{d\zeta^2}-rac{1-4L^2}{4\zeta^2}+\kappa^4\zeta^2+2\kappa^2(J-1)
ight)\phi_J(\zeta)=M^2\phi_J(\zeta)$$

• Normalized eigenfunctions $\ \langle \phi | \phi
angle = \int d\zeta \, \phi^2(z)^2 = 1$

$$\phi_{n,L}(\zeta) = \kappa^{1+L} \sqrt{rac{2n!}{(n+L)!}} \, \zeta^{1/2+L} e^{-\kappa^2 \zeta^2/2} L_n^L(\kappa^2 \zeta^2)$$

Eigenvalues

$$\mathcal{M}_{n,J,L}^2 = 4\kappa^2\left(n+rac{J+L}{2}
ight)$$

$$\vec{\zeta}^2 = \vec{b}_\perp^2 x (1-x)$$

G. de Teramond, H. G. Dosch, sjb



Light meson orbital (a) and radial (b) spectrum for $\kappa=0.6$ GeV.





Effective mass from $m(p^2)$

Tandy, Roberts, et al

Prediction from AdS/QCD: Meson LFWF



• Light Front Wavefunctions: $\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$ off-shell in P^- and invariant mass $\mathcal{M}^2_{q\bar{q}}$



Boost-invariant LFWF connects confined quarks and gluons to hadrons

week ending 24 AUGUST 2012



AdS/QCD Holographic Wave Function for the ρ Meson and Diffractive ρ Meson Electroproduction

Light-Front Perturbation Theory for pQCD

$$T = H_I + H_I \frac{1}{\mathcal{M}_{initial}^2 - \mathcal{M}_{intermediate}^2 + i\epsilon} H_I + cdots$$

- "History": Compute any subgraph only once since the LFPth numerator does not depend on the process — only the denominator changes!
- Wick Theorem applies, but few amplitudes since all $k^+ > 0$.
- J_z Conservation at every vertex
- Unitarity is explicit
- Loop Integrals are 3-dimensional
- hadronization: coalesce comoving quarks and gluons to hadrons using light-front wavefunctions $\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$

$$\sum_{initial} S^{z} - \sum_{final} S_{z} \mid \leq n$$
 at order g^{r}
$$\int_{0}^{1} dx \int d^{2}k_{\perp}$$
 K. Chiu, sjb

Connection to the Linear Instant-Form Potential





Harmonic Oscillator $U(\zeta) = \kappa^4 \zeta^2$ LF Potential for relativistic light quarks

A.P. Trawinski, S.D. Glazek, H. D. Dosch, G. de Teramond, sjb





Remarkable Features of Líght-Front Schrödínger Equation

Dynamics + Spectroscopy!

- Relativistic, frame-independent
- •QCD scale appears unique LF potential
- Reproduces spectroscopy and dynamics of light-quark hadrons with one parameter
- Zero-mass pion for zero mass quarks!
- Regge slope same for n and L -- not usual HO
- Splitting in L persists to high mass -- contradicts conventional wisdom based on breakdown of chiral symmetry
- Phenomenology: LFWFs, Form factors, electroproduction
- Extension to heavy quarks

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

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Color Confinement, Hadron Dynamics, and Hadron Spectroscopy from Light-Front Holography and Superconformal Algebra


QCD Lagrangian

$$\mathcal{L}_{QCD} = -\frac{1}{4} Tr(G^{\mu\nu}G_{\mu\nu}) + \sum_{f=1}^{n_f} i\bar{\Psi}_f D_{\mu}\gamma^{\mu}\Psi_f + \sum_{f=1}^{n_f} i_f\bar{\Psi}_f\Psi_f$$

$$iD^{\mu} = i\partial^{\mu} - gA^{\mu} \qquad G^{\mu\nu} = \partial^{\mu}A^{\mu} - \partial^{\nu}A^{\mu} - g[A^{\mu}, A^{\nu}]$$

Classical Chiral Lagrangian is Conformally Invariant

Where does the QCD Mass Scale come from?

QCD does not know what MeV units mean! Only Ratios of Masses Determined

ode Alfaro, Fubini, Furlan:

Scale can appear in Hamiltonian and EQM without affecting conformal invariance of action!

Unique confinement potential!

• de Alfaro, Fubini, Furlan (dAFF)

$$\begin{aligned} G|\psi(\tau) > &= i\frac{\partial}{\partial\tau}|\psi(\tau) > \\ G &= uH + vD + wK \\ G &= H_{\tau} = \frac{1}{2} \left(-\frac{d^2}{dx^2} + \frac{g}{x^2} + \frac{4uw - v^2}{4}x^2 \right) \end{aligned}$$

Retains conformal invariance of action despite mass scale! $4uw-v^2=\kappa^4=[M]^4$

Identical to LF Hamiltonian with unique potential and dilaton!

$$\begin{bmatrix} -\frac{d^2}{d\zeta^2} + \frac{1-4L^2}{4\zeta^2} + U(\zeta) \end{bmatrix} \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$
$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L+S-1)$$

Dosch, de Teramond, sjb



- Identify with difference of LF time $\Delta x^+/P^+$ between constituents
- Finite range
- Measure in Double-Parton Processes

Retains conformal invariance of action despite mass scale!

Haag, Lopuszanski, Sohnius (1974)

Superconformal Quantum Mechanics $\{\psi,\psi^+\} = 1$ $B = \frac{1}{2}[\psi^+,\psi] = \frac{1}{2}\sigma_3$ $\psi = \frac{1}{2}(\sigma_1 - i\sigma_2), \quad \psi^+ = \frac{1}{2}(\sigma_1 + i\sigma_2)$ $Q = \psi^{+}[-\partial_{x} + \frac{f}{x}], \quad Q^{+} = \psi[\partial_{x} + \frac{f}{x}], \quad S = \psi^{+}x, \quad S^{+} = \psi x$ $\{Q, Q^+\} = 2H, \{S, S^+\} = 2K$ $\{Q, S^+\} = f - B + 2iD, \ \{Q^+, S\} = f - B - 2iD$ generates conformal algebra [H,D] = i H, [H, K] = 2 i D, [K, D] = - i K $Q \simeq \sqrt{H}, S \simeq \sqrt{K}$

Superconformal Quantum Mechanics

Baryon Equation $Q \simeq \sqrt{H}, S \simeq \sqrt{K}$

Consider
$$R_w = Q + wS;$$

w: dimensions of mass squared

$$G = \{R_w, R_w^+\} = 2H + 2w^2K + 2wfI - 2wB \qquad 2B = \sigma_3$$

Retains Conformal Invariance of Action

Fubini and Rabinovici

New Extended Hamíltonían G ís díagonal:

$$G_{11} = \left(-\partial_x^2 + w^2 x^2 + 2wf - w + \frac{4(f + \frac{1}{2})^2 - 1}{4x^2}\right)$$

$$G_{22} = \left(-\partial_x^2 + w^2 x^2 + 2wf + w + \frac{4(f - \frac{1}{2})^2 - 1}{4x^2}\right)$$

Identify $f - \frac{1}{2} = L_B$, $w = \kappa^2$
Firmulas of C : $M^2(n, I) = 4w^2(n + I + 1)$

Eigenvalue of G: $M^2(n, L) = 4\kappa^2(n + L_B + 1)$

LF Holography

Baryon Equation

Superconformal Quantum Mechanics

$$\begin{pmatrix} -\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}(L_{B} + 1) + \frac{4L_{B}^{2} - 1}{4\zeta^{2}} \end{pmatrix} \psi_{J}^{+} = M^{2}\psi_{J}^{+} \\ \begin{pmatrix} -\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}L_{B} + \frac{4(L_{B} + 1)^{2} - 1}{4\zeta^{2}} \end{pmatrix} \psi_{J}^{-} = M^{2}\psi_{J}^{-} \\ M^{2}(n, L_{B}) = 4\kappa^{2}(n + L_{B} + 1) \qquad \text{S=1/2, P=+} \\ Meson Equation \qquad \lambda = \kappa^{2} \\ \begin{pmatrix} -\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}(J - 1) + \frac{4L_{M}^{2} - 1}{4\zeta^{2}} \end{pmatrix} \phi_{J} = M^{2}\phi_{J} \\ M^{2}(n, L_{M}) = 4\kappa^{2}(n + L_{M}) \qquad \text{S=0, P=+} \\ M^{2}(n, L_{M}) = 4\kappa^{2}(n + L_{M}) \qquad \text{S=0, P=+} \\ \text{S=0, I=I Meson is superpartner of S=1/2, I=I Baryon} \end{cases}$$

Meson-Baryon Degeneracy for L_M=L_B+1









de Tèramond, Dosch, Lorce', sjb





Fit to the slope of Regge trajectories, including radial excitations

Same Regge Slope for Meson, Baryons: Supersymmetric feature of hadron physics

Fermionic Modes and Baryon Spectrum

[Hard wall model: GdT and S. J. Brodsky, PRL **94**, 201601 (2005)] [Soft wall model: GdT and S. J. Brodsky, (2005), arXiv:1001.5193]



From Nick Evans

• Nucleon LF modes

$$\psi_{+}(\zeta)_{n,L} = \kappa^{2+L} \sqrt{\frac{2n!}{(n+L)!}} \zeta^{3/2+L} e^{-\kappa^{2}\zeta^{2}/2} L_{n}^{L+1} \left(\kappa^{2}\zeta^{2}\right)$$
$$\psi_{-}(\zeta)_{n,L} = \kappa^{3+L} \frac{1}{\sqrt{n+L+2}} \sqrt{\frac{2n!}{(n+L)!}} \zeta^{5/2+L} e^{-\kappa^{2}\zeta^{2}/2} L_{n}^{L+2} \left(\kappa^{2}\zeta^{2}\right)$$

• Normalization

$$\int_0^\infty d\zeta \int_0^1 dx \psi_+^2(\zeta^2, x) = \int_0^\infty d\zeta \int_0^1 dx \psi_-^2(\zeta^2, x) = \frac{1}{2}$$

Quark Chíral Symmetry of Eígenstate!

• Eigenvalues

$$\mathcal{M}_{n,L,S=1/2}^2 = 4\kappa^2 \left(n + L + 1 \right)$$

• "Chiral partners"

$$\frac{\mathcal{M}_{N(1535)}}{\mathcal{M}_{N(940)}} = \sqrt{2}$$

Nucleon: Equal Probability for L=0, I

• Compute Dirac proton form factor using SU(6) flavor symmetry

$$F_1^p(Q^2) = R^4 \int \frac{dz}{z^4} V(Q, z) \Psi_+^2(z)$$

• Nucleon AdS wave function

$$\Psi_{+}(z) = \frac{\kappa^{2+L}}{R^2} \sqrt{\frac{2n!}{(n+L)!}} z^{7/2+L} L_n^{L+1} \left(\kappa^2 z^2\right) e^{-\kappa^2 z^2/2}$$

• Normalization $(F_1^p(0) = 1, V(Q = 0, z) = 1)$

$$R^4 \int \frac{dz}{z^4} \, \Psi_+^2(z) = 1$$

• Bulk-to-boundary propagator [Grigoryan and Radyushkin (2007)]

$$V(Q,z) = \kappa^2 z^2 \int_0^1 \frac{dx}{(1-x)^2} x^{\frac{Q^2}{4\kappa^2}} e^{-\kappa^2 z^2 x/(1-x)}$$

• Find

$$F_1^p(Q^2) = \frac{1}{\left(1 + \frac{Q^2}{M_{\rho}^2}\right) \left(1 + \frac{Q^2}{M_{\rho'}^2}\right)}$$

with $\mathcal{M}_{\rho_n}^2 \to 4\kappa^2(n+1/2)$



Space-Like Dirac Proton Form Factor

• Consider the spin non-flip form factors

$$F_{+}(Q^{2}) = g_{+} \int d\zeta J(Q,\zeta) |\psi_{+}(\zeta)|^{2},$$

$$F_{-}(Q^{2}) = g_{-} \int d\zeta J(Q,\zeta) |\psi_{-}(\zeta)|^{2},$$

where the effective charges g_+ and g_- are determined from the spin-flavor structure of the theory.

- Choose the struck quark to have $S^z = +1/2$. The two AdS solutions $\psi_+(\zeta)$ and $\psi_-(\zeta)$ correspond to nucleons with $J^z = +1/2$ and -1/2.
- For SU(6) spin-flavor symmetry

$$F_1^p(Q^2) = \int d\zeta J(Q,\zeta) |\psi_+(\zeta)|^2,$$

$$F_1^n(Q^2) = -\frac{1}{3} \int d\zeta J(Q,\zeta) \left[|\psi_+(\zeta)|^2 - |\psi_-(\zeta)|^2 \right],$$

where $F_1^p(0) = 1$, $F_1^n(0) = 0$.



Using SU(6) flavor symmetry and normalization to static quantities



Superconformal Algebra

2X2 Hadronic Multiplets

Bosons, Fermions with Equal Mass!



Superconformal Algebra 4-Plet



New Organization of the Hadron Spectrum

	Meson (P(C) N		Baryon			Tetraquark					
	q-cont	JF(0)	Name	q-cont	Jr	Name	q-cont	$J^{r(0)}$	Name		
	qq	0-+	$\pi(140)$						(800)		
	$\bar{q}q$	1+-	$h_1(1170)$	[ud]q	$(1/2)^+$	N(940)	[ud][ūd]	0++	$\sigma(500)$		
	qq	2-+	$\eta_2(1645)$	[ud]q	$(3/2)^{-}$	$N_{\frac{3}{2}}$ (1520)	[ud][ud]	1-+			
	$\bar{q}q$	1	$\rho(770), \omega(780)$	_	_		—			L	
$(\)$	$\bar{q}q$	2++	$a_2(1320), f_2(1270)$	(qq)q	$(3/2)^+$	$\Delta(1232)$	$(qq)[\bar{u}d]$	1++	$a_1(1260)$	\square	
	qq	3	$\rho_3(1690), \ \omega_3(1670)$	(qq)q	$(3/2)^{-}$	$\Delta_{\frac{1}{2}}(1700)$	(qq)[ud]	1-+	$\pi_1(1600)$		
	$\bar{q}q$	4++	$a_4(2040), f_4(2050)$	(qq)q	$(7/2)^+$	$\Delta_{\frac{7}{2}+}(1950)$	$(qq)[\bar{u}\bar{d}]$				
	$\bar{q}s$	0-	K(495)	_	_	_	_		_		
	$\bar{q}s$	1+	$\bar{K}_{1}(1270)$	[ud]s	$(1/2)^+$	Λ(1115)	$[ud][\bar{s}\bar{q}]$	0+	$K_0^*(1430)$		
	$\bar{q}s$	2-	$K_2(1770)$	[ud]s	$(3/2)^{-}$	$\Lambda(1520)$	$[ud][\bar{s}\bar{q}]$	1-			
	$\bar{s}q$	0-	K(495)	_		_	_		_		
	$\bar{s}q$	1+	$K_1(1270)$	[sq]q	$(1/2)^+$	$\Sigma(1190)$	$[sq][\bar{s}\bar{q}]$	0++	$a_0(980)$		
									$f_0(980)$		
	$\bar{s}q$	1-	$K^{*}(890)$	_	_	_	—	_	_	L	
\Box	āq	2+	$K_{2}^{*}(1430)$	(sq)q	$(3/2)^+$	$\Sigma(1385)$	$(sq)[\bar{u}d]$	1+	$K_1(1400)$	D	
	$\bar{s}q$	3-	$K_{3}^{*}(1780)$	(sq)q	$(3/2)^{-}$	$\Sigma(1670)$	$(sq)[\bar{u}d]$	2-	$K_2(1820)$		
	$\bar{s}q$	4+	$K_{4}^{*}(2045)$	(sq)q	$(7/2)^+$	$\Sigma(2030)$	$(sq)[\bar{u}d]$		_		
	88	0-+	$\eta'(958)$				—		_		
(88	1+-	$h_1(1380)$	[sq]s	$(1/2)^+$	$\Xi(1320)$	$[sq][\bar{s}\bar{q}]$	0++	$f_0(1370)$	\bigcirc	
									$a_0(1450)$		
	88	2-+	$\eta_2(1870)$	sq s	$(3/2)^{-}$	$\Xi(1620)$	sq sq	1-+			
	88	1	$\Phi(1020)$	_			_		_		
	88	2^{++}	$f'_{2}(1525)$	(sq)s	$(3/2)^+$	$\Xi^{*}(1530)$	$(sq)[\bar{s}\bar{q}]$	1++	$f_1(1420)$		
									$a_1(1420)$		
	88	3	$\Phi_{3}(1850)$	(sq)s	$(3/2)^{-}$	$\Xi(1820)$	$(sq)[\bar{s}\bar{q}]$		_		
	88	2++	$f_2(1640)$	(ss)s	$(3/2)^+$	$\Omega(1672)$	$(ss)[\bar{s}\bar{q}]$	1+	$K_1(1650)$		
	Meson				Ranvon 7			otroquark			
			7 1 1	Daryon I			eu aqual K				

M. Níelsen, sjb

Dosch, de Teramond, sjb

Supersymmetry across the light and heavy-light spectrum



Dosch, de Teramond, sjb

Supersymmetry across the light and heavy-light spectrum



Heavy charm quark mass does not break supersymmetry

Superpartners for states with one c quark

	Me	eson		Bary	yon	Tetraquark		
q-cont	$J^{P(C)}$	Name	$q ext{-cont}$	J^P	Name	q-cont	$J^{P(C)}$	Name
$\bar{q}c$	0-	D(1870)						
$\bar{q}c$	1+	$D_1(2420)$	[ud]c	$(1/2)^+$	$\Lambda_c(2290)$	$[ud][\bar{c}\bar{q}]$	0^+	$\bar{D}_{0}^{*}(2400)$
$\bar{q}c$	2^{-}	$D_J(2600)$	[ud]c	$(3/2)^{-}$	$\Lambda_c(2625)$	$[ud][\bar{c}\bar{q}]$	1-	
$\bar{c}q$	0-	$\bar{D}(1870)$						
$\bar{c}q$	1+	$O_1(2420)$	[cq]q	$(1/2)^+$	$\Sigma_c(2455)$	$[cq][\bar{u}\bar{d}]$	0^{+}	$D_0^*(2400)$
$\bar{q}c$	1-	$D^{*}(2010)$			_ \			
$\bar{q}c$	2^{+}	$D_2^*(2460)$	(qq)c	$(3/2)^+$	$\Sigma_{c}^{*}(2520)$	$(qq)[\bar{c}\bar{q}]$	1+	D(2550)
$\bar{q}c$	3^{-}	$D_3^*(2750)$	(qq)c	$(3/2)^{-}$	$\Sigma_{c}(2800)$	$(qq)[\bar{c}\bar{q}]$		
$\bar{s}c$	0^{-}	$D_s(1968)$			_			
$\bar{s}c$	1+	$D_{s1}(2460)$	[qs]c	$(1/2)^+$	$\Xi_c(2470)$	$[qs][\bar{c}\bar{q}]$	0^+	$\bar{D}_{s0}^{*}(2317)$
$\bar{s}c$	2^{-}	$Q_{s2}(\sim 2860)?$	[qs]c	$(3/2)^{-}$	$\Xi_c(2815)$	$[sq][\bar{c}\bar{q}]$	1-	
$\bar{s}c$	1-	$D_{s}^{*}(2110)$	$\backslash -$					
$\bar{s}c$	2^{+}	$D_{s2}^{*}(2573)$	(sq)c	$(3/2)^+$	$\Xi_{c}^{*}(2645)$	$(sq)[\bar{c}\bar{q}]$	1+	$D_{s1}(2536)$
$\bar{c}s$	1+	$Q_{s1}(\sim 2700)?$	[cs]s	$(1/2)^+$	$\Omega_c(2695)$	$[cs][\bar{s}\bar{q}]$	0^{+}	??
$\bar{s}c$	2^{+}	$D_{s2}^* (\sim 2750)?$	(ss)c	$(3/2)^+$	$\Omega_c(2770)$	$(ss)[\bar{c}\bar{s}]$	1+	??
M. Níelsen, sjb				pr	edictions	beautiful agreement!		

Dosch, de Teramond, sjb

Supersymmetry across the light and heavy-light spectrum



Heavy bottom quark mass does not break supersymmetry



Channel

Regge slope for heavy-light mesons, baryons: increases with heavy quark mass



SELEX (3520 ± 1 MeV) $J^P = \frac{1}{2}^- |[cd]c >$ Two decay channels: $\Xi_{cc}^+ \to \Lambda_c^+ K^- \pi^+, pD^+ K^-$

SELEX Collaboration / Physics Letters B 628 (2005) 18-24



 $\Xi_{cc}^+ \rightarrow pD^+K^-$ mass distribution from Fig. 2(a) with high-statistics measurement of random combinatoric background computed from event-mixing.

Gaussian fits for $\Xi_{cc}^+ \to \Lambda_c^+ K^- \pi^+$ and $\Xi_{cc}^+ \to pD^+K^-$ (shaded data) on same plot.

Underlying Principles

- Poincarè Invariance: Independent of the observer's Lorentz frame
- Quantization at Fixed Light-Front Time ~ au
- Causality: Information within causal horizon
- Light-Front Holography: $AdS_5 = LF(3+1)$

 $z \leftrightarrow \zeta$ where $\zeta^2 = b_{\perp}^2 x(1-x)$



- Single fundamental hadronic mass scale κ: but retains the Conformal Invariance of the Action (dAFF)!
- Unique color-confining LF Potential! $U(\zeta^2) = \kappa^4 \zeta^2$
- Superconformal Algebra: Mass Degenerate 4-Plet:

Meson $q\bar{q} \leftrightarrow \text{Baryon } q[qq] \leftrightarrow \text{Tetraquark } [qq][\bar{q}\bar{q}]$

Partonic Correlations Yeravan September 25, 2018

Color Confinement, Hadron Dynamics, and Hadron Spectroscopy from Light-Front Holography and Superconformal Algebra



Universal Hadronic Decomposition

$$\frac{\mathcal{M}_{H}^{2}}{\kappa^{2}} = (1 + 2n + L) + (1 + 2n + L) + (2L + 4S + 2B - 2)$$
• Universal quark light-front kinetic energy
Equal:
Virial
Virial
Heorem
• Universal quark light-front potential energy
$$\Delta \mathcal{M}_{LFFE}^{2} = \kappa^{2}(1 + 2n + L)$$
• Universal quark light-front potential energy
$$\Delta \mathcal{M}_{LFPE}^{2} = \kappa^{2}(1 + 2n + L)$$
• Universal Constant Contribution from AdS
and Superconformal Quantum Mechanics
$$\Delta \mathcal{M}_{spin}^{2} = 2\kappa^{2}(L + 2S + B - 1)$$

hyperfine spin-spin



Use counting rules to identify composite structure

Lebed, sjb

Running Coupling from Modified AdS/QCD Deur, de Teramond, sjb

Consider five-dim gauge fields propagating in AdS₅ space in dilaton background $arphi(z) = \kappa^2 z^2$

$$e^{\phi(z)} = e^{+\kappa^2 z^2} \qquad S = -\frac{1}{4} \int d^4 x \, dz \, \sqrt{g} \, e^{\varphi(z)} \, \frac{1}{g_5^2} \, G^2$$

Flow equation

$$\frac{1}{g_5^2(z)} = e^{\varphi(z)} \frac{1}{g_5^2(0)} \quad \text{or} \quad g_5^2(z) = e^{-\kappa^2 z^2} g_5^2(0)$$

where the coupling $g_5(z)$ incorporates the non-conformal dynamics of confinement

- YM coupling $\alpha_s(\zeta) = g_{YM}^2(\zeta)/4\pi$ is the five dim coupling up to a factor: $g_5(z) \to g_{YM}(\zeta)$
- $\bullet\,$ Coupling measured at momentum scale Q

$$\alpha_s^{AdS}(Q) \sim \int_0^\infty \zeta d\zeta J_0(\zeta Q) \,\alpha_s^{AdS}(\zeta)$$

Solution

$$\alpha_s^{AdS}(Q^2) = \alpha_s^{AdS}(0) e^{-Q^2/4\kappa^2}.$$

where the coupling α_s^{AdS} incorporates the non-conformal dynamics of confinement

Bjorken sum rule defines effective charge
$$\alpha_{g1}(Q^2)$$
$$\int_0^1 dx [g_1^{ep}(x,Q^2) - g_1^{en}(x,Q^2)] \equiv \frac{g_a}{6} [1 - \frac{\alpha_{g1}(Q^2)}{\pi}]$$

- Can be used as standard QCD coupling
- Well measured
- Asymptotic freedom at large Q²
- Computable at large Q² in any pQCD scheme
- Universal β_0 , β_1



Analytic, defined at all scales, IR Fixed Point

AdS/QCD dilaton captures the higher twist corrections to effective charges for Q < 1 GeV

$$e^{\varphi} = e^{+\kappa^2 z^2}$$

Deur, de Teramond, sjb





Process-independent strong running coupling

Daniele Binosi,¹ Cédric Mezrag,² Joannis Papavassiliou,³ Craig D. Roberts,² and Jose Rodríguez-Quintero⁴



Universality of Generalized Parton Distributions in Light-Front Holographic QCD Guy F. de Téramond,¹ Tianbo Liu,^{2,3} Raza Sabbir Sufian,² Hans Günter Dosch,⁴ Stanley J. Brodsky,⁵ and Alexandre Deur²



Nonperturbative strange-quark sea from lattice QCD, light-front holography, and meson-baryon fluctuation models

Raza Sabbir Sufian,¹ Tianbo Liu,^{1,2,*} Guy F. de Téramond,³ Hans Günter Dosch,⁴ Stanley J. Brodsky,⁵ Alexandre Deur,¹ Mohammad T. Islam,⁶ and Bo-Qiang Ma^{7,8,9}

Nuclear physics in soft-wall AdS/QCD: Deuteron electromagnetic form factors

Thomas Gutsche, Valery E. Lyubovitskij, Ivan Schmidt, Alfredo Vega

We present a high-quality description of the deuteron electromagnetic form factors in a soft-wall AdS/QCD approach. We first propose an effective action describing the dynamics of the deuteron in the presence of an external vector field. Based on this action the deuteron electromagnetic form factors are calculated, displaying the correct 1/Q¹⁰ power scaling for large Q² values. This finding is consistent with quark counting rules and the earlier observation that this result holds in confining gauge/gravity duals. The Q² dependence of the deuteron form factors is defined by a single and universal scale parameter K, which is fixed from data.

arXiv:1501.02738 [hep-ph]
Poincarè Invariance: Independent of the observer's Lorentz frame: Quantization at Fixed Light-Front Time τ

Underlying Principles

- Causality: Information within causal horizon: Light-Front
- Light-Front Holography: AdS₅ = LF (3+1) $z \leftrightarrow \zeta$ where $\zeta^2 = b_{\perp}^2 x(1-x)$



- Introduce Mass Scale κ while retaining the Conformal Invariance of the Action (dAFF)
- Unique Dilaton in AdS₅: $e^{+\kappa^2 z^2}$
- Unique color-confining LF Potential $~U(\zeta^2)=\kappa^4\zeta^2$
- Superconformal Algebra: Mass Degenerate 4-Plet:

Meson $q\bar{q} \leftrightarrow \text{Baryon } q[qq] \leftrightarrow \text{Tetraquark } [qq][\bar{q}\bar{q}]$

Partonic Correlations Yeravan September 25, 2018



Features of LF Holographic QCD

- Color Confinement, Analytic form of confinement potential
- Massless pion bound state in chiral limit
- QCD coupling at all scales
- Connection of perturbative and nonperturbative mass scales
- Poincare' Invariant
- Hadron Spectroscopy-Regge Trajectories with universal slopes in n, L
- Supersymmetric 4-Plet: Meson-Baryon Tetraquark Symmetry
- Light-Front Wavefunctions
- Form Factors, Structure Functions, Hadronic Observables
- •OPE: Constituent Counting Rules
- Hadronization at the Amplitude Level
- Analytic First Approximation to QCD

Many phenomenological tests

• Systematically improvable: Basis LF Quantization (BLFQ)

Partonic Correlations Yeravan September 25, 2018



Invariance Principles of Quantum Field Theory

- Polncarè Invariance: Physical predictions must be independent of the observer's Lorentz frame: Front Form
- Causality: Information within causal horizon: Front Form
- Gauge Invariance: Physical predictions of gauge theories must be independent of the choice of gauge
- Scheme-Independence: Physical predictions of renormalizable theories must be independent of the choice of the renormalization scheme — Principle of Maximum Conformality (PMC)

 Mass-Scale Invariance: Conformal Invariance of the Action (DAFF)

Partonic Correlations Yeravan September 25, 2018



The Renormalization Scale Ambiguity for Top-Pair Production Eliminated Using the 'Principle of Maximum Conformality' (PMC)



BLM/PMC: Scheme-Independent, same as Gell-Mann-Low in pQED

Top quark forward-backward asymmetry predicted by pQCD NNLO within 1 σ of CDF/D0 measurements using PMC/BLM scale setting

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"One of the gravest puzzles of theoretical physics"

DARK ENERGY AND THE COSMOLOGICAL CONSTANT PARADOX

A. ZEE

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$$(\Omega_{\Lambda})_{QCD} \sim 10^{45}$$

$$(\Omega_{\Lambda})_{EW} \sim 10^{56}$$

$$\Omega_{\Lambda} = 0.76(expt)$$

Extraordinary conflict between the conventional definition of the vacuum in quantum field theory and cosmology

Elements of the solution: (A) Light-Front Quantization: causal, frame-independent vacuum (B) New understanding of QCD "Condensates" (C) Higgs Light-Front Zero Mode Two Definitions of Vacuum State

Instant Form: Lowest Energy Eigenstate of Instant-Form Hamiltonian

 $H|\psi_0>=E_0|\psi_0>, E_0=\min\{E_i\}$

Eigenstate defined at one time t over all space; Acausal! Frame-Dependent

Front Form: Lowest Invariant Mass Eigenstate of Light-Front Hamiltonian

$$H_{LF}|\psi_0\rangle_{LF} = M_0^2|\psi_0\rangle_{LF}, M_0^2 = 0.$$

Frame-independent eigenstate at fixed LF time τ = t+z/c within causal horizon

Frame-independent description of the causal physical universe!



Light-Front vacuum can símulate empty universe

Shrock, Tandy, Roberts, sjb

- Independent of observer frame
- Causal
- Lowest invariant mass state M= o.
- Trivial up to k+=o zero modes-- already normal-ordering
- Higgs theory consistent with trivial LF vacuum (Srivastava, sjb)
- QCD and AdS/QCD: "In-hadron"condensates (Maris, Tandy Roberts) -- GMOR satisfied.
- QED vacuum; no loops
- Zero cosmological constant from QED, QCD, EW

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Light-Front Pion Valence Wavefunctions

 $S_{\bar{u}}^z + S_d^z = +1/2 - 1/2 = 0$



Angular Momentum Conservation

$$J^z = \sum_i^n S_i^z + \sum_i^{n-1} L_i^z$$

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Maris, Roberts, Shrock, Tandy, sjb

Ward-Takahashí Identíty for axíal current

GMOR satisfied, no VEV

 $P^{\mu}\Gamma_{5\mu}(k,P) + 2im\Gamma_5(k,P) = S^{-1}(k+P/2)i\gamma_5 + i\gamma_5 S^{-1}(k-P/2)$

$$S^{-1}(\ell) = i\gamma \cdot \ell A(\ell^2) + B(\ell^2) \qquad m(\ell^2) = \frac{B(\ell^2)}{A(\ell^2)}$$



Identify pion pole at $P^2 = m_\pi^2$

$$P^{\mu} < 0 |\bar{q}\gamma_{5}\gamma^{\mu}q|\pi > = 2m < 0 |\bar{q}i\gamma_{5}q|\pi >$$
$$f_{\pi}m_{\pi}^{2} = -(m_{u} + m_{d})\rho_{\pi}$$

Revised Gell Mann-Oakes-Renner Formula in QCD

$$\begin{split} m_{\pi}^2 &= -\frac{(m_u + m_d)}{f_{\pi}^2} < 0 |\bar{q}q| 0 > & \text{current algebra:} \\ m_{\pi}^2 &= -\frac{(m_u + m_d)}{f_{\pi}} < 0 |i\bar{q}\gamma_5 q| \pi > & \text{QCD: composite pion} \\ & \text{Bethe-Salpeter Eq.} \\ & \text{No-VEV!} \end{split}$$

vacuum condensate actually is an "in-hadron condensate"



Maris, Roberts, Tandy

PHYSICAL REVIEW C 82, 022201(R) (2010)

New perspectives on the quark condensate

Stanley J. Brodsky,^{1,2} Craig D. Roberts,^{3,4} Robert Shrock,⁵ and Peter C. Tandy⁶ ¹SLAC National Accelerator Laboratory, Stanford University, Stanford, California 94309, USA ²Centre for Particle Physics Phenomenology: CP³-Origins, University of Southern Denmark, Odense 5230 M, Denmark ³Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA ⁴Department of Physics, Peking University, Beijing 100871, China ⁵C.N. Yang Institute for Theoretical Physics, Stony Brook University, Stony Brook, New York 11794, USA ⁶Center for Nuclear Research, Department of Physics, Kent State University, Kent, Ohio 44242, USA (Received 25 May 2010; published 18 August 2010)

We show that the chiral-limit vacuum quark condensate is qualitatively equivalent to the pseudoscalar meson leptonic decay constant in the sense that they are both obtained as the chiral-limit value of well-defined gauge-invariant hadron-to-vacuum transition amplitudes that possess a spectral representation in terms of the current-quark mass. Thus, whereas it might sometimes be convenient to imagine otherwise, neither is essentially a constant mass-scale that fills all spacetime. This means, in particular, that the quark condensate can be understood as a property of hadrons themselves, which is expressed, for example, in their Bethe-Salpeter or light-front wave functions.

Partonic Correlations Yeravan September 25, 2018



Quark and Gluon condensates reside within hadrons, not vacuum

Casher and Susskind Maris, Roberts, Tandy Shrock and sjb

- Bound-State Dyson Schwinger Equations
- AdS/QCD
- Implications for cosmological constant --Eliminates 45 orders of magnitude conflict

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 $\begin{array}{l} (\Omega_{\Lambda})_{QCD} \sim 10^{45} \\ (\Omega_{\Lambda})_{EW} \sim 10^{56} \end{array} \qquad \Omega_{\Lambda} = 0.76 (expt) \end{array}$

$$(\Omega_{\Lambda})_{QCD} = 0 \qquad (\Omega_{\Lambda})_{EW} = 0$$

Central Question: What is the source of Dark Energy? $\Omega_{\Lambda} = 0.76(expt) \qquad \text{Higgs Zero-Mode Curvature?}$

Advantages of the Dirac's Front Form for Hadron Physics Poincare' Invariant

Physics Independent of Observer's Motion

- \bullet Measurements are made at fixed τ
- Causality is automatic
- Structure Functions are squares of LFWFs
- Form Factors are overlap of LFWFs
- LFWFs are frame-independent: no boosts, no pancakes!

Penrose, Terrell, Weisskopf

- Same structure function measured at an e p collider and the proton rest frame
- No dependence of hadron structure on observer's frame
- Jz Conservation, bounds on ΔLz Chiu, sjb
- LF Holography: Dual to AdS space
- LF Vacuum trivial -- no vacuum condensates!

Roberts, Shrock, Tandy, sjb



Some References for Light-Front Holography

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Color Confinement and Supersymmetric Properties of Hadron Physics from Light-Front Holography



with Guy de Tèramond, Hans Günter Dosch, C. Lorce, K. Chiu, R. S. Sufian, A.Deur, T. Liu Yerevan, Armenia 24-28 September 2018

Correlations in Partonic and Hadronic Interactions 2018

Some Key QCD Issues in Electroproduction

• Intrinsic Heavy Quarks at high x;

$$s(x) \neq \bar{s}(x)$$

- Role of Color Confinement in DIS
- Hadronization at the Amplitude Level
- Leading-Twist Lensing: Sivers Effect
- Diffractive DIS
- Static versus Dynamic Structure Functions
- Origin of Shadowing and Anti-Shadowing
- Is Anti-Shadowing Non-Universal: Flavor Specific?
- Nuclear Correlations and Effects
- $\mathbf{I} < \mathbf{X} < \mathbf{A}$
- Is Momentum Sum Rule Correct?