Future perspectives of TMD phenomenology

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A tribute to the Armenian contribution

TMD theoretical developments 1990s

Kotzinian NPB 441 (1995)

$$\begin{split} &\frac{1-\epsilon}{2Q^2} l_{\mu\nu} W^{\mu\nu} \\ &= \frac{1}{2} \left(H_{11}^{(0)} + H_{22}^{(0)} \right) + \epsilon H_{00}^{(0)} + \sqrt{2\epsilon(1+\epsilon)} Re H_{01}^{(0)} \cos \phi_h^l + \frac{\epsilon}{2} \left(H_{11}^{(0)} - H_{22}^{(0)} \right) \cos 2\phi_h^l \\ &+ \lambda \sqrt{2\epsilon(1-\epsilon)} Im H_{01}^{(0)} \sin \phi_h^l \\ &- S_{(\gamma h)}^1 \left[\sqrt{2\epsilon(1+\epsilon)} Re H_{021}^{(S)} \sin \phi_h^l + \epsilon Re H_{121}^{(S)} \sin 2\phi_h^l \right] \\ &- S_{(\gamma h)}^2 \left[\frac{1}{2} \left(H_{112}^{(S)} + H_{222}^{(S)} \right) + \epsilon H_{002}^{(S)} + \sqrt{2\epsilon(1+\epsilon)} Re H_{012}^{(S)} \cos \phi_h^l + \frac{\epsilon}{2} \left(H_{112}^{(S)} - H_{222}^{(S)} \right) \cos 2\phi_h^l \right] \\ &- S_{(\gamma h)}^3 \left[\sqrt{2\epsilon(1+\epsilon)} Re H_{023}^{(S)} \sin \phi_h^l + \epsilon Re H_{123}^{(S)} \sin 2\phi_h^l \right] \\ &+ \lambda S_{(\gamma h)}^1 \left[-\sqrt{1-\epsilon^2} Im H_{121}^{(S)} + \sqrt{2\epsilon(1-\epsilon)} Im H_{021}^{(S)} \cos \phi_h^l \right] \\ &+ \lambda S_{(\gamma h)}^3 \left[-\sqrt{1-\epsilon^2} Im H_{123}^{(S)} + \sqrt{2\epsilon(1-\epsilon)} Im H_{023}^{(S)} \cos \phi_h^l \right], \end{split}$$

Decomposition of SIDIS cross section

Kotzinian, Mulders PLB 406 (1997)

Concept of weighted asymmetries

 $A_T(x, y, z; |S_T|) \equiv \frac{\int d\phi^l \int d^2 P_{h\perp} \frac{|P_{h\perp}|}{zM_h} \sin(\phi_s^\ell + \phi_h^\ell) \left(d\sigma^\uparrow - d\sigma^\downarrow \right)}{\int d\phi^l \int d^2 P_{h\perp} (d\sigma^\uparrow + d\sigma^\downarrow)}$

HERMES long. asymmetry 2000

Probably the first measurement related to polarized TMDs



FIG. 3. Target-spin analyzing powers in the sin ϕ moment as a function of transverse momentum, for π^+ (squares) and π^- (circles). Error bars show the statistical uncertainties and the band represents the systematic uncertainties.

HERMES Sivers asymmetry 2005

More than 650 citations, the most cited paper by HERMES and the third by HERA!



PHYSICAL REVIEW LETTERS

week ending 14 JANUARY 2005

PRL 94, 012002 (2005)

Single-Spin Asymmetries in Semi-Inclusive Deep-Inelastic Scattering on a Transversely Polarized Hydrogen Target

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(The HERMES Collaboration)

about 10% Armenian

The TMD acronym

My personal recollection is that the acronym was coined by Harut before 2008

I think however that Italians were among the "fist followers" [see D[erek] Sivers https://www.youtube.com/watch?v=fW8amMCVAJQ]

see, e.g., Anselmino, Boglione, D'Alesio, Kotzinian, Murgia, Prokudin,Turk, PRD75 (07) Pasquini, Cazzaniga, Boffi, PRD 78 (08) Bacchetta, Conti, Radici, PRD 78 (08)

Apart from the acronym, Harut was and still is one of the main actors of the TMD program at



Harut recently complained about the ambiguous use of the acronym

To avoid possible troubles, this is what I taught at the TMD Collaboration Summer School in Philadelphia in 2017

About names and acronyms

TMD as an adjective stands for Transverse Momentum Dependent

 TMD as a noun stands for Transverse Momentum Distribution and it is usually meant to encompass both transverse-momentum-dependent PDFs and Fragmentation Functions (FFs)

Younger generations

Theory of Fragmentation Functions

Analysis Coordinator of COMPASS experiment





Back to TMD perspectives...

In the last (great) workshop in Yerevan (2009) I was asked to speak about TMD perspectives

IN FIVE TO TENYEARS

- COMPASS and HERMES and JLab: all structure functions (including unpolarized cross sections), weighted asymmetries, multidimensional binning *K. Hafidi, M. Aghasyan talks*
- BELLE and BABAR unpolarized cross sections and Collins
 (also for etas and kaons)
 I. Garzia's talk
- RHIC: jet-jet, photon-jet, hadron-hadron correlations at forward rapidities
- Fermilab: improvement of unpolarized Drell-Yan

X

X

Polarized Drell-Yan at COMPASS and RHIC
 C. Quintans, L. Bland talks

IN FIVE TO TENYEARS

- Details and subtleties in TMD factorization sorted out
- Understanding of TMD factorization (or lack of it) in hadronhadron collisions to hadrons
- Evolution equations of all TMDs known
- Numerical implementation of evolution functions up and working

IN FIVE TO TENYEARS

- More and improved fits (including evolution, more data, different approaches)
- Parametrization that can describe both TMDs and GPDs
- Coarse and qualitative 3D pictures

X

• Hints about orbital angular momentum

The unpolarized TMD as case study

Available data



Quark unpol. TMD: extractions

		Framework	HERMES	COMPASS	DY	Z production	N of points
PAST	KN 2006 <u>hep-ph/0506225</u>	NLL/NLO	×	×	~	~	98
	Pavia 2013 <u>arXiv:1309.3507</u>	No evo	~	×	×	×	1538
	Torino 2014 <u>arXiv:1312.6261</u>	No evo	<pre>(separately)</pre>	(separately)	×	×	576 (H) 6284 (C)
	DEMS 2014 <u>arXiv:1407.3311</u>	NNLL/ NLO	×	×	~	~	223
	EIKV 2014 <u>arXiv:1401.5078</u>	NLL/LO	1 (x,Q²) bin	1 (x,Q²) bin	~		500 (?)
PRESENT	Pavia 2016 <u>arXiv:1703.10157</u>	NLL/LO		~	~		8059
	SV 2017 <u>arXiv:1706.01473</u>	NNLL/ NNLO	×	×	•		309

First global fit of TMDs



It's the dawn of TMD global fits era

... but there's still a lot of climbing to be done

Mean transverse momentum squared





CAVEAT: intrinsic transverse momentum depends on TMD evolution "scheme" and its parameters. Not the best quantity to consider.

Functional form of TMDs

We go beyond a simple Gaussian.

In scalar diquark model, the two components would be $f_1 + g_1$ and $f_1 - g_1$

Scimemi and Vladimirov tried

$$f_{NP}(b) = \frac{\cosh\left(\left(\frac{\lambda_2}{\lambda_1} - \frac{\lambda_1}{2}\right)b\right)}{\cosh\left(\left(\frac{\lambda_2}{\lambda_1} + \frac{\lambda_1}{2}\right)b\right)} \qquad \qquad f_{NP}(z,b) = \exp\left(\frac{-\lambda_2 z b^2}{\sqrt{1 + z^2 b^2 \frac{\lambda_2^2}{\lambda_1^2}}}\right)$$

Transverse size in momentum space

Bacchetta, Delcarro, Pisano, Radici, Signori, arXiv:1703.10157



The fact that it goes to zero at x=1 is built in, but the sharp decrease is coming from data. However, it could still be an artefact of x and Q² correlations

Problems and open questions

Improvement in perturbative accuracy



talk by A. Vladimirov at EIC UG meeting 2017

Problems with normalisation



Problems with matching in SIDIS

N. Sato's talk at SPIN 2018 and arXiv:1808.04396



At high q_T , the collinear formalism should be valid, but large discrepancies are observed

Problems with matching in SIDIS

N. Sato's talk at SPIN 2018 and arXiv:1808.04396



The discrepancies could be largely resolved by sharply modifying the gluon collinear fragmentation function

Problems with matching in DY

E866/NuSea F. Piacenza's talk at SPIN 2018, article in preparation $pp \rightarrow \mu^+\mu^- X^- \sqrt{s} = 38.8 \ GeV$



Cannot be fixed by changing fragmentation functions... Higher twist? QED radiation?

Problems with flavour structure

Signori, Bacchetta, Radici, Schnell JHEP 1311 (13)



The bottom message is: there is room for flavour dependence, but we don't control it well



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TMDs impact on precision observables





All analyses assume that TMDs are not flavour dependent. What happens if they are?

$$m_W = 80370 \pm 7 \text{ (stat.)} \pm 11 \text{ (exp. syst.)} \pm 14 \text{ (mod. syst.)} \text{ MeV}$$

 $= 80370 \pm 19$ MeV,

 $m_{W^+} - m_{W^-} = -29 \pm 28$ MeV.

Tabs. I and II for 5 representative sets out of the 30 "Z-equivalent" sets. The former table lists the values of The APP state in Eq. (2) for each of the 5 considered **DSERVables** shows the corresponding shifts induced in M_W when applying our analysis to the m_T^{-} , p_T^{-} distributions for the former, arXiv: 1807.02101 We tried some judicious choices of flavour dependent widths and checked

-		_	1	1			1		1			
Set	u_v	d_v	Setts	; U	$l_v d$	s a	v^{v}	a	s	d_s	s	
1	0.34	0.26	Q .4	6	3 45	60 .	26	302.	46	0at90	DW3Ø	edium, large
2	0.34	0.46	Q .5	6	B 4:	30.	466.	al'	56	Darr	5 0v5la	rge, narrow
3	0.55	0.34	6 .3	B	555	5 0 .	394.	300 ;	33	brige	,0n 30 1	row, large
4	0.53	0.49	₿.3	Ø	<u>53</u> 2	20.	49 .	D .	87	large	,0 , 52	lium, narrow
5	0.42	0.38	9 .2	9	4 25	59 .	38	207.2	29	hēđi	0^{27}	arrow. large

TABLE I: Values of the Δg_{W} parameter in Eq. (2) for the flavors $a = u_v, d_v, w_{T} = p_{T\ell} = p_{T\ell} = p_{T\ell} = 0$ for the G_{ℓ} of taking into account the flavors $a = u_v, d_v, w_{T} = p_{T\ell} = p_{T\ell} = p_{T\ell} = 0$ for the G_{ℓ} of taking into account the flavors of $u_{T} = u_v + d_v$. flavour dependence of -2 3 -1 1 0 As expected, the shifts induced by the analysis second lead to errors in the determination of the W -2 3 -1 -4 9 -2 0 0 4 -4 mass 5 -1 -3 $\mathbf{0}$ 4

¹ Our analysis is performed on 30 bins in the interval [60, 90] GeV for m_{π} and on 20 bins in the interval [30, 50] CeV for m_{π} .

Future experimental data



JLab data

Patrizia's talk and M. Mirazita's talk at SPIN 2018



(But remember that we want them in multidimensional binning)

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Evolution of TMD fragmentation funct.

Bacchetta, Echevarria, Signori, Radici, arXiv:1508.00402



Measurements of TMD multiplicities will be really useful!

Future experimental data



To better test the formalism, we would need more data covering the same x range and spanning over a large range in Q^2 .

In five to ten years

- TMD multiplicities for pions and kaons, off protons and deuterons, from COMPASS and JLab
- Drell-Yan and Z measurements from CERN, RHIC, FermiLab (COMPASS with pions)
- TMD multiplicities for pions and kaons in e^+e^- from BELLE and BES
- Better understanding and control of higher-order QCD corrections
- More flexible functional forms, flavour dependence, at least two or three alternative extractions
- Use TMDs for something else (W mass... comparison with lattice... Wigner distributions...)
- READY TO USE EIC DATA